SLAC-PUB-3128
MAY 83

New Tests for Quark and Lepton Substructure

ESTIA J. EICHTEN
Fermi National Accelerator Laboratory, Batavia, IL 60510
and
Enrico Fermi Institute, University of Chicago,
Chicago, IL 60637

AND

KENNETH D. LANE†
Ohio State University, Columbus, OH 43210

AND

MICHAEL E. PESKIN*
Stanford Linear Accelerator Center, Stanford, CA 94305

ABSTRACT

If quarks and leptons are composite at the energy scale Λ, the strong forces binding their constituents induce flavor-diagonal contact interactions, which have significant effects at reaction energies well below Λ. Consideration of their effect on Bhabha scattering produces a new, stronger bound on the scale of electron compositeness: $\Lambda > 750$ GeV. Collider experiments now being planned will be sensitive to $\Lambda \sim 1$-5 TeV for both electrons and light quarks.

†Work supported in part by the Department of Energy under Contract No. DOE/ER/01545-326.

*Work supported in part by the Department of Energy under Contract No. DE-ACOB-76SF00515.

The proliferation of quarks and leptons has naturally led to the speculation that they are composite structures, bound states of more fundamental constituents which are often called "preons".[1] Many authors have proposed models of such composite structure, but no obviously correct or compelling model has yet emerged. There is not even consensus on the most fundamental aspect of quark and lepton substructure - the value of the mass scale Λ which characterizes the strength of preon-binding interactions and the physical size of composite states. It is therefore important to devise experiments which probe this potential substructure as deeply as possible and which, at the same time, test the widest possible variety of models. In this Letter, we identify new observable consequences of quark and lepton substructure which do just that.[2] An immediate result of these is that existing Bhabha scattering measurements imply that $\Lambda > 750$ GeV for the electron, a factor of 5 larger than previous lower bounds.

Before spelling out our tests, let us review what is known about Λ. At present, high-energy cross sections are well explained by the standard $SU(3)\otimes SU(2)\otimes U(1)$ gauge theory with elementary quarks and leptons. If these fermions are composite, then Λ is much larger than their masses, completely unlike the situation in nuclear and hadron physics. However, 't Hooft has argued that gauge theories of preon-binding quite naturally produce composite fermions much less massive than the binding scale provided certain

symmetry constraints are satisfied.[3] Since the energy $\Lambda_{EW} = O(1\ \text{TeV})$ at which electroweak symmetry is broken is the lowest new dynamical scale we foresee, we expect $\Lambda \geq \Lambda_{EW}$.

Modifications of gauge-field (γ, Z^0,etc.) propagators and vertices with fermions occur in any preon model, though their precise form is model-dependent. In a favored parametrization,[4] one simply multiplies the gauge propagator by a form factor $F(q^2) \cong 1+q^2/\Lambda^2$. Measurements of $e^+e^- \to \psi\bar{\psi}(\psi=e,\mu,\tau,q)$ up to $\sqrt{s} = 35$ GeV at PETRA have excluded photon form factors for $\Lambda \lesssim 100$-200 GeV.[5] Composite fermions also possess new contact interactions generated by constituent exchange. These four-fermion interactions have strength $\pm g^2/\Lambda^2$, where g is an effective strong coupling constant analogous to the ρ-coupling $g_\rho^2/4\pi \cong 2.1$. If contact interactions mediate flavor-changing processes such as $K_L^0 \to \mu e$, D^0-$\bar{D}^0$ and K^0-$\bar{K}^0$ mixing, the lower limits on Λ range from ~ 30 TeV to ~ 800 TeV.[6] While these bounds are impressive, it is possible to construct composite models in which some[6] or all[7] of the dangerous flavor-changing interactions are absent. In summary, the only relatively model-independent constraints are the much lower ones deduced from the PETRA measurements.[8]

