Construction and Testing of the SLD Čerenkov Ring Imaging Detector*

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Abstract

We report on the construction of the Čerenkov Ring Imaging Detector (CRID) for the SLD experiment at the SLAC Linear Collider and the testing of its components. We include results from testing the drift boxes, liquid radiator trays, and mirrors for the barrel CRID. We also discuss development of the support systems essential for the operation of the CRID: gas and liquid recirculator systems and monitoring.

Introduction

A large Čerenkov Ring Imaging Detector (CRID) [1, 2] is currently under construction for the SLD experiment at the SLAC Linear Collider (SLC). The barrel CRID will provide almost complete particle identification over the central 70% of the solid angle (figure 1). By making use of

*Work supported by the Department of Energy, contracts DE-AC03-76SF00515 and DE-AT03-79ER70023 and by the National Science Foundation under Grants PHY88-13669 and PHY88-13018. both liquid and gaseous radiators, $\pi/K/p$ separation will be possible up to about 30 GeV/c, and e/π separation up to about 6 GeV/c. Endcap CRID detectors are currently under design; in this paper we restrict our discussion to the barrel.

The barrel is constructed with a geometry of ten azimuthal sectors on each side of the mid-plane. In each sector, there are two drift boxes, each coupled to a multiwire proportional chamber sensitive to single electrons. or "single-electron detector" [3]. The drift box and detector combination serves to image Čerenkov light produced in the radiators. A Cerenkov photon with wavelength less than 230 nm entering a drift box can ionize the drift gas mixture containing about 0.1% TMAE (tetrakis[dimethylamino]ethylene), giving a single photoelectron. (To be transparent to photons of this short wavelength, the drift box windows are made of fused silica, i.e., quartz.) The resulting photoelectron drifts up to 1.27 meters in a uniform electric field which is nearly parallel to the magnetic field in a time projection geometry. It is detected by a plane of ninety-three 7 μ m diameter carbon sense fibers, spaced at a pitch of 3.2 mm, in the detector. The electronics used to read out the detector are discussed elsewhere [5]. The three production coordinates of the electron are measured from the drift time, the position of the fiber giving the signal, and charge division along the fiber. The production coordinates are used together with track-

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Figure 1: Principal Components of the SLD CRID Barrel system.

ing information from the SLD central drift chamber [1] to determine the Čerenkov angle and thus particle mass.

The liquid radiator is a 1 cm thick layer of C_6F_{14} , contained in a G-10 tray with a quartz window. The quartz window is on the side away from the beam axis, allowing the Čerenkov radiation produced in the liquid to reach a drift box for detection. There are forty trays in the finished barrel, one corresponding to each drift box.

The gas radiator is obtained by filling the CRID vessel with gaseous C_5F_{12} . Čerenkov photons produced in the gas are focused on the drift boxes by spherical mirrors mounted on the outer wall of the CRID vessel. Each drift box detects light from two rows of five mirrors. The 400 mirrors in the barrel give complete coverage.

The CRID is operated at 40° C in order to keep the C_5F_{12} , which boils at 30° C, in the gaseous phase and to prevent condensation of TMAE in the drift box gas. The drift gas is saturated with TMAE at 28° C by bubbling the gas through liquid TMAE maintained at that temperature.

Production Status

The mechanical structure, or vessel, of the barrel CRID has been installed in SLD. All 400 mirrors have been mounted, aligned, and installed in the vessel. A final survey of the mirrors positions and orientations is in progress. Before it was moved into SLD, the vessel was also fitted with the radiator gas temperature control system and central HV distribution plane.

About 80% of the drift boxes for the barrel CRID are finished, and assembly of the single-electron detectors is complete. Almost two thirds of the liquid radiator trays are finished. Finished drift boxes, detectors, and radiator trays are stored in a clean room with the temperature con-



Figure 2: Transmission of a completed drift box. The circles are the measured values. The triangles are the estimate from the transmission of the individual windows, including losses from the field-shaping wires.

trolled to 37° C and a relative humidity of less than 10%. The room is equipped so that the interior of the drift boxes and radiator trays can be purged with dry, boil-off nitrogen.

