

How Resonance-Continuum Interference Changes 750 GeV Diphoton Excess: Signal Enhancement and Peak Shift

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ABSTRACT: The new scalar resonance contribution to the 750 GeV diphoton excess observed at the LHC 13 TeV necessarily interferes with the continuum background in the $gg \rightarrow \gamma\gamma$. The interference has two considerable effects: (1) enhancing or suppressing diphoton signal rate due to the imaginary-part interference and (2) distorting resonance shape due to the real-part interference. From the best-fit study of two benchmark models (two Higgs doublets with ~ 50 GeV widths and a singlet scalar with 5 GeV width, both extended with vector-like fermions), we find that the resonance contribution to the 750 GeV excess can be enhanced by a factor of 2(1.6) for 3(6) fb signal rate and the 68%(95%) CL best-fit mass range can shift by 1–4 (any $\mathcal{O}(1)$) GeV. If the best-fit excess rate decreases with future data, the interference effects will become more significant. The inevitable interferences can also provide a consistency check of a resonance hypothesis, whether or not future precision shape measurements confirm a Breit-Wigner shape or discover interesting deviations.

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1 Introduction

Recently, mild excesses in diphoton invariant mass distribution have been observed in both ATLAS [1] and CMS [2] experiments at Large Hadron Collider (LHC) 13 TeV running. The excesses are 3.6σ and 2.6σ significant from Standard Model (SM) hypothesis, respectively, and are found to prefer a new resonance at around 750 GeV decaying to diphotons [1–3]. The excesses at LHC 13 are currently not completely inconsistent with no significant excesses at LHC 8 TeV data, e.g. [3], and more data are needed to confirm or disfavor the resonance interpretation. The tantalizing hint of a new resonance triggered various theoretical proposals [5–73] allegedly regarded to fit the 750 GeV excess rate $\sim \mathcal{O}(1)$ fb. Also, both a narrow and a somewhat broad resonance with $\Gamma \sim \mathcal{O}(10)$ GeV can fit the data similarly well [1, 3].

The interference between a resonance and the SM continuum background, however, is inevitable [74–81] but has been ignored so far. The interference can have two considerable effects (see, e.g. Ref. [80]):

1. Enhancing or suppressing diphoton signal rate,

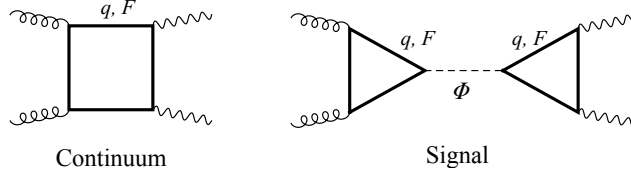


Figure 1. Representative Feynman diagrams of the interfering continuum background (left) and a scalar resonance signal (right) in the $gg \rightarrow \gamma\gamma$.

2. Distorting resonance shape.

The effects can be especially sizable if the resonance width is at least comparable to experimental resolutions or bin sizes, $\Gamma \gtrsim 5$ GeV. For the 125 GeV SM Higgs boson, for example, even though it is narrow, the resulting peak-shift is ~ 70 MeV [76, 78] and will be comparable to the pole-mass measurement uncertainty soon (currently ~ 490 MeV [82]). For a 750 GeV gg -fused scalar resonance with $\mathcal{O}(1)$ fb diphoton rate, the resonance-continuum interference is generally large: the resonance-squared $S \sim \mathcal{O}(1)$ fb and the $gg \rightarrow \gamma\gamma$ continuum background $B \sim 0.2$ fb/40 GeV naively generate $2\sqrt{SB}/S \sim (30 - 90)\%$ relative interference effect. The interference is particularly large in the diphoton channel because the scalar resonance contribution is two-loop suppressed while the interfering continuum background is only one-loop as shown in Fig. 1, so that the above naive estimation of the relative interference is generally loop-factor enhanced [80, 81].

The two main interference effects are induced by different relative phases between the resonance and the continuum processes. The real-part interference ($\propto \cos \phi$ as will be defined and discussed) induces either peak-dip or dip-peak pattern added to a resonance peak, hence distorting the resonance shape from a pure resonance peak. On the other hand, the imaginary-part interference ($\propto \sin \phi$) simply rescales the resonance peak, hence enhancing or suppressing the resonance peak. The non-zero phase is generated when some particles running in loops are lighter than 375 GeV.

In this paper, we investigate each interference effect on the current 750 GeV excess data by considering two benchmark models that exhibit maximally enhanced signals (from the purely imaginary-part interference) or maximally distorted resonance shape (from the purely real-part interference) in the $gg \rightarrow \gamma\gamma$ process. We first describe our method of calculating resonance shapes including interferences in Sec. 2.1 and the diphoton datasets and best-fit analysis method in Sec. 2.2. The two benchmark models are introduced and our main results are discussed in Sec. 3 and Sec. 4. Then we conclude and discuss prospects in Sec. 5.

2 Formalism and Analysis Method

2.1 Diphoton Rate and Resonance Shape

We consider a scalar resonance in the $gg \rightarrow \gamma\gamma$. It interferes with the one-loop continuum backgrounds shown in Fig. 1. The total differential cross section including the interference is written as

$$\begin{aligned} \frac{d\sigma}{dm_{\gamma\gamma}} &= \frac{d\sigma_{\text{bg}}}{dm_{\gamma\gamma}} + \frac{d\sigma_{\text{sig}}}{dm_{\gamma\gamma}} \\ &= \frac{2}{m_{\gamma\gamma}} \mathcal{L}_{gg}\left(\frac{m_{\gamma\gamma}^2}{s}\right) \left[\hat{\sigma}_{\text{bg}}(m_{\gamma\gamma}^2) + \hat{\sigma}_{\text{sig}}(m_{\gamma\gamma}^2) \right], \end{aligned} \quad (2.1)$$

where $\mathcal{L}_{gg}(x) = \int_x^1 dy (x/y) f_{g/p}(y) f_{g/p}(x/y)$ is the gg parton luminosity (we use CT10NNLO PDF set [95]) and $\hat{\sigma}_{\text{bg},\text{sig}}$ are the parton-level cross sections. The signal cross-section $\hat{\sigma}_{\text{sig}}$, the deviation from the SM background, consists of the resonance-squared and the resonance-continuum interference [80],

$$\hat{\sigma}_{\text{sig}}(\hat{s}) = \frac{M^4}{(\hat{s} - M^2)^2 + M^2\Gamma^2} \left[\left(1 + \frac{2\Gamma}{RM} s_\phi \right) \hat{\sigma}_{\text{res}} + \frac{2(\hat{s} - M^2)c_\phi}{M^2} \hat{\sigma}_{\text{int}} \right], \quad (2.2)$$

