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Shifts in the Properties of the Higgs Boson from Radion Mixing *

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Abstract

We examine how mixing between the Standard Model(SM) Higgs boson, h, and the radion of the Randall-Sundrum model modifies the expected properties of the Higgs boson. In particular we demonstrate that the total and partial decay widths of the Higgs, as well as the $h \to gg$ branching fraction, can be substantially altered from their SM expectations, while the remaining branching fractions are modified less than $\lesssim 5\%$ for most of the parameter space volume.

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The Randall-Sundrum(RS) model[1] offers a potential solution to the hierarchy problem that can be tested at present and future accelerators[2]. In this model the SM fields lie on one of two branes that are embedded in 5-dimensional AdS space described by the metric $ds^2 = e^{-2k|y|}\eta_{\mu\nu}dx^\mu dx^\nu - dy^2$, where k is the 5-d curvature parameter of order the Planck scale. To solve the hierarchy problem the separation between the two branes, r_c, must have a value of $kr_c \sim 11-12$. That this quantity can be stabilized and made natural has been demonstrated by a number of authors[3] and leads directly to the existence of a radion(r), which corresponds to a quantum excitation of the brane separation. It can be shown that the radion couples to the trace of the stress-energy tensor with a strength Λ of order the TeV scale, *i.e.*, $\mathcal{L}_{eff} = -r\ T^\mu_\mu/\Lambda$. (Note that $\Lambda = \sqrt{3}\Lambda_\pi$ in the notation of Ref.[2].) This leads to gauge and matter couplings for the radion that are qualitatively similar to those of the SM Higgs boson. The radion mass (m_r) is expected to be significantly below the scale Λ implying that the radion may be the lightest new field predicted by the RS model. One may expect on general grounds that this mass should lie in the range of a few $\times 10$ GeV $\leq m_r \leq \Lambda$. The phenomenology of the RS radion has been examined by a number of authors[4] and in particular has been recently reviewed by Kribs[5].

On general grounds of covariance, the radion may mix with the SM Higgs field on the TeV brane through an interaction term of the form

$$S_{rH} = -\xi \int d^4x \sqrt{-g_w} R^{(4)}[g_w] H^\dagger H\,, \tag{1}$$

where H is the Higgs doublet field, $R^{(4)}[g_w]$ is the Ricci scalar constructed out of the induced metric g_w on the SM brane, and ξ is a dimensionless mixing parameter assumed to be of order unity and with unknown sign. The above action induces kinetic mixing between the 'weak eigenstate' r_0 and h_0 fields which can be removed through a set of field redefinitions and rotations. Clearly, since the radion and Higgs boson couplings to other SM fields differ this mixing will induce modifications in the usual SM expectations for the Higgs decay widths and branching fractions.

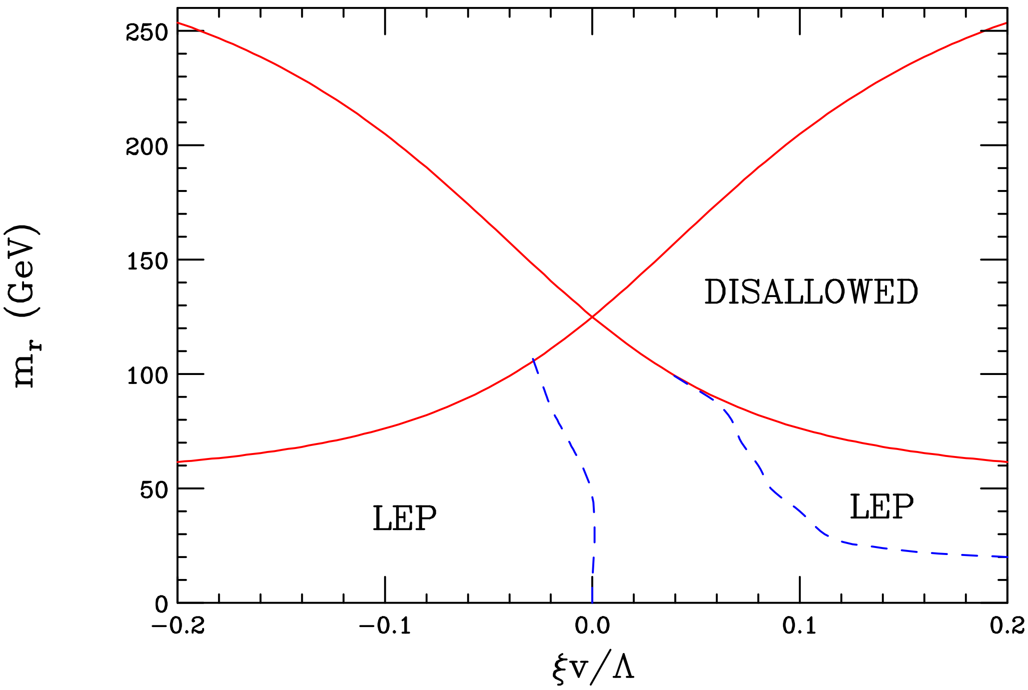

Figure 1: Constraint on the mass of the radion assuming $m_h = 125$ GeV as a function of the product $\xi v/\Lambda$ as described in the text. The disallowed region lies between the solid curves. The region excluded by LEP searches assuming $v/\Lambda = 0.2$ is also shown and are labeled by 'LEP'.

