## CHAPTER 5. ELECTRON AND POSITRON SOURCES

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## 5. ELECTRON AND POSITRON SOURCES

### 5.1 THE ELECTRON SOURCE

The electron source for the collider is required to provide two intense single bunches for each cycle of collider operation. One of these bunches is used to create the positrons for the next cycle of collider operation, while one provides the electrons to interact with the existing positrons in the same cycle of collider operation. Each of these bunches must result in $5 \times 10^{10}$ electrons being captured and damped in the damping ring. Normally, longitudinally polarized electrons will be delivered by a photoemission electron gun. This polarization is transformed into transverse polarization (parallel to the damping ring magnetic field) in the linac-to-ring beamline, and back to longitudinal polarization for reinjection into the linac in the ring-to-linac beamline. Alternatively, unpolarized electrons may be delivered by a thermionic gun.

### 5.1.1 The Laser Gun for Polarized Electron Beams

This gun utilizes photoemission from a semiconductor cathode to produce a longitudinally polarized electron beam. The short optical pulses which cause the photoemission are chopped by a Pockels cell optical pulse chopper from a much longer optical pulse provided by a flashlamp pumped dye laser. The laser operates at a wavelength corresponding to the energy of the minimum direct bandgap of the semiconductor. The chopped optical pulses are circularly polarized by a pulsed Pockels cell optical retarder, and this circular polarization, responsible for the longitudinal polarization of the electrons, may be reversed randomly on a pulse-to-pulse basis, providing considerable reduction in potential polarization dependent systematic effects.

The semiconductor cathode is activated by exposure to monolayer quantities of cesium and oxygen. This activated surface requires a good UHV environment for long lifetime, typically $10^{-10}$ torr or better.

SLAC has considerable experience operating polarized electron guns of this type on the linac. ${ }^{1}$ An electron gun suitable for use on the collider has been built
and is operational in a laboratory setup. Considerable developmental work has been done to simplify the operational aspects of this gun. Current photocathodes of GaAs provide a maximum polarization of only $50 \%$. A program is underway to develop alternate photocathodes to deliver much higher beam polarizations, with a goal of $\geq 90 \%$. We anticipate that such photoemitters could be incorporated into the existing gun structure without significant change.

Figure 5.1.1.1
Overall Schematic of Photocathode Gun Showing Side View and Downbeam View


| Table 5.1.1.1 Laser Gun Specifications |  |
| :--- | :--- |
| Electrons (charge) per pulse | $1 \times 10^{11}(16 \mathrm{nC})$ |
| Repetition Rate | 180 Hz |
| Pulse width | 2.5 ns FWHM |
| Cathode voltage | 200 kV |
| Anode voltage | Ground |
| Gun insulation | Low pressure SF $\mathbf{6}_{6}$ |
| Cathode area (diameter) | $1.8 \mathrm{~cm}^{2}(1.5 \mathrm{~cm})$ |
| Electron pulse time stability <br> w.r.t. RF zero crossing | $\sigma=10$ ps |
| Cathode material | Cesium activated GaAs |
| Cathode current | Space charge limited to 15 A peak |
| Pulse-to-pulse <br> amplitude stability | $\sigma \leq 0.5 \%$ |
| Electron polarization | $50 \%$ |
| Calculated invariant <br> emittance area | $1 \times 10^{-3} \pi m_{0} c \mathrm{~cm}$ |

### 5.1.2 The Thermionic Gun ${ }^{2}$

While thermionic emission does not provide the polarized electrons, guns of this type are far less complex than photoemission guns, and thus could provide a high reliability backup gun for the polarized guns. A high current thermionic gun has been constructed, based upon a dispenser cathode-grid assembly developed by EIMAC for use in high power planar triodes. This gun has delivered 20 ampere pulses at 175 kV during tests with a 20 nsec pulser. Using an avalanche transistor pulser, the gun can deliver 8 amp pulses with a FWHM of 2 nsec . The calculated gun emittance is $3 \times 10^{-3} \pi m_{0} c-c m$. This gun has been in routine single pulse operation at 150 kV on the collider injector development project for over 18 months without any major difficulty. Recently the pulser was modified ${ }^{3}$ to provide two single bunches per cycle.

Figure 5.1.2.1 Side View of Thermionic Gun


| Table 5.1.2.1 Thermionic Gun Specifications |  |
| :--- | :--- |
| Type | Triode with intercepting mesh <br> control grid |
| Electrons per pulse (charge) | $1 \times 10^{11}(16 \mathrm{nC})$ |
| Repetition rate | 360 Hz |
| Pulse width | 2.5 ns FWHM |
| Cathode voltage | 175 kV |
| Anode voltage | Ground |
| Gun insulation | Air |
| Cathode area (diameter) | $2 \mathrm{~cm}{ }^{2}(1.6 \mathrm{~cm})$ |
| Cathode-to-grid-spacing | $150 \mu \mathrm{~m}$ |
| Grid pulse height | 500 V |
| Electron pulse time stability <br> w.r.t. RF zero crossing | $\sigma=10$ ps |
| Cathode material | Dispenser type |
| Cathode current | Space charge limited to 20 A peak; <br> Maximum of 9 A peak for 2.5 ns pulse |
| Pulse-to-pulse <br> amplitude stability | $\sigma \leq 0.5 \%$ |
| Calculated invariant <br> Emittance area | $3 \times 10^{-3} \pi m 0 c \mathrm{~cm}$ |

### 5.1.3 The Buncher ${ }^{4}$

While it appears likely that photoemission cathodes could directly generate an electron pulse containing $5 \times 10^{10}$ electrons lasting only about 100 psec , computer simulations indicate that the transient and space-charge phenomena associated with such a bunch are so severe that it might not be possible to bunch and capture $100 \%$ of this charge into a single S-band bunch in the linac. Since these problems decrease significantly as the electron bunch becomes longer, we have decided to employ a subharmonic buncher consisting of two cavities located between the electron gun and the S-band linac. In the present scheme, initial bunching is accomplished at the 16 th subharmonic ( 178.5 MHz ) of the linac RF frequency. This allows the electron pulse from the gun to be several nsec long, thus eliminating most $o^{\prime}$ ' he ill effects associated with the higher charge density of the 100 psec bun h. The subharmonic buncher will deliver an electron pulse of approximately 150 psec length at 150 kV to a standard SLAC injector section. Subharmonic bunching with 178.5 MHz cavities has been used at SLAC to achieve single $S$-band bunches containing more than $1 \times 10^{11}$ electrons.

Figure 5.1.3.1
Schematic of Existing 178.5 MHz Subharmonic Buncher'System


### 5.1.4 The Injector ${ }^{5}$

The bunch from the subharmonic bunching section enters a standard SLAC injector section. This section is composed of a standard 3-meter accelerator section with an attached $10-\mathrm{cm}$ traveling-wave buncher. At the output of this section, the beam has a $30-\mathrm{psec}$ bunch length and an energy of about 50 MeV . The beam then enters an achromatic bending system which can be used to alter the longitudinal bunch shape. Further acceleration will then attain the energy spectrum necessary for damping ring injection. Solenoidal focusing is used along the entire length of the subharmonic buncher and the injector section. The cathode itself will be in a magnetic-field-free region in order to avoid emittance degradation. For injection into Section 1-2 along with the positrons, the electron beam is accelerated to about 200 MeV by four conventional accelerator sections interspersed with instrumentation. A quadrupole FODO array with a $\beta$ of four m beginning at the injector accelerator section is used to focus this beam.

| Assumptions concerning gun beam |  |
| :---: | :---: |
|  |  |
| Peak current | 15 A |
| Pulse shape | Gaussian |
| Pulse width | $\sigma=1.0 \mathrm{~ns}$ |
| Energy | 200 keV |
| $\Delta E / E$ | $\pm 0.2 \%$ or less |
| Emittance area | $5 \times 10^{-3} \pi m_{0} \mathrm{c}-\mathrm{cm}$ or less |
| Cathode area | $1.8 \mathrm{~cm}^{2}$ |
| Repetition rate | up to 180 pps |
| Gun lenses |  |
| Focal length | 11.5 cm |
| Peak field | 586 gauss |
| Sub-harmonic buncher |  |
| Frequency | 178.5 MHz |
| Type | re-entrant cylindrical cavity |
| Gap length | 3.8 cm |
| Bore | 5 cm |
| $Q_{L}$ | $4 \times 10^{3}$ |
| Shunt impedance | Rs $-V^{2} / 2 P=2.2 M \Omega$ |
| Pulse length | $30 \mu \mathrm{~s}$ |
| Drive power (typical) | 2 kW |
| Gap voltage | 60 kV |
| Drift distances | 1 m and 40 cm |
| Solenoid field | 200 to 600 Gauss |
| Matching solenoid field | $\sim 1 \mathrm{kG}$ |
| S-band Buncher |  |
| Type | Travelling wave DLWG |
| Length | 10.5 cm |
| Field strength | $20 \mathrm{kV} / \mathrm{cm}$ |
| Buncher solenoid field | 1.2 kG |
| Accelerator Section |  |
| Gap to buncher | 2 cm |
| Length of accelerator section | $3 \mathrm{~m}$ |
| Accelerating field <br> Solenoid field | $\begin{aligned} & 150 \mathrm{keV} / \mathrm{cm} \\ & 1 \text { to } 2 \mathrm{kG} \end{aligned}$ |
| Beam injected into sector 1-2 |  |
| Current | $5 \times 10^{10}$ electrons/bunch <br> ( $8 \mathrm{nCoulomb} / \mathrm{bunch}$ ) |
| Bunch length ( $\sigma_{z}$ ) | $2 \mathrm{~mm}$ |
| Energy | 200 MeV (at BAS I) |
| Emittance area | $3 \times 10^{-2} \pi m_{0} c \mathrm{~cm}$ |

5.1.5 State of Development ${ }^{6}$

A complete single-bunch injection system composed of the thermionic gun (Fig. 5.1.5.1), two subharmonic bunchers operating at 178.5 MHz (Fig. 5.1.5.2), the standard injection section, and the instrument section, has been in operation at the front end of the linac since spring, 1981. This injection system meets the SLC specifications for intensity and emittance.

Figure 5.1.5.1 The Thermionic Gun


Figure 5.1.5.2: Subharmonic Buncher


### 5.2 THE POSITRON SOURCE

In order to optimize the collider luminosity, the positron source is required to produce sufficient numbers of positrons to saturate the intensity capability of the linac. Since the intensity of the scavenger pulse is subject to the same linac capability, the positron system is required to deliver at least one positron into the damping ring for each electron incident on the positron target. This requirement implies more than an order-of-magnitude improvement over the existing source which, at 7 GeV incident electron energy, normally has a yield of $7 \%$ but, under optimum conditions, has been tuned to $15 \%$. The needed improvement will result from (a) increasing the incident electron energy to 33 GeV , and (b) using a collection system with a larger admittance.

The overall positron production system is shown in Figs. 5.2.0.1-3. and component parameters are summarized in Table 5.2.0.1. A single bunch of electrons (the scavenger pulse) is extracted from the linac at the $2 / 3$ point by means of a fast kicker magnet and a pair of septum magnets. The extracted electron beam is transported through a $24^{\circ}$ bend and into the positron vault where it is focused to the desired beam spot size at the target. Positrons are collected and accelerated up to 200 MeV in a $35-\mathrm{ft}$ length of linac accelerator section. They are then transported back to the beginning of the linac for further acceleration to 1.2 GeV , whereupon they are sent to the damping ring. Achieving the desired number of positrons depends on the efficiency of the collection and transport systems. The optimization of the yield is constrained by the minimum allowable beam size as determined by the target strength.

