CP Violation Results from B Decays at BABAR

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In the present paper we review recent experimental results from the BABAR experiment concerning the measurement of the CKM angles. A particular highlight is given to the novel independent determination of the angle \( \alpha \) from \( B^0 \to a_1(1260)^\pm \pi^\mp \) and to the recent full-luminosity updates of several angle \( \gamma \) measurements.

INTRODUCTION

The measurement of CP violation (CPV) in B meson decays provides crucial tests of the Standard Model and of the Cabibbo-Kobayashi-Maskawa (CKM) matrix \cite{1}. The angle \( \beta \) is experimentally measured with a precision of \( O(10) \) in \( B^0 \to (\pi^0)K^0 \) \cite{2} and is not covered in this paper. The determination of angles \( \alpha \) and \( \gamma \) still suffers from larger experimental uncertainties. We review results from BABAR, including the measurement of \( \alpha \) in \( B^0 \to \pi^+\pi^- \) and in the novel decay mode \( B^0 \to a_1(1260)^\pm \pi^\mp \), and full-luminosity updates of several angle \( \gamma \) measurements.

EXPERIMENTAL TECHNIQUES

The results are based on data collected with the BABAR detector \cite{3} at the PEP-II asymmetric-energy \( e^+e^- \) collider, at a center-of-mass energy near the \( T(4S) \) resonance.

The \( B \) meson is kinematically characterized by the variables \( \Delta E \equiv E_B - \frac{1}{2} \sqrt{s} \) and \( m_{ES} \equiv \sqrt{s/4 - |\mathbf{p}_B|^2} \), where \( (E_B, \mathbf{p}_B) \) is the \( B \) four-momentum vector expressed in \( T(4S) \) rest frame. The total integrated luminosity corresponds to about \( 468 \times 10^6 \) \( B\bar{B} \) pairs.

Background arises primarily from random combinations of particles in \( e^+e^- \to q\bar{q} \) events \( (q = u, d, s, c) \), and is discriminated against \( B\bar{B} \) events by using event shape variables, combined into multivariate “shape” classifiers, that are indicated with \( SC \) in the following.

MEASURING \( \alpha \)

The CKM angle \( \alpha \) is measured in \( b \to u\bar{u}d \) transition, exploiting the interference between the decay of mixed and unmixed \( B^0 \) mesons. The signal \( B \) meson \( (B_{CP}) \) is reconstructed into its decay to a CP-eigenstate, accessible from both \( B^0 \) and \( B^0 \). From the remaining particles in the event, we reconstruct the decay vertex of the other \( B \) meson \( (B_{tag}) \) and identify its flavor, through the analysis of its decay products \cite{4}.

The distribution of the difference \( \Delta t \equiv t_{CP} - t_{tag} \) of the proper decay times of \( B \) mesons into CP-eigenstates, such as \( \rho^+\rho^- \), is given by

\[
\begin{equation}
\begin{aligned}
f_q^{\rho\rho}(\Delta t) &= \frac{e^{-|\Delta t|/\tau}}{4\pi} \left\{ 1 - q_{tag} \left[ C \cos(\Delta m_d \Delta t) - S \sin(\Delta m_d \Delta t) \right] \right\},
\end{aligned}
\end{equation}
\]

where \( \tau \) is the mean \( B \) lifetime, \( \Delta m_d \) the \( B^0 - \bar{B}^0 \) mixing frequency, and \( q_{tag} = +1 \) (\(-1\)) if the \( B_{tag} \) decays as a \( B^0 \) (\( \bar{B}^0 \)). The parameters \( S \) and \( C \) describe mixing-induced and direct CPV, respectively. Considering just tree level contributions to the process, \( S = \sin(2\alpha) \) and \( C = 0 \). However, non negligible penguin (loop) amplitudes may contribute to \( b \to u\bar{u}d \) transitions. The different strong and weak phase of the penguin amplitudes may give rise to direct CPV \( (C \neq 0) \), and modify \( S \) into \( S = \sin(2\alpha_{eff})\sqrt{1-C^2} \), where \( \alpha_{eff} = \alpha - \Delta \alpha \), with \( \Delta \alpha \neq 0 \). However, \( \Delta \alpha \) may be extracted via an isospin SU(2) or a flavor SU(3) analysis of the decay.

