

Observation of orbital order in the half-filled $4f$ Gd compound

H. Jang,¹ B. Y. Kang,² B. K. Cho,^{2,*} M. Hashimoto,¹ D. Lu,¹ C. A. Burns,^{1,3} C.-C. Kao,⁴ and J.-S. Lee^{1,†}

¹Stanford Synchrotron Radiation Lightsource, SLAC National Accelerator Laboratory, Menlo Park, California 94025, USA

²School of Materials Science and Engineering, Gwangju Institute of Science and Technology, Gwangju 61005, Korea

³Dept. of Physics, Western Michigan University, Kalamazoo, Michigan, 49008, USA

⁴SLAC National Accelerator Laboratory, Menlo Park, California 94025, USA

(Dated: October 27, 2016)

Half-filled electron systems, even with the maximized spin angular moment, have been given little attention because of their zero-orbital angular moment according to Hund's rule. Nevertheless, there are several measurements which show evidence of a non-zero orbital moment as well as spin-orbit coupling. Here we report for the first time the orbital order in a half-filled $4f$ -electron system GdB_4 , using the resonant soft x-ray scattering at Gd $M_{4,5}$ -edges. Furthermore, we discovered that the development of this orbital order is strongly coupled with the antiferromagnetic spin order. These results clearly demonstrate that even in half-filled electron systems, the orbital angular moment can be an important parameter to describe material properties, and may provide significant opportunities for tailoring new correlated electron systems.

PACS numbers: 71.70.Ej, 75.25.Dk, 75.25.-j, 78.70.Ck

According to Hund's rule [1], half-filled electron configurations, such as $3d^5$, $4f^7$, lead to the zero-orbital angular momentum, $\mathbf{L} = 0$, while the spin angular momentum \mathbf{S} , is maximized. In this sense, no spin-orbit coupling is expected in half-filled electron systems. For this reason, therefore, an atomic element which has the half-filled electron configuration has been uninteresting for synthesis of strongly correlated electron systems. On the other hand, nowadays, the spin-orbit coupling interaction has been found to be an important parameter to design emergent materials such as topological insulators and superconductors, in addition to magnetic applications [2–7]. In particular, most exotic phenomena [2] in strongly correlated electron systems and correspondingly their complex phase diagrams [3, 4] are attributed to an inter-coupling between several degrees of freedom, including the spin-orbit coupling interaction [5–7]. To reinforce this interaction in material design, researchers have employed elements which have a large spin and/or orbital moment. In this context, half-filled systems have been given little attention even with the maximized spin moment, because of their zero-orbital moment.

Interestingly, in the last decade, many studies have demonstrated experimentally that the orbital angular momentum in half-filled systems does not follow Hund's rule [8–13]. For example, the non-zero d orbital moment in Fe^{3+} (d^5 configuration) systems, *e.g.*, GaFeO_3 and $\epsilon\text{-Fe}_2\text{O}_3$ compounds, has been observed by soft x-ray magnetic circular dichroism measurements at the Fe L -edge [8, 9]. In the Ca-doped BiFeO_3 compound, the anisotropically reconstructed Fe^{3+} d orbital band was observed by resonant soft x-ray scattering [10]. Similarly, a considerable f orbital moment contribution in the rare earth metal Gd ($4f^7$ configuration) has been observed by using time-resolved x-ray experiments [11, 12]. These experimental demonstrations imply that the orbital angular

momentum is a non-negligible degree of freedom. Moreover, this implication makes a connection with previous findings [14, 15] which observed a change in the g -factor from the free electron value and a splitting of the Gd^{3+} ($4f^7$) Hund's rule ground state manifold, revealing that the spin-orbit coupling effect, even in a half-filled system, is an important parameter. In this context, the half-filled electron systems can be regarded as the strongly correlated electron systems. Several authors have considered the correlated effect theoretically [16, 17]. Nevertheless, the half-filled electron system is still at an early stage in being accepted as an important issue in materials science. This is because no experimental demonstration of spin and orbital degree of freedom under the non-zero spin-orbit coupling has been performed in half-filled electron systems. Therefore, exploring static spin and orbital orders as well as their intercoupling behavior are critical to determine whether a half-filled system can be used to help design new materials.