where $s_\phi = \sin \phi$ and $c_\phi = \cos \phi$, and we factor out Breit-Wigner (BW) parts. We define $\hat{\sigma}_{\text{res},\text{int}}$ and the relative phase ϕ in terms of phase-space integrated squared amplitudes ($\mathcal{M}_i = \mathcal{A}_i e^{i\phi_i}$)

$$\hat{\sigma}_{\text{int}} e^{i\phi} \equiv \frac{1}{32\pi\hat{s}} \int d\cos\theta^* \sum \mathcal{A}_{\text{bg}} \mathcal{A}_{\text{res}} e^{i(\phi_{\text{res}} - \phi_{\text{bg}})}, \quad (2.3)$$

and similarly for $\hat{\sigma}_{\text{res}}$. The summation is over helicity and color indices, and θ^* is the scattering angle in the c.m. frame. We introduce a key parameter R , defined by

$$R \equiv \frac{\hat{\sigma}_{\text{res}}}{\hat{\sigma}_{\text{int}}} \approx \frac{\mathcal{A}_{\text{res}}}{\mathcal{A}_{\text{bg}}}, \quad (2.4)$$

which measures the relative size of interference.

For a narrow resonance (whose width is not much larger than experimental resolutions or bin sizes), the real part interference, the term proportional to c_ϕ in Eq. (2.2), is washed out after the integration over $m_{\gamma\gamma}$. Since the invariant mass distribution is highly accumulated near the resonance peak, we can consider parameters R , ϕ and parton luminosity as constant values. Then the total signal rate with the interference

effect, defined as σ_{mNWA} , is obtained as [80]

$$\sigma_{\text{mNWA}} = \sigma_{\text{NWA}} \cdot C = \left[\frac{M\pi}{\Gamma} \mathcal{L}_{gg}(M^2/s) \hat{\sigma}_{\text{res}}(M^2) \right] \cdot C, \quad \text{for a narrow resonance,} \quad (2.5)$$

where $C = (1 + \frac{2\Gamma}{RM} s_\phi)$ quantifies the strength of the imaginary-part interference. Note that the terms inside the square bracket corresponds to the usual total rate in the narrow width approximation (NWA), production cross section times branching ratio. The subscript mNWA represents *modified* NWA. It is useful to express $\hat{\sigma}_{\text{sig}}(\hat{s})$ in terms of σ_{mNWA} which is measured in experiments:

$$\hat{\sigma}_{\text{sig}}(\hat{s}) = \frac{\Gamma M^3 / (\pi \mathcal{L}_{gg}(M^2/s))}{(\hat{s} - M^2)^2 + M^2 \Gamma^2} \left[\frac{2(\hat{s} - M^2)}{M^2} \frac{c_\phi}{RC} + 1 \right] \cdot \sigma_{\text{mNWA}}. \quad (2.6)$$

This is our resonance shape function for a narrow resonance.

For a broad resonance, with $\Gamma \gtrsim 50$ GeV, we now need to take into account the $m_{\gamma\gamma}$ dependence of R , ϕ and parton luminosity; they are not constant in $m_{\gamma\gamma}$ anymore in broad resonance region. We redefine the total rate σ_{mNWA} for a broad resonance by integrated differential rate, Eq. (2.2), around the resonance mass M :

$$\sigma_{\text{mNWA}} = \int_{M-\Delta}^{M+\Delta} dm_{\gamma\gamma} \left[\frac{d\sigma_{\text{sig}}}{dm_{\gamma\gamma}} \right]_{\text{peak}}, \quad \text{for a broad resonance.} \quad (2.7)$$

We set $\Delta = 100$ GeV for our broad resonance example. We also use the following ratio

$$K_{\text{intf}} = \frac{\sigma_{\text{mNWA}}}{\sigma_{\text{prod}} \cdot \text{Br}_{\gamma\gamma}} \quad (2.8)$$

to quantify the strength of the imaginary-part interference for a broad resonance. This K_{intf} factor is approximately equal to the C factor for a narrow resonance in Eq. (2.5).

The resonance shape function is parameterized not only by usual mass M , width Γ and the total rate σ_{mNWA} but also by the relative interference phase ϕ . R is not a completely independent parameter as shall be discussed. The purely real-part (imaginary-part) interference corresponds to $\phi = 0, 180^\circ$ ($\phi = \pm 90^\circ$). The real-part interference induces peak-dip or dip-peak structure in addition to a BW peak while the imaginary-part interference either enhances or reduces the BW peak or convert the peak to a BW dip (without associated peak) [80]. Thus, the purely real-part interference can most significantly change the resonance shape from a BW peak while the purely imaginary-part interference can most significantly enhance the signal rate (or peak height). These two effects are our main topics; we will study two benchmark models for each of them.

It is hard to carry out a model-independent best-fit analysis including interference effects based on Eq. (2.2) and Eq. (2.6). The interference depends not only on M , Γ , σ_{mNWA} , which are usually chosen in model-independent analysis without interference effects, but also on ϕ and R . In particular, R is correlated with σ_{mNWA} , which is hard to obtain the analytic relation. In this regard we use two benchmark models to numerically discuss the interference effects. For the (purely) real-part interference, we consider a singlet model which introduces a CP-odd SM singlet scalar with a minimal set of vector-like quarks and vector-like leptons: see Sec. 3. For the (purely) imaginary-part interference, Type II 2HDM with vector-like leptons is to be studied: Sec. 4.

There is an important assumption in our implementation of higher-order corrections. We first normalize the total rate without interferences (equivalent to multiplying the correction factor to $\hat{\sigma}_{\text{res}}$) to the result obtained by HIGLU fortran package [88] which includes next-to-next-to-leading-order QCD and next-to-leading-order EW contributions. Then we multiply the same correction factor to the interference term $\hat{\sigma}_{\text{int}}$, as no results are available. Although this assumption approximately accounts for higher-order corrections to the total rate, it implies that $R = \hat{\sigma}_{\text{res}}/\hat{\sigma}_{\text{int}}$ does not receive appreciable higher-order corrections. This may not be an unreasonable assumption since higher-order corrections to the resonance-squared and the resonance-continuum interference can be similar, hence cancelling out in their ratio R . In any case, both the purely real-part and the purely imaginary-part interferences approximately grow with $1/R$. Thus, any corrections to R would directly affect what we discuss in this paper.

2.2 Dataset and Method

In order to quantitatively study interference effects on the 750 GeV diphoton excess data, we perform a Poissonian likelihood analysis to find the best fit. The dataset is from the latest LHC 8 and 13 TeV diphoton resonance search data at around $m_{\gamma\gamma} = 750$ GeV from both ATLAS and CMS experiments. We read in the predicted backgrounds and observed data from the reported plots in Refs. [1, 2, 83, 84]. The total uncertainty in each bin is assumed to be $2(1.5) \times$ statistical uncertainty for LHC 13(8) TeV data.