Our new tests for substructure are based on two observations: <u>First,</u> in any model in which one or both chiral components of the fermion ψ is composite, there must occur flavor-diagonal, helicity-conserving contact interactions of the form[9]

$$\mathcal{L}_{\psi\psi} = (g^2/2\Lambda^2)[\eta_{LL}\bar{\psi}_L\gamma_\mu\psi_L\bar{\psi}_L\gamma^\mu\psi_L + \eta_{RR}\bar{\psi}_R\gamma_\mu\psi_R\bar{\psi}_R\gamma^\mu\psi_R$$
$$+ 2\eta_{RL}\bar{\psi}_R\gamma_\mu\psi_R\bar{\psi}_L\gamma^\mu\psi_L] \quad . \qquad (1)$$

In our construction of Eq. (1), we assume that the standard $SU(3)\otimes SU(2)\otimes U(1)$ gauge theory is correct and that $\Lambda \geq \Lambda_{EW}$.[10] Then ψ_L and ψ_R are distinct species and there is no reason why $\mathcal{L}_{\psi\psi}$ should conserve parity. We define Λ in Eq. (1) such that the strong coupling $g^2/4\pi=1$ and the largest $|\eta_{ij}|=1$. Color indices, if any, are suppressed in Eq. (1). Second, if some kinematic region of ψ-ψ elastic scattering is, in the standard theory, controlled by a gauge coupling $\alpha_\psi<<1$, then $\mathcal{L}_{\psi\psi}$ produces interference terms in the cross section of order $(4\pi\alpha_\psi/q^2)^{-1}(g^2/\Lambda^2) = q^2/\alpha_\psi\Lambda^2$ relative to the standard-model contribution.[11] This model-independent effect overwhelms the $O(q^2/\Lambda^2)$ contribution of form factors.

We apply our tests below to high-energy Bhabha scattering $(\alpha_\psi=\alpha)$ and to jet production at high transverse momentum (p_T) in hadron-hadron colliders $(\alpha_\psi=\alpha_{QCD}(q^2))$. It is also important to consider the model-dependent possibility that distinct fermions ψ_1 and ψ_2 have some constituents in common. Then an interaction such as (1) exists, with roughly the same strength, and will modify cross sections for $\psi_1\bar{\psi}_1\to\psi_2\bar{\psi}_2$, $\psi_1\psi_2\to\psi_1\psi_2$ and their $SU(2)_W$ transforms. As an example, we shall consider $e^+e^-\to\mu^+\mu^-$.

Bhabha Scattering. The unpolarized beam cross section, including γ and Z^0 exchanges and $\mathcal{L}_{\psi\psi}$ with $\psi=e$, is given by

$$d\sigma/d(\cos\theta) = (\pi\alpha^2/4s)[4A_0 + A_-(1-\cos\theta)^2 + A_+(1+\cos\theta)^2] \quad ;$$

$$A_0 = (\frac{s}{t})^2|1 + \frac{g_R g_L}{e^2}\frac{t}{t_z} + \frac{\eta_{RL}t}{\alpha\Lambda^2}|^2 \quad ,$$

$$A_- = |1 + \frac{g_R g_L}{e^2}\frac{s}{s_z} + \frac{\eta_{RL}s}{\alpha\Lambda^2}|^2 \quad , \qquad (2)$$

$$A_+ = \frac{1}{2}|1 + \frac{s}{t} + \frac{g_R^2}{e^2}(\frac{s}{s_z} + \frac{s}{t_z}) + \frac{2\eta_{RR}s}{\alpha\Lambda^2}|^2$$
$$+ \frac{1}{2}|1 + \frac{s}{t} + \frac{g_L^2}{e^2}(\frac{s}{s_z} + \frac{s}{t_z}) + \frac{2\eta_{LL}s}{\alpha\Lambda^2}|^2 \quad .$$

In Eq. (2), $t=-s(1-\cos\theta)/2$, $s_z=s-\mu_z^2+i\mu_z\Gamma_z$ and $t_z=t-\mu_z^2+i\mu_z\Gamma_z$; $g_R/e=\tan\theta_W$ and $g_L/e=-\cot 2\theta_W$.[12]

A useful way to search experimentally for electron substructure is to plot the fractional deviation

$$\Delta_{ee}(\cos\theta) = \frac{d\sigma/d(\cos\theta)|_{meas.}}{d\sigma/d(\cos\theta)|_{EW}} - 1 \quad , \qquad (3)$$

where $d\sigma/d(\cos\theta)|_{EW}$ is given by Eq. (2) with $\Lambda=\infty$. Since Δ_{ee} must vanish in the forward direction, the measured cross section can be normalized there to the electroweak value.