Component Testing

Drift boxes

In order for the drift boxes to detect Čerenkov photons from the radiators with high efficiency, three items are critical to their performance. The quartz windows must have good transmission for wavelengths short enough to ionize the TMAE in the drift box. Second, the resulting photoelectrons must have a long lifetime (~ 100-200 μ s) against reattachment, giving them time to drift to the detector. Finally, the drift field must be uniform so that the production point of an electron can be reconstructed with millimeter accuracy from the signal observed by the single-electron detector.

The UV transparency of all quartz sheets arriving from the manufacturer is measured in a vacuum test station. This device allows a UV beam to be scanned over the whole 125 cm \times 30 cm surface of an assembled window or a completed drift box. Figure 2 shows the transmission measured at one point for a completed drift box. It is in good agreement with the transmission predicted from the measurements of the individual windows and losses due to the field-shaping wires on the quartz surface, indicating that the assembly process does not degrade the transmission of the windows. Measurements taken at many positions on the box indicate good uniformity of transmission over the surface.

Finished drift boxes are fitted with a single-electron de-



Figure 3: The test apparatus for the measurement of photoelectron drift distortion, electron lifetime, and drift gas UV transparency. The distortion measurements of Figures 4-6 were made with this apparatus.

tector and tested in a high voltage set-up which closely simulates the electrostatic configuration of the actual CRID vessel. Here we test the high voltage performance only. The detector has first been tested in a fixture described in [3].

The drift box and detector assembly then undergoes further testing in the apparatus of figure 3. The assembly is placed in a Lucite test enclosure and connected to a gas system which provides drift gas either with or without TMAE. The Lucite enclosure can be filled with Freon to help the drift box hold high voltage and to test for Freon leakage into the drift volume. For these tests we use the Freon, $C_2Cl_2F_4$. The output gas can be monitored for electron lifetime [7] and UV transparency. UV absorption in the 160-230 nm region is a sensitive check of the TMAE concentration.

To maintain a long lifetime against electron reattachment, the drift gas must be of high purity. Not only must pure gas be delivered to a drift box by the gas system, but it must remain pure inside the box. (The gas flow is about one volume change every two hours.) To prevent the electronegative C_5F_{12} of the gas radiator from getting into a drift box, each finished drift box is leak checked with a helium leak detector to a sensitivity of 10^{-5} cm³/s. The materials used in the construction of the drift box have been carefully selected to minimize reaction with the TMAE in the drift gas (see [6]).

The electron lifetime monitor on the output gas has been useful for comparing relative lifetime but has been difficult to calibrate absolutely. We have recently measured the electron lifetime directly in a drift box. The apparatus in figure 3 includes a pulsed N₂ laser (wavelength = 337 nm) mounted on an x-y transport allowing it to be pointed into any region of the drift box. The laser shines into the drift box through a quartz window and ionizes the drift gas via a double-photon process. By firing the laser into the box at several different drift distances and examining the number of observed electrons versus drift distance, we



Figure 4: Difference between the expected and measured x position for different values a x across the width of CRID drift box number 6. Data are shown for four different drift distances, ranging from 11 to 121 cm.

have measured a lifetime of approximately 100 μ s [8]. This is satisfactory for drift box operation.

The laser can also be used to measure distortions in the trajectories of drifting electrons by comparing the reconstructed light spots with their known positions (see [8]). (Although each drift box is equipped with gating electrodes to suppress the entry of positive ions into the drift volume [4], measurements done so far have been without gating.) Distortions in the x-y plane, the plane parallel to the windows, are measured by firing the laser vertically into the drift box. The laser intensity is adjusted so that the mean number of photoelectrons produced by one pulse is ~ 0.1 . At this intensity, events with multiple photoelectrons are rare. Upon reaching the detector, the electron produces a signal on a single wire. Taking many laser pulses in the same position, the hits are distributed over two to four wires because of diffusion and the finite spot size from the laser. This allows the mean position to be determined to better than the wire pitch. Comparing the reconstructed wire address with the address that is predicted from the laser position, assuming a uniform drift field, measures the extent to which the drift field is not uniform.

Distortion measurements have been made in this way on several drift boxes. Figure 4 shows distortions measured across the width of a drift box, for several drift distances. The plot shows the measured position minus the predicted position, in mm. The distortions are largest near the G-10 side walls, for x near 0 and 300 mm, but even there they do not exceed 2 mm.