To assess such a critical advance, we employ GdB_4 as a model system. GdB_4 has a primitive tetragonal structure ($P4/mbm$) with four formula units per lattice point as the basis (Fig. 1(a))[18]. In the Gd metal case, the anisotropic $5d$ band is strongly coupled with the $4f^7$ band via intra-atomic exchange [25], leading to an f -orbital contribution in the half-filled Gd [11, 12]. On the other hand, the ionic state becomes Gd^{3+} in Gd-compounds like GdB_4 , indicating a half-filled $4f^7$ configuration with an empty $5d$ band. However, through inter-atomic hybridization between the Gd $5d$ and B $2p$ states [26], the $5d$ orbital state undergoes a distortion, leading to the anisotropic $5d$ band [26, 27]. As a consequence, the f orbital contribution in GdB_4 is expected to be similar to Gd metal. Moreover, rare earth tetraborides (RB_4) with $\mathbf{L} \neq 0$ rare-earth elements show various ordering phenomena [28–31], such as antiferromagnetic (AFM)

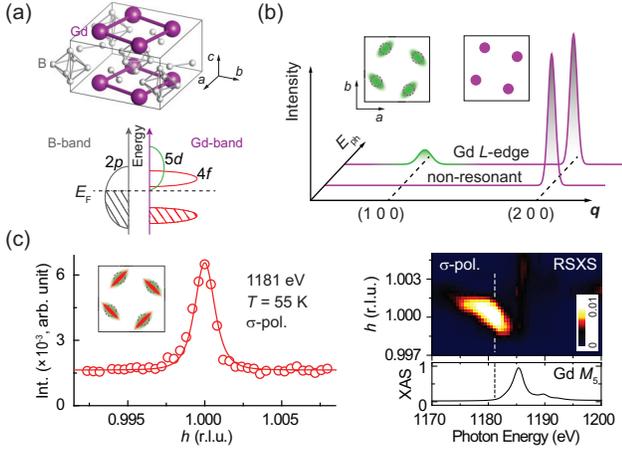


FIG. 1. (color online) (a) Crystal and schematic-band structures of GdB₄. (b) Schematic diffraction patterns along $q = (h 0 0)$ at non-resonant/resonant Gd L -edge (d band). Insets represent Gd positions in the ab -plane (right panel) and anisotropic $5d$ orbitals (left panel). (c) The ATS (1 0 0) reflection (left panel). The inset in the left panel represents an overlay of the anisotropic $5d$ (green colored ellipse) and $4f$ (red colored ellipse) orbitals. (Right panel) Top: the ATS resonant profile. Bottom: the XAS spectrum of GdB₄.

order, quadrupolar order, and structural transitions, as well as their inter-coupling effect. In this work, we employ resonant soft x-ray scattering (RSXS) technique on GdB₄. With exploring f -band through dipole-transition ($3d \rightarrow 4f$), we directly address non-zeroed f orbital behavior even with Gd³⁺ (*i.e.*, half-filled $4f^7$ configuration) and investigate its coupling with the AFM spin order above and below the Néel temperature, $T_N = 42$ K.

As a first step, the f orbital behavior of GdB₄ was explored using RSXS at the Gd M_5 -edge. All RSXS experiments were carried out at beamline 13-3 of the Stanford Synchrotron Radiation Lightsource. Four Gd atoms are located at the $4g$ symmetry points in the crystalline ab -plane of a tetragonal unit cell and B atoms fill the space between Gd atoms (Fig. 1(a)). Beyond the inter-atomic hybridization between the Gd $5d$ -B $2p$ states [26], with the intra-atomic coupling the distorted $5d$ band would also hybridize with the Gd $4f$ band (Fig. 1(a)). Therefore, these intra- and inter-atomic mixing of the Gd $5d$ band lead to an anisotropic f -band tensor of the x-ray susceptibility. A forbidden reflection, namely an anisotropic tensor susceptibility (ATS) reflection [31, 32], is allowed in resonant x-ray scattering (similar to Jahn-Teller assisted orbital behavior in manganites [4]). Figure 1(b) schematically shows the ATS reflection of the Gd $5d$ band at 55 K, revealing a forbidden reflection at $q(H K L) = (2n + 1 0 0)$ seen at the resonant Gd L -edge ($2p \rightarrow 5d$) [31], which originates from the distorted d orbital band within the crystal structure (left inset in Fig. 1(b)). Since a resonance at the Gd $M_{4,5}$ -