To make unique predictions in this scenario we need to specify four parameters: the masses of the *physical* Higgs and radion fields, $m_{h,r}$, the mixing parameter ξ and the ratio v/Λ, where v is the vacuum expectation value of the SM Higgs $\simeq 246$ GeV. Clearly the ratio v/Λ cannot be too large as Λ_π is already bounded from below by collider and electroweak precision data[2]; for definiteness we will take $v/\Lambda \leq 0.2$ and $-1 \leq \xi \leq 1$ in what follows although larger absolute values of ξ and wider ranges of v/Λ have been entertained in the literature. The values of the two physical masses themselves are not arbitrary. When we require that the weak basis mass-squared parameters of the radion and Higgs fields be real, as is required by hermiticity, we obtain an additional constraint on the ratio of the physical radion and Higgs masses which only depends on the product $|\xi|\frac{v}{\Lambda}$. Explicitly one finds that either $\frac{m_r^2}{m_h^2} \geq 1+2\sin^2\rho+2|\sin\rho|\sqrt{1+\sin^2\rho}$ or $\frac{m_r^2}{m_h^2} \leq 1+2\sin^2\rho-2|\sin\rho|\sqrt{1+\sin^2\rho}$ where $\rho = \tan^{-1}(6\xi\frac{v}{\Lambda})$. This disfavors the radion having a mass too close to that of the Higgs when there is significant mixing; the resulting excluded region is shown in Fig. 1. These constraints are somewhat restrictive; if we take $m_h = 125$ GeV and $\xi\frac{v}{\Lambda} = 0.1(0.2)$ we find that either $m_r > 205(254)$ GeV or $m_r < 76(61)$ GeV. This lower mass range for the radion is somewhat disfavored by direct LEP searches as also can be seen from Fig. 1. Here we have assumed that $v/\Lambda = 0.2$ and converted the LEP Higgs search bounds[6] into ones for the radion using the appropriate set of rescaling factors. While most of the region is excluded for this values values of the parameters, the parameter space for a light radion is certainly not closed. Furthermore, as we decrease the assumed value of v/Λ, the size of the allowed region grows since the radions couplings to the Z are rapidly decreasing. One may argue, however, that a somewhat more massive radion is more likely than one in the remaining allowed region in the lower part of Fig. 1.

Let us now turn our attention to the properties of the Higgs boson in this model. Following the notation of Giudice *et al.*[4], the coupling of the physical Higgs to the SM fermions and massive gauge bosons $V = W, Z$ is now given by

$$\mathcal{L} = \frac{-1}{v}(m_f\bar{f}f - m_V^2 V_\mu V^\mu)[\cos\rho\cos\theta + \frac{v}{\Lambda}(\sin\theta - \sin\rho\cos\theta)]h\,, \tag{2}$$

where the angle ρ is given above and θ can be calculated in terms of the parameters ξ and v/Λ and the physical Higgs and radion masses. Denoting the combinations $\alpha = \cos\rho\cos\theta$ and $\beta = \sin\theta - \sin\rho\cos\theta$, the corresponding Higgs coupling to gluons can be written as $c_g\frac{\alpha_s}{8\pi}G_{\mu\nu}G^{\mu\nu}h$ with $c_g = \frac{-1}{2v}[(\alpha + \frac{v}{\Lambda}\beta)F_g - 2b_3\beta\frac{v}{\Lambda}]$ where $b_3 = 7$ is the $SU(3)$ β-function and F_g is a well-known kinematic function of the ratio of masses of the top quark to the

