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# On Testing V-A in $\Lambda_b$ Decays<sup>\*</sup>

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## ABSTRACT

We comment on a recent suggestion by Amundson, Rosner, Worah and Wise to test the chirality of the *b*-quark decay coupling via polarized  $\Lambda_b$  baryons produced in  $e^+e^- \to Z \to \Lambda_b + X$ . We study the effect of contributions from an amplitude in which a right-handed *b* to *c* current couples to a V-A lepton current.

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b-quark decays are known to be governed by very weak couplings, which in the Standard Model are given by the two tiny mixing angles  $\theta_{23}$ ,  $\theta_{13}^{-1}$ . These decays are therefore very sensitive to new kinds of interactions, and in particular to right-handed couplings  $^2$ . Recently we described a viable  $SU(2)_L \times SU(2)_R \times U(1)$  model, in which b-quarks decay purely right-handedly  $^3$ . We pointed out that measurement of the lepton asymmetry in  $B \to D^* \ell \nu^4$  cannot distinguish our model from the Standard Model, since in this process the lepton current in our model is dominantly V+A  $^5$ . A nice test, which may distinguish between the exchange of left- and right-handed gauge bosons in b decays was suggested by Amundson, Rosner, Worah and Wise  $^6$ . These authors noted that the electron spectrum from highly polarized  $\Lambda_b$ 's produced in  $e^+e^- \to Z \to \Lambda_b + X$  is significantly harder for V-A than for V+A quark and lepton currents. Thus, ongoing experiments at LEP, in which leptons from  $\Lambda_b$  decay were already observed  $^7$ , may offer an early test of the model.

In this brief note we wish to study the effect of  $W_L-W_R$  mixing, which exists in our model in addition to  $W_R$  exchange<sup>3</sup>, to see how much it can affect the electron spectrum calculated in <sup>6</sup>. Also, as a model-independent study and to demonstrate another version of right-handed b to c couplings, we will first calculate the electron spectrum for a V+A quark coupling, assuming that the lepton current is purely left-handed. Such a possibility is outside the parameter range of the model of <sup>3</sup>, since it corresponds to decays due to  $W_L-W_R$  mixing alone. This case, which seems to be one's first guess of what right-handed b couplings might be, was recently excluded by the  $B\to D^*\ell\nu$  data <sup>4</sup>. Our purpose of including a discussion of this case is to show that also in the case of  $\Lambda_b$  decays it can be most easily distinguished from other cases studied here.

We use the physics of the heavy quark symmetry presented in <sup>68</sup> and the

notations of  $^6$  to describe the lepton spectrum in terms of the free-quark decay  $b \to ce^-\overline{\nu}_e$ . For a V+A b to c coupling and a V-A lepton current the normalized electron decay distribution in the b rest frame is given by  $^9$ :

$$\frac{1}{\Gamma} \frac{d^2 \Gamma}{dx d(\cos \psi)} = \frac{6x^2 (1 - \zeta)^2}{f(m_c^2 / m_b^2)} (1 - x)(1 - P\cos \psi) , \qquad (1)$$

$$x \equiv 2E^*/m_b \; , \qquad \zeta \equiv m_c^2/[m_b^2(1-x)] \; .$$

f(y) is a well-known phase-space function, and  $E^*$ ,  $\psi$  are the energy of the electron and its angle with respect to the b-quark polarization. The polarization is almost complete, P=-0.93. The boost from the b rest-frame to the Z frame, in which  $Z\to b\bar{b}$  occurs, is described in  $^6$ . For P=-1 we find from Eq.(1) the electron energy spectrum shown in Fig. 1(c). We used the values of  $m_b=5$  GeV,  $m_c=1.66$  GeV,  $E_b=45$  GeV from  $^6$ , and chose a minimum electron transverse momentum of  $p_T^{min}=0.8$  GeV/c. This spectrum should be compared with the two spectra of  $^6$  using the same momentum cut, shown in Fig.1(a), Fig.1(b), which describe the cases of V-A and V+A quark and lepton currents, respectively. The difference is striking. The distribution of Fig.1(c) peaks at a considerably higher energy (14 GeV instead of 7-9 GeV) and involves many fewer low energy electrons than the two other distributions.

