

## Spin Physics Experiments at SLAC

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### Abstract

Some recent results on the  $g_1$  and  $g_2$  spin structure functions of the proton and neutron are presented. New data on the inclusive photoproduction of negative hadrons from transversely polarized protons and deuterons are shown for the first time. Plans for future experiments using polarized photon beams are discussed.

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# 1 Introduction

During the conference talk, I presented recent SLAC results for the  $g_1$  and  $g_2$  spin structure functions, and related sum rules, as well as preliminary results for the spin asymmetry in inclusive hadron photoproduction. Because the  $g_1$  and  $g_2$  results are published (or soon will be), I will focus on the photoproduction results, hoping to stimulate some theoretical interest in the interpretation.

The second half of my talk concerned three future experiments at SLAC, all involving polarized photon beams. The physics prospects from these approved experiments is discussed in the second part of this paper.

## 2 Results on $g_1$

The final results from SLAC E155 have now been published[1]. The structure functions  $g_1^p$  and  $g_1^n$  were measured in a single experimental setup over the large kinematic range  $0.014 < x < 0.9$  and  $1 < Q^2 < 40 \text{ GeV}^2$  using deep-inelastic scattering of 48 GeV longitudinally polarized electrons from polarized protons and deuterons. The data indicate that the  $Q^2$  dependence of  $g_1^p$  ( $g_1^n$ ) at fixed  $x$  is very similar to that of the spin-averaged structure function  $F_1^p$  ( $F_1^n$ ). Simple empirical fits to the data are given by

$$\frac{g_1^p}{F_1^p} = x^{0.700}(0.817 + 1.014x - 1.489x^2)(1 - \frac{0.04}{Q^2}) \quad (1)$$

$$\frac{g_1^n}{F_1^n} = x^{-0.335}(-0.013 - 0.330x + 0.761x^2)(1 + \frac{0.13}{Q^2}). \quad (2)$$

From an NLO QCD fit to all available data, E155 finds that the difference of first moments  $\Gamma_1^p - \Gamma_1^n = 0.176 \pm 0.003 \pm 0.007$  at  $Q^2 = 5 \text{ GeV}^2$ , in agreement with the Bjorken sum rule prediction of  $0.182 \pm 0.005$ . Using the same NLO pQCD fit, the quark singlet contribution in the  $\bar{M}S$  scheme is  $\Delta\Sigma = 0.23 \pm 0.04(\text{stat}) \pm 0.06(\text{syst})$  at  $Q^2 = 5 \text{ GeV}^2$ , confirming earlier indications that quarks carry only a small fraction of the spin of the nucleon.

### 3 The $g_2$ structure function

The recent (1999) experiment SLAC E155x made the best measurements of  $g_2$  for the proton and deuteron to date. The final results should be published soon; the results presented here are still preliminary. We used the 120 Hz SLAC electron beam with a longitudinal polarization of  $(83 \pm 3)\%$  at energies of 29.1 and 32.3 GeV and a typical current of 25 nA. We used transversely polarized  $\text{NH}_3$  and  $^6\text{LiD}$  targets as sources of polarized protons (average polarization 75%) and deuterons (average polarization 20%). Scattered electrons were detected in three independent spectrometers centered at  $2.75^\circ$ ,  $5.5^\circ$ , and  $10.5^\circ$ . Electrons in each spectrometer were separated from pions using gas Cherenkov counters and segmented electromagnetic calorimeters. Tracking was done with scintillator hodoscopes.

Since the results for  $g_2$  in the three spectrometers and two beam energies are reasonably consistent with the  $Q^2$  dependence of the twist-two  $g_2^{WW}$  model, they are averaged together using this assumption to produce the averaged values shown in Fig. 1.

The proton results are clearly different than zero, and exhibit an  $x$ -dependence similar to that of the  $g_2^{WW}$  model. There appear to be statistically significant differences from the  $g_2^{WW}$  model, possibly indicating non-zero twist-3 contributions. The data are in qualitative agreement with the bag model calculation of Stratmann [4] and the chiral soliton calculation of Gamberg and Weigel [5], but are considerably more negative than the model of Song [6]. The deuteron data have larger errors than the proton data, but also indicate significantly negative values at high  $x$ , and are in qualitative agreements with  $g_2^{WW}$ , Stratmann [4], and Gamberg and Weigel [5].