The fit ranges considered in this paper are

$$\begin{aligned} m_{\gamma\gamma} &= \{630, 830\} \text{ GeV} && \text{for ATLAS 13 (3.2/fb), CMS 13 (2.6/fb), CMS 8 (19.7/fb),} \\ m_{\gamma\gamma} &= \{642, 835\} \text{ GeV} && \text{for ATLAS 8 (20.3/fb).} \end{aligned} \tag{2.9}$$

We choose ATLAS 8 data bins closest to 630 and 830 GeV. The range is somewhat broad so that we can consider a broad resonance as well. CMS 13 dataset is divided into CMS EBEB 13 and CMS EBEE 13 categories depending on which parts of detectors

identify photons. We consider them as independent datasets. Fiducial signal efficiencies are taken from the experimental references and Ref. [3].

We carry out a χ^2 -fit to all the data bins within the range, and take the total change of χ^2 compared to the SM-fit (background-only), $\Delta\chi^2 = \chi^2 - \chi_{\text{SM}}^2$, as a measure of how well the model fits the data. Our SM-fit (background-only) results are:

$$\chi_{\text{SM}}^2 = 7.02, 4.93, 17.77, 1.52, 16.65, \quad (2.10)$$

for ATLAS 13, CMS EBEB 13, CMS EBEE 13, ATLAS 8, and CMS 8, respectively. The results are, of course, sensitive to the assumption of total uncertainties. As will be discussed, although CMS EBEE 13 and CMS 8 show worst fits, these data do not strongly support a 750 GeV resonance – various excesses and deficits around 750 GeV are not significantly fitted better with new resonance contributions. However, ATLAS 13 and CMS EBEB 13 data are fitted better with a new resonance at around 750 GeV. Our read-in data and model-independent fit results without interferences approximately agree with those in Ref. [3]; assuming a BW peak with both fixed $\Gamma = 5, 40$ GeV and with varying Γ .

3 Singlet Model: Real-Part Interference

3.1 Singlet Model

Consider a CP-odd SM-singlet scalar $\Phi = A$, coupling to vector-like quarks $Q \equiv Q^{7/6} = (\mathbf{3}, \mathbf{2}, 7/6)$ and vector-like leptons $L \equiv L^{3/2} = (\mathbf{1}, \mathbf{2}, 3/2)$

$$\mathcal{L} \ni \frac{1}{2}M_\Phi^2\Phi^2 + \sum_Q (s_Q\Phi + M_Q)\bar{Q}\gamma_5 Q + \sum_L (s_L\Phi + M_L)\bar{L}\gamma_5 L, \quad (3.1)$$

where $s_{Q,L}$ are real Yukawa couplings, $M_{\Phi,Q,L}$ mass eigenvalues, $N_{Q,L}$ number of fermions, and $q_{Q,L}$ electric charges. We choose $Q^{7/6}$ and $L^{3/2}$ from the minimal matter list [85] – the list of new particles that can eventually decay to SM particles – since they have the largest electric charges. We consider A , but H shall also exhibit similar effects.

In the quark sector, we introduce a single vector-like Q with fixed parameters

$$M_Q = 1 \text{ TeV}, N_Q = 2, s_Q = 0.2. \quad (3.2)$$

We still have enough lepton sector free parameters that we can use to fit the data and to illustrate interference effects.

In the lepton sector, we consider

$$M_L = 400 \text{ GeV}, N_L = 6, s_L \text{ is varied.} \quad (3.3)$$

The sign of the Yukawa s_L determines the sign of the relative phase: $s_L \rightarrow -s_L$ approximately changes the relative phase $\phi \rightarrow \pi + \phi$. It is an approximate relation because Q also contributes to the $\Phi \rightarrow \gamma\gamma$ part although it is subdominant to the L contribution. We will compare the results with positive and negative s_L (as well as with the results without any interference accounted for) to see how the best-fit changes with interference effects.

Another important parameter is the width. In the above model, the width is typically too small ($\lesssim 1 \text{ GeV}$) to make interference effects apparent in current experiments; Φ mainly decays to loop-induced gg and $\gamma\gamma$

$$\Gamma(\Phi \rightarrow gg) = \frac{\alpha_S^2}{128\pi^3} \frac{M_\Phi^3}{M_Q^2} \left| \sum_Q s_Q A_{1/2}^\Phi \left(\frac{M_\Phi^2}{4M_Q^2} \right) \right|^2, \quad (3.4)$$

$$\Gamma(\Phi \rightarrow \gamma\gamma) = \frac{\alpha^2}{256\pi^3} M_\Phi^3 \left| \sum_{f=Q,L} N_C q_f^2 \frac{s_f}{M_f} A_{1/2}^\Phi \left(\frac{M_\Phi^2}{4M_f^2} \right) \right|^2, \quad (3.5)$$

where loop functions $A_{1/2}^{\Phi=H,A}$ are defined as in Ref. [89], and other signals such as $Z\gamma$, ZZ , WW are currently well below their LHC 8 sensitivities. If such a narrow resonance falls within a single experimental bin, the real-part interference (although itself is independent on the width) is cancelled out; in addition, the imaginary-part interference is small since it is directly proportional to the width as $C - 1 \propto \Gamma$. Thus, to illustrate possible impacts of interference effects, we assume a bigger constant width

$$\Gamma_\Phi = 5 \text{ GeV}. \quad (3.6)$$

It is easy to add extra hidden decay modes of Φ , not constrained at all, to make so. If the assumed width were much bigger than the true width, diphoton signal will be suppressed; but if the true width is bigger, the interference will become more relevant. Also, if the N_L were smaller, although one can still have almost 100% $\text{BR}(\gamma\gamma)$, the total width decreases and the interference effects will be less significant. Meanwhile, for $M_L \leq M_\Phi/2$, the decays into vector-like leptons dominate and the diphoton signal becomes too suppressed. Although such light leptons can change the phase ϕ and introduce different interference effects, we cannot fit the diphoton excess data well and do not discuss this possibility further.