edges results from the $3d \rightarrow 4f$ dipole transition, the ATS reflection seen by RSXS indicates an anisotropic f -band in the GdB₄ chemical environment. As we expected, the $4f$ -bands ATS (1 0 0) reflection is clearly observed by the RSXS measured in the paramagnetic state (Fig. 1(c)). The ATS reflection shows resonant enhancement at the photon energy (E_{ph}) ~ 1181 eV. Note that the resonant energy position is slightly lower than the maximum seen in Gd x-ray absorption spectroscopy (XAS). The ATS resonance persists up to room temperature [18]. This indicates that the Gd f orbital band is anisotropic, *i.e.*, $L \neq 0$, via the Gd $4f$ - $5d$ hybridization, even in the half-filled configuration.

At $T > T_N$ where no long-range ordered spin contribution (*i.e.*, AFM order) in GdB₄ [33] is expected, we found clear $4f$ orbital anisotropy due to Gd $4f$ - $5d$ hybridization. To explore coupling between the anisotropic $4f$ orbital and the AFM phase, the RSXS measurements were performed below T_N . Figure 2(a) shows the AFM Gd spin configuration and the RSXS experimental geometry [18]. The scattering plane lies exactly in the crystalline ac -plane. In the crystal's ab -plane, a non-collinear arrangement of Gd spins [34] leads to AFM order, creating scattering intensity at the wavevector $q(H K L) = (1 0 0)$, which is same as the ATS reflection q . With the scattering geometry shown in Fig. 2(a), therefore the (1 0 0) AFM represents Gd spin components projected along the b -axis. Further considering photon polarizations, with $\psi = 0^\circ$ the (1 0 0) AFM scattering intensity can only detect the AFM order through the π' - π channel, $\sim |-\cos \psi|_\pi^2$ (details in Eq.S11 [18]). In other words, it is impossible to detect any magnetic signal with an incident σ polarization at $\psi = 0^\circ$. In this sense, with the incident π polarization, we could observe a clear AFM intensity at $T = 28$ K (Fig. 2(b)). Interestingly, a clear resonant reflection is observed in the σ polarization (see the inset of Fig. 2(b)). Note that the observed intensity in the σ

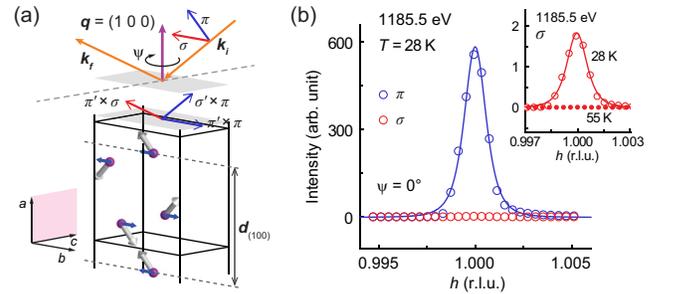


FIG. 2. (color online) (a) The RSXS experimental geometry ($\psi = 0^\circ$) and its polarization configuration, and Gd AFM spin structure. k_i (k_f) denotes incident (scattered) photons. Blue arrows in the spin structure represent the projected Gd spins along the b -axis. $d_{(100)}$ indicates the real space periodicity of (1 0 0) wavevector. (b) (1 0 0) reflections below T_N ($T = 28$ K) in the π channel. The inset shows the σ channels.