Table 5.2.0.1 Positron Source Specifications

| EXTRACTION |  |
| :--- | :--- |
| Electron Scavenger Pulse | 33 GeV |
| Energy | $5.0 \times 10^{10} \mathrm{e}^{-} / \mathrm{pulse}$ |
| Intensity | 0.6 mm |
| Size $(1 \sigma)$ | $264 \mathrm{Joules} / \mathrm{pulse}$ |
| Pulse energy | 180 Hz |
| Pulse rate | 47 kW |
| Power |  |
| Target | $90 \% \mathrm{Ta}-10 \% \mathrm{~W}$ |
| Material | 6 radiation lengths $=24 \mathrm{~mm}$ |
| Length | $53 \mathrm{~J} / \mathrm{pulse}$ |
| Energy deposited in target | $380^{\circ} \mathrm{C}$ |
| Pulse temperature rise | $580^{\circ} \mathrm{C}$ |
| Max. pulse temp. | $32,000 \mathrm{psi}$ |
| Max. compressive stress | 9 kW |
| Power deposition | $200^{\circ} \mathrm{C}$ |
| Steady-state temp. |  |
| Positron Beam at Target | $2-20 \mathrm{MeV}$ |
| Energy range | $2 \mathrm{~mm} \times 2.5 \mathrm{MeV} / \mathrm{c}=0.01 \mathrm{~m}-\mathrm{radians}$ |
| Transverse emittance (Invariant) | 2.5 |
| Yield (e $\mathrm{e}^{+} / \mathrm{e}^{- \text {in. }}$ ) | 1.21 GeV |
| Beam Properties at End of Sector 1 | $2 \% \mathrm{full}$ |
| Energy | $4.2 \times 10^{-6} \mathrm{~m}-\mathrm{radians}$ |
| Energy spread |  |
| Transverse emittance |  |

Figure 5.2.0.1



### 5.2.1. Yield Calculations

It is important to determine where gains in positron yield can be made and what the limits are due to the nature of the electromagnetic cascade. Yields have been calculated for various beam energies and target materials using the EGS computer program, ${ }^{7}$ and some of the results are shown in Fig. 5.2.1.1. The following conclusions have been drawn from these studies: (a) Materials with large atomic number are more efficient. For example, tungsten gives twice the yield of copper under the same conditions at their shower maxima. (b) Denser materials localize the shower to a smaller region. This results in a smaller emittance and hence better collection efficiency, but it also produces higher physical stress and pulse temperature rise. (c) Positrons are produced preferentially in the $2-20 \mathrm{MeV}$ range. (d) In dense materials, the positrons come from a small spot, typically 0.5 mm radius, but with large angles, typically $\pm 20^{\circ}$. (e) The yield is nearly proportional to the incident electron energy. (f) The number of positrons produced is more than sufficient, provided they can be captured and accelerated as desired. Using the parameters for the collection optics of the existing SLAC source, we calculate a yield of $15 \%$, in good agreement with the measured yield. The intended factor of 3 increase in admittance in each transverse plane of the collecting optics, and the higher incident electron energy, result in a total calculated yield of $2.5 \mathrm{e}^{+} /$incident $\mathrm{e}^{-}$.

Figure 5.2.1.1
Positron yield vs depth. The longitudinal distribution of positrons in a tungsten target per incident 33 GeV electron is shown for maximum positron energies of $5,10,20,50$ and 100 MeV .


Figure 5.2.1.2
A schematic representation of positron source


### 5.2.2 Target

Beam tests indicated that a target composed of either W-26Re (tungsten $26 \%$ rhenium) or Ta-10W (tantalum - $10 \%$ tungsten) material would be suitable for our application. Both materials were analyzed on an engineering basis using their basic material properties and how these would affect pulse temperature rise, heat transfer, and thermal stress. The stress was compared to the material's tensile strength and fatigue strength. A slight advantage for Ta-10W over W26Re was found. In as much as Ta-10W is also more readily available, less expensive and easier to work with, it is the obvious material of choice. Ta-10W material properties are shown in Table 5.2.2.1.

Calculations indicate that the single pulse temperature rise will be $380^{\circ} \mathrm{C}$ for a 0.6 mm incident beam radius ( 1.3 mm at shower peak). This, when added to the expected steady state temperature, results in a peak temperature of $580^{\circ} \mathrm{C}$. The resulting stress from each thermal pulse will be $\approx 32,000$ psi which is well below the expected yield strength of 60,000 psi. Fatigue characteristics for pure tantalum indicate an infinite number of cycles at 34000 psi . To avoid interaction between adjacent pulses and the resulting compounding of stress, the target will be rotated, thus distributing the power deposited over a larger area. Figure 5.2.2.2 shows the prototype target assembly. The target wheel, inside a vacuum chamber, is spun at about 2 Hz to distribute the power. The rotating motion is attained utilizing a commercially available feedthrough which uses a ferromagnetic fluid to seal the drive shaft. The target is cooled by water passages in the drive shaft.

| Table 5.2.2.1 |  |
| :--- | :---: |
| TARGET MATERIAL PROPERTIES |  |
| TANTALUM - 10\% TUNGSTEN |  |
|  |  |
| Atomic Number | $73 / 74$ |
| Density, gm/cm ${ }^{3}$ | 16.9 |
| Specific Heat, cal/gm ${ }^{\circ} \mathrm{C}$ | 0.035 |
| Radiation Length, $\mathrm{gm}^{\prime} / \mathrm{cm}^{2}$ | 6.814 |
| Thermal Conductivity cal/s-cm- ${ }^{\circ} \mathrm{C}$ | 0.185 |
| Tensile Strength, $\mathrm{lb} / \mathrm{in}^{2}$ | 71000 |
| Melt Temperature, ${ }^{\circ} \mathrm{C}$ | 3035 |
|  |  |

Figure 5.2.2.1
Positron target energy deposition. The radial distribution of energy deposition from an electromagnetic shower cascade is shown at 6 radiation lengths in tungsten for 33 GeV incident electrons. The left-hand scale is the energy density normalized to the incident beam energy. The right-hand scale shows the pulse temperature rise for $3.7 \times 10^{10}$ incident electrons, corresponding to 194 Joules of beam energy. The histogram is the Monte Carlo calculation for a point incident beam. The curves are for Gaussian incident beam with $\sigma=0.4$ and 0.8 mm .



Figure 5.2.2.2

### 5.2.3 Transverse and Longitudinal Acceptance

The positron beam emerging from the target will enter a focusing solenoid system. The goal is to transform a $2 \mathrm{~mm} \times 2.5 \mathrm{MeV} \sigma_{x} \times \sigma_{p_{x}}$ transverse emittance into a $8 \mathrm{~mm} \times 0.6 \mathrm{MeV}$ emittance, which corresponds to the accelerator aperture with a superimposed 5 kG solenoidal field. This is accomplished with a pseudoadiabatically changing solenoidal field where the initial-to-final field ratio is the square of the radii ratio. ${ }^{8}$ Thus a 80 kG initial field is desired. We believe fields approaching this can be achieved with a pulsed magnetic field utilizing a flux concentrator. ${ }^{9}$ The flux concentrator can conveniently be made to produce the desired field taper; it produces field only where needed for the particle orbits, and provides to some extent, a correctly shaped collimator to stop unwanted shower particles.

The positrons emerging from the target have a longitudinal spatial extent which is nearly identical to the extent of the incident electron bunch, which is roughly 2 mm long. The emergent positrons have a large spread in energy from roughly 2 to 20 MeV with most concentrated at low energies. At the lower energies the positrons are sufficiently non-relativistic that they slip in phase, so that after the positrons have been accelerated in the booster to about 200 MeV there is a distinct energy-phase angle correlation, with the positrons that emerged from the target with low energies lagging in phase. Additional phase slip occurs for positrons that leave the target with non-zero transverse momentum because the resulting helical motion has a longer path length than straight motion down the axis.

Two design considerations for the positron source have been incorporated to minimize the phase slip. The first is to minimize drift regions before acceleration and the second is to accelerate initially with a high field accelerator section. A 1.5 m accelerator section with acceleration at $50 \mathrm{MeV} / \mathrm{m}$ will be used.

The positron capture efficiency has been calculated by ray tracing positrons through a target solenoidal field of 10 kG , the pulsed solenoidal field "flux concentrator" magnet, and into a solenoid of 5 kG axial field enclosing the RF accelerating cavity. Positrons were taken from the output of EGS for an inci-
dent 33 GeV electron beam having a nominal beam radius of 0.6 mm standard deviation both vertically and horizontally. Figure 5.2.3.1(a) shows the momentum spectrum ( $1 / \mathrm{N}_{e_{-}}$) $\mathrm{d} \mathrm{N}_{e_{+}} / \mathrm{dp}$ of positrons produced where $\mathrm{N}_{e_{--}}$is the number of electrons incident on the target, and $\mathrm{N}_{e^{+}}$is the number of positrons produced. Because of bandwidth limitation in the later transport system, only positrons below $20 \mathrm{MeV} / \mathrm{c}$ momentum are considered; the integral of this plot is $31 \mathrm{e}^{+}$ per incident $\mathrm{e}^{-}$. Figure $5.2 .3 .1(\mathrm{~b})$ shows the momentum spectrum of positrons accepted by the geometrical aperture of 18 mm diameter; the integral of this spectrum is $5.5 \mathrm{e}^{+} / \mathrm{e}^{-}$. Figure 5.2 .3 .2 shows the distribution $\left(1 / N_{e^{-}}\right) \mathrm{dN}_{e^{+}} / \mathrm{dp}_{t}$ where $p_{t}$ is the momentum transverse to the nominal beam direction at the exit of the target.The curves show the yield for all positrons and positrons accepted geometrically. The primary loss at high $p_{t}$ is due to the limited transverse acceptance.

The effect of the phase slip is illustrated in Figure 5.2.3.3 where the geometrically accepted positrons have their final momentum plotted against the phase lag $\Delta t$, which is the difference between the time a positron reaches the end of the accelerator section and that of a relativistic positron with straight motion down the axis. The momentum acceptance of the damping ring ( $2 \%$ ) will limit $\Delta t$ to less than 22 ps . Wake fields or beam loading may further reduce $\Delta t$. For the present study, we have assumed a total width of $\Delta t=15 \mathrm{ps}$ which corresponds to the positrons in the first three columns of Fig. 5.2.3.3. This additional requirement reduces the positron efficiency. The in-phase positrons within $\Delta t=15$ ps are shown in Figure 5.2.3.1 and 5.2.3.2. The integral of this system is $2.5 \mathrm{e}^{+}$per incident $\mathrm{e}^{-}$.

Many studies ${ }^{10}$ have been made varying magnetic field, drift spaces, and accelerating fields to other values which design and operating criteria may require. For these studies the acceptance for in-phase positrons varies between 2 and 3.4 positrons per incident electron. These studies do not take into account effects due to space charge and wake fields. Longitudinal wake fields in the 1.5 m , section have been estimated ${ }^{11}$ to cause 2 MeV lower energy of the tail of the positron bunch than at the head. Transverse wakes are on the order of 0.1 $\mathrm{MeV} / \mathrm{c}$ per mm offset of the bunch centroid.

This scheme does not use deceleration followed by acceleration which is known to lead to a more compact $E_{f}-\psi_{f}$ distribution in phase space and to the capture of lower $E_{i}$ particles. ${ }^{11}$ In the deceleration mode many of the particles spend some time in the booster accelerator at low energies where they are more susceptible to the adverse effects of space charge and wake field forces. More detailed calculations are needed to identify an optimum capture strategy. Eventually the optimum operation will be based on experimental tests. We will try to make the hardware compatible with operation in higher acceleration or deceleration mode.

Figure 5.2.3.1


Figure 5.2.3.2


Figure 5.2.3.3


### 5.2.4.1 Design Parameters for the Capture Accelerator Section:

Choice of Structure and Iris Diameter: A preliminary study was made to compare the predicted performances of uniformly periodic Constant Gradient and Constant Impedance Disk Loaded Waveguide Structures, operated Traveling-Wave with SLED, and Standing-Wave and Recirculating (no SLED). The results of this study, reported in CN-282 dated 17 Oct' 84, showed that, for a given iris aperture diameter, the Constant Impedance (CI) Disk Loaded Waveguide (DLWG) structure would give the highest energy gain, and would also be simpler to build and operate. The expected energy gain increased rapidly as the iris aperture diameter decreased. However, separate calculations showed that the beam-induced wake-field effects increased even more rapidly with decreasing aperture diameter. As a compromise, it was decided to set the iris aperture to 23 mm diameter.