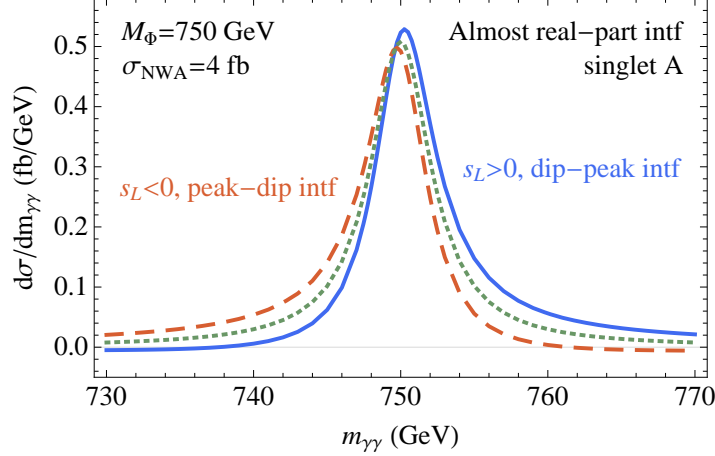


Figure 2. Example diphoton resonance shapes with $s_L > 0$ (dip-peak interference, **blue-solid**), $s_L < 0$ (peak-dip interference, **red-dashed**), and no interference (**green-dotted**) for the same mass $M_\Phi = 750$ GeV and the NWA rate $\sigma_{\text{NWA}} \simeq 4$ fb. The relative phase $\phi \simeq 8.3^\circ(+\pi)$ for $s_L > 0(< 0)$ induces almost purely real-part interference, and the resulting peak shifts and long tails affect best-fit analysis. The small imaginary-part interference also makes true observable mNWA rates σ_{mNWA} and peak heights slightly different. We set $|s_L| \simeq 1.5$ and $\Gamma_\Phi = 5$ GeV.

An important feature of the singlet scalar model is that the relative phase is small:

$$\phi \simeq \begin{cases} 8.3^\circ & \text{for } s_L > 0; \\ 188.3^\circ & \text{for } s_L < 0, \end{cases} \quad (3.7)$$

which induces almost purely real-part interference. This is the case in which resonance shape is maximally distorted from pure BW shape (and the peak location is maximally shifted), for the given total rate. The small but non-zero phase is generated from the SM quark loops in $gg \rightarrow \gamma\gamma$ background box diagrams.

3.2 Results – Singlet Model

In Fig. 2 we show an example of the SM-singlet scalar resonance shapes for $s_L > 0$ (blue-solid) and $s_L < 0$ (red-dashed) with full interference effects. For comparison, we also show the resonance shape without any interferences taken into account (green-dotted). All three cases have the same NWA rates and the width Γ_Φ . But $s_L > 0(< 0)$ induces a small dip-peak (peak-dip) interference pattern added to the BW peak, so that a long tail toward a high (low) invariant mass region appears and the peak shifts toward the same direction. As a result, the best-fit results change, even though the NWA rates, masses and widths are all the same. We quantify such interference effects in

this subsection. The small but non-zero imaginary-part interference, Eq. (4.9), actually makes σ_{mNWA} (true observable rate slightly different from the NWA rate; see Eq. (2.5)) and the peak heights slightly different among the three shapes.

Fig. 3 shows the best-fit results to ATLAS 13 (left) and CMS EBEB 13 (right) datasets individually, for a singlet scalar $\Phi = A$ model with $s_L > 0$ (upper) and $s_L < 0$ (lower). For comparison, we also show the results without any interferences accounted for (dashed). These datasets are the ones that most strongly prefer the existence of a 750 GeV resonance, and the interference effect does not change the preference of the resonance existence; the data fit much better with a new resonance around 750 GeV even with interferences. Comparing the upper panels for $s_L > 0$ with the lower panels for $s_L < 0$, we find that the 68% CL best-fit mass parameter is shifted by about 1–4 GeV while a much bigger shift $\mathcal{O}(1)$ GeV is expected for the 95% CL region or for weaker couplings s_L . Meanwhile, similar magnitudes of couplings are preferred regardless of interference effects. For $s_L > 0$ with dip-peak interference, the peak shifts toward high-mass region and the high-mass region is more accumulated (see Fig. 2); consequently, somewhat smaller masses are preferred compared to the $s_L < 0$ case (and to the case without interferences).

The interference effects are still apparent, even after including all other LHC 13 and LHC 8 datasets that do not strongly prefer the existence of an additional resonance. This is shown in Fig. 4; the 68% CL best-fit regions again shift by about 1–4 GeV and a bigger shift is expected for the 95% CL region or for weaker couplings s_L . The preference of an additional resonance also still exists with interference effects.

There is a noticeable tendency that interference effects become stronger with a weaker s_L , as can be deduced from a wider best-fit mass shift with a weaker s_L in Fig. 3 and Fig. 4. This is a general result of interference; the real-part interference approximately grows with $1/R \sim \mathcal{A}_{\text{bg}}/\mathcal{A}_{\text{res}}$ amplitude ratio, which measures the background-resonance interference contribution compared to the resonance-squared contribution. If future data prefer to a weaker signal, the interference effects will be larger and more important.

Finally, we briefly compare best-fit results to various datasets. that compared to the ATLAS 13 result in Fig. 3, the CMS EBEB 13 prefers to a resonance with a slightly higher mass and weaker coupling. But the preferences of a new resonance around 750 GeV from both data are consistent with each other. Including LHC 8 datasets in Fig. 4 significantly prefers to a weaker coupling and actually worsens the best-fit (total $|\Delta\chi_{\text{min}}^2|$ in the right panel decreased from the left panel). This may imply that the LHC 8 datasets do not strongly favor the resonance contribution. Future data can only clarify the origin of the excess.