Distortions in the direction normal to the window, parallel to the wires in the detector, can also be investigated. This is done by pivoting the laser so that it shines into the box at an angle, typically 45°. The electrons are observed



Figure 5: Use of an angled laser track to investigate z distortions. (a) Reconstructed z position as a function of drift time (the zero for the time coordinate is arbitrary). A line with slope predicted from the laser angle and the drift velocity is superimposed on the data. (b) The residuals to the line shown in part a. The curve shows the residuals predicted by an electrostatic simulation. These results are for the normal operating potentials.

on a few neighboring wires, as before, but now the drift time varies depending on the depth (z) of ionization. The apparent z, as observed at the detector, is reconstructed by charge division. We divide the measured drift time of the observed signals into 20 ns bins, and in each bin we calculate the mean measured z. Figure 5a shows data on the resulting mean z versus drift time. We fit the data to a straight line with the slope fixed by the known laser angle and measured drift velocity. The residuals (or distortions) are shown in figure 5b. The measurements we show here are taken at the center of the detector, *i.e.*, equidistant from the two side walls. Similar measurements 1.5 cm from the side wall show distortions of up to 5 mm in the corners of the drift volume, with 2 mm being a more typical value.

One important source of possible distortions in the drift field is a mismatch of potentials in the transition region between the drift volume and the detector. We have investigated this by varying the potential of the drift box (keeping the gradient of the drift field within the box constant at 400 V/cm) while keeping the detector at the standard operating voltages. The results shown in figures 4 and 5 were obtained at the optimum voltages as determined both from



Figure 6: Effect of potential matching on z reconstruction. These plots are analogous to figure 5b. They result from decreasing (a) and increasing (b) the potential difference between the drift box and the detector by 200 V. The curves are the predictions of the electrostatic simulation.

the data and from electrostatic simulations [9]. The residuals predicted from the electrostatic simulation are shown as a curve on figure 5b. Repeating the measurement for drift box potentials 200 V lower and higher, results in the residuals shown in figure 6. The data have been fit to a line with the same slope as before. The data agree nicely with the distortions predicted by the electrostatic simulation.

Liquid radiator trays

Operation of the liquid radiator trays requires good UV transmission of the liquid C_6F_{14} and quartz window. The quartz is checked upon receipt from the manufacturer in the same way as the drift box quartz.

We have built an apparatus to measure the UV transmission of liquid C_6F_{14} . This device is being used to conduct tests of how long the liquid retains transmission in the presence of the tray materials. We have begun by studying the transmission of the C_6F_{14} before exposure to the tray components. The liquid contains dissolved gases, particularly oxygen, that must be removed to obtain good transmission. The C_6F_{14} test samples are degassed by several cycles of freezing in liquid nitrogen, purging the container with boil-off N_2 gas, and allowing the frozen liquid to thaw and equilibrate with the gas in the container. We have observed good transmission in the purified samples, although we are still studying the systematic limitations of our transmission measurements. When the tests are finished, the UV spectrometer will be moved to the SLD detector. There it will monitor the liquid at various stages as it is recirculated through the trays, including the output of any single tray.

It is essential that the liquid trave be strong enough to support the weight of the liquid and other smaller pressure differentials for the several-year life of the experiment. Two sample trays loaded with C_6F_{14} have been tested under slowly increasing pressure for many months [10]. The tests were performed at the CRID operating temperature of 40° C. The results are summarized in figure 7. The first tray tested had a 3 mm thick window of ordinary float glass. The second test tray had a 4 mm thick quartz window, the same thickness used in the production travs. For each tray, the window deflection was monitored throughout the test. The figure shows the observed window deflection for each tray as a function of the applied pressure, up to the breaking point. By comparing the deflections to those predicted by simple models of uniformly loaded plates with known boundary conditions [11], we have estimated the maximum stress in the window at the breaking point. In the case of the glass window, the deflections were well represented by a model in which the sides and rib of the tray serve as simple supports. The thicker quartz window is more closely reproduced by a model in which the rib is a fixed support. In both cases, the window failed at a stress of approximately 48 MPa and in the region of maximum stress as predicted by the model. The quartz tray, having a thicker window, sustained a higher applied pressure. The quartz tray failed at a pressure of six times the largest load borne by any of the trays under normal operation.