### 4 Inclusive Photoproduction of Hadrons

The double spin asymmetry for inclusive photoproduction of hadrons by longitudinally polarized electrons impinging on longitudinally or transversely polarized targets was measured in E155 and E155x respectively. Because the scattered electron was not detected, the production is determined by scattering at very small angles, and the beam essentially consists of circularly polarized real photons. The E155 results have been published[8], and show small but significant double spin asymmetries for a proton target, with larger

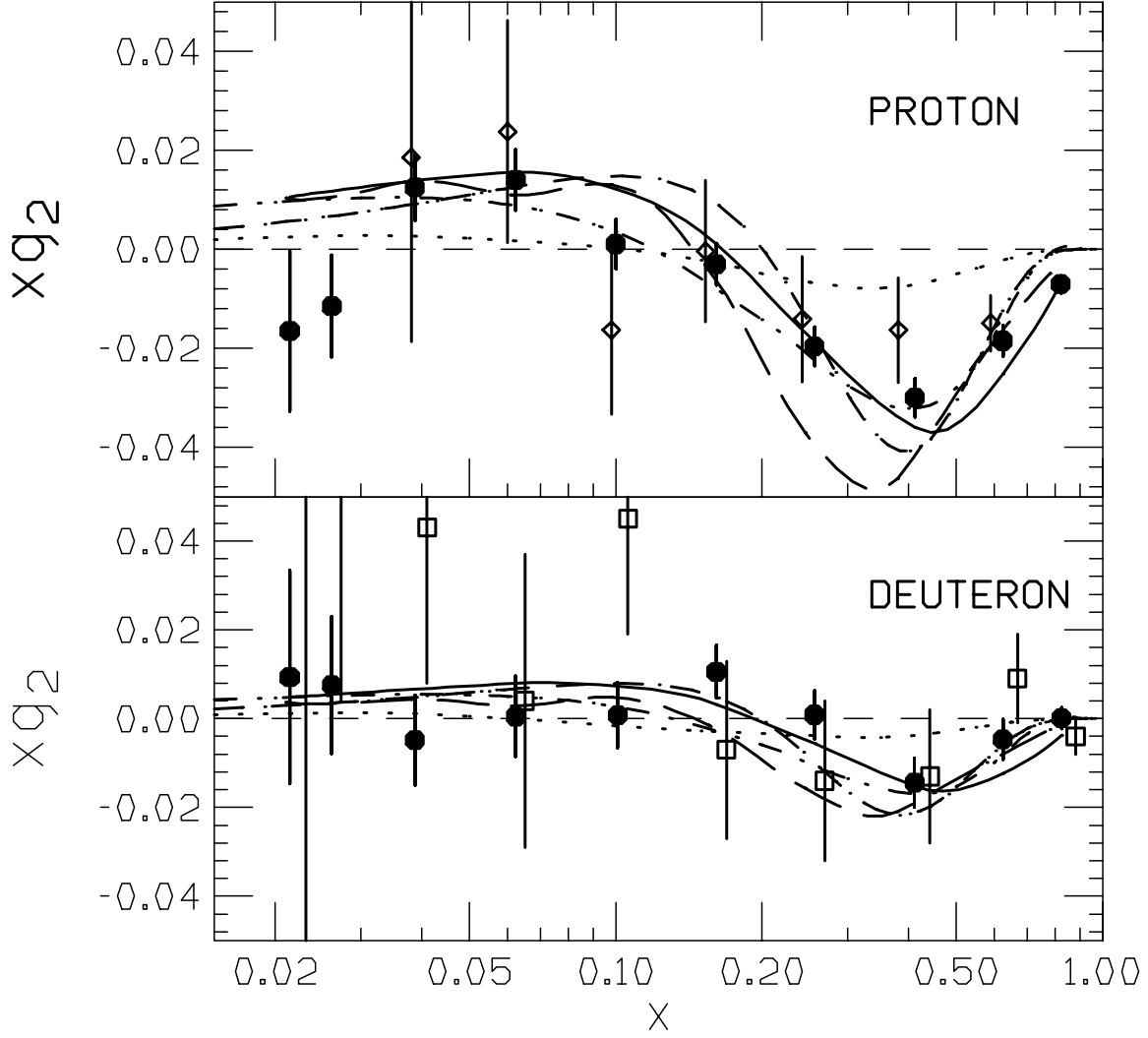


Figure 1: The structure function  $xg_2$  averaged over the three spectrometers for E155x (solid circles, preliminary), and data from E143 [2] (diamonds) and E155 [3] (stars). The errors are statistical; the systematic errors are negligible. Also shown is  $g_2^{WW}$  at the average  $Q^2$  of this experiment at each value of  $x$  (solid curves) and the calculations of Stratmann [4] (dot-dash), Gamberg and Weigel [5] (short dash), Song [6] (dot), and Wakamatsu[7] (long dash).