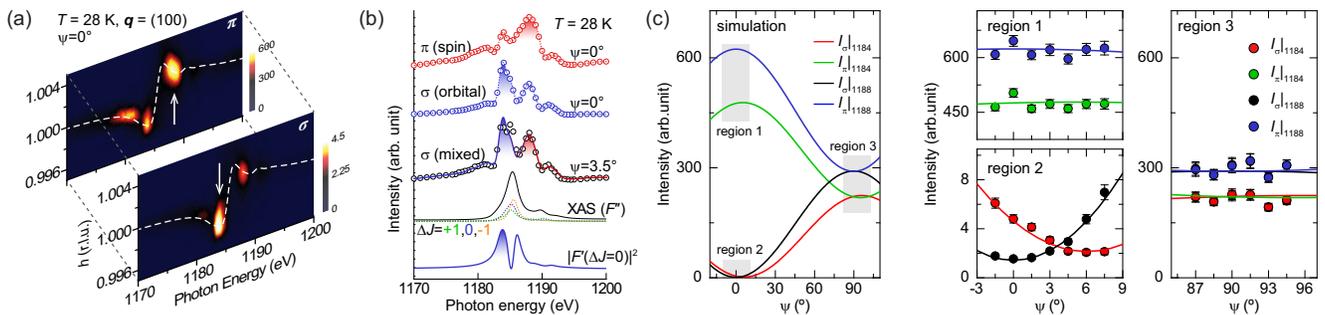


FIG. 3. (color online) (a) Two-dimensional maps of the $(1\ 0\ 0)$ resonant profiles in the σ and π channel. The arrows indicate the strongest resonant positions. The white dashed lines indicate slight variations of the wavevector due to the real part (F') of the atomic form factor. (b) Integrated resonant profiles in the π channel ($\psi = 0^\circ$) and the σ channels ($\psi = 0^\circ$ and 3.5°) as a function of the photon energy. At $\psi = 3.5^\circ$, the black-solid-line is the estimation. XAS spectrum is displayed with three dotted lines ($\Delta J = 0, \pm 1$) [18]. The bottom spectrum shows the squared atomic form factors with $\Delta J = 0$. (c) Simulated azimuthal angle dependences with different photon energies and polarizations (left panel). Comparison between simulated (lines) and experimental (circles) results (middle and right panels). The three regions are denoted as grey areas in the left panel.

polarization is about two orders weaker than that in the π polarization. Hence, we infer that such an unexpected resonance in the σ channel must be associated with either the ATS reflection or something else (*e.g.*, orbital order). Importantly, GdB_4 does not undergo any structure transition as a function of temperature [35] (and See Supplemental Material [18]), indicating a conservation of the crystal symmetry above and below T_N . For this reason, no significant temperature dependence of the ATS reflection is expected. As compared in the inset of Fig. 2(b), however, the σ channel intensity below T_N is much stronger ($\sim \times 100$) than the pure ATS intensity above T_N . Thus, this σ channel feature can be distinguished from the ATS reflection.

In order to investigate the origin of this σ channel feature, the resonant energy profiles of $\mathbf{q} = (1\ 0\ 0)$ in both the σ and π channels were monitored (Fig. 3(a)). Those two resonant profiles are not identical, showing a different intensity distribution as a function of E_{ph} . Note that the slight variation (the white dashed lines) of the \mathbf{q} vector as a function of E_{ph} is proportional to a change in the real part of the atomic form factor [18, 36]. In the π channel case (*i.e.*, the AFM reflection), the resonant maximum is at $E_{ph} = 1188$ eV. Also, there are two relatively weak multiplets of Gd^{3+} at 1181 and 1184 eV. On the other hand, the resonant maximum in the σ channel (1184 eV) is 4 eV below the maximum for the magnetic order. This proves that the σ channel's feature is distinct from the magnetic order. It is worth noting that the small intensity spot in the σ channel that exists at 1188 eV might come from the imperfect incident x-ray polarization (98 %). Moreover, the σ resonant profile is quite different from the ATS one shown in Fig. 1(d). In parallel, this energy selective sensitivity is also supported by analysis of the Gd XAS which consists of a combination of three terms due to the dipole selection

rule $\Delta J = 0, \pm 1$ [18, 37]. In particular, the quadrupole term is proportional to the absorption cross-section for $\Delta J = 0$; *i.e.*, $F'(\Delta J = 0) + iF''(\Delta J = 0)$ [18, 38, 39]. As shown in Fig. 3(b) bottom, the maximum resonance in the scattering profile in the σ channel is qualitatively represented by the $|F'(\Delta J = 0)|^2$ spectrum, revealing that the σ profile originates from the quadrupole moment. In this sense, we clearly exclude both the magnetic order and the pure ATS reflection as candidates for the origin of this σ channel feature.