Choice of Length: The length of DLWG was initially taken to be 1 m . However, increasing the iris diameter to 23 mm lowered the expected energy gain to 60 MeV for a 1 m section. After looking at the tradeoffs involved in using a longer section (especially the decrease in input cavity field at beam time, due to the increased time-decay of the incoming SLED pulse), it was decided to increase the length to 1.5 m .

## Parameters:

| Length | 1.5 | m |
| :--- | :---: | :--- |
| Iris Aperture diameter | 23.0 | mm |
| Number of Cavities (including couplers) | 42 |  |
| Normalized group velocity, $v_{g} / c$ | 0.0124 |  |
| Shunt Impedance | 57.0 | $\mathrm{M} \Omega / \mathrm{m}$ |
| $Q$ | 13600 |  |
| Attenuation per unit length | 0.177 | nepers |
| Total attenuation | 0.266 | nepers |
| Filling Time | 403 | ns |
| Energy (assuming 50 MW klystron, SLEDed |  |  |
| in with 0.97 db waveguide feed loss) | 82.8 | MeV |
| Peak SLED Power into First Cavity | 256 | MW |
| Peak Axial Field in First Cavity | 71.9 | $\mathrm{MV} / \mathrm{m}$ |
| Axial Field in First Cavity at Beam Time | 54.3 | $\mathrm{MV} / \mathrm{m}$ |
| Axial Field in Last Cavity at Beam Time | 55.1 | $\mathrm{MV} / \mathrm{m}$ |
| Average Axial Field at Beam Time | 55.2 | $\mathrm{MV} / \mathrm{m}$ |

Description of RF Feed to Capture Section: The layout of the RF System feeding the SLC Positron Source is shown in Fig. 5.2.4.1. Power from a SLEDed 50 MW klystron installed at Station 20-3C travels through approximately 131 ft of waveguide to the 1.5 m section. The attenuation is about 0.97 db . Some system parameters follow. Average power values assume 120 pps operation.
Klystron average RF power output ..... 30.0 ..... kW
Power dissipation in SLED cavities ..... 9.1 ..... kW
SLED Power Gain Factor ..... 6.4
SLED Accelerator Voltage Gain Factor
(for 1.5 m section with 403 ns filling time) ..... 2.21
Power into rectangular waveguide ..... 20.9 ..... kW
Power into 1.5 m DLWG section ..... 16.7 ..... kW
Power dissipation (RF) in 1.5 m section ..... 7.9 ..... kW
Power dissipation in first cavity of section ..... 210 ..... W
Power dissipation in last cavity of section ..... 125 ..... W
Power dissipation in each of two output loads ..... 4.9 ..... kW
5.2.4.2 Note on Radiation Power Dissipation in 1.5 m Section: EGS computations have shown that radiation generated in the positron target will deposit 12.6 kW of power in the 1.5 m section. The radial temperature drop across the first disk in the section (which can be water-cooled only at its outer periphery, and absorbs a total of 0.51 kW from radiation and $R F$ losses) will be about $40^{\circ} \mathrm{F}$. Preliminary calculations ${ }^{3}$ indicate that the radial and hoop stresses will be well below the yield values for annealed copper, and disk 'dishing' will not be a problem. Detuning due to an increase in the iris aperture diameter will be offset to first order by the effect of an increase in the disk wall thickness. More detailed calculations are in progress.

With a total flow of 27.4 gpm in 12 parallel circuits, the output water temperature will be $4.9^{\circ} \mathrm{F}$ above the input. A feedback system will set the inlet water temperature so that the average metal temperature is held at $113^{\circ} \mathrm{F}$.
5.2.4.3 The Constant Gradient Sections: The 1.5 m capture section will be followed by three standard 3.05 m constant gradient sections, mounted on a standard girder. The first section will be about 0.5 m downbeam of the capture section. The space normally occupied by the second section on a standard girder
will be used for instrumentation. The arrangement, including the waveguide feed from a second 50 MW SLEDed klystron at station $20-4 \mathrm{~A}$ is also shown in Fig. 5.2.4.1. Half the klystron power goes to the first section, and one fourth the power goes to each of the second and third sections. The length of rectangular waveguide run to each section is about 200 ft , introducing an attenuation of about 1.5 db . Parameters are as follows:
Power dissipation in SLED cavities ..... 7.2 kW
SLED Power Gain Factor ..... 6.2
SLED Accelerating Voltage Gain Factor
(for 3.05 m section with 830 ns filling time) ..... 1.78
Power dissipation in rectangular waveguide ..... 6.7 kW
Power dissipation in first CG DLWG section ..... 5.5 kW
Power dissipation in each of 2 nd and 3 rd
CG DLWG sections ..... 2.7 kW
Power dissipation in each load for these sections ..... 1.3 kW
Peak SLED power into first section ..... 110Energy gain in first section81 MeV
Energy gain in each of 2 nd and 3 rd sections ..... 57 ..... MeV

Figure 5.2.4.1:
Waveguide Layout for SLC Positron Source


### 5.2.5 Transport Systems

The beam transport systems consist of an extraction line to bring the 33$\mathrm{GeV} \mathrm{e}^{-}$scavenger pulse from the linac to the target, the east turn which takes the output of the booster accelerator to the return line in the linac housing, the return line, the west turn and reinjector, and the Sector 1 focusing array. Each of these systems must be sufficiently isochronous so to not spread the Sband bunch unduly. In addition, the total path length of the system from the scavenger extraction kicker in the $\mathrm{e}^{-}$damping ring to the injection kicker in the $e^{+}$ring must be a multiple of the damping ring circumference in order to place the positrons in the correct bucket of the $\mathrm{e}^{+}$ring. The east turn will be built as a movable "trombone" to allow small adjustments ( $\pm 8 \mathrm{~cm}$ ) to be made.
5.2.5.1 Extraction Line: The extracted beam is first deflected down by a kicker at Sector 19, girder 4d. At girder 7 of Sector 19, a horizontal and vertical pair of Lambertson septum dipoles deflect the beam down and to the north. The main portion of the $24^{\circ}$ bend is provided by a FOD0 array of 9 quadrupoles and 8 bend magnets. In order to match the desired elevation and slope of the $\mathrm{e}^{+}$ systems, the plane of the bending arc is rolled by about $4.4^{\circ}$ and an additional vertical bend magnet just ahead of the arc provides the correct initial pitch angle. Two additional quadrupoles midway between the Lambertson septa and the vertical bend aid in matching the beam optics.


Fiqure 5.2.5.1.2


```
Figure 5.2.5.1.4 TURTLE plots for the Extraction Line. The initial particle distribution was generated by K. Bane( ) for a case in which blowup of the beam by transverse wakefields was suppressel by Landau damping resulting from the introduction of energy spread. Energy \(=33 \mathrm{GeV}\). No aberrations in magnets.
```

Figure 5.2.5.1.4.1 Horizontal phase plot at beginning of Extraction Line.

TWO DIMENSIONAL PLOT OF $X$ ' (mr) VS $X(m m)$


TOTAL NUMBER OF ENTRIES $=1000$ INCLUDING UNDERFLOW AND OVERFLOW

|  | UNDERFLOW | OVERFLOW |
| :--- | ---: | ---: |
| ACROSS | 0 | 0 |
| DOWN | 0 | 0 |

SUM OF SQUARES $=93070$.
CENTER $=\quad-0.017 \quad$ RMS HALF WIDTH $=0.146$
CENTER $=0.000 \quad$ RMS HALF WIDTH $=0.002$
CORRELATION $=0.0268$

NO 1 TWO DIMENSIONAL PLOT OF

| $X$ IN MM | 0.0 | $M$ | FROM THE TARGET |
| ---: | :--- | :--- | :--- |
| $X$ IN MR | 0.0 | $M$ | FROM THE TARGET |

Figure 5.2.5.1.4.2 Vertical phase plot at beginning of Extraction Line.


Figure 5.2.5.1.4.3 Longitudinal phase plot at beginning of Extraction Line.


Figure 5.2.5.1.4.4 Horizontal phase plot at Positron Target.
$\therefore$ TWO DIMENSIONAL PLOT OF $X$ ( mr ) VS $\mathrm{X}(\mathrm{mm})$

|  |  |  | -0.500 | 0.000 | 0.500 totals |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
|  |  |  | I** | **---**---** | *---** |  |  |
| -2.000 |  | -1.800 | I |  |  | I | 0 |
| -1.800 |  | -1.600 | I |  |  | I | 0 |
| -1.600 |  | -1.400 | I |  |  | I | 0 |
| -1.400 |  | -1.200 | I |  |  | I | 0 |
| -1.200 | T0 | -1.000 | I | 1 |  | I | 1 |
| -1.000 | то | -0.800 | I | 1 |  | I | 1 |
| -0.800 | T0 | -0.600 | I |  |  | I | 0 |
| -0.600 | T0 | -0.400 | I | 1 |  | I | 1 |
| -0.400 | T0 | -0.200 | I | 11 |  | I | 2 |
| -0.200 | T0 | 0.000 | I | 16B\$\$3761 |  | I | 151 |
| 0.000 | то | 0.200 | I | 58V\$\$ZA | 1 | I | 200 |
| 0.200 | T0 | 0.400 | I | 3C\$\$\$G51 |  | I | 157 |
| 0.400 | T0 | 0.600 | I | $1 \mathrm{CNTA73}$ |  | I | 85 |
| 0.600 | T0 | 0.800 | I | 6KY\$A61 |  | I | 115 |
| 0.800 | T0 | 1.000 | I | 12BTW0531 |  | I | 108 |
| 1.000 | T0 | 1.200 | I | 2S\$\$\$8 |  | I | 170 |
| 1.200 | то | 1.400 | I | 126 |  | I | 9 |
| 1.400 |  | 1.600 | I |  |  | I | 0 |
| 1.600 |  | 1.800 | I |  |  | I | 0 |
| 1.800 | T0 | 2.000 | I |  |  | I | 0 |
|  |  |  | I**---**---**---**---** |  |  |  |  |
|  |  |  | I |  |  | I |  |
|  |  |  | I |  |  | I |  |
|  |  |  | I | 1221 |  | I |  |
|  |  |  | I | 6368572 |  | I |  |
|  |  | totals | I 00001511144462010000 |  |  | I | 1000 |

TOTAL NUMBER OF ENTRIES = $\quad 1000$ INCLUDING UNDERFLOW AND OVERFLOW

|  | UNDERFLOW | OVERFLOW |
| :--- | ---: | ---: |
| ACROSS | 0 | 0 |
| DOWN | 0 | 0 |

SUM OF SQUARES $=33414$.
CENTER $=-0.041$ RMS HALF WIDTH $=0.071$

| CENTER $=$ |
| :--- |
| CORRELATION $=$ |
| $=$ |$\quad-0.364$ RMS HALF WIDTH $=0.433$

NO 7 TWO DIMENSIONAL PLOT OF

| X IN MM | 68.625 M | FROM THE TARGET |
| ---: | :--- | :--- |
| X ' IN MR | 68.625 M | FROM THE TARGET |

Figure 5.2.5.1.4.5 Vertical phase plot at Positron Target.


Figure 5.2.5.1.4.6 Longitudinal phase plot at Positron Target.