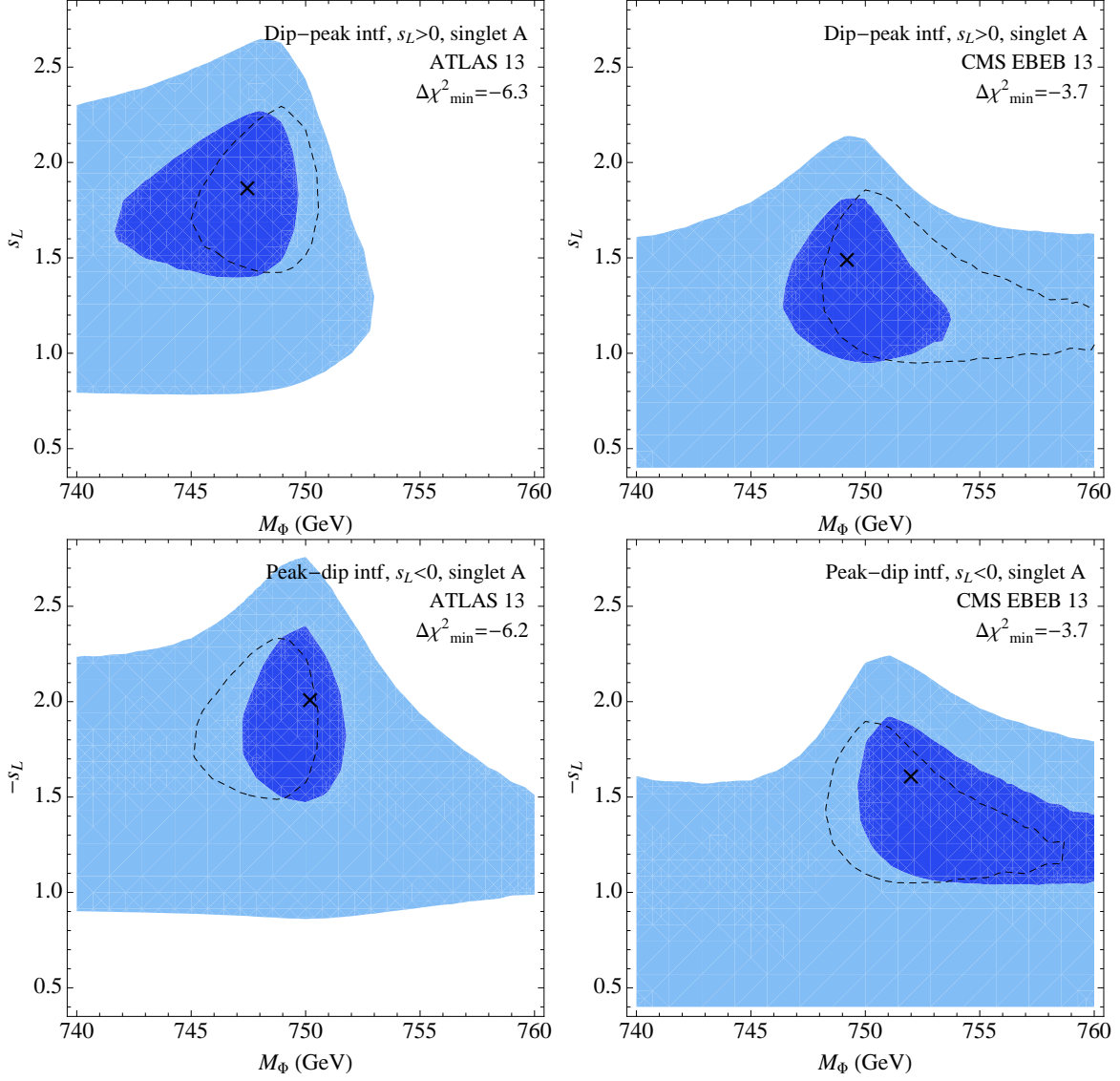


Figure 3. The 68% CL(darker blue) and 95% CL(lighter blue) preferred regions for CP-odd singlet A . $s_L > 0$ (**upper**) and $s_L < 0$ (**lower**) can be compared with each other (and with dashed lines for the 68% CL results without interferences accounted for) to see interference strength. $\Gamma_\Phi = 5$ GeV. Fit is performed for $m_{\gamma\gamma} = 630 - 830$ GeV from ATLAS 13 (**left**, $\chi_0^2 = 7.02$) and CMS EBEB 13 (**right**, $\chi_0^2 = 4.93$) datasets. The best-fit mass ranges with positive and negative s_L differ by $\mathcal{O}(1)$ GeV and the difference is bigger with weaker s_L . The best-fit $\Delta\chi_{\min}^2$ compared to the SM fit χ_0^2 is also shown.

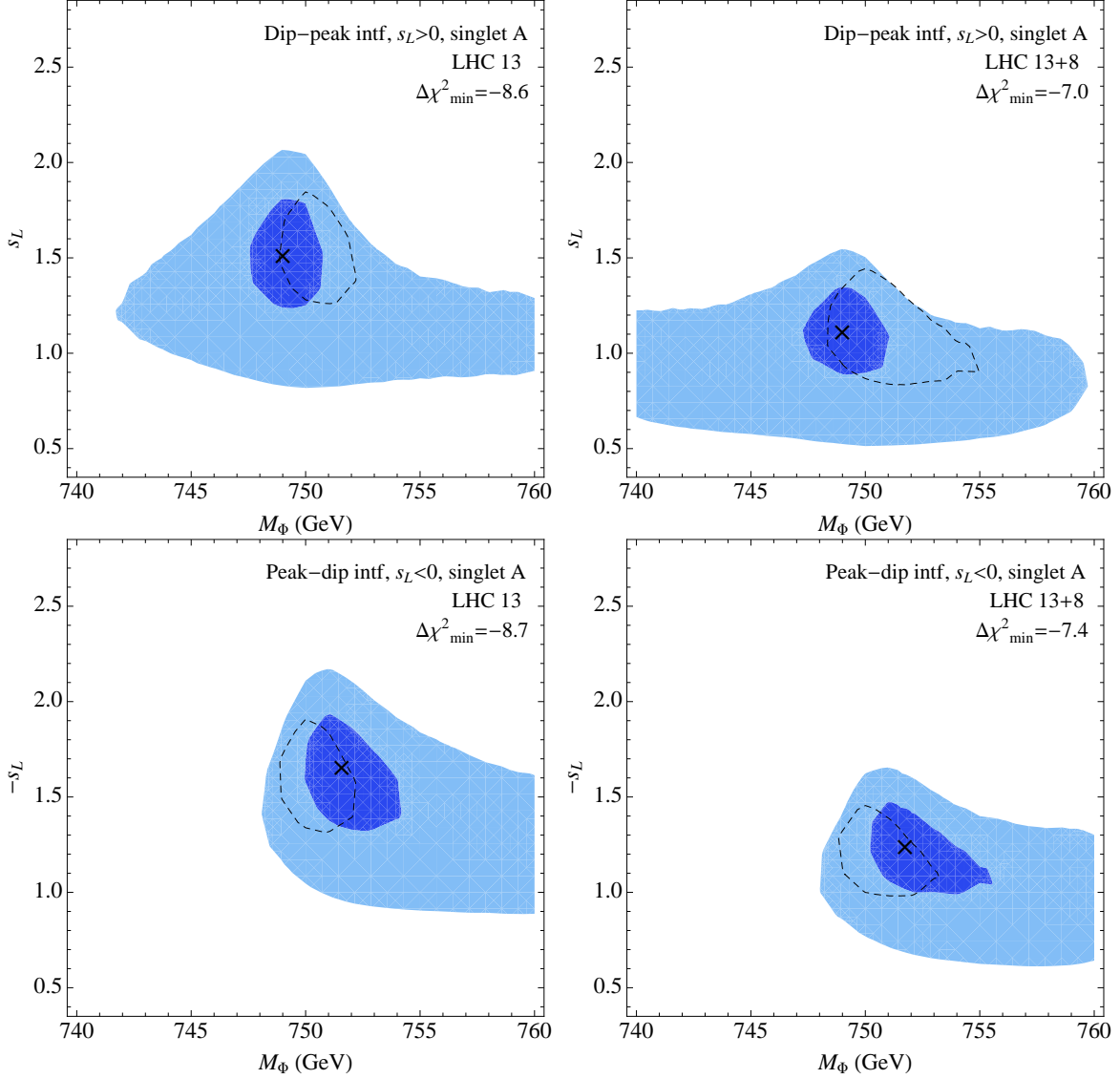


Figure 4. The 68% CL(darker blue) and 95% CL(lighter blue) preferred regions for CP-odd singlet A . $s_L > 0$ (**upper**) and $s_L < 0$ (**lower**) can be compared with each other (and with dashed lines for the 68% CL results without interferences accounted for) to see interference strength. $\Gamma_\Phi = 5$ GeV. Fit is performed for $m_{\gamma\gamma} = 630 - 830$ GeV from LHC 13 (**left**, $\chi_0^2 = 29.7$) and LHC 13+8 (**right**, $\chi_0^2 = 47.9$) datasets. The best-fit mass ranges with positive and negative s_L differ by $\mathcal{O}(1)$ GeV and the difference is bigger with weaker s_L . The best-fit $\Delta\chi_{\min}^2$ compared to the SM fit χ_0^2 is also shown.