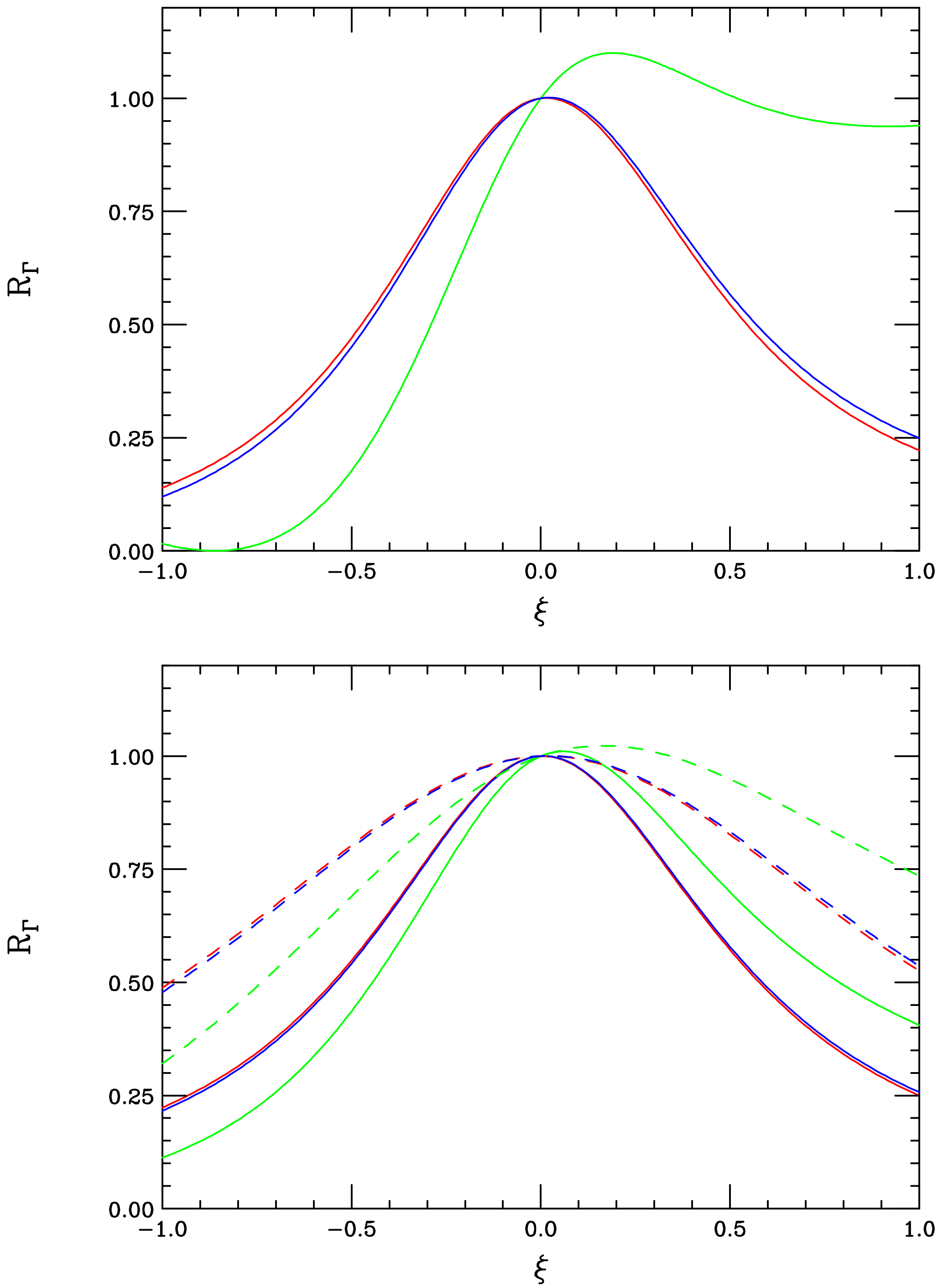

Figure 2: Ratio of Higgs widths to their SM values, R_Γ, as a function of ξ assuming a physical Higgs mass of 125 GeV: red for fermion pairs or massive gauge boson pairs, green for gluons and blue for photons. In the top panel we assume $m_r = 300$ GeV and $v/\Lambda = 0.2$. In the bottom panel the solid(dashed) curves are for $m_r = 500(300)$ GeV and $v/\Lambda = 0.2(0.1)$.

physical Higgs. Similarly the physical Higgs couplings to two photons is now given by $c_\gamma \frac{\alpha_{em}}{8\pi} F_{\mu\nu}F^{\mu\nu}h$ where $c_\gamma = \frac{1}{v}[(b_2+b_Y)\beta\frac{v}{\Lambda} - (\alpha+\frac{v}{\Lambda}\beta)F_\gamma]$, where $b_2 = 19/6$ and $b_Y = -41/6$ are the $SU(2)\times U(1)$ β-functions and F_γ is another well-known kinematic function of the ratios of the W and top masses to the physical Higgs mass. (Note that in the simultaneous limits $\alpha \to 1,\ \beta \to 0$ we recover the usual SM results.) From these expressions we can now compute the change of the various decay widths and branching fractions of the SM Higgs due to mixing with the radion.

Fig. 2 shows the ratio of the various Higgs widths in comparison to their SM expectations as functions of the parameter ξ assuming that $m_h = 125$ GeV with different values of m_r and $\frac{v}{\Lambda}$. We see several features right away: (i) the shifts in the widths to $\bar{f}f/VV$ and $\gamma\gamma$ final states are very similar; this is due to the relatively large magnitude of F_γ while the combination b_2+b_Y is rather small. (ii) On the otherhand the shift for the gg final state is quite different since F_g is smaller than F_γ and b_3 is quite large. (iii) For relatively light radions with a low value of Λ the Higgs decay width into the gg final state can come close to vanishing due to a strong destructive interference between the two contributions to the amplitude for values of ξ near -1. (iv) Increasing the value of m_r has less of an effect on the width shifts than does a decrease in the ratio $\frac{v}{\Lambda}$.

The deviation from the SM expectations for the various branching fractions, as well as the total width, of the Higgs are displayed in Fig. 3 as a function of the mixing parameter ξ. We see that the gluon branching fraction and the total width may be drastically different than that of the SM. As we will see below the former will affect the Higgs production cross section at the LHC. However, the $\gamma\gamma$, $f\bar{f}$, and VV, where $V = W, Z$ branching fractions receive small corrections to their SM values, of order $\lesssim 5-10\%$ for almost all of the parameter region. Observation of these shifts will require the accurate determination of the Higgs