TWO DIMENSIONAL PLOT OF DP/P (\%) VS DL (mm)


TOTAL NUMBER OF ENTRIES $=1000$ INCLUDING UNDERFLOW AND OVERFLOW
UNDERFLOW OVERFLOW

| ACROSS | 0 | 2 |
| :--- | :--- | ---: |
| DOWN | 0 | 0 |

SUM OF SQUARES $=35872$.
CENTER $=-0.627$ RMS HALF WIDTH $=0.755$
CENTER $=0.497$ RMS HALF WIDTH $=0.457$
CORRELATION $=-0.0122$
NO 9 TWO DIMENSIONAL PLOT OF
DL IN MM 68.625 M FROM THE TARGET
DP/P IN PC $68.625 \mathrm{M} \quad$ FROM THE TARGET

Figure 5.2.5.1.5 TURTLE plots showing effects of sextupole fields induced by dipole excitation of the QD quadrupoles. The sextupole field is assumed to be induced by dipole excitation to correct for vertical misalignments of magnitude $|d y|=1 \mathrm{~mm}$. Cases (a) and (b) show the extreme effects for different distributions of the sign of $\therefore \quad$ the sextupole in the four QD's. Compare to Figure 5.2.5.1.4.5 which is the case of no sextupole. The effects on horizontal and longitudinal phase space is minimal.

Figure 5.2.5.1.5.1 Signs of the sextupoles in the QD's is +--+.

TWO DIMENSIONAL PLOT OF $\quad Y^{\prime}(\mathrm{mr})$ VS $\quad Y$ (mm)


TOTAL NUMBER OF ENTRIES $=1000$ INCLUDING UNDERFLOW• AND OVERFLOW
UNDERFLOW OVERFLOW

| ACROSS | 0 | 0 |
| :--- | :--- | :--- |
| DOWN | 0 | 0 |

SUM OF SQUARES $=61804$.

| CENTER $=$ | $0.181 \quad$ RMS HALF WIDTH $=$ | 0.211 |
| :--- | :---: | :---: | :---: |
| CENTER $=$ | $0.064 \quad$ RMS HALF WIDTH $=$ | 0.064 |
| CORRELATION $=$ | 0.9613 |  |

NO 8 TWO DIMENSIONAL PLOT OF

| Y IN MM | 68.625 M | FROM THE TARGET |
| ---: | :--- | :--- |
| $\mathrm{Y}^{\prime}$ IN MR | 68.625 M | FROM THE TARGET |

Figure 5.2.5.1.5.2 Signs of the sextupoles in the $Q D^{\prime}$ s is -++-.

5.2.5.2 Booster: The beam transport system around the booster accelerator consists of solenoid and then quadrupole focusing. A uniform 5 kG solenoidal field extends from the target through the high gradient 1.5 m section to the end of the first normal gradient 3 m section. There a transition is made to a FOD0 quadrupole array. The 3 m section following the solenoid is devoted to instrumentation and a magnetic charge separator to stop the electrons.

Figure 5.2.5.2.1

Figure 5.2.5.2.2 TURTLE plots at end of first 10-ft accelerator section in positron booster. Initial number of electrons on target $=327$. Relative yield $=2.7 \mathrm{e}+\mathrm{e}$ - .

Figure 5.2.5.2.2.1 Horizontal phase plane.


Figure 5.2.5.2.2.2 Vertical phase plane.


Figure 5.2.5.2.2.3 Longitudinal phase plane.

TWO DIMENSIONAL PLOT OF DP/P (\%) VS DL (mm)

5.2.5.3 East Turn \& $24^{\circ}$ Bend: Two $90^{\circ}$ dipoles turn the positron beam around $180^{\circ}$ in a vertical plane and back along the $24^{\circ}$ housing. In order to satisfy the isochronous condition, the integral of the dispersion must average to zero in the bend dipoles. This is accomplished by two small $10^{\circ}$ bends on each side of the main $180^{\circ}$ bend. A weak quadrupole between each of the $10^{\circ}$ bends permits adjustment of the isochronism. Four $6^{\circ}$ dipoles are used to bring the beam into the return line. This system is also made isochronous, by quadrupoles which suitably control the dispersion function. A sextupole at the symmetry point of the $180^{\circ}$ turn corrects the second order non-isochronism.

Fiqure 5.2.5.3.1



Figure 5.2.5.3.5

5.2.5.4 Return Line: The transport from Sector 19 to Sector 1 is done with a FODO array of eight quadrupoles per sector of nominal phase advance of $90^{\circ} / \mathrm{cell}$. Dipole correctors, 4 vertical and 4 horizontal per sector, are located at respective maxima in the $\beta$ function to provide steering and compensation for the deflection due to the earth's magnetic field or other stray fields.

Figure 5.2.5.4.1


5.2.5.5 West Turn and Inflector: This transport matches from the return line to a $180^{\circ}$ turn similar to the east turn but in the horizontal plane. The $\mathrm{e}^{+}$ beam is brought down to the linac by a $12^{\circ}$ bend and recombined with the $\mathrm{e}^{-}$ beam by an upward $12^{\circ}$ bend. The combined west turn and inflection system is made isochronous by slightly over compensating the path length dispersion of the west turn.


E+ RETURN SLCPRTN 3-11-85

Figure 5.2.5.5.3


Figure 5.2.5.5.4 TURTLE plois at reinjection into Sector 1. Initial number of electrons on target $=327$. Relative yield $=2.2 \mathrm{e}+/ \mathrm{e}^{-}$. Transmission through Return Line $=81 \%$.

Figure 5.2.5.5.4.1 Horizontal phase plane.

TWO DIMENSIONAL PLOT OF $X$ ( $m$ r) VS $X(m m)$


TOTAL NUMBER OF ENTRIES $=\quad 729$ INCLUDING UNDERFLOW AND OVERFLOW

|  | UNDERFLOW | OVERFLOW |
| :--- | ---: | ---: |
| ACROSS | 0 | 0 |
| DOWN | 3 | 1 |

SUM OF SQUARES = 6155.

| CENTER $=$ | $0.167 \quad$ RMS HALF WIDTH $=$ | 2.256 |
| :--- | :--- | :--- |
| CENTER $=$ | $0.049 \quad$ RMS HALF WIDTH $=$ | 1.817 |
| CORRELATION $=$ | -0.4230 |  |

NO 22 TWO DIMENSIONAL PLOT OF

| X IN MM | 2033.382 M | FROM THE TARGET |
| ---: | :--- | :--- |
| X ' IN MR | 2033.382 M | FROM THE TARGET |

Figure 5.2.5.5.4.2 Vertical phase plane.

TWO DIMENSIONAL PLOT OF $\quad Y^{\prime}(\mathrm{mr})$ VS $Y(\mathrm{~mm})$


Figure 5.2.5.5.4.3 Longitudinal phase plane.

5.2.5.6 Sector 1 Focusing: In order to constrain the large $\mathrm{e}^{+}$emittance to the available aperture and overpower transverse wake fields, an en'ergy-graded FODO array is used. Seventy-two quadrupoles of maximum strength 15 kG (gradient $\times$ length) have been placed around the linac structure. In order to minimize interference with existing structures, the spacings of quadrupoles deviate slightly from that of an ideal graded FODO array due to linac physical constraints.

Figure 5.2.5.6.1



Figure 5.2.5.6.2



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### 5.2.6 Flux Concentrator:

The flux concentrator is a pulsed magnet producing an axially tapered magnetic field with a maximum value of 6 T . The magnet immediately follows the target and acts as a short focal length lens in order to collect the highly divergent positrons.

A cross section of the magnet is shown schematically in Fig. 5.2.6.1. The body of the magnet core is a copper slug 120 mm long with a radius of 37 mm . The throat of the core tapers from a radius of 3.5 mm with a slope of 0.233 to about 31 mm . The coil consists of a total of 13 turns of water cooled copper conductor, and is wound in two halves so that both the outside ends remain at ground potential while the high voltage is applied at the center. This is done to reduce breakdown problems as well as to reduce the length of the water passage.

The magnet is powered by a resonant, pulsed supply; the resonant frequency of the charging capacitor and the magnet is of the order of 15 kHz . The supply delivers a half sine-wave at a repetition rate of 180 Hz with a peak current of 32 kA and a peak potential of 2 kV for a total input power to the magnet of about 80 kW . For test purposes a SLAC test lab klystron modulator will be modified to power the magnet. Design for a normal pulsed supply is under way.

The power dissipated in the magnet is estimated to about 50 kW of electrical power plus a significant fraction of the beam power. The electrical dissipation has been estimated from measurements of the Q of a prototype magnet. Estimates of the beam power dissipation are still in progress; we have assumed 25 kW for design purposes, which should be conservative.

Removing the total amount of heat from the magnet by water cooling presents no particular problems; the main concern is in local heating of the magnet core. It is estimated that about 4 kW due to eddy current heating will be dissipated in the first 5 mm of the throat of the core throat.

Instantaneous temperature rise in the coil is of the order of $1^{\circ} \mathrm{C}$, and exterior of the core is of the order of $5^{\circ}$ and present no problem. The throat of the core is estimated to have a maximum instantaneous temperature rise of $40^{\circ} \mathrm{C}$. This
may present long term fatigue problems; tests are required.
Prototype magnets operating at reduced current and repetition rate have been constructed to measure magnetic fields and make general design tests. A full scale prototype magnet is being built to test the design at full current. These will be especially useful to measure heat dissipation in the core.

The coil is mechanically constrained by a surrounding band to contain the radial force per unit length of coil of $5 \times 10^{4} \mathrm{nt} / \mathrm{m}$ at full current. The mechanical constraint poses no particular problems, but it must not form a shorted turn which would decrease the magnetic field in the concentrator. A mechanical design for this exists.

Figure 5.2.6.1

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## CHAPTER 6. DAMPING RING AND BUNCH COMPRESSOR

6.1 INTRODUCTION
6.1.1 Choice of Parameters
6.1.2 General Features and Layout
6.2 DAMPING RING
6.3 LINAC-TO-RING TRANSPORT (LTR)
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## 6. DAMPING RING AND BUNCH COMPRESSOR

### 6.1 INTRODUCTION

The primary purpose ${ }^{1}$ of damping rings is to reduce the transverse phase spaces of positron and electron beams to suitably small values at the injection end of a linear collider such that their bunches, when accelerated to high energies, can be made to collide with high luminosity. Further, longitudinal bunch length control must be provided to minimize wake field effects in the accelerator. Finally, it is desirable to provide a means of preserving and manipulating the spin polarization of the electron beam.

The following description pertains primarily to the first SLC damping ring which was made operational in 1983. the second ring is identical in concept and major features, but differs in some details. These dissimilárities are due to differing requirements for the two beams and some improvements that were indicated by experience with the first ring.

### 6.1.1 Choice of Ring Parameters

Taking as given ${ }^{2}$ that positrons are created at the converter with an invariant emittance of about $\epsilon=5 \pi(\mathrm{MeV} / \mathrm{c}) \mathrm{mm}$, and that the design luminosity requires (before certain emittance growths are accounted for) an invariant emittance of $0.015 \pi(\mathrm{MeV} / \mathrm{c}) \mathrm{mm}$, a transverse phase space reduction of about 300 is needed. The damping time must be in the millisecond region, since the Collider will operate at 180 Hz . The interplay between damping ring energy $E_{D}$, bending radius $r_{D}$, field $B_{D}$, the number of bunches per ring $n_{B}$, the number of rings $n_{M}$, the time between bunches $t_{B}$, the total beam current $i_{D}$ (assuming $5 \times$ $10^{10}$ particles/bunch) has been parameterized ${ }^{3}$ as a function of bending magnet length for a separated function lattice and is shown in Fig. 6.1.1.

Combining this parameterization with the physically realizable values of magnetic fields, gradients and practical rise and fall times for injection and extraction kickers leads to the ring parameters shown in Table 6.1.1. Two bunches circulate in each ring in diametrically opposite buckets. For the positron case,
the two buckets are filled and emptied on alternating accelerator pulses. Hence, the positrons spend 11.1 ms or 3.6 damping times in the ring. For the electron case, less damping is required, so two bunches are simultaneously injected and extracted each 5.6 ms , the second bunch being used to subsequently generate positrons for refilling the emptied positron bucket.