Another candidate, orbital order, is natural to consider. To check this possibility, the azimuthal angle dependences in both the π and the σ polarizations were performed. Similar to deriving the AFM intensity, firstly we can describe the orbital scattering intensity in the incident σ polarization at $\psi = 0^\circ$, $\sim |\cos \theta \cos \psi|_\sigma^2$ (details in Eq.S15 [18]). Accordingly, we infer that the signal in the incident σ polarization at $\psi = 0^\circ$ is from the orbital. If $\psi \neq 0^\circ$ (or $\neq 90^\circ$), however we must consider that the scattering intensities are from both the AFM and orbital ordering intensities because all components are non-zero. Thereby, we combined the both orders, and describe the total scattering intensities in the σ and π as follows [18]:

$$I_\sigma(\psi) = |-\mathcal{S} \cos \theta \sin \psi + \mathcal{O} \cos \theta \cos \psi|^2 \quad (1)$$

$$I_\pi(\psi) = |\mathcal{S} \cos \theta \sin \psi + \mathcal{O} \cos \theta \cos \psi|^2 + |-\mathcal{S} \cos \psi|^2,$$

where \mathcal{S} and \mathcal{O} is a scattering factors which are proportional to the ordered spin and orbital moment. In this context, with a small variation from $\psi = 0^\circ$, we can easily monitor a mixed feature containing both spin and orbital orders. For example (see Fig. 3(b)), it clearly shows the mixed feature (black circles) between the spin and orbital even with the $\psi = 3.5^\circ$. Using a linear combination of both the σ and π intensities at the $\psi = 0$ (*i.e.*, $I_\sigma(0) \rightarrow$ orbital and $I_\pi(0) \rightarrow$ spin, and experimentally $\mathcal{S} \gg \mathcal{O}$), we produced the mixed feature (black solid line in

Fig. 3(b)), which is well matched with the experimental azimuthal data.

Moreover, we expand the azimuthal angle study at the specifically tuned E_{ph} , because the orbital (1184 eV) and spin (1188 eV) resonances are energetically distinguishable. Figure 3(c) summarized the simulated azimuthal angle dependences (left panel) of Eq. (1) depending on the photon energy and polarizations, as well as comparison in experimental data (middle and right panels), resulting in a good agreement. Note that details of parameters are in Ref.[18]. On the one hand, overall spin's azimuthal dependences at the σ and the π channel are $\sim |\sin(\psi)|^2$ and $\sim |-\cos(\psi)|^2$, respectively. On the other hand, as shown in Fig. 3(c), overall orbital response still looks like $\sim |\sin(\psi)|^2$ even with the σ channel at 1184 eV. This is because the spin ordering signal is predominant even at the orbital resonant energy. However, the actual $I_{\sigma|1184\text{eV}}$ (*i.e.*, orbital response) is non-zero at $\psi = 0$ (see the middle-bottom panel of Fig. 3(c)). The minimum $I_{\sigma|1184\text{eV}}$ locates at $\psi \sim 6^\circ$. According to Eq. (1), this shifted minimum indicates an existence of the non-zero $\mathcal{O}_{1184\text{eV}}$, resulting in an interference effect between the AFM and orbital orders [18]. From these findings, therefore we conclude that the origin of the σ channel is mainly Gd 4*f* orbital order.