We have used Eqs. (2) to calculate Δ_{ee} at $\sqrt{s}=35$ GeV for the cases in which $\mathcal{L}_{ee}$ reduces to the coupling of two left-handed (LL), right-handed (RR), vector (VV) and axial-vector (AA) currents. In the LL model, e.g., $\eta_{LL}=\pm 1$,

$\eta_{RR}=\eta_{RL}=0$. The results are shown in Fig. 1 for values of Λ such that $|\Delta_{ee}|\approx 3$-5% over a wide angular range, consistent with the PETRA measurements.[5] Several comments are in order: (1) For $s\ll\mu_z^2$, the RR model is indistinguishable from LL, because the parity-violating Z^0-terms are negligible there. (2) Greater sensitivity to Λ occurs when both left- and right-handed electron components are composite and have common constituents ($|\eta_{RL}|\approx 1$). (3) Even greater sensitivity to the space-time structure of $\mathcal{L}_{ee}$ may be obtained by using polarized e^+,e^- beams.[13] (4) The PETRA measurements imply the bounds

$$\Lambda(LL,RR) > 750\text{ GeV} \quad ; \Lambda(VV,AA) > 1500\text{ GeV} \quad . \tag{4}$$

Most other physically reasonable models will give bounds lying between these two.

Experiments at higher energy e^+e^- colliders will probe even deeper into the electron. Figure 2 shows Δ_{ee} at $\sqrt{s}=100$ GeV for the same four models. We chose Λ in each case so that $|\Delta_{ee}|\approx 5$-8% over a large angular range. Note the distinctive effects of parity-violating Z^0-terms. We expect that high-luminosity Z^0-factories will be able to set the limits $\Lambda(LL,RR) > 2$ TeV and $\Lambda(VV,AA) > 5$ TeV.[2,13]

$\bar{q}q$ and qq Hard Scattering. The most general $SU(3)\otimes SU(2)\otimes U(1)$-invariant contact interaction involving only light quarks $q_{L,R}=(u,d)_{L,R}$ contains 12 independent helicity-conserving terms.[13] Here, we consider only the simple case

of the product of two left-handed color- and isospin-singlet currents:

$$\mathcal{L}_{qq} = \pm(g^2/2\Lambda^2)\bar{q}_L\gamma_\mu q_L\bar{q}_L\gamma^\mu q_L \quad . \tag{5}$$

We have calculated the cross section for high-p_T jet production using lowest-order QCD and the interaction (5). The contributions of light quark and gluon jets were included. The results are shown in Fig. 3 for $\bar{p}p$ and pp collisions at $\sqrt{s}=2$ TeV. We assume an effect is detectable if it gives a deviation from the expected QCD shape that is at least a factor of two and amounts to at least 100 events/yr. Then, for a $\bar{p}p$ collider with annual integrated luminosity of 10^{37}cm^{-2}, the limit $\Lambda>1.0$ TeV can be set for the interaction (5). The corresponding limit for a pp collider with 10^{40}cm^{-2} is $\Lambda>1.5$-2.0 TeV. More immediately, the CERN $\bar{p}p$ collider, with integrated luminosity 10^{36}cm^{-2}, can limit $\Lambda\gtrsim 250$ GeV for this interaction.