		$SU(3) \times SU(2) \times U(1)_Y$
$L_L = \begin{pmatrix} E_L \\ D_L \end{pmatrix}$	$L_R = \begin{pmatrix} E'_R \\ D'_R \end{pmatrix}$	$(\mathbf{1}, \mathbf{2}, -\frac{3}{2})$
E_R	E'_L	$(\mathbf{1}, \mathbf{1}, -1)$
D_R	D'_L	$(\mathbf{1}, \mathbf{1}, -2)$

Table 1. The contents and quantum numbers of vector-like leptons in the VLL-2HDM model. The electric charges of the doublet components are $(-1, -2)$.

4 VLL-2HDM: Imaginary-Part Interference

4.1 VLL-2HDM Model

We consider the Type II two-Higgs-Double-Model (2HDM) in the alignment limit extended with extra vector-like leptons (VLL). We first summarize the (heavy) Higgs sector and then introduce the VLL sector.

The Higgs sector consists of three neutral Higgs bosons, h , H (scalar), A (pseudo-scalar) and two charged Higgs bosons H^\pm . In the alignment limit, the h is the 125 GeV SM Higgs boson, and the heavier Higgs bosons, H and A , are our focus in this paper. To be consistent with electroweak precision data and to explain the diphoton excess, we consider a degenerate heavy Higgs bosons

$$M_\phi = M_H = M_A = 750 \text{ GeV} . \quad (4.1)$$

In the alignment limit with small $t_\beta \sim 1$ as we will focus in this paper¹, the 2HDM alone cannot explain the diphoton excess. It is mainly because the heavy Higgs bosons dominantly decay to the top pair (the decays to ZZ and WW are forbidden in the alignment limit), and the relevant the diphoton branching ratio is only $\text{Br}(\phi \rightarrow \gamma\gamma) = 7.8(8.7) \times 10^{-6}$ for $\phi = H(A)$, leading to too small signal rates $\sigma(pp \rightarrow H, A \rightarrow \gamma\gamma) = 0.012 \text{ fb}$. Thus, we extend the model by extra VLL to achieve the needed ~ 400 enhancement of diphoton signal.

We now introduce VLLs, $L_L, E_R, D_R, E'_L, D'_L$, of which the quantum numbers are summarized in Table 1. Note that the electric charges of $E^{(\prime)}$ and $D^{(\prime)}$ are -1 and -2 , respectively. All of the VLLs in Table 1 are imbedded in one family. The Lagrangian

¹The parameter space of $t_\beta \gtrsim 25$ at $M_\phi = 750 \text{ GeV}$ is excluded by the 8 TeV LHC data [90].

of the VLLs in Type II 2HDM is

$$\begin{aligned}
-\mathcal{L} = & Y_D \bar{L}_L H_1 D_R + Y'_D \bar{L}_R H_1 D'_L + Y_E \bar{L}_L \tilde{H}_2 E'_R + Y'_E \bar{L}_R \tilde{H}_2 E'_L \\
& + \left[M \bar{L}_L L_R + M_E \bar{E}'_L E_R + M_D \bar{D}'_L D_R + \text{h.c.} \right].
\end{aligned} \tag{4.2}$$

The mass matrix in the basis of (E, E') is

$$\mathcal{M}_E = \begin{pmatrix} M & \frac{1}{\sqrt{2}} Y_E v_2 \\ \frac{1}{\sqrt{2}} Y'_E v_2 & M_E \end{pmatrix}. \tag{4.3}$$

We have similar form of \mathcal{M}_D by changing $Y_E^{(\prime)} \rightarrow Y_D^{(\prime)}, v_2 \rightarrow v_1, M_E \rightarrow M_D$. We focus on the no-mixing case, which is possible if $M_E \gg M, Y_E v_2$ and $M_D \gg M, Y_D v_1$. Then the light masses of E and D are degenerate as $M_{E_1} = M_{D_1} = M$. The heavy masses are $M_{E_2} = M_E$ and $M_{D_2} = M_D$, which suppresses the contribution from E_2 and D_2 . We do not consider the mass M below $M_\Phi/2$ since the new decay channels of $H/A \rightarrow E\bar{E}/D\bar{D}$ raise the total width quickly. We also assume that $Y_E = Y'_E$ and $Y_D = Y'_D$ for simplicity.

The Yukawa terms for the VLLs in the mass eigenstate basis become

$$\begin{aligned}
-\mathcal{L}_{\text{Yukawa}} = & -\frac{1}{t_\beta} y_E H (\bar{E}_1 E_1 + \bar{E}_2 E_2) + t_\beta y_D H (\bar{D}_1 D_1 + \bar{D}_2 D_2) \\
& -i \frac{1}{t_\beta} y_E A (\bar{E}_1 \gamma_5 E_1 + \bar{E}_2 \gamma_5 E_2) - i t_\beta y_D A (\bar{D}_1 \gamma_5 D_1 + \bar{D}_2 \gamma_5 D_2) \\
& + y_E h (\bar{E}_1 E_1 + \bar{E}_2 E_2) + y_D h (\bar{D}_1 D_1 + \bar{D}_2 D_2),
\end{aligned} \tag{4.4}$$

where $y_E = s_\beta Y_E / \sqrt{2}$ and $y_D = c_\beta Y_D / \sqrt{2}$.