Considering the 4*f* orbital configurations in other RB_4 compounds [40, 41], this orbital order might be quadrupolar order. In the other RB_4 compounds with $L \neq 0$ rare earth elements, it is known that the 4*f* quadrupolar order is closely correlated with the in-plane AFM magnetic structures through spin-orbit coupling [31]. For example, an antiferro-quadrupolar (AFQ) order at $\mathbf{q} = (1\ 0\ 0)$ in DyB_4 forms only when the in-plane non-collinear spin components of the AFM order develop [31]. At higher temperature, the AFM spin component is parallel to the out-of-plane direction, and no AFQ order appears. Similarly, we can associate the newly observed Gd 4*f* orbital order at $\mathbf{q} = (1\ 0\ 0)$ with the AFQ order, if we find a similar relationship between the in-plane AFM order and the orbital order. Such relationship was investigated by the temperature dependencies of both the AFM order and the orbital orders (Fig. 4(a)). We found that an onset temperature of the orbital order coincides with T_N , indicating that the newly observed Gd 4*f* orbital order is AFQ order.

While the AFM order shows the typical power-law behavior, however, the temperature behavior of the AFQ is quite distinguishable from that of the AFM. To understand the order parameter shapes, both the AFM and AFQ orders were fit using a power law $[1 - (T/T_c)]^{2\beta}$. The AFM order is well described with T_c and 2β given by 41.8 ± 1 K and 0.78 ± 0.02 , respectively. We note that the estimated β is very similar to the value in the three-dimensional Heisenberg magnet model [42], which is consistent with the non-collinear spin configuration in GdB_4 [34]. On the other hand, in the AFQ case, the 2β is

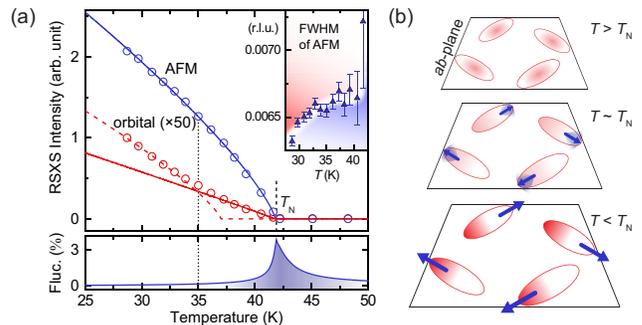


FIG. 4. (color online) (a) Temperature dependences of the integrated intensities of the AFM order (at 1188 eV, π) and orbital order (at 1184 eV, σ) with power-law fits and the simulated spin fluctuation. The inset displays the FWHM (*i.e.*, $1/\text{correlation-length}$) change in the AFM order, as a function of temperature. (b) Schematic picture of the evolutions of both spin (arrows) and orbital (ellipses) orders in the Gd 4*f* band as a function of temperature.

estimated to be 0.97 ± 0.02 , more than 20% slower than the AFM order development. However, this 2β value does not fit below $T \sim 35$ K, due to the relatively faster growth of the AFQ intensity. The order parameter in the low temperature region was fit by $2\beta = 0.79 \pm 0.02$ which is similar to the AFM order, as well as the AFQ order in DyB_4 [31]. These fitting results indicate that the AFQ order in temperature region $\sim 35 \text{ K} < T < 42$ K is somewhat different from the AFQ order below 35 K. To understand this region, we simulated the Gd spin fluctuations using a Ginzburg-Landau free energy model [18]. In GdB_4 case, the simulation result is shown in bottom panel, Fig. 4(a). Below T_N , the spin fluctuation is continuously decreasing. Interestingly, the vanishing point of the fluctuation occurs at $T \sim 35$ K, revealing that the slow development of the AFQ order is likely related with spin fluctuations. When the AFM order is fully developed, the AFQ order becomes reinforced. In other words, this AFQ order is induced by the AFM order through spin-orbit coupling. Furthermore, the AFM order below 35 K is affected by the reinforced AFQ order, resulting in the increased AFM correlation-length (inset in Fig. 4(a)). These implications are schematically summarized in Fig. 4(b). We note that the alignment of the orbital and magnetic orders also provides a possible reason for the AFM pattern found in this system, due to its spin-orbit coupling, and that the AFQ pattern in Fig. 4(b) used the same pattern as DyB_4 [31] because the in-plane AFM patterns in both GdB_4 [34] and DyB_4 [31] are identical.