asymmetries observed for positive hadrons than for negative hadrons. The asymmetries for the deuteron target are consistent with zero. The data were compared[9] to the photoproduction limit of a Regge model fit[10] to the world data on  $g_1$ , assuming that the asymmetry of low transverse momentum hadrons at forward angles is representative of the total cross section asymmetry. Reasonably good agreement was found for both the proton and deuteron. This may have interesting implications for the convergence of the GDH sum rule.

The preliminary data for negative hadrons from transversely polarized protons from E155x are presented in Fig. 2. The positive hadron asymmetry was not measured in E155x. The raw experimental asymmetries have been corrected for target polarization, beam polarization, and dilution factor (fraction of target nucleons that are polarized), but no average bremsstrahlung de-polarization correction has been applied (nor was it for the E155 results). It can be seen that the proton target asymmetries are clearly non-zero for momenta below 15 GeV, especially at 2.75 degrees. The deuteron asymmetries are consistent with zero, with larger errors than the proton data due the lower average target polarization. The magnitude of the proton asymmetries are fairly similar to those observed with the longitudinally polarized target, at about 0.005. This is somewhat surprising, since a) the target polarization component parallel to the photon direction was only about 4%; b) there was no coincident electron to define a direction for a single spin asymmetry. The only direction is the plane defined by the photon and target polarizations (and our measurements are in this plane). Parity violating effects are expected to be negligible due to frequent target polarization direction reversals. I do not know of any theoretical calculations for this double spin asymmetry. I hope these data will inspire such calculations to be carried out.

## 5 Future: $\Delta\sigma^{\gamma N}(k)$ and the High Energy Contribution to the GDH Sum Rule

An experiment (E159[11]) has recently been approved at SLAC to measure  $\Delta\sigma^{\gamma N}(k)$ , the helicity-dependent total photo-absorption cross section, for photon energies  $5 < k < 40$  GeV, on both proton and deuteron targets. Our first goal is to complement our extensive set of measurements of  $g_1$  at