The partial decay widths of $\Phi = h, H, A$ in the VLL-2HDM are

$$\begin{aligned}
\Gamma(\Phi \rightarrow gg) = & \frac{G_F \alpha_s^2 M_\Phi^3}{64 \sqrt{2} \pi^3} \left| \sum_q \hat{y}_q^\Phi A_{1/2}^\Phi(\tau_q) \right|^2, \\
\Gamma(\Phi \rightarrow \gamma\gamma) = & \frac{G_F \alpha_e^2 M_\Phi^3}{128 \sqrt{2} \pi^3} \left| \sum_q \hat{y}_q^\Phi N_c Q_q^2 A_{1/2}^\Phi(\tau_q) + \sum_\ell \hat{y}_\ell^\Phi Q_\ell^2 A_{1/2}^\Phi(\tau_\ell) + \mathcal{A}_{\gamma\gamma, \text{VLL}}^\Phi \right|^2,
\end{aligned} \tag{4.5}$$

where $\tau_f = M_\Phi^2 / (4m_f^2)$, the relative Yukawa couplings normalized by the SM values are $\hat{y}_{t,b,\tau}^h = 1$, $\hat{y}_t^{H,A} = \mp 1/t_\beta$ and $\hat{y}_{b,\tau}^{H,A} = t_\beta$ for Type II in the aligned 2HDM, and the loop functions $A_{1,1/2}^{H/A}(\tau)$ are referred to Ref. [89]. The VLL contributions $\mathcal{A}_{\gamma\gamma, \text{VLL}}^\Phi$ in

Eq. (4.5) are given as

$$\mathcal{A}_{\gamma\gamma, \text{VLL}}^\Phi = \sum_{\text{VLL}} \sum_{i=1,2} \left[Q_{E_i}^2 \frac{\hat{y}_t^\Phi y_{E_i}^v}{M_{E_i}} A_{1/2}^\Phi(\tau_{E_i}) + Q_{D_i}^2 \frac{\hat{y}_b^\Phi y_{D_i}^v}{M_{D_i}} A_{1/2}^\Phi(\tau_{D_i}) \right]. \quad (4.6)$$

In order to greatly enhance the $H/A \rightarrow \gamma\gamma$ partial decay width through VLL loop one needs multiple number of VLL families. In the following analysis we introduce 3 VLL families. We vary M from 375 GeV to 600 GeV and $y_{E,D}$ from -4π to 4π .

The final comment is on the constraint from the Higgs precision data. As shown Eq. (4.5), the VLL loop also contributes to $h \rightarrow \gamma\gamma$, which is already very limited by the 8 TeV LHC data. If two Yukawa couplings y_D and y_E are tuned as

$$y_D = -\frac{Q_E^2}{Q_D^2} y_E = -0.25 y_E, \quad (4.7)$$

new contribution to the Higgs precision data vanishes if E and D are degenerate in mass M . If $t_\beta = 1$, the cancellation of the VLL contributions to $h \rightarrow \gamma\gamma$ equally happens to the $A \rightarrow \gamma\gamma$ decay. Since the A diphoton signal is usually larger than the H one if no cancellation occurs, we choose $t_\beta = 0.7$ in the analysis. Other exclusion limits from $Z\gamma$ [91], $b\bar{b}$ [92], $\tau^+\tau^-$ [93], and jj [94] channels at the 8 TeV LHC are satisfied in the parameter space under consideration.

4.2 Results – VLL-2HDM Model

We first discuss the total widths of H and A , both of which are dominated by the $t\bar{t}$ decay channel. Using the running top quark mass $m_t(\mu = 750 \text{ GeV}) = 147 \text{ GeV}$ [96], we have $\Gamma_{H(A)} = 46(58) \text{ GeV}$. Since the degenerate H and A do not interfere, we treat them as BW peaks. We perform a minimum χ^2 analysis (see Sec. 2.2) and find the best-fit signal rates to the LHC 13+8 datasets

$$\sigma(pp \rightarrow \Phi \rightarrow \gamma\gamma) = \begin{cases} 6.5 \pm 2.5 \text{ fb} & (68\% \text{CL}) \\ 6.5_{-3.5}^{+4.5} \text{ fb} & (95\% \text{CL}) \end{cases}, \quad (4.8)$$

which are in agreement with Ref. [3].

In our scenario of VLL-2HDM the relative interference phase corresponds to almost imaginary interference:

$$\phi \simeq \begin{cases} 90^\circ & \text{for } y_E > 0; \\ -90^\circ & \text{for } y_E < 0. \end{cases} \quad (4.9)$$

The reasons are as follows. The complex phase from the continuum background amplitude is minor [81]. But the production part $gg \rightarrow H/A$ is dominated by top quark loop

and the loop function generate large complex phase: $77^\circ(91^\circ)$ for $M_H(M_A) = 750 \text{ GeV}$. In addition, the decay part $H/A \rightarrow \gamma\gamma$, dominated by VLL loop contribution, is also real since $M_L > M_{H,A}/2$ in our scenario. Depending on the sign of Yukawa coupling y_E , the whole complex phase is changed by π . It maximally enhances the signal rate for $\phi \approx 90^\circ$ (constructive interference) and maximally suppress the signal rate for $\phi \approx -90^\circ$ (destructive interference).

Figure 5 shows our results in the parameter space (y_E, M) for the VLL-2HDM: the K_{intf} in Eq. (2.8) (upper panels) and the allowed parameter space by the 750 GeV diphoton excess data (lower panels). It is of great interest that quite large interference effects (large $K_{\text{intf}} - 1$) appear around the measured total signal rate, as shown in Figs. 5(a) and (b). For $y_E > 0$, $K_{\text{intf}} > 1$ and thus constructive interference occurs: the interference effect can make even factor 2 for the 3 fb total rate. Within the 68% CL best-fit signal rate the interference effect ranges from 40% to 80% when $y_E > 0$. For $y_E < 0$, $K_{\text{intf}} < 1$ so that destructive interference occurs: in order to explain the signal rate, we need quite large magnitude of y_E and thus very limited parameter space is allowed. Figures 5(c) and (d) show that the allowed parameter space significantly change by including the interference effect. For comparison, we show the allowed parameter region without including interference effects. With positive y_E and $M = 400 \text{ GeV}$, for example, required y_E for the signal rate 6.5 fb is reduced from ~ 7.5 to ~ 5.5 by including the interference effects. Equivalently, the change of required number of VLL family is from 3 to 4. In all, interference effects have significant implications on the underlying physics.

$M[\text{GeV}]$	y_E	$\phi^H[^\circ]$	$\phi^A[^\circ]$	K_{intf}^H	K_{intf}^A	K_{intf}^{H+A}	$\sigma_{\text{mNWA}}^{H+A}[\text{fb}]$
457	2	99	123	2.6	5.9	3.5	1
413	4	93	108	1.6	3.0	2.0	3
400	6	91	104	1.3	2.1	1.6	6
385	-5	-96	-88	0.38	0.20	0.32	1
395	-8	-95	-86	0.54	0.21	0.43	3

Table 2. Numerical values for ϕ , K_{intf} , and σ_{mNWA} for H , A and the total. The benchmark parameter points are chosen to yield total signal rates of 1, 3, 6 fb. VLL-2HDM model.