In summary, we demonstrate for the first time the existence of orbital order in the half-filled 4*f*-electron system, GdB_4 , using resonant soft x-ray scattering. We also reveal that the orbital order is strongly coupled with the antiferromagnetic spin order. Considering these findings as well as the intrinsic advantage such as the maximized

spin moment, the half-filled electron system may be appreciated as a new ingredient for materials design. The current work highlights a compelling new opportunity for tailoring emergent functionality in the strongly correlated electron systems.

We thank Warren E. Pickett, Ki Bong Lee and Young S. Lee for valuable discussions and comments. All synchrotron studies were carried out at the SSRL, a Directorate of SLAC and an Office of Science User Facility operated for the US DOE Office of Science by Stanford University. J.-S.L. acknowledges support by the Department of Energy, Office of Basic Energy Sciences, Materials Sciences and Engineering Division, under contract DE-AC02-76SF00515. C.A.B. was supported by the U.S. Department of Energy, Office of Basic Energy Sciences, Division of Materials Sciences and Engineering, under Award DE-FG02-99ER45772. B.Y.K. and B.K.C were supported by National Research Foundation of Korea (NRF) grants funded by the Korean government (MSIP; Grants No. 2011-0028736 and No. 2013K000315).

* chobk@gist.ac.kr

† jslee@slac.stanford.edu

- [1] C. Kittel, *Introduction to Solid State Physics* (John Wiley and Sons, Inc. New York, 7th ed, 1996).
- [2] L. A. Wray, S.-Y. Xu, Y. Xia, Y. S. Hor, D. Qian, A. V. Fedorov, H. Lin, A. Bansil, R. J. Cava, and M. Z. Hasan, *Nat. Phys.* **6**, 855 (2010).
- [3] S.-W. Cheong and M. Mostovoy, *Nat. Mater.* **6**, 13 (2007).
- [4] E. Dagotto and Y. Tokura, *MRS Bulletin* **33**, 1037 (2008).
- [5] A. D. Caviglia, M. Gabay, S. Gariglio, N. Reyren, C. Cancellieri, and J.-M. Triscone, *Phys. Rev. Lett.* **104**, 126803 (2010).
- [6] V. Galitski and I. B. Spielman, *Nature* **494**, 49 (2013).
- [7] A. Manchon, H. C. Koo, J. Nitta, S. M. Frolov, and R. A. Duine, *Nat. Mater.* **14**, 871 (2015).
- [8] J.-Y. Kim, T. Y. Koo, and J.-H. Park, *Phys. Rev. Lett.* **96**, 047205 (2006).
- [9] Y.-C. Tseng, N. M. Souza-Neto, D. Haskel, M. Gich, C. Frontera, A. Roig, M. van Veenendaal, and J. Nogués, *Phys. Rev. B* **79**, 094404 (2009).
- [10] H. Jang, G. Kerr, J. S. Lim, C.-H. Yang, C.-C. Kao, and J.-S. Lee, *Sci. Rep.* **5**, 12402 (2015).
- [11] M. Wietstruk, A. Melnikov, C. Stamm, T. Kachel, N. Pontius, M. Sultan, C. Gahl, M. Weinelt, H. A. Dürr, and U. Bovensiepen, *Phys. Rev. Lett.* **106**, 127401 (2011).
- [12] C. E. Graves *et al.*, *Nat. Mater.* **12**, 293 (2013).
- [13] G. van der Laan, E. Arenholz, A. Schmehl, and D. G. Schlom, *Phys. Rev. Lett.* **100**, 067403 (2008).
- [14] B. G. Wybourne, *Phys. Rev.* **148**, 317 (1966).
- [15] H. A. Buckmaster and Y. H. Shing, *Phys. Status Solidi A* **12**, 325 (1972).