$e^+e^-\to\mu^+\mu^-$. If the electron and the muon have one or more constituents in common, the helicity-conserving terms in their contact interaction are

$$\mathcal{L}_{e\mu} = \frac{g^2}{\Lambda^2}\sum_{i,j=L,R}\eta'_{ij}\bar{e}_i\gamma_\lambda e_i\bar{\mu}_j\gamma^\lambda\mu_j \quad , \tag{6}$$

where Λ, g^2 and η'_{ij} are normalized as in Eq. (1). The fractional deviation $\Delta_{e\mu}$ from the electroweak cross section for $e^+e^-\to\mu^+\mu^-$ has the following properties: (1) Existing measurements at $\sqrt{s}$=35 GeV are consistent with $|\Delta_{e\mu}| \leq$ 6-8% over a wide angular range.[5] This corresponds to Λ>1.4 TeV for LL and RR models and to Λ>2.2 TeV for VV and AA. These bounds on lepton substructure are stronger, but more model-dependent, than those in Eq. (4). Muon decay and ν_μ-e elastic scattering give $\Lambda\geq$6 TeV for a LL isovector interaction and $\Lambda\geq$2 TeV for a LL isoscalar interaction. (2) For $\sqrt{s}<<\mu_z$, $\Delta_{e\mu}(LL)=\Delta_{e\mu}(RR) \propto (1+\cos\theta)^2$. Also, $\Delta_{e\mu}(VV)$=constant, while $\Delta_{e\mu}(AA) \propto \cos\theta$; these effects could be hidden by a normalization error and by the Z^0-induced asymmetry, respectively. (3) Because γ and Z^0 appear only in the s-channel, the beam energy can be tuned to enhance the effect of particular space-time structures in $\Delta_{e\mu}$. When $Re(1+g_ig_js/e^2 s_z) = 0$, the η'_{ij}-contribution is negligible, $\propto s/\Lambda^4$, while the fractional deviations due to other couplings is greater than at nearby energies. This occurs at $\sqrt{s_{LL}}$ = 77.4 GeV, $\sqrt{s_{RR}}$ = 82.2 GeV and $\sqrt{s_{RL}}$ = 115.9 GeV.

Finally, comparable limits on other flavor-nondiagonal interactions can be obtained from existing data on deep-inelastic ν_μ-nucleon scattering, $\bar{p}p\to\mu^+\mu^-X$ at $\sqrt{s}$=2 TeV and e-p collisions at Q^2=(100 GeV)^2.

We have shown that flavor-diagonal contact interactions induced by preon-binding forces significantly alter hard-scattering cross sections at energies well below Λ. Searches for these effects are the most sensitive model-independent tests of quark and lepton substructure.

If $\Lambda\approx$1-5 TeV, deviations from the standard model will soon be observable. The coming generation of multi-TeV colliders should be able to detect substructure up to $\Lambda\approx$10-50 TeV. But, if Λ is only a few TeV, the implications for experiments at these colliders will be more profound. In particular, if $\Lambda\approx$2 TeV, the Bhabha cross section at a 1 TeV×1 TeV linear e^+e^- collider[2] would be $\propto 1/\Lambda^2\approx$0.1 nb, or about 5000 units of R.

This work was begun at the 1982 DPF Summer Study in Snowmass. We thank several participants for stimulating discussions, particularly I. Hinchliffe, H. Kagan, J. Leveille, D. Pellett, M. Perl and H. Wiedemann. K.L. thanks the CERN and Fermilab theory groups for their hospitality during the course of this research.

REFERENCES

1. Two recent reviews are: M.E. Peskin, in Proceedings of the 1981 International Symposium on Lepton and Photon Interactions at High Energy, W. Pfeil, ed., p880 (Bonn, 1981); L. Lyons, Oxford University Publication 52/82 (June 1982).

2. Preliminary accounts of our work appear in the following subgroup reports of the Division of Particles and Fields Summer Study on Elementary Particle Physics and Future Facilities, R. Donaldson, ed. (Fermilab, 1982): M. Abolins, et al.; H. Kagan; F. Bulos, et al..