In Table 2, we present the numerical values for ϕ , K_{intf} , and σ_{mNWA} . The benchmark parameter points are chosen to yield total signal rates of 1, 3, 6 fb. In order to see the individual interference effects, we show ϕ and K_{intf} for H and A separately. For both H and A , the relative interference phase is about $\pm 90^\circ$: almost purely imaginary interference occurs. K_{intf}^H and K_{intf}^A show that the interference effects are larger for A

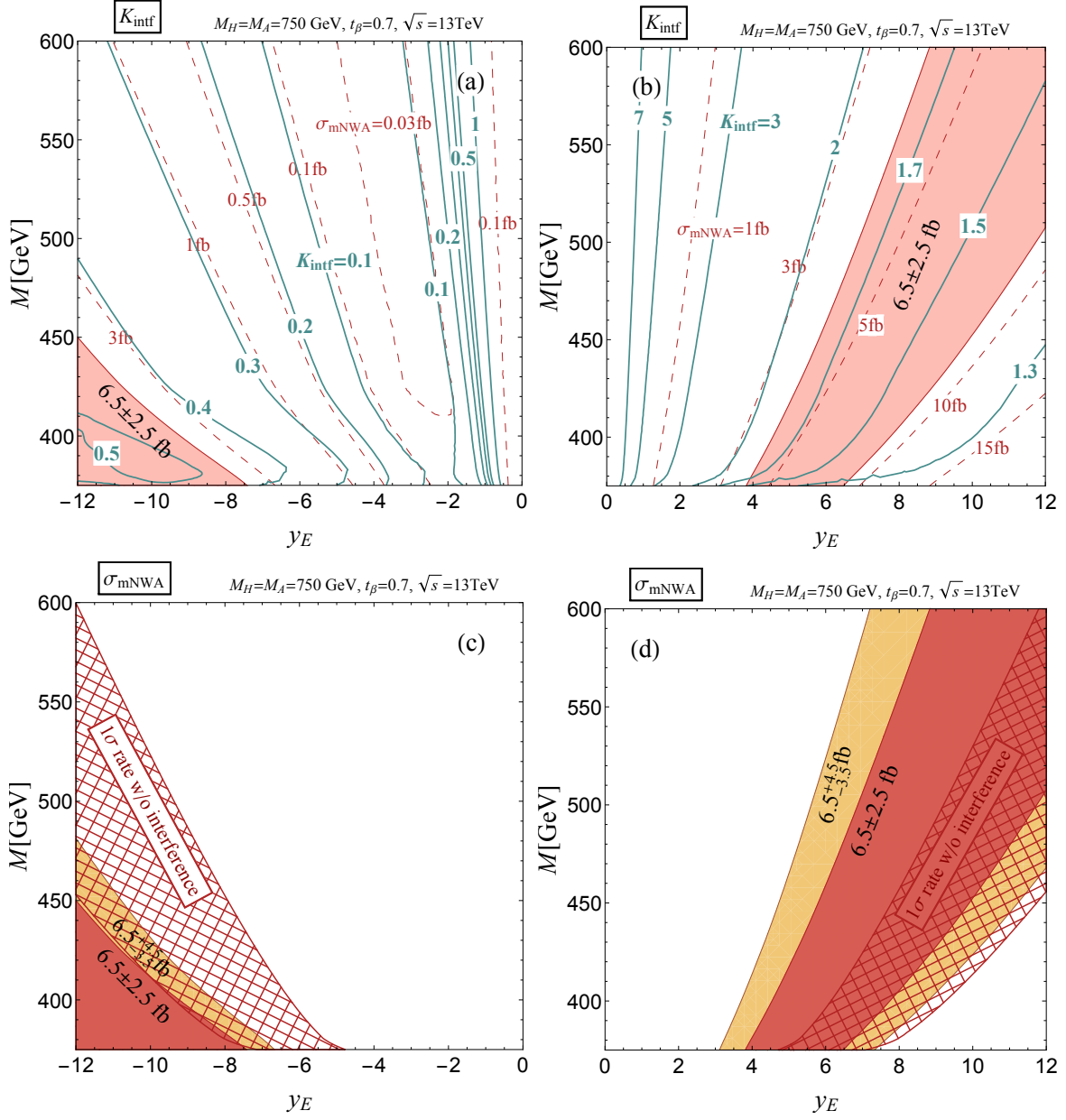


Figure 5. (Upper): Contour plots for K_{intf} (solid lines) and σ_{mNWA} (dashed line) in the (y_E, M) plane of VLL-2HDM model. (Lower): The 68% CL (1σ , darker orange) and 95% CL (lighter orange) best-fit regions are shown. For comparison, the 68% CL results without interferences accounted for are also shown as hatched regions. (Left): $y_E < 0$ induces signal suppression. (Right): $y_E > 0$ induces signal enhancement.

than for H . This is attribute to different loop functions and thus different M_Φ positions for the vanishing real part of the corresponding loop functions. One crucial result is that the interference effects become larger with decreasing signal rate $\sigma_{\text{mNWA}}^{H+A}$. For 1 fb signal rate, for example, the enhancement factor due to the interference can be as large as a factor of three. It is very possible that the current signal rate is fluctuated up and the future precision measurement may lead to lower signal rate. Then the interference effects become crucial.

5 Conclusions and Discussions

We have investigated the impacts of the resonance-continuum interference in the $gg \rightarrow \gamma\gamma$ process on the the recently observed 750 GeV diphoton excess. The two most important interference effects – maximal signal enhancement from the purely imaginary-part interference and maximal shape distortion from the purely real-part interference – have been studied in two benchmark models. First, a CP-odd singlet scalar (extended with vector-like fermions) represents the purely real-part interference case, and it predicts that the 68%(95%) CL best-fit mass range shifts by 1–4 (any $\mathcal{O}(1)$) GeV. The shift is expected to be larger with a weaker coupling parameter space, which will be more preferred if the excess rate decreases in the future. Second, the heavy Higgs bosons in the two-Higgs-doublet-model (extended with vector-like leptons) represent the purely imaginary-part interference case, and the diphoton resonance signal is found to be enhanced or suppressed by a factor of 2(1.6) for 3(6) fb signal rate. Again, the effect is bigger for a weaker coupling parameter space.

Although our results are obtained with benchmark models, any scalar resonance in the $gg \rightarrow \gamma\gamma$ process with similar widths and total rates would exhibit similar sizes of interference effects; and the relative phase ϕ between the resonance and the continuum will determine the type of interference effects. For the given diphoton rate and the phase ϕ , the total width is the most important parameter. If the width is much smaller than the current resolution $\sim \mathcal{O}(1)$ GeV, the real-part interference will cancel out and the imaginary-part interference will be small in proportion to the width. If a resonance is very broad, a careful study of resonance shape including its $m_{\gamma\gamma}$ -dependence shall be carried out, regardless of interference effects, based on our formalism and method presented in this paper.

The future precision shape measurements and interpretations taking into account the resonance-continuum interference can provide important information and consistency check of a new resonance. One can not only test a BW resonance hypothesis but also measure ϕ (and the rate, mass, width). Such precision observables, in particular ϕ , can subsequently be interpreted in terms of the properties of new particles running

in loops. Remarkably, if any noticeable deviations from a BW shape can be fit well with the real-part interference, this would just be another convincing evidence of a new resonance.

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