Figure 6.1.1 Parameterization for Design
$\rho$ : bending radius
$E_{D}$ : energy
$B_{D}$ : bending field
$i_{D}$ : total beam current ( $5 \times 10^{10}$ particles/bunch)
$t_{B}$ : time between bunches
$n_{D}=n_{B} n_{m}$ : total number of bunches
$n_{B}$ : number of bunches per damping ring
$n_{m}$ : number of damping rings



### 6.1.2 General Features and Layout

The first damping ring is located on the south side of the west end of the linac near the juncture of Sector 1 and Sector 2. The geometry of this complex is shown in Figure 6.1.2.1. A second ring, now under construction and not shown in the figure, will lie on the north side of the linac. The plan view of this second ring will be almost a perfect mirror image of the south ring with the linac beam line being the axis of symmetry. The insert in the figure depicts how particles of opposite charge are extracted and reinjected into the accelerator in the presently limited space between linac sectors. The scheme employs two horizontally bending dipole magnets. The upstream magnet is the extraction dipole which bends electrons (positrons) to the south (north) into the entrance of the linac-to-Ring (LTR) transport line. For stability during collider operation this magnet will operate dc, but for ring tests it is pulsed to provide interlacing with beams for the other linac programs. The downstream or reinjection magnet is designed for dc operation only. It, again, bends returning beams of opposite charge in opposite directions so that they both exit in a direction collinear with the center line of the linac wave guide. When this magnet is operating the the damping rings become the only sources of beams for the linac sectors downstream of Sector 1.

The south ring and those parts of its transport lines outside of the linac housing lie in a common level plane at an elevation 36 cm below the elevation defined by the height of the undeflected linac beam at the center of the extraction-reinjection region. This elevation for the intial design was chosen to allow the option of building the second ring above the first. The north ring and its corresponding transport lines will also lie in a common level plane, but at an elevation 4 cm above the previously defined point, an optimal height for minimizing vertical deflections in the second order optical design of the return line.

Both rings reside in concrete vaults some 35 ft . below grade. The south transport lines are housed in eight-ft diameter corrugated steel tunnels, while those on the north are hosed in ten-ft square concrete tunnels.

## Figure 6.1.2



### 6.2 DAMPING RING

The global design parameters for the damping ring have been previously given in Table 6.1.1. Figure 6.2 .1 shows the ring and important components in greater detail. In this figure the two injection and two extraction septa magnets are labeled with the prefix $S B$, the two kicker magnets with the prefix $K$, and the two RF cavities with prefix RFC. The nominal design betatron and dispersion functions of the ring are shown in Figure 6.2.2. A complete listing of the Twiss parameters at the exit of each element as generated by the machine modeling code COMFORT ${ }^{4}$ is given in Table 6.2.1. The position of all orbit correctors ${ }^{5-6}$ and beam position monitors are included in this deck for the purpose of modeling closed orbit corrections. The bend magnet cross section is shown in Figure 6.2.3. The field in the 2 cm gap is 19.8 kG and the effective length is 32 cm . The good field region where $|\delta B / B|<0.1 \%$ extends over $\pm 1.0 \mathrm{~cm}$ in the horizontal plane.

Quadrupoles which bear a common name in the COMFORT deck are bussed in six families so that their strengths can be varied to change tunes and the optical properties of the insertion. It is intended to run the machine on the coupling resonance $\nu_{x}-\nu_{y}=4$ to minimize the beam emittance. The cross section of a typical laminated quadrupole is shown in Figure 6.2.4. The bore of the quadrupoles in the arc lattice is 25.8 mm . The bore of the four insertion quadrupoles and the four matching quadrupoles is 51.5 mm and 35.4 mm , respectively.

Sextupoles are introduced at the ends of each bending magnet by pole face shaping to correct the natural chromaticities from $\approx-9$ to $\approx+0.5$. In addition, four discrete sextupoles ( $S_{F} \approx S_{D} \leq 25 \mathrm{~m}^{-2}$ ) powered in pairs are used to trim the chromaticity.

Figure 6.2.1 Damping Ring Components


Figure 6.2.2 Nominal Betatron and Dispersion Functions


## Table 6.2.1



Table 6．2．1 Continued

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Figure 6.2.3


SECTION THRU 19.7 KG EEND MAGNET

Figure 6.2.4


TYPICAL DAMPING RING QUAD
DIM.'A'\{ $\begin{cases}\text { FOCUSING }-5.950 \prime \prime & \text { ( } 16 \text { TOTAL) } \\ \text { DEFOCUSING }-3.375^{\prime \prime} & \text { ( } 18 \text { TOTAL) }\end{cases}$

### 6.3 LINAC-TO-RING TRANSPORT LINE (LTR)

The unique geometry of the LTR transport line evolved early in the design of the damping ring complex, in response to the need for beam polarization control while simultaneously matching geometric constraints determined by 1) the choice to combine the extraction and injection region in the linac and 2) the optical requirements of the more sensitive designs of the ring (DR) and return line (RTL).

To preserve and manipulate the spin polarization of the electron beam, the optics and geometry of the transport lines are so chosen that solenoidal spin rotators ( $B_{z} \ell \approx 6.3 T-m$ ) can be inserted in regions of zero dispersion and at odd $\pi / 2$ multiples of spin precession angle. At 1.21 GeV this occurs at bend angles of $(2 n+1) \times 32.8^{\circ}$. Thus in LTR after $164^{\circ}$ of bend the longitudinally polarized beam extracted from the linac has precessed to become transversly polarized in the plane of bend. The solenoidal field then precesses the spin direction to be normal to the bend plane of the ring so that polarization will be preserved during damping. The transfer matrix for a solenoidal field can be represented by a focusing element nested between a rotation and inverse rotation about the beam axis, thus a requirement that the dispersion be zero at the rotator solenoid avoids mixing horizontal and vertical dispersion caused by rotation. The induced correlation in the $x-y$ projection of the beam phase space can be made transparent by causing $\beta_{x}=\beta_{y}$ and $\alpha_{x}=\alpha_{y}$ in the solenoid. Such a provision will be a minor future modification to LTR.

Having provided for polarization control, the need to match to a geometry prescribed by the RTL optics (where a smooth monotonic $\int(\eta / \rho) \delta s$ is desirable) and the ring resulted in the unusual geometry and optics of LTR.

A path length constraint was also applied to synchronize positron and electron beams.

Figures 6.3.1 and 6.3.2 show the betatron and dispersion functions for LTR. The future spin rotator solenoid will be located near the point where $s=40 \mathrm{~m}$. Figure 6.3.3 shows the mechanical apertures superimposed on the envelope of a beam for which $\epsilon_{x}=\epsilon_{y}=1.0 \mathrm{~mm}-\mathrm{mr}$ and $\delta=1.0 \%$.

The transport line contains 23 horizontal bend magnets identical in cross section to the ring bend magnets (see Fig. 6.2.3). These magnets are powered in series from a single source and provide $200.7^{\circ}$ of bend. Three vertical bends also powered in series are used to bring the beam to the elevation of the level plane of the damping ring, 36 cm below the height of the extraction point. All of the LTR line outside of the linac housing also lies in this plane.

The main windings of the 45 quadrupoles are powered in series while the first eight and the last ten are equipped with trim windings for correction and tuning. In the two central straight sections the quadrupoles are arranged in a FODO array with $45^{\circ}$ phase shift per cell. The $\beta$ function is not now matched in these regions but this loss of elegance has not effected performance. The remaining quadrupoles ( $\sim 1 / 2$ ) were used in the design to match to the linac and damping ring machine functions.

Extraction from the linac is achieved by pulsing a laminated core horizontal bending magnet (r.t. $\approx 3 \mathrm{~ms}$, rate $\geq 20 \mathrm{~Hz}$ ) which bends the electron(positron) beam south (north) by $7.0^{\circ}$ (see Fig. 6.3.4). The beam after passing through a protection collimator and a ten port vacuum chamber then enters a field free aperture in the yoke of the downstream reinjection magnet. A small correcting dipole has been placed immediately upstream of this aperture to correct for any field leakage into the aperture. Vertical bending magnets not shown in the figure then deflect the beam into the level plane which is common to the remainder of LTR, the ring and most of RTL. When the upstream magnet is not pulsed the undeflected beam enters the good field region of the reinjection magnet, thus this magnet must remain off for straight ahead operation.

## Diagnostics

The LTR transport line is equipped with diagnostic instrumentation for measuring the beam size, intensity, energy spread and orbit distortions. Control, status and appropriate output signals are interfaced through CAMAC to a microprocessor dedicated to LTR. The microprocessor is linked through SLCNET to the VAX residing at the Main Control Center (see I\&C Section 7.).

The energy spectrum of the beam being delivered from Sector 1 is measured
with a 20 foil spectrum monitor (see Fig. 6.3.5) placed 3 m downstream of the extraction point. The tungsten foils are 0.8 mm wide with 2.0 mm spacing giving a resolution of $0.4 \%$ ( $\delta E / E$ per foil). In the axial direction the foils are thick ( $\approx 3$ r.l.) thus the unamplified signals are suitable input for a simple multichannel charge integrator with video graphics display. Charge integration and hold time are controllable through the VAX. During typical operation, the display is updated at 2 Hz for a 10 Hz beam of $<10^{10}$ electrons per pulse. The spectrum monitor foils are supported on the front face of a movable water cooled collimator. The collimator has two positions. When in one position, all foils are in the beam and the collimator serves as a 4 kW dump for tuning Sector 1. When the collimator and its foils are moved to the alternate position the 8 central foils are removed from the beam which now passes through a 16 mm diameter hole defining the energy spread to be $\pm 1 \%$, matching the acceptance of the ring.

Three beam current measuring toroids are located in LTR. The first is placed upstream of the energy defining aperture and measures the intensity of the beam delivered from Sector 1. The second toroid, downstream of the collimator measures the beam intensity within the energy acceptance of the ring. A third toroid at the end of LTR measures the transmission efficiency of the beam line. Signals from these last two toroids are to be integrated and compared during each machine pulse providing a means of limiting component damaging beam loss to an acceptable level (100-200 W).

A Faraday cup can be remotely inserted into the beam downstream of the last toroid as an independent measurement of the charge delivered to the ring.

The beam orbit is measured by 19 beam position monitors (BPM) placed at intervals of approximately one-quarter of a betatron wavelength. Each beam position monitor has four strip line detectors, one in each quadrant. The four signals (one from each electrode) are electronically summed then digitized for processing by the SLC VAX computer (see Fig. 6.3.6). Beam position information is then presented in either numeric or graphic form. All beam position monitors are not read out during the passage of a single beam pulse. Instead, a
system of CAMAC-controlled solid state one pole-ten throw relays sequentially switch signals to the processor during the inter-pulse period. Three of these relays with three processors presently allow a sequential scan of the beam position monitors in sets of three. The scan rate is limited by software processing time and it requires approximately 3 seconds for a scan of all LTR BPM's. This slow scan rate has not presented any operational difficulty however, since orbit correction is usually done with low intensity beams having good pulse to pulse stability. The measured orbits reproduce with rms deviations of $\approx 100 \mu \mathrm{~m}$.

## Figure 6.3.1

LTR Betatron Functions

Figure 6.3.2
LTR Eta Functions

LTR SOUTH: LTRS413 DECK: LTRS413 FINAL : DJ13




Figure 6.3.3 LTR Beam and Vacuum Chamber Clearance


Figure 6.3.4 Extraction and Re-Injection at Girder 1-9



Figure 6.3.5 Spectrum Monitor Foils and Collimator


Figure 6.3.6 Beam Position Monitor Signal Processing


### 6.4 RING-TO-LINAC TRANSPORT LINE (RTL)

The first and second order optics for the damping ring to linac transport line are designed ${ }^{7}$ to preserve the damped transverse emittance while simultaneously compressing the bunch length of the beam to that length required for reinjection into the linac. This design, including provisions for future control of beam polarization, is described here.