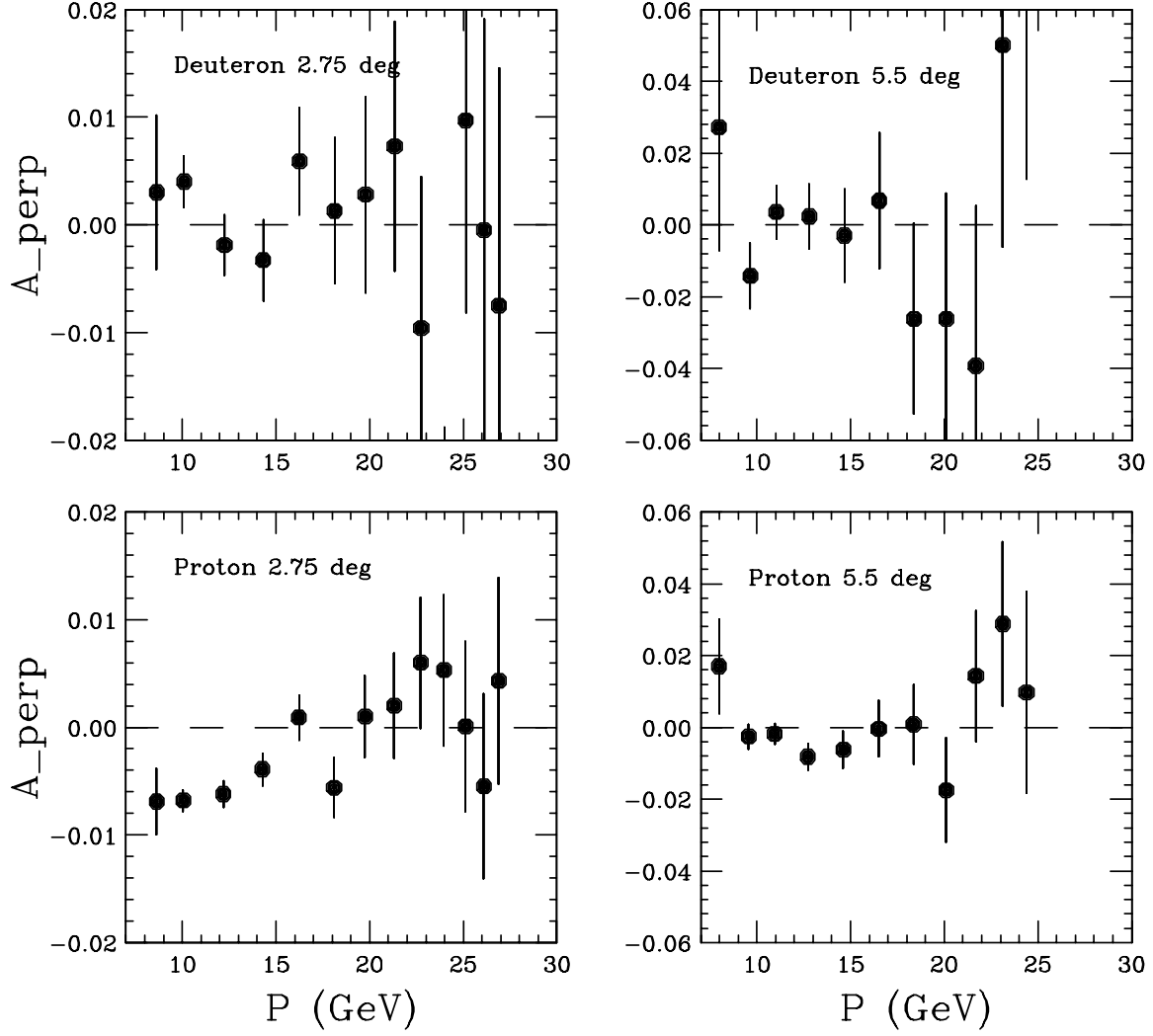


Figure 2: The double spin asymmetry for inclusive negative hadron photo-production from transversely polarized protons and deuterons as a function of hadron momentum at scattering angles of 2.75 and 5.5 degrees, averaged over beam energies of 29 and 32 GeV.

$Q^2 > 0$  with the anchor points at  $Q^2 = 0$ , useful for global fitting[10, 13] and understanding the low- $x$  behavior. Our second goal is to test the convergence of the GDH sum rule [12],

$$\int_{k_\pi}^{\infty} \frac{dk}{k} \Delta\sigma^{\gamma N}(k) = \frac{2\pi^2\alpha\kappa^2}{M^2} \quad (3)$$

where  $M$  and  $\kappa$  are the nucleon mass and anomalous magnetic moment, and  $k_\pi$  is the threshold energy needed to produce at least one pion. Early indications from measurements in the resonance region are that the sum rule may already be over-saturated, requiring a sign change to  $\Delta\sigma^{\gamma N}(k)$  for convergence.

We will use an untagged coherent bremsstrahlung beam to create a high flux of circularly polarized photons. With coherent bremsstrahlung, a set of high intensity spikes is generated by proper orientation of a diamond crystal radiator. With longitudinally polarized electrons, the incoherent bremsstrahlung photons are circularly polarized, with the polarization maximal at the endpoint. The coherent photons are elliptically polarized: the circular component is almost identical to that for incoherent photons. The coherent peak polarization also has a linear component which will cancel in the measurement of  $\Delta\sigma^{\gamma N}(k)$ , but will allow for the measurement of possibly interesting azimuthal asymmetries.

For targets, we will use polarized  $\text{NH}_3$  and  $\text{ND}_3$  as sources of polarized protons and neutrons. Polarized deuterons to first order allow measurements of the isovector combination  $(n+p)/2$ , with small corrections for the deuteron  $D$ -state, shadowing, and nuclear coherent hadron production. An extension to this proposal could use a polarized  $^3\text{He}$  target to verify the consistency of  $\Delta\sigma^{\gamma n}(k)$  for the neutron as extracted from either deuterium or  $^3\text{He}$ . The detector is a simple calorimeter optimized to measure  $> 98\%$  of all hadronic interactions, and to reject electromagnetic backgrounds.