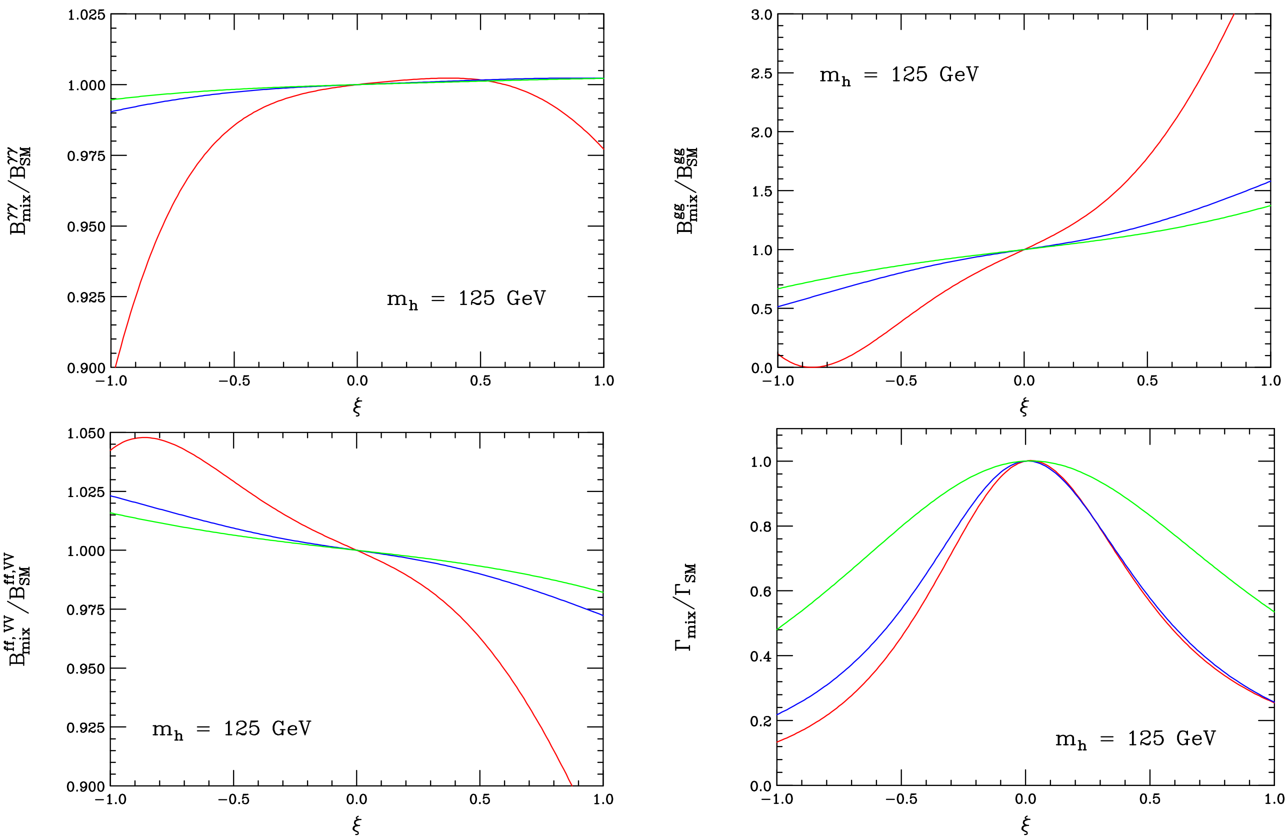

Figure 3: The ratio of the Higgs branching fraction into $\gamma\gamma$, gg, $f\bar{f}$, and VV final states as labeled, as well as for the total width, with radion mixing to that of the SM as a function of ξ. The red (blue; green) curves correspond to the choice $m_r = 300$ GeV, $v/\Lambda = 0.2$ (500, 0.2; 300, 0.1).

branching fractions obtainable at an e^+e^- Linear Collider[7] from which constraints on the radion model parameter space may be extracted. These small changes in the ZZh and $hb\bar{b}$ couplings of the Higgs boson can also lead to small reductions in the Higgs search reach from LEPII. This is shown in Fig. 4 for several sets of parameters; except for extreme cases this reduction in reach is rather modest.

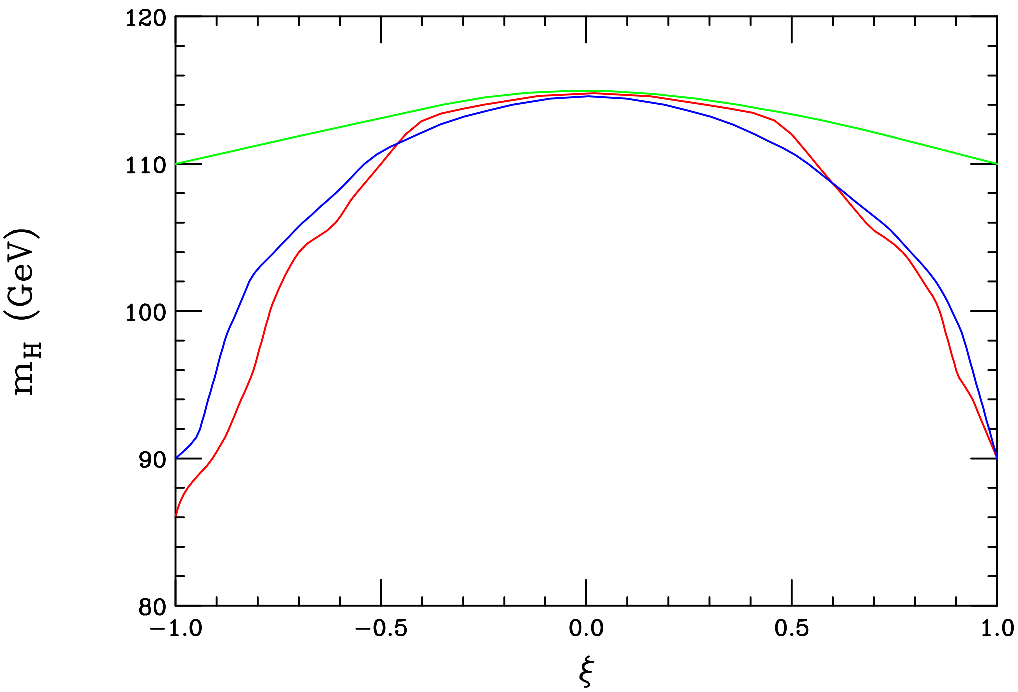

Figure 4: Lower bound on the mass of the Higgs boson from direct searches at LEP as a function of ξ including the effects of mixing. The red (blue; green) curves correspond to the choice $m_r = 300$ GeV, $v/\Lambda = 0.2$ (500, 0.2; 300, 0.1).

At the LHC the dominant production mechanism/signal for the light Higgs boson is via the gluon-gluon fusion through a triangle graph with subsequent decay into $\gamma\gamma$. Both the production cross section and the subsequent $\gamma\gamma$ branching fraction are modified by mixing as shown in Fig. 5. This figure shows that the Higgs production rate in this mode at the LHC is always reduced in comparison to the expectations of the SM due to the effects of mixing. For some values of the parameters this reduction can be by more than an order of magnitude which could seriously hinder Higgs discovery via this channel at the LHC.