The design parameters for the SLC damping ring (DR) have been described elsewhere. ${ }^{8}$ Those parameters important to the description of the RTL transport line are listed in Table 6.4.1 and were taken as given input conditions for the transport line. At the output of the transport line it is important to have control of the bunch length to reduce both the transverse emittance growth and the energy spread due to transverse and longitudinal wake fields as the particle bunches are accelerated through the linac. A delicate balance must be maintained to minimize these effects since the transverse emittance growth decreases with shorter bunch length while the opposite is true for the energy spread. Calculation of the effects of these fields ${ }^{9}$ have shown that an rms bunch length on the order of 1.0 mm is suitable for the acceleration of $5 \times 10^{10}$ particles/bunch. To ensure that adequate control of this parameter can be provided the choice was made to design the RTL beam line such that compression of the 6.0 mm rms DR bunch length to a minimum rms length of 0.5 mm will be achieved.

| Table 6.4.1 Damping Ring Parameters |  |
| :--- | :--- |
| Energy | 1.21 GeV |
| Equilibrium emittance | $9.1 \times 10^{-9} \pi$ rad-m |
| Extracted beam emittance | $1.2 \times 10^{-8} \pi$ rad-m |
| Equilibrium-related energy spread | $7.4 \times 10^{-4}$ |
| Equilibrium bunch length | 5.9 mm |

## Compression

The bunch length compression in RTL is accomplished by first passing the particle bunches through an S-band accelerating section at $0^{\circ}$ central phase angle immediately after extraction from the damping ring. This maneuver introduces a correlation between the position of a particle in the bunch and its energy (see Fig. 6.4.1). Compression then occurs due to energy dependent'path length differences proportional to the momentum compaction factor $\alpha$ (see Fig. 6.4.2).

It can be seen from Fig. 6.4.1 that the slope of the required correlation between energy and bunch length is very nearly defined by the damping ring energy spread $\delta(0)$ and the chosen rms bunch length $\ell(1)$ of 0.5 mm after compression. Hence the energy spread at the end of RTL will be given by $\delta(1) \approx[\delta(0) / \ell(1)] \ell(0) \approx 0.009$. The required peak accelerating voltage $V_{\mathrm{RF}}$ then follows from $\delta(1)=\left(V_{\mathrm{RF}} / E_{R}\right) \sin (2 \pi \ell(0) / \lambda)$ where $E_{R}$ is the ring operating energy of 1210 MeV and $\lambda$ is the $S$-band wavelength of 0.105 m . Thus, $V_{\mathrm{RF}}$ is $\approx 30 \mathrm{MeV}$ which in practice can be easily supplied by a single $3-\mathrm{m}$ disc loaded waveguide identical to those used in the linac. The momentum compaction factor given by $\ell(1) / L \delta(0)=\alpha \approx 1 / \nu^{2}$ can be used to estimate that the RTL should span approximately two betatron wavelengths.

## Optical Design

The RTL transport line will be required to have an energy acceptance greater than $\pm 2 \delta(1)$. For such a large energy spread control of chromatic aberrations becomes a primary concern. The concept of the second-order magnetic optical achromat ${ }^{10}$ was used as a design guide though strict adherence to its principles


Figure 6.4.1


Figure 6.4.2
was precluded because of the need to fit the transport line into an existing linac housing. The theoretical principles of the second-order achromat are described elsewhere. ${ }^{11}$ Here we describe its application to the RTL design. The beam line consists of two stages (see Fig. 6.4.2) each characterized by a transformation matrix equal to the identity matrix. Stage one consists of four identical cells, each cell containing two bending magnets, two quadrupoles, and two sextupoles. The strength of the two quadrupoles and two drift distances are determined by the constraint that the overall transformation be the identity with a total phase shift of $2 \pi$. Then, since there are four identical cells the second-order geometric terms vanish due to symmetry. And, as shown by K. Brown, all second-order chromatic terms can be removed by properly setting the strength of the two sextupoles.

Stage two is not so neatly characterized. The placement of the bending magnets in this stage is dominated by the geometry of the RTL tunnel and the linac housing. With needed reverse bends and vertical as well as horizontal bends $n$-fold symmetry could not be maintained for $n>2$. Instead, the stage is composed of two identical "cells" giving "half wave symmetry." All components in this stage, including those special magnets used to re-inject into the linac are imaged by the negative of the identity matrix onto an identical component in the other "cell," assuring that overall dispersion and second-order geometric aberrations vanish. Thus, stage two consists of six independent pairs of bend magnets, five independent pairs of quadrupoles and four independent pairs of sextupoles.

Another unusual feature of stage two can be observed in Fig. 6.4.2. It will be noted that the first half of this stage has been split and stage one has been inserted, a maneuver that was made possible by the stage one transformation being the identity. This "nested achromat" was introduced to take advantage of the large horizontal bend angle in the first two bending magnets of stage two to obtain a practical fit to the existing housing.

With the bend magnet strengths thus constrained by the geometry, six parameters, namely four of the quadrupole strengths and two pairs of intervening
drifts, are determined by the constraint that the overall transformation of stage two be the identity. One pair of imaged quadrupoles are used to match to a periodic $\eta$-function in the nested stage one, providing a smooth bunch length compression in this section.

## Matching the Linac Lattice

Figures 6.4.3a-c illustrate the resulting matching functions after completion of the first-order fitting described above and a match has been made to the linac lattice. This latter matching is achieved by inserting four quadrupoles between the damping ring extraction optics and the upstream end of the compressor waveguide. These quadrupoles are used to match to the desired linac lattice at the beginning of the two achromats just downstream of the compressor waveguide where there is an identical image of the linac entry point. Second-order effects in the region before the compressor waveguide are negligible because of the small momentum spread of the extracted beam.

## Second-Order Correction in Stage Two

The mixture of horizontal and vertical bends in stage two results in mixed $x, y$ dispersion at the sextupole sites. In order to avoid the introduction of crossplane chromatic aberrations, this means that the four sextupole pairs must be rotated axially with respect to the beam coordinate system. That was done in conjunction with variation of the strengths of the four sextupole pairs to zero the $T_{i 66}, i=1,2,3,4$ and the $T_{i j 6}, i=1,2, j=3,4$, elements of the second-order transfer matrix.

It was found empirically that the foregoing prescription reduced all chromatic aberrations in stage two to insignificant levels. It may not, however, be generally applicable because we are not dealing with a theoretically perfect second-order achromat.

## Bunch Length

The progressive compression of a bunch as it moves through the RTL beam line is shown in Fig. 6.4.4a. Most of the compression occurs relatively smoothly in stage one. There is some unwanted but unavoidable fluctuation of the length

Figure 6.4.3: (a) $\beta$ vs. $s(m)$; (b) $\eta$ vs. $s(m)$; (c) Envelope vs. $s(m)$; where $E=\sqrt{\epsilon \beta+(\eta \delta)^{2}}$, for $\epsilon_{z}=\epsilon_{y}=0.01 \mathrm{~mm}-\mathrm{mrad}$ and $\delta=2 \%$

of the fully compressed bunch in stage two. Figure 6.4.4b shows the correlation of the bunch length and the momentum spread as given by the $\boldsymbol{r}_{56}$ correlation term of the TRANSPORT sigma matrix. The locations in the second stage where this term is less than zero correspond to over-compression. At the point of injection into the linac this term is equal to zero as intended.

## Bunch Lengthening Effects

Figure 6.4.1a-b indicate a linear correlation for the particle position and its energy. This would only be true for $2 \pi \ell(0) \ll 1$. Antisymmetric bunch lengthening occurs when the small angle approximation is not satisfied. A second effect which is symmetric about the bunch center is also introduced by the second order momentum dependent path length differences. Both these effects have been examined ${ }^{12}$ for this design and can be adequately compensated by trimming the compressor waveguide voltage and phase angle.

Coherent bunch lengthening due to the longitudinal wake field of the bunch has also been examined and shown to be negligible.

## Spin Polarization of the Electron Beam

It is expected that an electron beam from a polarized source will be injected into the DR with transverse spin polarization perpendicular to the plane of the DR. One of the design goals for the RTL electron beam is to be able to rotate the spin polarization to any arbitrary direction. ${ }^{13}$ This will be done with the aid of suitably disposed solenoid magnets. In order to avoid disruption of the basic optics of the RTL beam, solenoids can only be located in places where there is no dispersion and where the beam is "round," (i.e., $\beta_{x}=\beta_{y} ; \alpha_{x}=\alpha_{y}=0$ ). Two solenoids are needed, and for arbitrary control the spin precession angle in the RTL-bends between them should be an odd multiple of $90^{\circ}$. There are only two possible locations which satisfy these criteria, one just before entry into the achromats and the other just after reinjection into the linac. The spin precession angle in the intervening beam line, the whole RTL beam, is almost exactly $270^{\circ}$ at 1.2 GeV .

Figure 6.4.4: (a) Bunch length vs. $s(m)$ for $V_{R F}=33 \mathrm{MeV}$; final rms length is 0.5 mm . (b) $\mathrm{r}_{56} \ell-\delta$ correlation) vs. $\mathrm{s}(\mathrm{m})$.


## Orbit Correction

The placement of beam position monitors and orbit correcting magnets in RTL was determined by the same method ${ }^{14}$ used for the LTR beam line.

The construction of ten beam lines, each with randomly displaced components were simulated and the efficacy of various correction schemes were tested. The criteria for judging performance included a one milliradian bend as a practical upper limit for the strength of the correctors. Each component including the beam position monitors themselves were independently misaligned using untruncated gaussian distributions. It is expected that in practice such errors would not be truly gaussian but would be truncated, thus this estimate is conservative.

The alignment tolerances used for these simulations and construction specifications were:

| for all <br> $(\mathrm{mm})$ | for quads <br> $(\mathrm{mr})$ | for bends <br> $(\mathrm{mr})$ |
| :---: | :---: | :---: |
| $\delta x=0.2$ | $\delta x^{\prime}=0.5$ | $\delta x^{\prime}=0.1$ |
| $\delta y=0.2$ | $\delta y^{\prime}=1.0$ | $\delta y^{\prime}=0.1$ |
| $\delta z=2.0$ | $\delta z^{\prime}=0.5$ | $\delta z^{\prime}=0.1$ |

The results of this study are shown in Figures 6.4.5-6.4.7.

Figure 6.4.5a


Figure 6.4.5b
RTL TRANSPORT LINE


Figure 6.4.6a

Figure 6.4.6b


CORRECTED HORIZØNTAL TRAJECTØRIES


6-42

RTL TRANSPORT LINE

Figure 6.4.7a

Figure 6.4.7b



### 6.5 SUBSYSTEMS AND COMPONENTS

### 6.5.1 Damping Ring RF System ${ }^{15}$

The linear collider project at SLAC contains two damping rings to reduce the emittance of short electron or positron bunches which contain $5 \times 10^{10}$ particles per bunch. Two of these bunches are stored at a time and then extracted for acceleration in the collider. The RF system is subject to strong transients in beam loading. A computer model is used to optimize capture while minimizing RF power. The introduction of phase jump in the RF drive at injection time together with offsets in the tuning loops of the RF cavities when no beam is stored allows optimum performance under heavy beam load conditions. The RF system ( 800 kV at 714 MHz ) for the electron damping ring has been built, tested and installed, and is being tested with beam.

## Introduction

The linear collider project at SLAC (SLC) contains two damping rings to reduce the emittance of the $\mathrm{e}^{+}$and $\mathrm{e}^{-}$bunches. This is necessary to attain high luminosity at collision. The bunches of each charge type are stored in separate damping rings. Each ring is a small high-field storage ring in which radiation damping takes place. The RF system is required to replace the energy lost by synchrotron radiation. The parameters of the RF system are given in Table 6.5.1.1. At the 1.21 GeV energy of the stored beams the radiation per turn is a modest 93 keV . However, each injected bunch contains $5 \times 10^{10}$ particles and this induces parasitic mode losses and strong transient beam loading. A theoretical treatment of beam loading effects indicates that optimum performance can be obtained by introducing a phase jump at injection together with an offset in the tuning of the cavity under no beam conditions. All particles which arrive with an energy spread of $\pm 1 \%$ will be captured with the RF system preset in such a fashion. No rapid tuning is necessary, which greatly simplifies the system, but the offset in the tuning loop has to be applied during the time when no beam is stored to keep the cavities in the proper tune condition.

| Table 6.5.1.1 RF System Parameters |  |  |
| :--- | :---: | :--- |
| RF Frequency | 714 | MHz |
| Harmonic Number | 84 |  |
| Available Klystron Power | 50 | kW |
| Total Cavity Shunt Impedance | 35 | $\mathrm{M} \Omega$ |
| Unloaded Q | 24000 |  |
| Cavity Coupling Coefficient | 2.5 |  |
| Cavity Tuning Angle | $-64^{\circ}$ |  |
| Injection Phase Angle | $78^{\circ}$ |  |
| Phase Jump at Injection | $-50^{\circ}$ |  |
| Energy of Stored Beam | 1.21. | GeV |
| Energy Loss per Turn | 93 | keV |
| Parasitic Mode Loss per Turn (full current) | 48 | keV |
| Circulating Current per Bunch | 70 | mA |
| Storage Time Between Pulses (electrons) | 5.56 | msec |
| Storage Time Between Pulses (positrons) | 11.11 | msec |
|  |  |  |

Two RF cavities with two cells each provide the required gap voltage. Amplitude and phase feedback loops stabilize the RF fields in the cavities and thus give the beams a precise timing when they are ejected into the linear accelerator.