Mirrors

The mirrors for the barrel CRID [12] are approximately 30 cm \times 30 cm, with a radius of curvature 95.8 or 101.6 cm. There are ten slightly different shapes of mirrors. Each of the ten mirrors directed on a drift box requires a shape specific to its position. The mirrors need to have an angular reflection error of less than 1 milliradian in order to make a negligible contribution to the position error of the photoelectrons. They also need to have reflectivity of at least 85% in the wavelength range between 180 and 220 nm.

The optical distortion of each mirror has been measured photographically prior to applying the reflective coating by imaging a mask placed in front of the mirror. The apparatus is such that a 1 milliradian reflection error corresponds to an error of approximately 1 mm on the negative. Only 3.5% of the mirror substrates failed this test and had to be replaced.

To obtain good reflectivity, the mirrors need both good surface smoothness and a good coating. To keep light loss due to scattering to a few percent, the mirror substrates are polished to an average surface roughness of ≤ 3 nm



Figure 7: Pressure tests of sample trays. The solid circles are for a tray with a 3 mm thick glass window. The two dotted lines show predictions for the model with simple support at the rib (steeper curve) and fixed support at the rib (lower curve). The open squares and dashed lines show the data and predictions, respectively, for a tray with a 4 mm thick quartz window. The normal operating range for the trays is 1.7-6.1 kPa.

(rms). The surface profile of approximately 70% of the mirrors was checked at Applied Magnetics Co.¹ with a surface profilometer. Seven percent of the mirrors failed the test and were replaced. Reflectivity of the coated mirrors relative to a standard sample was measured using a custom instrument (see [12]). Figure 8 shows the average reflectance of all the mirrors at eight wavelengths in the range 160-230 nm. All of the mirrors satisfied the reflectivity requirements.

Fluid Systems and Monitoring

The CRID fluid delivery and monitoring systems are presently being assembled and tested. A comprehensive on-line system will continually monitor the purity, UV transparency, electron lifetime, and mixing of the various fluids so that problems may be detected before the physics performance of the CRID is degraded. A more complete description of the monitor system can be found in ref. [13].

The 40° C operating temperature of the CRID is maintained using Kapton-insulated heater pads attached to the surfaces of the vessel. The heaters are controlled according to the output of temperature sensors interspersed with the heaters. Current to the heaters is digitally switched using an AC zero-crossing solid state relay. The relay is followed by a rectifier and a smoothing capacitor to provide "softstart" DC current for use inside the SLD magnetic field. The total available heating power is 15 KW at 150 volts



Figure 8: Average UV reflectance measured for the production mirrors for the barrel CRID. Also shown is the minimum acceptable reflectance: 80% at 160 nm, rising to 85% above 180 nm.

DC. All of the gases and liquids circulating through the CRID will be preheated and delivered in heated lines.

Three gases are necessary for the operation of the CRID. The gas inside the drift box is ethane, bubbled through TMAE at 28° C to achieve a concentration of about 0.1%. It is then heated to 40° C for delivery to the drift boxes. The drift boxes are built with double side walls, separated by a purge space [14], to help keep the electronegative radiator gas from entering the drift box or the gas laden with reactive TMAE from escaping into the vessel volume. The purge volume is flushed with pure methane. Finally, the vessel is filled with gaseous C_5F_{12} . The cost of the C_5F_{12} requires that it be recirculated. After flowing through the vessel, the gas is passed through filters to remove water, oxygen, and particulates before flowing through the vessel again. In addition, pure N2 is used as a purge gas before or after opening the vessel to air. After sealing the vessel, it is first purged with nitrogen until a satisfactory level of O_2 and H_2O is reached. The N_2 is then thermodynamically replaced with C_5F_{12} . Similarly, the C_5F_{12} will be replaced by nitrogen before opening the vessel. The exchange is monitored by measuring the density of the gas with a sonar device [15], both directly in the vessel and in the monitor system described below.

The drift gas, drift box purge gas, and radiator gas are all at atmospheric pressure. An analog system controls the pressure of each of the three systems by regulating the flow using a mass flow controller.² Adjustments are made based on the readings of capacitance manometers³ attached to each of the systems. Pressure differences between the drift box gas and radiator gas system must be small, no more than a few Torr, to prevent the pressure

¹Applied Magnetics Co., Goleta, CA 93117

²Model 258B, MKS Inc., Burlington, MA 01803, USA.