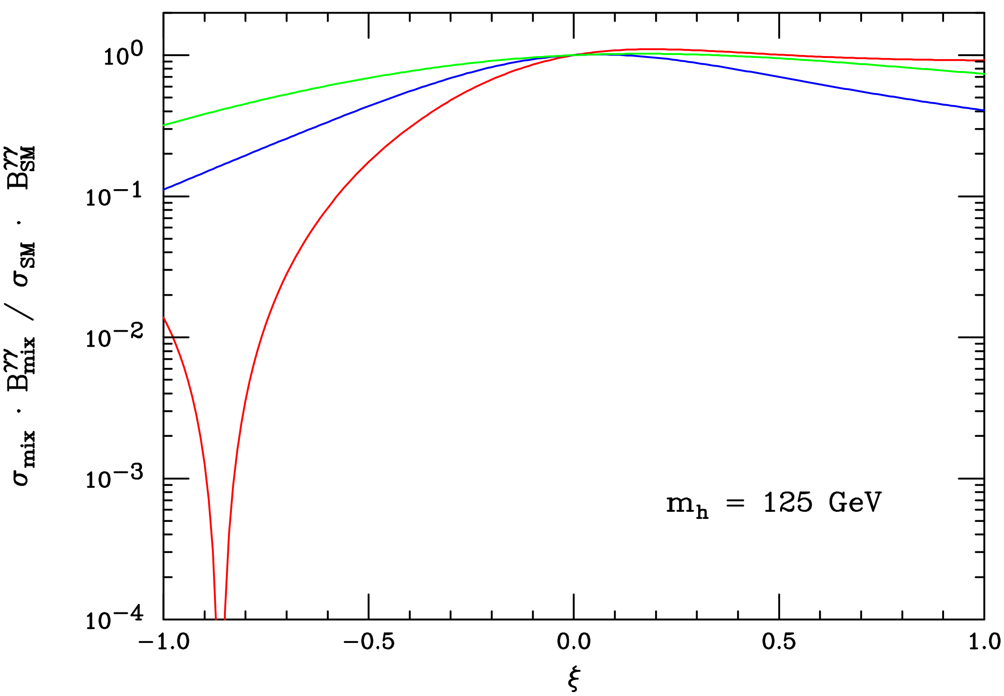

Figure 5: The ratio of production cross section times branching fraction for $pp \to h \to \gamma\gamma$ via gluon fusion with radion mixing to the SM expectations as a function of ξ. The Higgs mass is taken to be 125 GeV. The red (blue; green) curves correspond to the choice $m_r = 300$ GeV, $v/\Lambda = 0.2$ (500, 0.2; 300, 0.1).

In summary, we see that Higgs-radion mixing, which is present in some extra dimensional scenarios, can have a substantial effect on the properties of the Higgs boson. These modifications affect the widths and branching fractions of Higgs decay into various final states, which in turn can alter the expectations for Higgs production at both LEP and the LHC. For some regions of the parameters the size of these width and branching fraction shifts may require the precision of a Linear Collider to study in detail.