The expected errors are shown in Fig. 3 for both the proton and neutron, and for two data taking modes, one involving counting each hadronic interaction individually, and one where only the total flux of hadrons for each helicity state is measured. Even with the larger counting mode statistical errors, a very good determination can be made of both the magnitude and energy dependence of  $\Delta\sigma^{\gamma N}(k)$  for  $5 < k < 40$  GeV. By measuring with both proton and deuteron targets, the high energy contributions to both the isovector and isoscalar GDH sum rules can be determined. This will allow

tests of Regge-inspired models, which predict very different behavior for the isovector and isoscalar contributions, and will provide a baseline for studies of the polarized spin-structure functions measured with virtual photons.

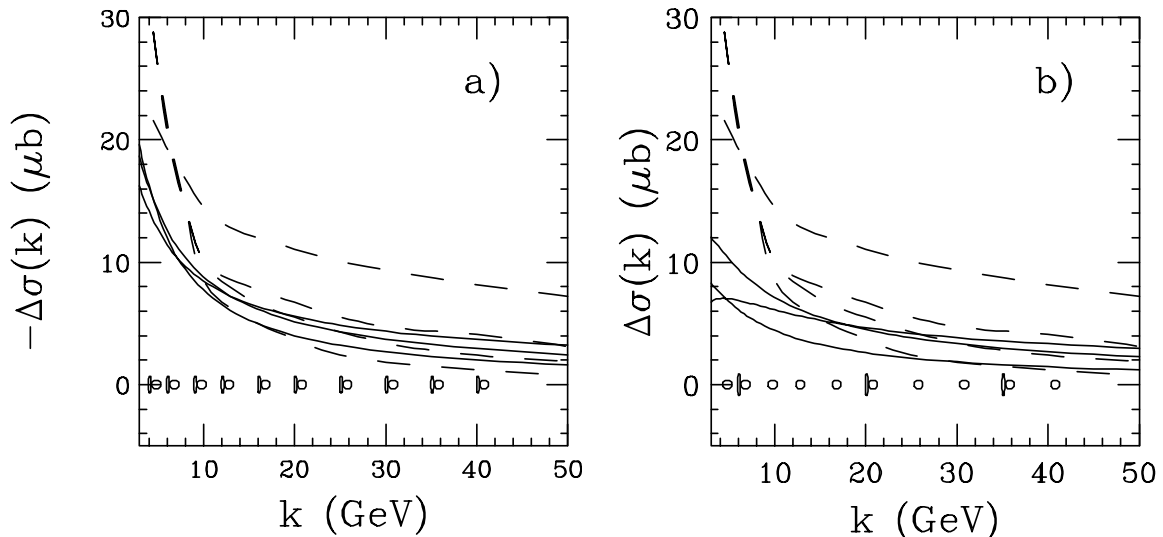


Figure 3: a) Projected proton error bars for  $\Delta\sigma^p(k)$  for E159 as a function of photon energy for the counting mode (rectangles) and flux integration mode (circles). The dashed curves are representative models from Ref. [13], the solid curves are from Ref. [10]; b) same but for the neutron as measured with  $\text{ND}_3$ .

## 6 Future: Polarized Charm Photoproduction and the Gluon Spin

Another approved SLAC experiment using the circularly polarized photon facility mentioned above is E161[14], designed to study the gluon spin structure of nucleons using open charm photoproduction. The measurements will utilize a  $^6\text{LiD}$  polarized target to measure the asymmetry  $A_{cc}$  in open charm photoproduction. This process is dominated by the photon-gluon fusion mechanism. The open charm signal will be measured by detecting the muons from charm decay at large  $p_T$ . This experiment will measure the asymme-



try  $A_{cc}$  over a range of energies and  $p_T$  sensitive to  $x$  from 0.1 to 0.3 with statistical precision of about 0.01. This is to be compared with the range of current theoretical models in which the values of  $A_{cc}$  differ by more than 0.1 and  $x\Delta g(x)$  differ by up to 0.3 in this  $x$  range.

Figure 4 shows the expected statistical error on  $A_{cc}$  as a function of  $p_T^\mu$  for  $5 < P_\mu < 10$  GeV for three incident photon coherent peak energies. The points are arbitrarily plotted at a value of zero. Also shown are the calculated asymmetries from a sample of gluon polarization models. The systematic errors of 10% of the value of the asymmetry (typical error 0.01) will be highly correlated point-to-point. There will be additional data for higher momentum muons. Our statistical errors are projected to be significantly smaller than the similar COMPASS experiment[15] at CERN, but our lower photon energies correspond to larger values of  $x$  for the gluons.

