

SLAC-PUB-7283
August 1996

First Measurement of the T-Odd Correlation Between the Z^0 Spin
and the Three-Jet Plane Orientation in Polarized Z^0 Decays to
Three Jets

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We present the first measurement of the correlation between the Z^0 spin and the event-plane orientation in polarized Z^0 decays into three jets in the SLD experiment at SLAC utilizing a longitudinally polarized electron beam. The CP-even and T-odd triple product $\vec{S}_Z \cdot (\vec{k}_1 \times \vec{k}_2)$ formed from the two fastest jet momenta, \vec{k}_1 and \vec{k}_2 , and the Z^0 polarization vector \vec{S}_Z , is sensitive to physics beyond the Standard Model. We measure the expectation value of this quantity to be consistent with zero and set 95% C.L. limits of $-0.022 < \beta < 0.039$ on the correlation.

Polarization is an essential tool in investigations of fundamental symmetries in particle physics. Parity violation was first discovered in β decays from polarized ^{60}Co , and T, CP and CPT violations were searched for using polarized neutrons¹ and positronium². The recent development of high-polarization electron sources based on strained-lattice GaAs photocathodes, in conjunction with the high luminosity achieved at the SLAC Linear Collider (SLC), has allowed production of highly polarized Z^0 bosons by e^+e^- annihilation, enabling investigations of symmetries at the Z^0 resonance.

For polarized Z^0 decays to three hadronic jets one can define the triple product $\vec{S}_Z \cdot (\vec{k}_1 \times \vec{k}_2)$, which correlates the Z^0 boson polarization vector \vec{S}_Z with the normal to the three-jet plane defined by \vec{k}_1 and \vec{k}_2 , the momenta of

* Work supported in part by Department of Energy contract DE-AC03-76SF00515 (SLAC) and National Science Foundation contract NSF-PHY95-10439 (Rutgers).

Presented at the annual Divisional Meeting (DPF 96) of the
Division of Particles and Fields of the American Physical
society, 10-15 August 1996, Minneapolis, MN.

the highest- and the second-highest energy jets, respectively. Here we report the first experimental study of this quantity.

The triple product $\vec{S}_Z \cdot (\vec{k}_1 \times \vec{k}_2)$ is even under C and P reversals, and odd under T_N , where T_N reverses momenta and spin vectors without exchanging initial and final states. Since T_N is not a true time-reversal operation, a non-zero value does not signal CPT violation and is possible in a theory that respects CPT invariance³.

The differential cross section for $e^+e^- \rightarrow q\bar{q}g$ for massless quarks may be written^{4,5}

$$\frac{1}{\sigma} \frac{d\sigma}{d\cos\omega} = \frac{9}{16} \left[\left(1 - \frac{1}{3} \cos^2\omega\right) + \beta A_Z \cos\omega \right], \quad (1)$$

where ω is the polar angle of the vector normal to the event plane with respect to the electron-beam direction, and A_Z is the polarization of the Z^0 boson along this same direction ($A_Z = (P_{e^-} - A_e)/(1 - P_{e^-} \cdot A_e)$, where P_{e^-} is the electron-beam polarization, defined to be negative for a left-handed beam, and $A_e = 2v_e a_e / (v_e^2 + a_e^2)$ with v_e and a_e the electroweak vector and axial vector coupling parameters of the electron, respectively). With $|\beta A_Z|$ representing the magnitude⁶, the second term is proportional to the T_N -odd triple product, and appears as a forward-backward asymmetry of the event-plane normal relative to the Z^0 polarization axis. The sign and magnitude of this term are different for the two beam helicities.

Recently Brandenburg, Dixon, and Shadmi have investigated Standard Model T_N -odd contributions of the form $\vec{S}_Z \cdot (\vec{k}_1 \times \vec{k}_2)$ at the Z^0 resonance⁵. The triple product vanishes identically at tree level³, but non-zero contributions arise from higher-order processes; these contributions are found to be very small at the Z^0 resonance and yield values of the correlation parameter $|\beta| \lesssim 10^{-5}$. Because of this background-free situation, measurement of the cross section (1) is sensitive to physics processes beyond the Standard Model that give $\beta \neq 0$.

The measurement was performed with the SLC Large Detector (SLD)⁷ using approximately 50,000 Z^0 decays into multi-hadrons collected in 1993 and 100,000 decays collected in 1994-95, for which the magnitude of the average electron-beam polarization was 0.63 and 0.77 respectively. In the present analysis⁸ the hadronic event selection and three-jet reconstruction were based on the topology of energy depositions in the liquid argon calorimeter, taking advantage of its large solid-angle coverage. The efficiency for selecting hadronic events was estimated to be $92 \pm 2\%$, with a background in the selected sample of $0.4 \pm 0.2\%$, dominated by $Z^0 \rightarrow \tau^+ \tau^-$ and $Z^0 \rightarrow e^+ e^-$ events.

To measure the triple-product correlation for $e^+e^- \rightarrow q\bar{q}g$, three-jet events

were selected and the three momentum vectors of the jets were reconstructed using the ‘‘Durham’’ jet algorithm⁹. Planar three-jet events were selected by requiring exactly three reconstructed jets to be found with a jet-resolution parameter value of $y_c=0.005$, the sum of the angles between the three jets to be greater than 358° , and that each jet contain at least two clusters. A total of 44,683 events satisfied these criteria. The jet energies were calculated by using the measured jet directions and solving the three-body kinematics assuming massless jets, and were then used to label the jets such that $E_1 > E_2 > E_3$.

For each event the reconstructed jet vectors were used to determine the vector normal to the jet plane and its polar angle ω , from which the measured distribution of $\cos\omega$ was derived. A bin-by-bin correction factor $\epsilon(|\cos\omega|)$, for detector acceptance and initial-state radiation, was determined from Monte Carlo simulations.

The $\cos\omega$ distribution is described by

$$\frac{1}{\sigma} \frac{d\sigma}{d\cos\omega} = \frac{9}{16} \left[\left(1 - \frac{1}{3} \cos^2\omega\right) + \beta A_Z (1 - 2 P_{mis}(|\cos\omega|)) \cos\omega \right], \quad (2)$$

where P_{mis} is the probability of incorrectly ordering the energy of the three jets such that the sign of the $\cos\omega$ term is incorrect.

We performed a maximum-likelihood fit of Eq. 2 simultaneously to the $\cos\omega$ distributions from the 1993 and 1994–1995 left- and right-handed event samples, with the relevant values of A_Z , and allowing the parameter β to vary. We found $\beta = 0.008 \pm 0.015$, where the error is statistical only. The T_N -odd contribution is consistent with zero within the statistical error and we calculate limits of⁸

$$-0.022 < \beta < 0.039 \quad @ \quad 95\% \quad C.L. \quad (3)$$

A number of systematic checks were performed. The analysis was performed on two Monte Carlo samples with simulated T_N -odd values; in both cases β was measured to be consistent with the input value within the statistical error. The dependence of the β value on the jet-finding algorithm and the jet-resolution parameter was examined. The analysis was also performed using only charged tracks measured in the central drift chamber. In each case the T_N -odd contribution was found to be consistent with zero within the statistical error.

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