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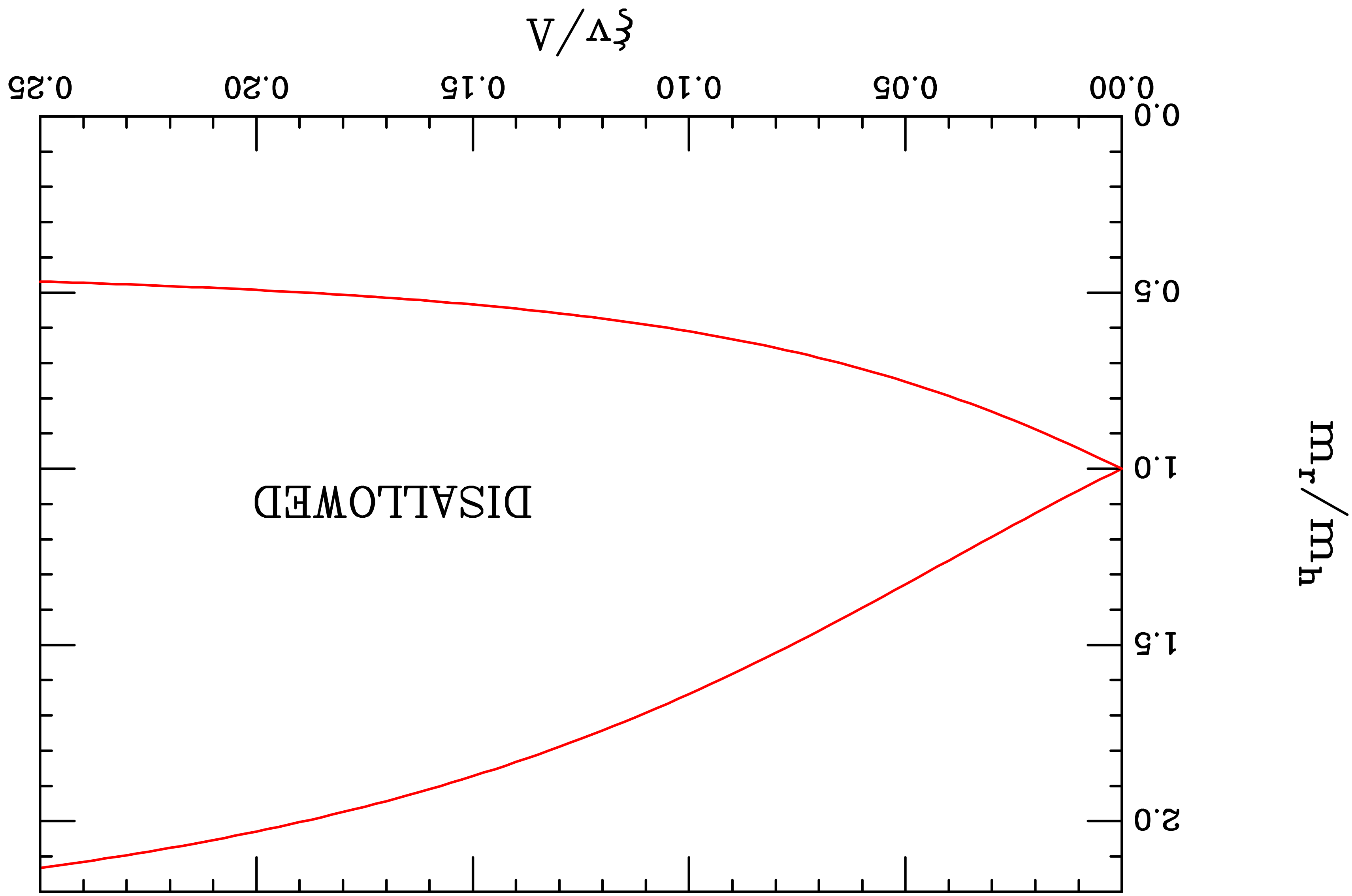
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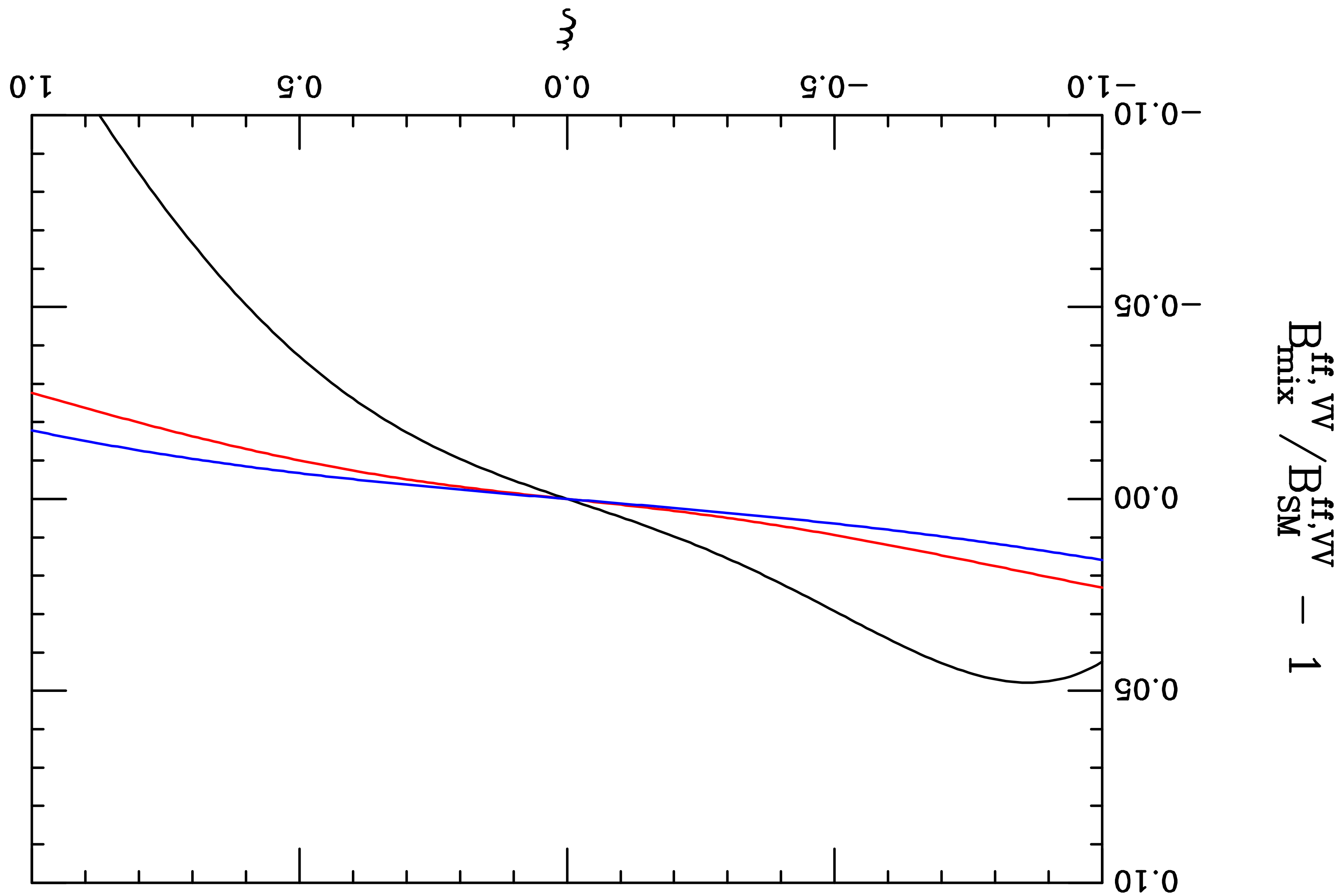