The E161 experiment will also measure the double spin asymmetry for elastic and inelastic photoproduction of closed charm ( $J/\psi$  particles). The latter may also yield interesting information on the gluon spin, if the relative contributions from color singlet and color octet mechanisms can be reliably modeled. Under study is the possibility of looking at the photon-gluon fusion process producing a pair of strange quarks, tagged by their decay to muons.

## 7 Future: Linear Polarization Asymmetry in Charm Photoproduction

A third planned experiment at SLAC, E160[16], will use unpolarized electrons to make coherent bremsstrahlung beams at 15, 25, and 35 GeV. These photons will have a fairly high degree of linear polarization. While the main goal of the experiment is to measure the  $A$ -dependence of  $J/\psi$  and  $\psi'$  quasi-elastic photoproduction, we will also measure the linear polarization single-spin asymmetry for nuclear coherent, quasi-elastic, and inelastic  $J/\psi$  photoproduction “for free”. At present, I do not know of any predictions for these asymmetries. However, there is a QCD prediction[17] for open charm photoproduction at our photon energies: the single-spin asymmetry is predicted to be large, at about 0.2, and unlike the cross section itself, is quite stable against higher order QCD corrections. We can identify open charm events in our muon spectrometer by single muons with transverse momenta near 1

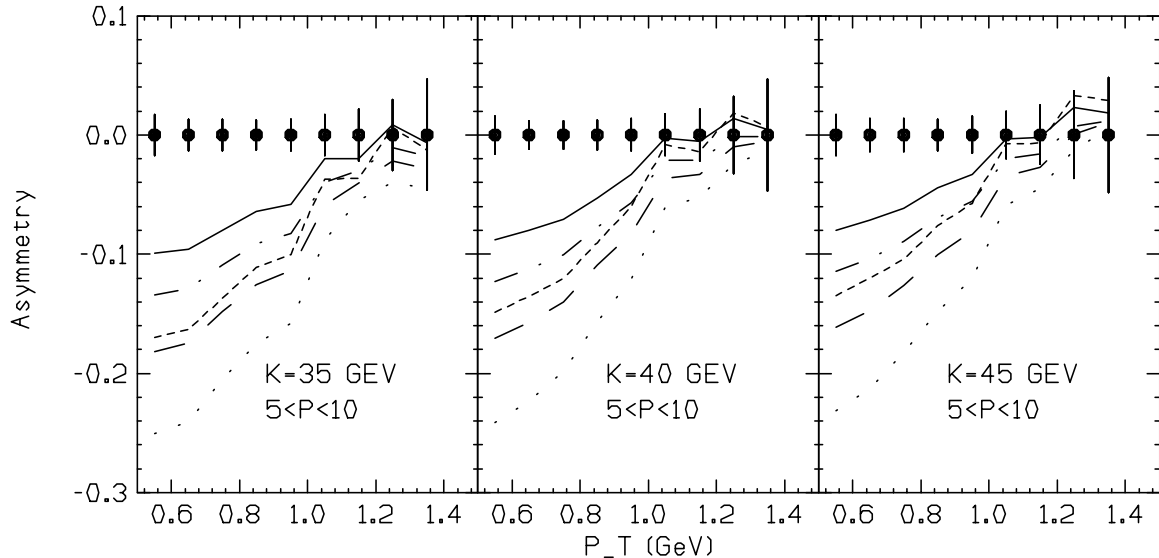


Figure 4: The E161 projected statistical errors on the asymmetry  $A_{cc}$  for open charm production as a function of  $p_T^\mu$  of the detected muon for  $5 < P_\mu < 10$  GeV. Also shown are asymmetries for several gluon polarization models.

GeV, where the backgrounds from  $\pi$ ,  $K$ , and  $J/\psi$  decays are the smallest, or as two like-sign muons. Preliminary estimates are that we can measure the linear polarization asymmetry with a statistical error of about 0.02 or better, which would provide a meaningful test of the prediction. It would be very nice if calculations could be done for the closed charm case: this may be a good way to learn about the relative importance of color singlet and color octet mechanisms.

## 8 Summary

The recent SLAC data on  $g_1$  and  $g_2$  have provided significant new information of the spin structure of the nucleon. The inclusive hadron photoproduction data from these experiments could also prove to be valuable, with more theoretical attention. Future experiments using polarized photon beams should provide insight into the role of gluon polarization in the nucleon, and the behavior of the spin structure functions in the limit of  $Q^2 \rightarrow 0$ .

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