- [16] R. Skomski, A. Kashyap, and A. Enders, *J. Appl. Phys.* **109**, 07E143 (2011).
- [17] T. Hotta, *J. Phys. Soc. Jpn.* **84**, 114707 (2015).
- [18] See Supplemental Material at <http://> for the experimental/simulation details and additional data, which includes Refs. [19–24].
- [19] B. D. Cullity and S. R. Stock, *Elements of X-ray Diffraction* (Prentice Hall, New Jersey, 3rd ed, 2001) Chap. 4.
- [20] J. P. Hannon, G. T. Trammell, M. Blume, and D. Gibbs, *Phys. Rev. Lett.* **61**, 1245
- [21] S. W. Lovesey and S. P. Collins, *Scattering and Absorption by Magnetic Materials* (Clarendon, Oxford, 1996) Chap. 5.
- [22] H. Wadati *et al.*, *J. Appl. Phys.* **106**, 083705 (2009).
- [23] E. Stavitski and F. M. F. de Groot, *Micron* **41**, 687 (2010).
- [24] P. C. Hohenberg and A. P. Krekhov, *Phys. Rep.* **572**, 1 (2015).
- [25] M. Colarieti-Tosti, S. I. Simak, R. Ahuja, L. Nordström, O. Eriksson, D. Åberg, S. Edvardsson, and M. S. S. Brooks, *Phys. Rev. Lett.* **91**, 157201(2003).
- [26] Z. P. Yin and W. E. Pickett, *Phys. Rev. B* **77**, 035135 (2008).
- [27] S. Ji, C. Song, J. Koo, K.-B. Lee, Y. J. Park, J. Y. Kim, J.-H. Park, H. J. Shin, J. S. Rhyee, B. H. Oh, and B. K. Cho, *Phys. Rev. Lett.* **91**, 257205 (2003).
- [28] D. Okuyama, T. Matsumura, H. Nakao, and Y. Murakami, *J. Phys. Soc. Jpn.* **74**, 2434 (2005).
- [29] D. Okuyama, T. Matsumura, T. Mouri, N. Ishikawa, K. Ohoyama, H. Hiraka, H. Nakao, K. Iwasa, and Y. Murakami, *J. Phys. Soc. Jpn.* **77**, 044709 (2008).
- [30] Z. Fisk, M. B. Maple, D. C. Johnston, and L. D. Woolf, *Solid State Commun.* **39**, 1189 (1981).
- [31] S. Ji, C. Song, J. Koo, J. Park, Y. J. Park, K.-B. Lee, S. Lee, J.-G. Park, J. Y. Kim, B. K. Cho, K.-P. Hong, C.-H. Lee, and F. Iga, *Phys. Rev. Lett.* **99**, 076401 (2007).
- [32] V. E. Dmitrienko, *Acta Crystallograp. Sec. A* **40**, 89 (1984).
- [33] B. K. Cho, J.-S. Rhyee, J. Y. Kim, M. Emilia, and P. C. Canfield, *J. Appl. Phys.* **97**, 10A923 (2005).
- [34] J. A. Blanco, P. J. Brown, A. Stunault, K. Katsumata, F. Iga, and S. Michimura, *Phys. Rev. B* **73**, 212411 (2006).
- [35] J. Fernández-Rodríguez, J. A. Blanco, P. J. Brown, K. Katsumata, A. Kikkawa, F. Iga, and S. Michimura, *Phys. Rev. B* **72**, 052407 (2005).
- [36] R. B. Leighton, *Principles of Modern Physics* (McGraw-Hill, New York, 1959), Chap. 12.
- [37] J. B. Goedkoop, B. T. Thole, G. van der Laan, G. A. Sawatzky, F. M. F. de Groot, and J. C. Fuggle, *Phys. Rev. B* **37**, 2086 (1988).
- [38] H. Ott, C. Schüßler-Langeheine, E. Schierle, A. Yu. Grigoriev, V. Leiner, H. Zabel, G. Kaindl, and E. Weschke, *Phys. Rev. B* **74**, 094412 (2006).
- [39] T. Nagao and J. Igarashi, *J. Phys. Soc. Jpn.*, **77**, 084710 (2008).
- [40] T. Hotta, *Rep. Prog. Phys.* **69**, 2061 (2006).
- [41] Y. Kuramoto, H. Kusunose, and A. Kiss, *J. Phys. Soc. Jpn.* **78**, 072001 (2009).
- [42] J. C. Le Guillou and J. Zinn-Justin, *Phys. Rev. Lett.* **39**, 95 (1977).