3. G. 't Hooft, in Recent Developments in Gauge Theories, G.'t Hooft, et al., eds. (Plenum Press, New York, 1980).

4. M.S. Chanowitz and S.D. Drell, Phys. Rev. Lett. 30, 807 (1973).

5. J.G. Branson, in Proceedings of the 1981 International Symposium on Lepton and Photon Interactions at High Energy, W. Pfeil, ed., p.279 (Bonn, 1981); R. Brandelik, et al., (Tasso Collaboration), Phys. Lett. 117B, 365 (1982).

6. I. Bars, in Proceedings of the XVII^{me} Rencontre de

Moriond, J. Tran Thanh Van, ed. (1982). Similar contact interactions and bounds on Λ occur in theories of extended technicolor; see E. Eichten and K. Lane, Phys. Lett. 90B, 125 (1980).

7. See, for example: H. Terazawa, Y. Chikashige, and K. Akama, Phys. Rev. D15, 480 (1977); L. Abbott and E. Farhi, Phys. Lett. 101B, 69 (1981); H. Fritzsch and G. Mandelbaum, Phys. Lett. 102B, 319 (1981).

8. The only stronger bound comes from the rapport between theory and experiment for (g-2) of the muon; this gives $\Lambda \gtrsim 870$ GeV. See R. Barbieri, L. Maiani and R. Petronzio, Phys. Lett. 96B, 63 (1980); S.J. Brodsky and S.D. Drell, Phys. Rev. D22, 2236 (1980).

9. Only if ψ is a Goldstone fermion of spontaneously broken global supersymmetry will its associated dimension-six interactions vanish. See W. Bardeen and V. Visnjić, Nucl. Phys. B194, 422 (1982).

10. By implication, we assume that G_F and $\sin^2\theta_W$ take the values determined from ν_μ-nucleon scattering by ignoring contact interactions. While this assumption is difficult to justify satisfactorily and models exist which violate it (see the last two papers in Ref. 7), it may soon be confirmed by the observation of Z^0 and $W^\pm$ in $\bar{p}p$ collisions.

11. The existence of such relatively large interference terms in Bhabha scattering has been noted by John Preskill, Invited Talk presented at the Vanderbilt Conference on "Novel Results in Particle Physics," May, 1982 (unpublished).

12. We use μ_Z=93.0 GeV, Γ_Z=2.92 GeV and $\sin^2\theta_W$=0.222. See W. Marciano and Z. Parsa, in Proceedings of the Cornell Z^0 Theory Workshop, M.E. Peskin and S.-H.H. Tye, eds. p127.

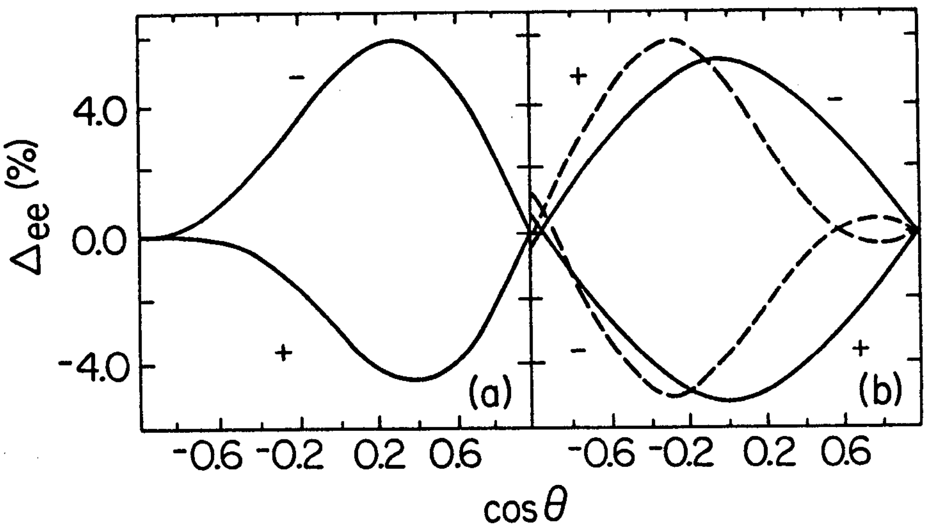
13. This will be discussed in more detail by E. Eichten, H. Kagan and K. Lane, in preparation.