To date the damping ring for $\mathrm{e}^{-}$particles has been constructed and is being commissioned.

## Parasitic Mode Losses

A bunch of charged particles circulating in a storage ring excites electromagnetic fields in the neighborhood of discontinuities in the vacuum chamber walls. In the damping ring there are discontinuities and transitions between vacuum chamber components. The bunch can, in addition, excite higher-order modes in the RF cavities. Field and modes excited by the bunch represent an energy loss which is in addition to the synchrotron radiation loss per turn. By adding together losses for all the individual components plus the higher-order mode losses
in the RF cavities, an estimate can be made for the total loss voltage, $V_{1}=Z_{1} I_{0}$, where $I_{0}$ is the circulating current per bunch and $Z_{1}$ is the loss impedance. For the damping ring

$$
Z_{1}(\mathrm{M} \Omega) \approx 1.9 \exp \left[-0.17 \sigma_{z}(\mathrm{~mm})\right]
$$

Note that the loss impedance depends strongly on the bunch length. At the design current of 70 mA per bunch, the loss voltage is $V_{1} \approx 48 \mathrm{kV}$ for the damped bunch length of 6 mm . This must be added to the synchrotron radiation loss of 93 kV to obtain the total voltage per turn that must be supplied by the RF system.

## Transient Beam Loading

In the damping ring the total stored charge is injected in one or two bunches on a single turn, rather than over many turns lasting many cavity filling times, as is usually the case for a storage ring. Thus transient beam loading effect might affect the RF voltage required to capture the injected bunches. It is not simple to write analytic expressions taking into account transient beam loading effects because of the large injected energy spread ( $\pm 1 \%$ ), and the resulting large-amplitude phase oscillations for particles at the extremes of this spread. It is straightforward, however, to write turn-by-turn expressions for the energy and phase deviations from the synchronous energy and phase for a relatively small number of superparticles distributed over the injected phase space. The recurrence relations expressing the energy deviation $\epsilon$ and phase deviation $\theta$ of the $i$ th particle on the $n$th turn in terms of quantities on the previous turn are:

$$
\begin{array}{r}
\theta_{n}^{(i)}=\theta_{n-1}^{(i)}+A \epsilon_{n}^{(i)} \\
\epsilon_{n}^{(i)}=\epsilon_{n-1}^{(i)}+V_{g_{n-1}}^{(i)}-V_{B_{n-1}}^{(i)}-V_{s}
\end{array}
$$

Here $A=2 \pi \alpha h / E_{0}$ where $\alpha$ is the momentum compaction factor, $h$ is the harmonic number and $\mathrm{e} E_{0}$ is the beam energy; $\mathrm{e} V_{s}$ is the synchrotron energy loss per turn, $V_{g}$ is the generator voltage component $V_{g}=V_{g R} \cos \psi \cos (\theta+\psi)$, where $\psi$ is the tuning angle and $V_{g R}$ is the generator voltage component at res-
onance; and $V_{B}$ is the beam loading voltage component given by an appropriate sum over the beam induced voltages of all of the individual super particles. In calculating this sum, the decay of the beam induced voltages between turns due to the finite cavity filling time, and the phase shifts due to phase oscillations, must be taken into account. The program DAMP ${ }^{16}$ has been written at SLAC to perform this calculation. ${ }^{17}$

Initial results indicated that a minimum klystron power of 44 kW would be required to capture uniform injected energy distribution with $\delta \epsilon_{\max }= \pm 12$ $\mathrm{MeV}, \sigma_{\epsilon}= \pm 7 \mathrm{MeV}$. This minimum power was obtained by adjusting the cavity coupling and tuning, and the injection phase and energy deviation of the central superparticle in a distribution which is uniform in energy up to $\delta \epsilon_{\max }$ with zero phase width. It was soon realized that the required klystron power could be reduced by making a jump in the generator phase at the moment of injection such that the transient in the generator voltage roughly cancels the transient beam induced voltage. A more precise cancellation of all transient beam loading effects could be obtained by jumping both the phase and amplitude of the generator voltage component, but a phase jump alone gives good results and the hardware and control problems are simpler.

Figures 6.5.1.1 and 6.5.1.2 give a typical example of the oscillations in energy and phase for 500 turns after injection for a phase jump in VG at injection of $-40^{\circ}$.

Details of the other RF and ring parameters are given in Ref. 18. Fig. 6.5.1.3 shows the distribution in phase and energy for the injected bunch and for the superparticles after 500 turns ( $\approx 60 \mu \mathrm{~s}$ ). Also shown is the damped distribution after several damping times ( $\tau_{d}=1.5 \mathrm{~ms}$ for phase oscillations). Note that the rms energy spread has been reduced by damping by about a factor of six.

By adjusting the magnitude of the phase jump, the cavity coupling coefficient, the cavity tuning angle and the injection phase angle, it is possible to reduce the klystron power required to about 30 kW and still capture the complete injected energy spread. Without a phase jump, the minimum klystron power is 44 kW . The maximum energy excursion is also reduced slightly, from
$\pm 15 \mathrm{MeV}$ to $\pm 13 \mathrm{MeV}$.

## RF System Description

The block diagram of the RF system is shown in Fig. 6.5.1.4. The design of the RF accelerator structure was optimized to achieve a peak accelerating potential in excess of 1 MV with the available 50 kW of RF power. Calculations with the LALA program yielded a cavity design using four re-entrant copper cells with slot coupling between cells. Due to limited straight section space in the storage ring the four cells were split into two assemblies (see Fig. 6.5.1.5) with two cells each, installed almost opposite each other in the ring. Each structure includes a ceramic window in the input waveguide and a slot coupling from the waveguide into one cell as well as a fixed and a moveable tuner. Two ports are provided for vacuum pumps. A shunt impedance of $R=V^{2} / P=17.4 \mathrm{M} \Omega$ was measured for each two-cell RF structure.

A waveguide magic tee is used to split the power feed to the two RF cavities. Since no counter rotating particles are stored in the damping rings the waveguide lengths and position of the cavities in the ring can be arranged such that reflected power from each cavity is combined in the magic tee's terminated port. Thus the requirement for a costly isolator to protect the klystron is eliminated. The klystron is a commercially-available TV tube with 50 kW CW output.

The RF system is stabilized by a total of four feedback loops to assure accurate ejection of the stored and damped beams back into the linear accelerator at the correct phase of the fields in the linear accelerator.

Each cavity is tuned by a feedback loop which compares the RF phase of the field in the cavity to the phase of the driving signal and uses the resulting error to operate the tuner via a stepping motor. Since the intermitantly stored beam is a strong detuning element to the cavity, the tuning loop is adjusted to provide optimum tuning, i.e., minimizing reflected power from the cavity with the beam stored. A pulsed offset is introduced when no beam is in the ring. This offset counteracts the large error which would otherwise be detected when the beam is not present and thus keeps the cavity in a state ready for the next injected beam.

In a similar fashion a pulsed offset is introduced in the main phase feedback loop during the "beam-off" time to rotate the phase of the field in the two cavities to a position where it can best capture the injected beam. The offset is removed at beam injection, and produces the phase jump discussed earlier in this paper. ${ }^{18}$ Immediately after injection the rapidly rising beam induced fields rotate the cavity fields to the desired phase. The main function of the phase feedback loop is to lock the vector sum of the field vectors in the two RF cavities to the input signal of the RF system. Long cable runs in this loop are temperature stabilized with the coax cable surrounded by a coaxial water jacket operating at $+45^{\circ} \mathrm{C} \pm 0.1^{\circ} \mathrm{C}$. Phase stable coax cables with $\pm 9 \mathrm{ppm} /{ }^{\circ} \mathrm{C}$ temperature coefficient and foam dielectric are used.

The gap voltage feedback loop sums the field amplitudes detected from samples of each cell and compares it to a fixed reference voltage. The resulting error signal is applied to a variable attenuator in the drive line to the $k$ klystron.

The amplitude and phase detectors, electronic attenuators and phase shifters used in the drive and feedback circuits had been developed for the PEP storage ring RF systems and are described in an earlier paper. ${ }^{18}$

## Present Status

The RF system as described is operational and the damping ring has stored current of $\approx 10 \mathrm{~mA}$ with a lifetime of close to one hour. The synchrotron frequency has been measured as a function of gap voltage and the "cold tested" parameter of the cavities have been verified with the beam. The heavy beam loading tests await the production of intense bunches from the linac.


Figure 6.5.1.2 Variations in the mean energy deviation, 'the rms energy deviations and the extreme energy deviations for 500 turns.


Figure 6.5.1.1 Variations in the mean ohase. the rms phase and the extreme phase excursions for 500 turns.


Figure 6.5.1.4 Block diagram of damping ring RF system



Figure 6.5.1.5

### 6.5.2 Vacuum System

The vacuum system for the first damping ring has designed ${ }^{19}$ requirements to 1) provide sufficient cooling capability to absorb 13 kW of synchrotron radiation power; 2) provide beam position monitor electrodes shielded from synchrotron radiation; 3) include ion clearing electric fields, while; 4) providing a base pressure of $\approx 5 \times 10^{-8}$ Torr in a low conductance system. The small size and close packed construction of the ring made meeting the requirement a unique engineering challenge.

Ion effects, if not dealt with, can be expected to affect the phase space of the beams in the electron ring. ${ }^{20}$ A two step attack on the problem is employed. Distributed ion pumps ( $150 \mathrm{l} / \mathrm{sec}$ ) are located in each of the 40 bending magnets to reduce pressure (after clean-up) into the $10^{-8}$ Torr range. Further, the partitioned vacuum chamber is designed to let the electric field of the pump electrodes leak into the beam region to cause the ions to migrate and be cleared. Clearing in the remaining circumference is provided by clearing electrodes and by biasing the beam position monitor strips with voltages up to $\pm 1 \mathrm{kV}$.

## Damping Ring (DR)

The ring is oval in shape with short straight sections near the end of each arc and two longer straight sections separating the arcs. Forty bend magnets with alternating focusing and defocusing quadrupoles define the arcs, while the straight sections are delineated by quadrupoles, sextupoles, kickers, septa and RF cavities.

The vacuum system consists of 34 metal chambers joined by conflat-type flanges with metal gaskets. The arcs are composed of sixteen standard cells which are, in many respects, similar to the standard bend chambers of SPEAR and PEP. Each arc chamber is contained within two bend magnets and focusing and defocusing quads. Sections through the bend magnets are aluminum extrusions containing distributed ion pumps. Sections through the quads are stainless steel tubing which contain clearing field electrodes, beam position monitors (BPM), water cooling and synchrotron masks. Each standard aluminum arc chamber also has stainless steel welded bellows at one end. At the ends of
each straight section are matching chambers (eight total) which are contained within a single bend magnet and various combinations of sextupole and quad magnets. The components of the matching chambers represent slight variations of the standard chamber to accommodate the matching cells lattice. The short straight sections near the ends of one (north) arc are formed by the RF cavities. The cavities are oxygen-free, electronic (OFE) copper assemblies, brazed in a hydrogen furnace. Near the ends of the other (south) arc are straight sections composed of alumina $\left(\mathrm{Al}_{2} \mathrm{O}_{3}\right)$ ceramic vacuum chambers clamped with the pulsed kicker magnets. Both long straight sections contain the septa, drift chambers, and instrumentation chambers: the optical monitor and horizontal and vertical beam scrapers. The septum magnets are contained with the vacuum chamber.