The possibility of a V+A b to c current coupled to a V-A lepton current was recently excluded by measurement of the forward-backward asymmetry in  $B \to D^*\ell\nu^4$ . This measurement favors the two cases in which both quark and lepton currents are either V-A or V+A<sup>5</sup>. The first case corresponds to the Standard Model, while the second one describes the dominant  $W_R$  exchange contribution in model<sup>3</sup>. For both cases the measured lepton angular distribution sets 95% C.L.

upper limits on the allowed rates coming from amplitudes in which the quarks and leptons couple with opposite chiralities. The form-factor-dependent limits on the ratio of rates of opposite and equal chiralities are at the level of 30%. The implication of these limits on the model<sup>3</sup> is a bound on the  $W_L - W_R$  mixing parameter,  $\zeta_g$ :

$$\left(\frac{\zeta_g}{\beta_g}\right)^2 < 0.30 , \qquad \beta_g \equiv \frac{g_R^2}{g_L^2} \frac{M_L^2}{M_R^2} .$$
 (2)

We use this constraint to study within model<sup>3</sup> the effect of  $W_L - W_R$  mixing on the lepton energy spectrum of  $\Lambda_b$  decay. Fig.1(d) describes the spectrum corresponding to  $(\zeta_g/\beta_g)^2 = 0.29$ , which is just below the limit (2). The peak of this distribution lies between the peaks of the V-A and V+A distributions, Fig.1(a) and Fig.1(b), respectively. That is, the effect of  $W_L - W_R$  mixing is to diminish the difference between the distributions of our model and the Standard Model. Nevertheless, the feature of a considerably lower high-energy electron tail persists in our model. An observation of a higher tail, as in Fig.1(a), would clearly favor the Standard Model.

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## FIGURE CAPTION

FIG.1. Distributions in electron laboratory energy for inclusive semileptonic  $\Lambda_b$  decays from  $e^+e^- \to Z \to \Lambda_b + X$ , with  $p_T^{min} = 0.8$  GeV/c. (a) Standard Model<sup>6</sup>, (b)  $W_R$  exchange<sup>6</sup>, (c) V+A quark coupling and V-A lepton coupling, (d) Model<sup>3</sup> with  $(\zeta_g/\beta_g)^2 = 0.29$ .

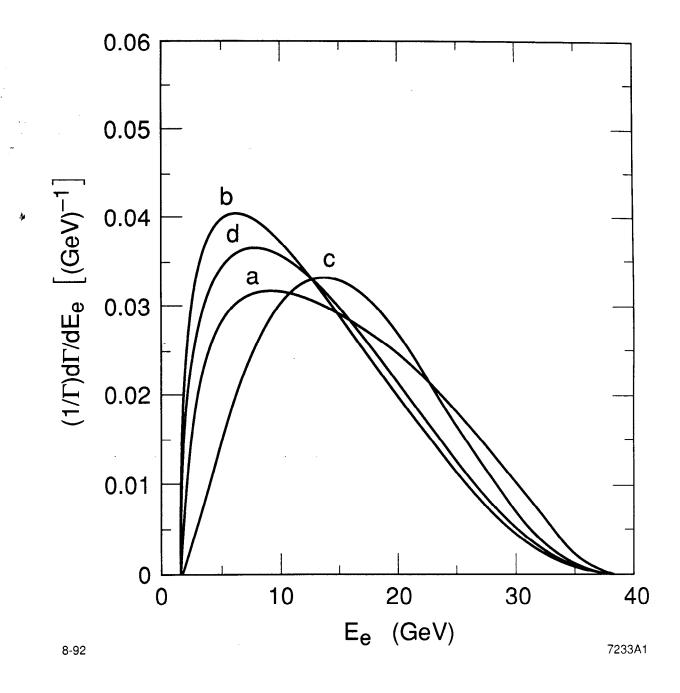


Fig. 1