FIGURE CAPTIONS

Fig. 1 $\Delta_{ee}(\cos\theta)$, in per cent, at $\sqrt{s}$=35 GeV. (a) The LL and RR models with Λ=750 GeV. (b) The VV model (solid lines) with Λ=1700 GeV and the AA model (dashed lines) with Λ=1400 GeV. The ± signs refer to the overall sign of the contact interaction in each case.

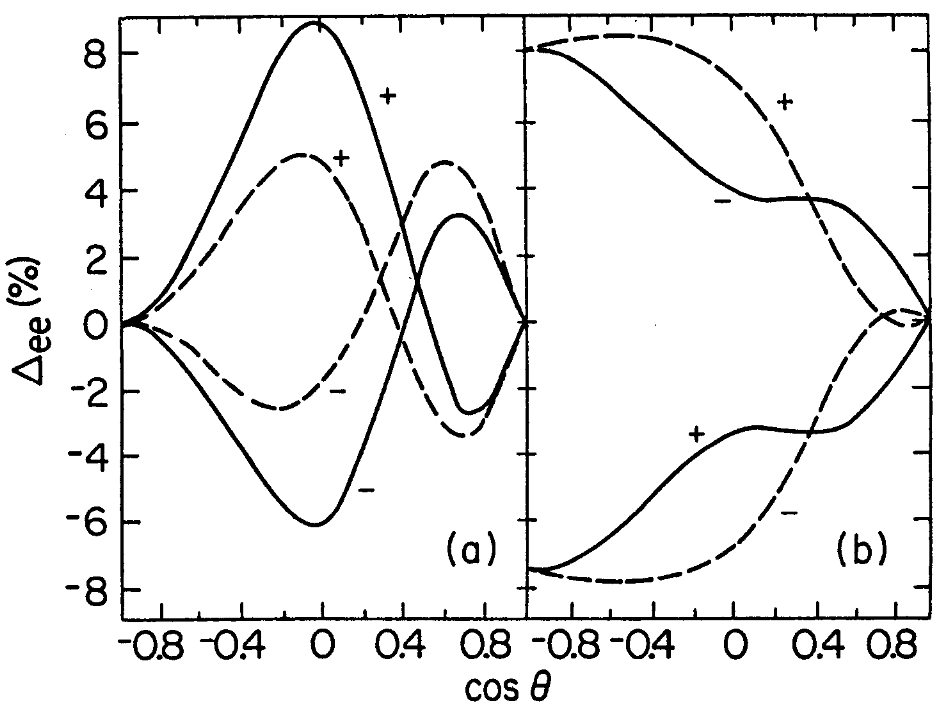
Fig. 2 $\Delta_{ee}(\cos\theta)$, in per cent, at $\sqrt{s}$=100 GeV. (a) The LL model (solid lines) and RR model (dashed lines) for Λ=2 TeV. (b) The VV model (solid) and AA model (dashed) for Λ=5 TeV. The ± signs have the same meaning as in Fig. 1.

Fig. 3 The jet production cross section (in picobarns/GeV) at rapidity y=0 vs. transverse momentum at $\sqrt{s}$=2 TeV in (a) $\bar{p}p$ collisions and (b) pp collisions for various Λ(in TeV). The solid and dashed lines in (b) refer, respectively, to the + and - signs in Eq. (5). Due to a cancellation near y=0, the interference is negligible in (a).



