Microwave Cold-Testing Techniques for the NLC^{*}

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Abstract

The R & D program for the Next Linear Collider Test Accelerator (NLCTA) includes the development of microwave techniques for testing Xband accelerating structures at different stages of design, manufacturing and assembly. During the design phase, short stacks were built and tested to finalize dimensions. Cell by cell measurements were performed on the NLCTA injector cells as a microwave quality control (QC) after manufacturing. The two injector sections were tuned using a moveable plunger and tested using a semiautomated system of bead-pull. Using this perturbation technique, we were able to map the amplitude and phase of the electric field on the axis throughout the assembled structures under traveling-wave conditions.

1 INTRODUCTION

The development of high-gradient X-band accelerator sections for the Next Linear Collider (NLC), requires low power rf testing techniques that ensure the precise and reliable performance of these sections. For this reason, the NLC Test Accelerator (NLCTA) project [1] provides for an R & D program to develop the low-power testing techniques. These characterizing techniques are aimed at supporting different development phases of the accelerator sections; from the design, through manufacturing, and finally testing of the complete accelerator section. In Fig. 1 we summarize the series of tests that can be used through the development phases of an accelerator section.

In this paper, we will outline some of the above testing techniques. Results of measurements performed on the two NLCTA injector sections are presented as examples of the implementation of these testing techniques.

2 MICROWAVE QC

In the Detuned Structure (DS) used for the NLCTA injector sections, the diameter of each cell (2b), the thickness of its iris (t) and the diameter of its aperture (2a), all vary progressively from cell to cell. The goal is to detune the dipole modes and prevent short range cumulative build up of wake fields, while maintaining quasi-constant gradient characteristics of the fundamental accelerating mode.



To verify the regularity of machining of successive cells we measured the resonant frequency of the individual cells in the configuration shown in Fig. 2. The results of the measurement for the 204 cells of the two injector sections are plotted in Fig. 3. By plotting the inverse of the measured frequency versus the cell radius b, one obtains a straight line fit which is a convenient way of verifying the (2b) dimension.



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3 TUNING AND TESTING INJECTOR SECTIONS

The use of brazing in the assembly of an accelerator section can result in deviations in the volume of the constituent cavities (cells). Thus it was required to tune individual cells in the injector sections. A stainless steel plunger was progressively pulled through the structure from the center of one cell to the next. The resulting phase shift was measured on a network analyzer at a frequency corresponding to the operating frequency 11.424 GHz and corrected for room temperature and dry N_2 under atmospheric pressure. In the two injector sections we tuned, the majority of the cells needed to be tuned by pulling their cylindrical walls outward to make up for the volume of the brazing fillet. The completed structure was finally checked using the bead-pull technique.



A. TW Perturbation Technique:

Mallory and Miller [2] were the first to use the bead-pull technique under traveling-wave conditions to measure the phase and field strength in accelerator structures. Steele [3] formulated the theoretical basis for the technique. We have built a semi-automated system [4] based on this technique to evaluate the tuning of X-band accelerators. Fig. 4 shows a schematic for our bead-pull system. A small metallic cylindrical bead is attached to a thin nylon string along the axis of the vertically-mounted accelerator section. The position of the bead is determined by the position of the carriage moving on a lead screw driven by a stepper motor. The computer which controls the motor also collects and processes the reflection coefficient data (S_{11}) from the microwave network analyzer. The field perturbation introduced by the bead located at a





position z can be modeled as a reactive discontinuity in a transmission line. The susceptance associated with this discontinuity is proportional to the energy perturbation; and hence to the square of the field at this location. For a small needle-like bead located on the axis of the accelerator structure the field that is perturbed is the electric field of the $2\pi/3$ mode.

$$E_{z}(z) = |E(z)|e^{-j\Theta_{E}(z)}$$

The difference between the unperturbed input reflection $S_{11}^{\ \mu}$ and the perturbed reflection $S_{11}^{\ p}$ is then proportional to the square of the field as follows:

$$\Delta S_{11}(z) = S_{11}^{p}(z) - S_{11}^{u}(z) = c |E(z)|^{2} e^{-2j\Theta_{E}(z)}$$
(1)
where c is a constant which depends on the input

where c is a constant which depends on the input power and the bead characteristics (material, shape, and volume).

B. Results of Characterization of NLCTA Intector Sections

Using the bead-pull system shown schematically in Fig. 4, we tested both sections of the NLCTA injector. We used a stainless steel cylindrical bead of diameter $d_b=0.625$ mm and a length $l_b=1$ mm supported on a monofiliment nylon string of diameter $d_s = 0.145$ mm. The bead was moved along the structure in steps of 0.5mm (approximately 17 steps per cell). At each bead position, the S₁₁ value used was the average of 256 measurements. The measurement time for the 90-cm structure was 90 minutes. The measurements were made at frequency corrected for the nylon string, the room temperature and dry N₂. The results of these measurements on both injector sections are shown in Fig. 5.

The three-fold symmetry of the accelerating $2\pi/3$ mode is manifested clearly. The variation in the reflection from the bead as it is pulled along the axis of each structure traces the hypotrochoid function. This

trace can be predicted theoretically for a given longitudinal electric field, Ez based on its space harmonics. Our previous work, [4] showed good agreement between theory and experiment. It is to be noted that the data shown in Fig. 5 represents the measured reflection after centering by subtracting a phasor corresponding to the reflection from the unperturbed structure, S_{11}^{u} according to Eq. (1). The amount of centering ($|S_{11}^{u}|=0.02$) for both sections was consistent with the VSWR=1.04 measured for both of them.

The electric field Ez can be obtained from the measured reflection (after centering) using Eq. (1). Ez for the second injector section is plotted as (a phasor) in Fig. 6. In Fig. 7 we show phase variations of Ez for the same the same section for the first eleven cells.



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