³Model 221, MKS Inc., Burlington, MA 01803, USA.

differences across the quartz drift box windows from cracking them. Similarly, deviations from atmospheric pressure must be small, so that pressure differences between the gas radiator and liquid radiator, which is separately referenced to atmosphere, do not jeopardize the liquid tray windows. The normal operating range for the gas systems is within 1 Torr of atmospheric pressure. The system has several levels of response to increasingly large pressure excursions, culminating in passive pressure relief through bubblers for deviations exceeding 5 Torr.

We have made extensive provision for monitoring the quality of the gases delivered by the gas systems and also the gases returning from the CRID. This is shown schematically in figure 9. The upper part of the figure shows the system for checking source methane and ethane gas and the radiator gas. They can be checked for density (using the sonar), UV transmission, water content, and oxygen content. Samples can also be extracted for analysis in a gas chromatograph. Similarly, gas containing TMAE from either of two TMAE bubblers and the output gas from the drift boxes and purge flow are examined by the system in the lower part of the drawing. Each drift box has a separate output line that can be selected. The output flow from the purge volumes is manifolded together to give one line for each half of the barrel. The gases connected to this part of the system can be checked in much the same way as the first group, except that the sonar is replaced by an electron lifetime monitor [7], which measures the electron attachment lifetime in the gas.

To reconstruct the production coordinate of an electron drifting over 1 m to an accuracy of 1 mm, the drift velocity needs to be known to 0.1%. The drift velocity will change during the day-to-day operation of the CRID, if for no other reason, then because of variations in atmospheric pressure. In addition to monitoring the pressure, the drift velocity will be measured directly in each drift box using fiber optics. Each box is equipped with 19 silica fibers 200 μm in diameter. They are all illuminated using a common 600 μ m fiber. The end of each fiber is held in a collimator glued to the quartz window on the outer (mirror) side of the box. UV light, from either a laser or a flash lamp, travels down the fibers and into the drift box, where it ionizes the TMAE. Comparing the arrival time of signals from fibers at different known positions on the box will enable the drift velocity to be determined. The fibers will also be used to follow any time-dependent distortions in the drift field. One possible source of such a distortion would be positive ions accumulating on the sides, or windows, of the drift box.

Liquid C_6F_{14} for the radiator trays is provided by a recirculation system. The system maintains the C_6F_{14} at 40° C throughout and uses only gravity to move the liquid through the trays. Spill-tanks (overflow devices) are used repeatedly at various levels to define and limit the liquid pressure heads at the trays and other critical points. The entire system is, at all levels, in contact with a buffer gas (purified N₂), which serves as a pressure reference and, in turn, communicates with the atmosphere through a C_6F_{14} condenser and bubblers. Backup blow-away bubblers are used to relieve the gas buffer system if there is a deviation from atmospheric pressure that could break the quartz windows in the trays. Pumps are used only to raise the liquid to the top of the system, whereupon it is dumped into a reservoir to begin its gravity descent.

Before each cycle through the trays, the liquid is passed through particulate filters, water absorbents, and gasremoving chemical columns.⁴ This guards against loss of UV transmission caused by buildup of O₂, H₂O, or other impurities leaching from the trays or lines. The capacity is sized such that 120 cc/s of fluid can be delivered with \lesssim 10 ppm of contaminants of any given type, after initial purification of the C₆F₁₄, *e.g.*, by distillation and degassing. Monitoring of all critical temperatures, pressures, flow rates, impurities and UV transmission are carried out continuously.

To monitor the fluid level in the trays to be sure that they fill completely, we have developed a simple technique using a miniature light reflection emitter/sensor.⁵ This device contains an infrared LED transmitter and a phototransistor. The conventional use of the sensor is to detect changes in the distance or reflectivity of surfaces in front of it, with the highest sensitivity at a distance of about 1 mm. The liquid C_6F_{14} is clear, and the bare sensor does not respond to being immersed in the liquid. However, we have obtained a nice signal to the presence of liquid if we optically couple the emitter and sensor with a drop of epoxy⁶ on the face of the device. Then, enough light is reflected from the surface of the drop to turn on the phototransistor. When the device is immersed in C_6F_{14} , less light is reflected at the interface between the epoxy and the liquid, reducing the current drawn by the transistor. The change in current, e.g., 1.4 mA dry to 0.4 mA wet, is easy to measure. Each tray is fitted with two of the epoxycoupled sensors, one each at the high and low corners of the tray, in the path of the fluid inlet and outlet. The low sensor will be used to confirm when flow has started to a tray, and the high sensor will check that the tray fills completely and remains full as liquid circulates.