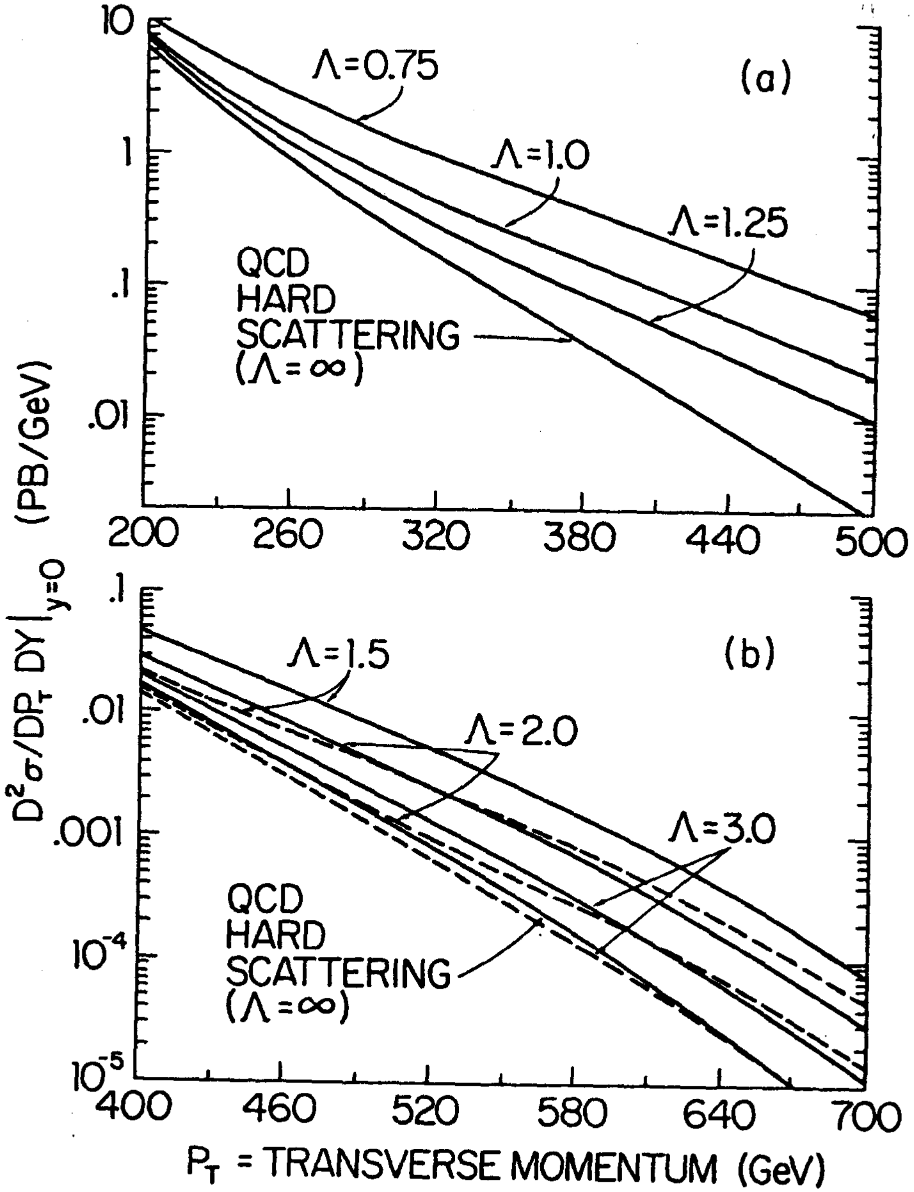
4.0
0.0
-4.0
Δee (%)
(a)
(b)
-0.6 -0.2 0.2 0.6
-0.6 -0.2 0.2 0.6
cos θ

Fig. 1



8
6
4
2
0
-2
-4
-6
-8
Δee (%)
(a)
(b)
-0.8 -0.4 0 0.4 0.8
-0.8 -0.4 0 0.4 0.8
cos θ

Fig. 2



10
1
.1
.01
(a)
Λ=0.75
Λ=1.0
Λ=1.25
QCD
HARD
SCATTERING
(Λ=∞)
200
260
320
380
440
500
(PB/GeV)
D²σ/DP_T DY|_{y=0}
.1
.01
.001
10⁻⁴
10⁻⁵
(b)
Λ=1.5
Λ=2.0
Λ=3.0
QCD
HARD
SCATTERING
(Λ=∞)
400
460
520
580
640
700
P_T = TRANSVERSE MOMENTUM (GeV)

Fig. 3