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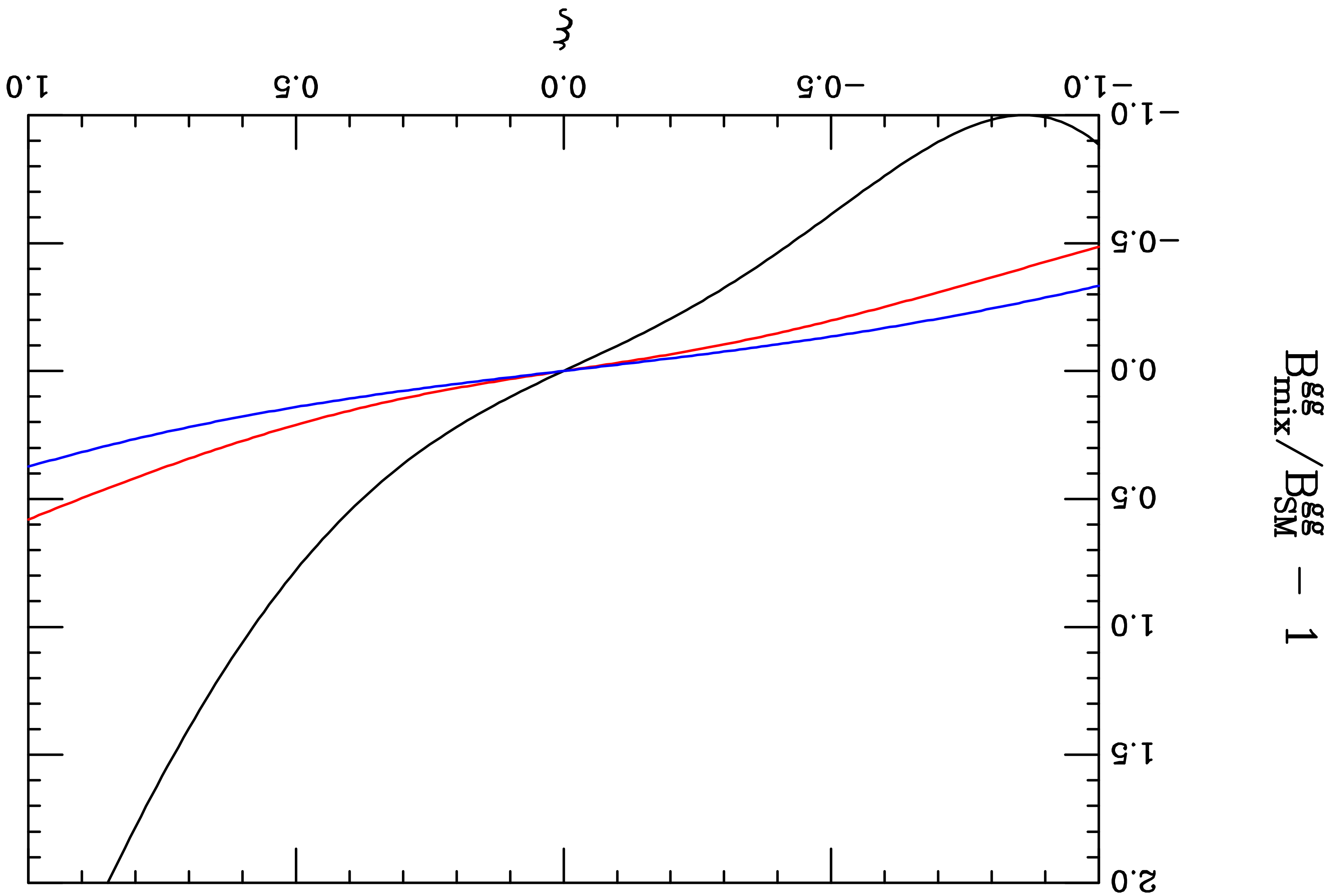
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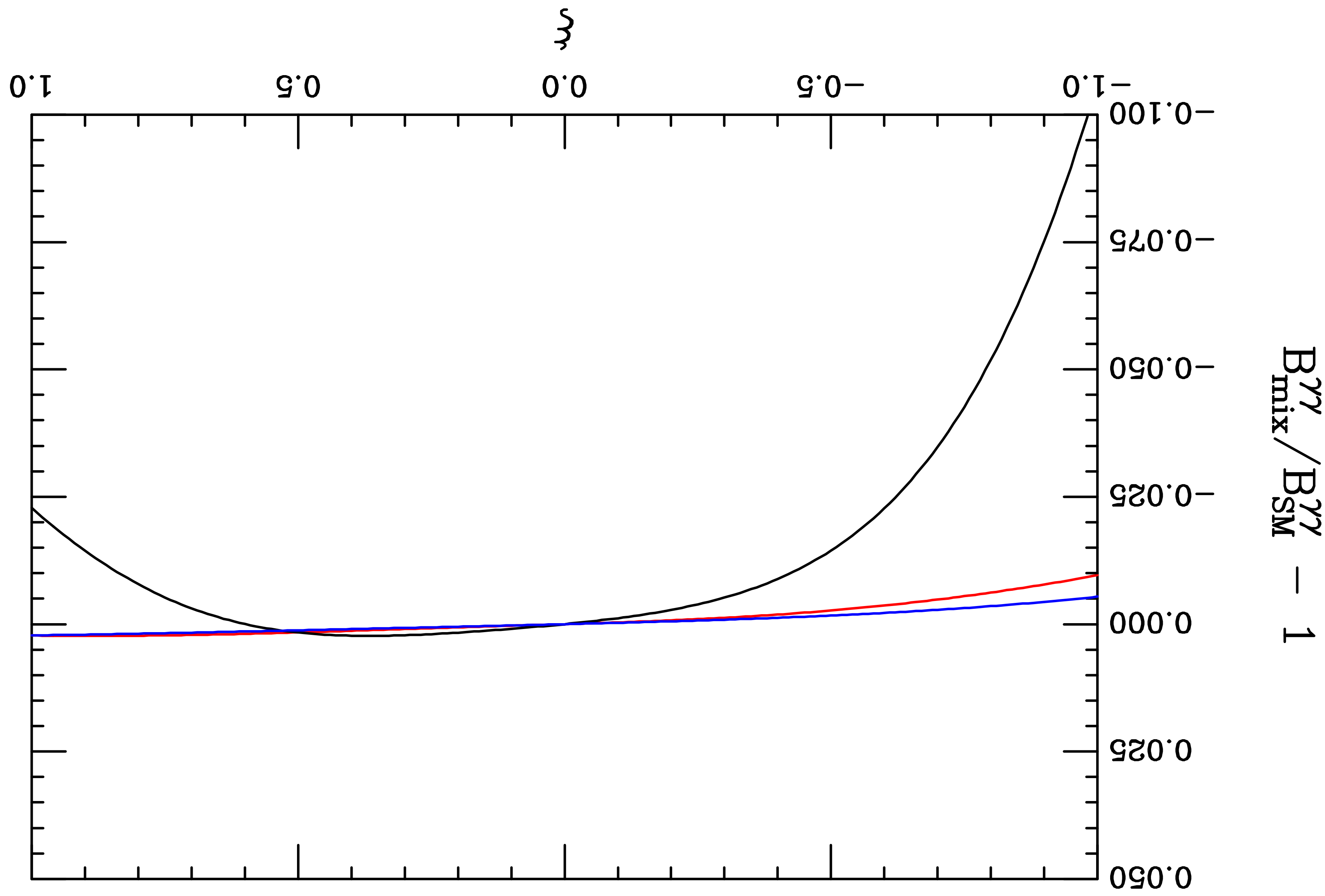