## Distributed Ion Pumps

The distributed pump for the damping ring is a diode sputter ion pump which operates in the 19.8 kG magnetic field of each bend magnet. The pump structure shown in Figure 6.5.2.1 consists of a stainless steel anode array sandwiched between two titanium cathode plates. The anode array is an assembly of $728,3 \mathrm{~mm}$ dia $\times 6 \mathrm{~mm}$ long stainless steel tubes. The cells are spot welded together to form an array of 91 rows of 8 cells each. The array structure is bent to conform to the bend radius of the vacuum chamber. The anode array is held at a potential of 3.5 kV which can be increased up to a maximum of 7 kV . The pump evacuates the beam chamber through a slot 6 mm wide. The slot limits the effective speed of the pump to $\approx 21 / \mathrm{sec}$ per cm of length. The stray electric field from the anode array combined with the magnetic field of the bending magnet produces a clearing field which removes ionized gas molecules from the bend magnet chamber.

Figures 6.5.2.2-6.5.2.4 show the performance of a prototype structure consisting of $207,3 \mathrm{~mm}$ dia $\times 5.7 \mathrm{~mm}$ cells. The solid curves were calculated from the empirical formulae of Malev and Trachtenberg. ${ }^{21}$

## Transport Lines

The transport lines are named the linac-to-Ring (LTR) line and the Ring-tolinac (RTL) line. The LTR consists of 81 vacuum chambers interconnected with various off-the-shelf bellows and fittings. The system is primarily aluminum (tubing and extrusions) to distribute the thermal effects of a miss-steered beam striking the wall. Various diameters of tubing are used to match the bores of magnets and size of the beams through each section of the line. Rectangular extrusions form the beampipe through bending magnets. Diagnostic instruments (profile monitors, spectrum analyzers, Faraday cups, beam stoppers, etc.) are mounted in standardized instrument containers and distributed along the LTR. The line is ion-pumped and pressure is monitored with nude Bayard-Alperttype ionization gauges. The line is isolated from the linac and damping ring by Viton ${ }^{\text {TM }}$ o-rings sealed, electropneumatically operated gate valves. Magnets and vacuum components are mounted and prealigned on concrete rafts prior to installation in the tunnel.

The RTL is similar to the LTR except for a compressor section which shortens the bunch length of beams removed from the damping ring. The compressor is a standard linac accelerator section with its own waveguide and klystron station. The RTL is ion-pumped and isolated from the linac and damping ring by gate valves.

In operation, the LTR and RTL maintain pressures better than the design criteria of $5 \times 10^{-8}$ Torr. The damping ring has a base pressure of $5 \times 10^{-9}$ Torr. Normal operation with high current beams may cause some degradation of the vacuum.
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Figure 6.5.2.1



Figure 6.5.2.2


Figure 6.5.2.3

Figure 6.5.2.4


### 6.5.3 Magnets and Power Supplies

The following lists (not yet completed) of magnets are in order as seen by the beam. The identification number in the leftmost column is that used in the engineering drawings. Power supplies are identified by the magnet string that they power. Further documentation is being prepared.
drimping rimg

| (Systea) |  |
| :---: | :---: |
| quadiupale | maghets |

Bysteal Location (L.T.R)

Table 6.5.3

## Date 01/18/84

Sheet 1 of 1
Listing by W.A.D-W/G.F.C


Table 6.5.3 Continued


Table 6.5.3 Continued

> DAMPING RING SUPPLY SYSTEM OUTPUT VOLTAGE/OUTPUT CURRENT VALUES

| Beam <br> Line <br> Section | P.S. Id. | Rating | Run val. 950 Mev | Loop Ohens | Run val. $1.216 \mathrm{ev}$ |
| :---: | :---: | :---: | :---: | :---: | :---: |
| LTR | PB-0 | 300V/200A | 210V/10A | 21.0 | 273V/13A |
| LTR | LTR-VB | $45 \mathrm{~V} / 400 \mathrm{~A}$ | 16V/255A | . 063 | 21V/326A |
| LTR | LTR-BEND | 600V/450A | $310 \mathrm{~V} / 2681$ | 1.16 | 453V/392A |
| LTR | LTR-QUAD | 600V/450A | $410 \mathrm{~V} / 120 \mathrm{~A}$ | 3.41 | 530V/155A |
| LTR | SE-O (ACME) | 100V/450A | 45V/255A | 1.76 | 58V/326A |
| LTR | QE13 | 20V/120A | NA |  | N/A |
| RING | RING BEND | 850V/450A | 540V/271A | 2.00 | 827V/415A |
| RING | R Q | 300V/200A | 170V/124A | 1.37 | 214V/156A |
| RING | R QF | 300V/200A | 150V/98A | 1.53 | 195V/127a |
| RING | R QDI | 300V/200A | 80V/74A | 1.08 | 118V/109A |
| RING | R QFI | $300 \mathrm{~V} / 200 \mathrm{~A}$ | 120V/100A | 1.20 | 178V/148A |
| RING | $R$ QFM | 45V/400A | 30V/79A | . 380 | 40V/104A |
| RING | K1-INS | 20 KV | 16.3 KV | N/A | 20400V |
| RING | K2-EXT | 20 KV | 18 KV | N/A | 23130 V |
| RING | SEPT-INS | 30V/2500A | 10V/2003A | . 005 | 13V/2600A |
| RING | SEPT-EXT | 30V/2500A | 10V/2038A | . 005 | 13V/2500A |
| RING | SEPT-INSIB | 336V/80A | 119/48A | . 230 | NA |
| RING | SEPT-EXTIB | 160V/50A | 11V/48A | 230 | NA |
| RING | B1 BYPASS | 9V/1000A | 8V/100A | . 030 | N/A |
| RING | B2 BYPASS | 9V/1000A | 7.5V/250A | . 030 | NA |
| RTL | RTL BEND | 600V/450A | 178V/249A | . 720 | 233V/326A |
| RTL | RTL VB | 150V/350A | 32V/214A | . 150 | $44 \mathrm{~V} / 294 \mathrm{~A}$ |
| RTL | RTL QUAD | 300V/200A | 155V/96A | 1.614 | N/A |
| RTL | RTL Q04A | 20V/125A | 12V/91A | . 130 | 15V/196 |
| RTL | RTL QF4A | 20V/125A | 11V/82A | .130 | 14V/104A |
| RTL | RTL Q04B | 20V/125A | 1V/8A | .130 | 1.3V/10A |
| RTL | RTL QF4B | 20V/125A | 8V/58A | .130 | 10V/74A |
| RTL | V842-vB0 | $95 \mathrm{~V} / 220 \mathrm{~A}$ | NA |  | 1288A |

6.5.3.1 Septa: The ring injection and extraction septa are of the current sheet type and can operate, in the high vacuum environment of the ring, since there are no vacuum windows. The pre-septum is 3 mm thick and carries 2400 Amps to produce 3 kG in a 1 cm gap. The post-septa are of two-turn construction to produce 8 kG in a 0.7 cm gap. Total deflection is $\approx 10^{\circ}$.
6.5.3.2 Kickers: The injection and extraction kickers are 33 cm long, 2.2 cm gap, ferrite loaded (Ferrox-cube 4C4) transmission lines of characteristic impedance $16.6 \Omega$. They transmit their field through ceramic vacuum tubes, the inside of which have been coated with a conducting layer ( $1.2 \Omega / \square$ of Kovar) to carry the high image currents of the beam. The time to establish peak fields of 1 kG is about 45 nsec , short enough to not disturb the second, circulating bunch. This time is a combination of the natural rise time of the pulser, the filling time of the ferrite and the shift introduced by eddy currents in the coating of the vacuum chamber. The 40 kV pulser (see Fig. 6.5.3.2.1)'consists of a resonantly-charged Blumlein discharged by means of a thyratron firing at rates up to 180 pps. Since the launch angle and position into the linac, which controls wake field effects, is related to the angle of the extraction kick, the pulse-to-pulse amplitude jitter must be held to better than $0.03 \%$. This tolerance has been met!

Figure 6.5.3.2.1 Kicker Magnet Pulser

Ferrite Loaded Magnet


## Inventory of Components

With the exception of those items common to both beams in the linac and the fact that the positron ring will not require clearing electrodes, the construction of the second ring and its transport will entail the duplication of the components listed in the following Tables:

| Table 6.5.4.1 Magnet and Power Supply Inventory |  |  |  |  |  |  |  |
| :--- | ---: | :---: | :---: | :---: | :---: | :---: | :---: |
| LOCATION |  | RING | LINAC TO | RING TO | LINAC | TOTAL | NO. OF |
| TYPE |  | RING | LINAC |  |  | POWER |  |
| BANDS | (dc) | 40 | 24 | 16 | 1 | $81^{\prime}$ | 3 |
| QUADS | (dc) | 50 | 45 | 23 | 1 | 119 | 8 |
| SEXTUPOLES | (dc) | 4 | 0 | 16 | 0 | 20 | 6 |
| KICKERS | (100 | ns) | 2 | 0 | 0 | 0 | 2 |
| PULSED MAGN. (8 | MS) | 0 | 0 | 0 | 1 | 1 | 2 |
| SEPTA |  | (dc) | 4 | 0 | 0 | 0 | 4 |
| VERTICAL | BENDS | (dc) | 0 | 3 | 4 | 1 | 8 |
| MAGNET | TOTALS | 100 | 72 | 59 | 4 | 235 | 23 |
| BIPOLAR | CORRECTORS | 0 | 25 | 35 | 0 | 60 | 64 |
| UNIPOLAR CORRECTORS | 40 | 24 | 16 | 1 | 80 | 64 |  |
| CORRECTOR | TOTALS | 40 | 49 | 52 | 1 | 140 | 128 |

Table 6.5.4.2 RF, Vacuum, I\&C Component Inventories

| LOCATION |  | LINAC TO | RING TO |  |  |
| :--- | :---: | :---: | :---: | :---: | :---: |
|  | RING | RING | LINAC | LINAC | TOTALS |
| RF Transmitter 50 kW | 1 | 0 | 0 | 0 | 0 |
| RF Cavities 714 MHz | 2 | 0 | 0 | 0 |  |
| S. Band Acc. Section |  |  | $10 \mathrm{ft} /$ | 0 | $10 /$ |
| \& Waveguide | 0 | 0 | 200 ft | 0 | 200 |
| Low Level Instrumentation | 1 | 0 | 1 | 0 | 2 |
| Length of Vac. Chamber | 111 ft | 229 ft | 180 ft | 0 | 520 |
| Number of Pumps | 40 DIP | 9 HIP | 9 HIP | 0 | 78 |
|  | 20 HIP | 9 HIP | 9 HIP |  |  |
| Number of Ion Gauges | 4 | 3 | 3 | 0 | 10 |
| Number of Valves | 0 | 3 | 3 | 0 | 6 |
| Beam Position Monitors | 26 | 19 | 22 | 0 | 67 |
| Clearing Electrodes | 28 | 0 | 0 | 0 | 28 |
| Moveable Scrapers, H.V. | 1 | 0 | 0 | 0 | 1 |
| Profile Monitors | 0 | 4 | 5 | 1 | 9 |
| Spectrum Monitors | 0 | 1 | 1 | 0 | 2 |
| Toroids | 0 | 3 | 3 | 1 | 7 |
| Gap Monitores | 0 | 3 | 3 | 0 | 6 |
| Synchrotron Light Monitor | 1 | 0 | 0 | 0 | 0 |
| Faraday Cups | 0 | 1 | 1 | 0 | 2 |

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