Conclusions

The production of the components for the SLD barrel CRID is well along, with all of the mirrors completed and installed. We have described several tests which are designed to ensure the reliable operation of the components in the assembled CRID. Fluid and monitoring systems for the CRID are complicated by the 40° C operating temperature, the importance of all subsystems closely tracking atmospheric pressure, the high purity required, and the ne-

⁴Oxisorb Filters, Messer Griesheim GMBH, D-4000 Düsseldorf 30, FRG.

⁵Catalog No. SFH900-2, Siemens Components, Inc.

⁶Equal parts by weight Epon 826 (Shell Co.) and Versamid 140 (Henkel Co.). Some other transparent epoxies also work.



Figure 9: Schematic of the barrel CRID gas monitor system.

cessity of recycling the C_5F_{12} gas and C_6F_{14} liquid. These systems are currently being assembled and tested.

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References

- [1] SLD Design Report SLAC-273 UC-34D, May 1984, and revisions.
- [2] D. Aston et-al., SLAC-Pub-4795, Nucl. Instr. and Meth., Sect. A, vol. 283, pp. 582-589, 1989.
- [3] D. Aston et al., SLAC-Pub-4875, Nucl. Instr. and Meth., Sect. A, vol. 283, pp. 590-595, 1989; see also J. Va'vra et al., SLAC-Pub-4432, IEEE Trans. Nucl. Sci., vol. NS-35, pp. 487-492, 1988, and ref. [4].
- [4] D. Aston et al., SLAC-Pub-4785, IEEE Trans. Nucl. Sci., vol. NS-36, pp. 276-282, 1989.
- [5] E. Spencer et al., SLAC-Pub-4404, IEEE Trans. Nucl. Sci., vol. NS-35, pp. 231-236, 1988, and E. Spencer et al., "A Low Noise Preamplifier System for Čerenkov Ring Imaging Detectors," to be submitted to Nucl. Instr. and Meth., Sect. A.
- [6] R. T. Rewick et al., SLAC-Pub-4415, Anal. Chem., vol. 60, p. 2095, 1988.
- [7] M. Cavalli-Sforza, "Electron Lifetime Measurements in the Disaster Monitor," CRID Internal Note No. 15, 1986.
- [8] P. Coyle, et al., "Electron Drift Properties of the First Production Drift Boxes," CRID Internal Note No. 62, 1990.
- [9] S. Dasu, "Electrostatic Simulation of CRID Detector," CRID Internal Note No. 63, 1990.
- [10] D. Coyne and D. Williams, "Strength and Breakage Tests of Liquid Radiator Trays," CRID Internal Note No. 61, 1990.
- [11] See, e.g., S. Timoshenko and S. Woinowsky-Krieger, Theory of Plates and Shells, second ed., New York: McGraw-Hill, 1959.
- [12] A. Bean et al., "Production of 400 Mirrors with High VUV Reflectivity for Use in a Ring Imaging Cherenkov Counter in SLD at SLAC," submitted to Nucl. Instr. and Meth., Sect. A.

- [13] P. Antilogus et al., SLAC-Pub-5105, "Monitor and Control Systems for the SLD Čerenkov Ring Imaging Detector," contributed paper to the IV Conference on Control Systems for Accelerators and Large Physics Experiments, Vancouver, Canada, October 29-November 3, 1989, to be published in Nucl. Instr. and Meth., Sect. A.
- [14] G. Hallewell et al., SLAC-Pub-4405, IEEE Trans. Nucl. Sci., vol. NS-35, pp. 398-403, 1988.
- [15] G. Hallewell et al., SLAC-Pub-4122, Nucl. Instr. and Meth., Sect. A, vol. 264, pp. 219-234, 1988.