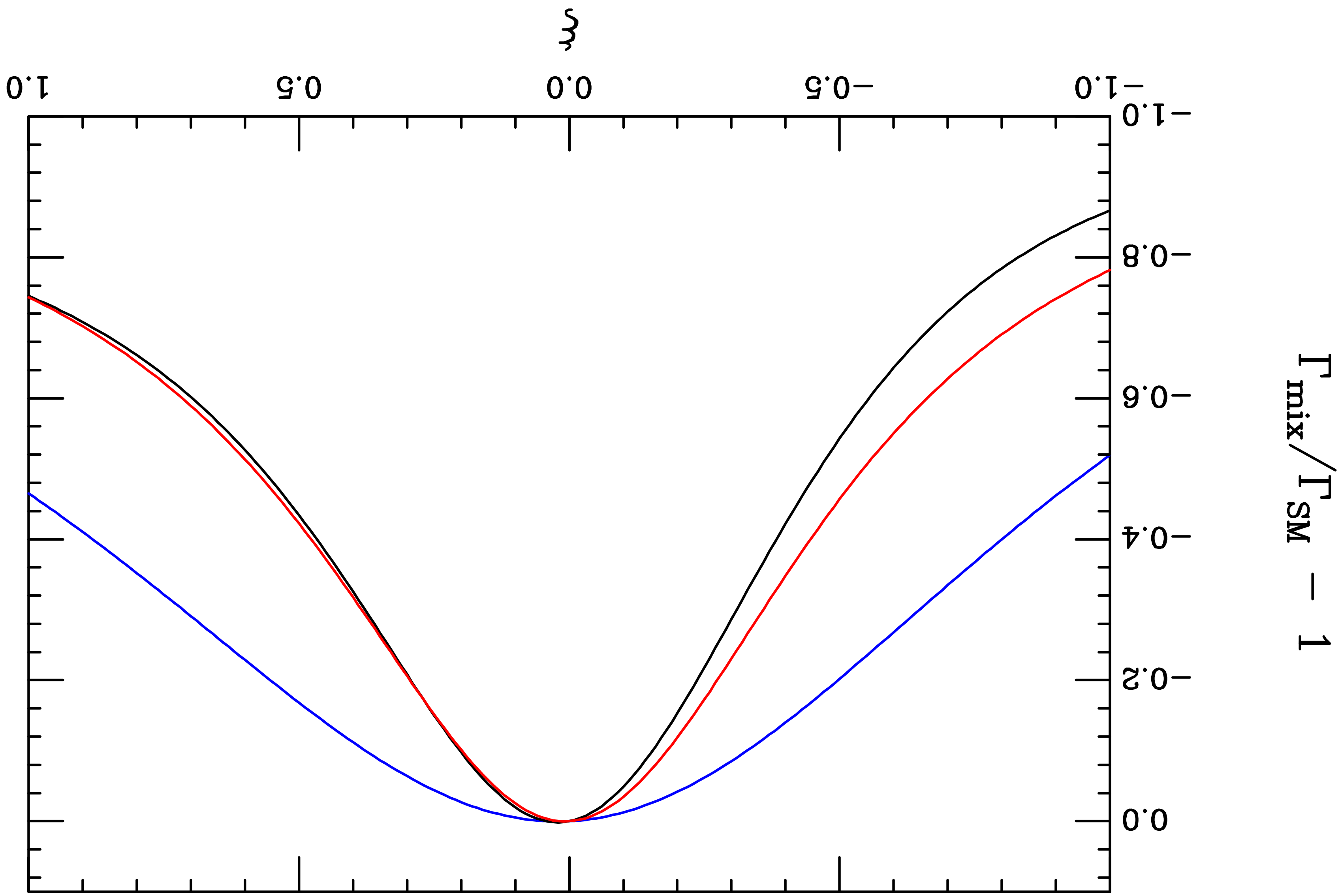

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0.0
0.2
0.4
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1.0
Γ_mix/Γ_SM − 1