Work supported in part by US Department of Energy contract DE-AC02-76SF00515.

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\textbf{α FROM } B \rightarrow \rho \rho

**Experimental Inputs**

An isospin SU(2) analysis is used to extract \( \alpha \) in \( B \rightarrow \rho \rho \) decays. The \( B \rightarrow \rho \rho \) decays are \( P \rightarrow V V \) transitions, where \( P \) (\( V \)) denotes a pseudoscalar (vector) meson. Hence, the decay is described by three different amplitudes, one for each helicity state, with different CP transformation properties [5]. The analysis of the angular distributions of the \( B \) meson decay products allows to extract the fraction \( f_L \) of longitudinal polarization. In the helicity formalism, the differential decay rate is

\[
\frac{1}{\Gamma} \frac{d^2\Gamma}{d\cos\theta_1d\cos\theta_2} \propto 4f_L \cos^2\theta_1 \cos^2\theta_2 + (1 - f_L) \sin^2\theta_1 \sin^2\theta_2,
\]

where \( \theta_1(2) \) is the helicity angle between the daughter \( \pi \) and the \( B \) recoil direction in the first (second) \( \rho \) rest frame. Since experimental measurements have shown the decay to be dominated by the longitudinal, CP-even polarization, a full angular analysis, that allows to separate the definite-CP contributions of the transverse polarization, is not needed.

Several inputs are needed to perform the SU(2) analysis of the \( B \rightarrow \rho \rho \) decay. They are: the time-dependent (TD) parameters and branching fraction (BF) of \( B^0 \rightarrow \rho^+\rho^- \) [6, 7], BF and direct CP-asymmetry \( (A_{CP}) \) of \( B^+ \rightarrow \rho^+\rho^0 \) [8], and TD parameters and BF of \( B^0 \rightarrow \rho^0\rho^0 \) [9]. In Table I we summarize the results of the different \( B \rightarrow \rho \rho \) analyses, and the number of \( B\bar{B} \) pairs used in each measurement. The \( B^+ \rightarrow \rho^+\rho^0 \) analysis has been updated using the final

\begin{table}[h]
\centering
\begin{tabular}{|c|c|c|c|c|c|}
\hline
Mode & \( BF \) & \( f_L \) & \( C(A_{CP}) \) & \( S \) & \( N_{B\bar{B}} \) \\
& \( (10^{-6}) \) & & & & \\
\hline
\( B^0 \rightarrow \rho^+\rho^- \) & 25.5 \pm 2.1 \pm 1.9 & 0.992 \pm 0.024 \pm 0.013 & 0.01 \pm 0.15 \pm 0.06 & -0.17 \pm 0.20 \pm 0.06 & 383 \\
\( B^0 \rightarrow \rho^0\rho^0 \) & 23.7 \pm 1.4 \pm 1.4 & 0.950 \pm 0.015 \pm 0.006 & -0.054 \pm 0.055 \pm 0.010 & - & 465 \\
\hline
\end{tabular}
\caption{Results (BF, \( f_L \), \( A_{CP} \), and \( S \)) for \( B \rightarrow \rho \rho \) analyses. The number of \( B\bar{B} \) pairs \( N_{B\bar{B}} \) used in each analysis is also reported.}
\end{table}

**BABAR** dataset. Signal yields, \( A_{CP} \) and \( f_L \) are extracted using a maximum likelihood (ML) fit to \( m_{CS} \), \( \Delta E \), \( SC \), and the masses and helicity angles of the \( \rho \) mesons. Multidimensional probability density functions are used to properly account for variables correlations in the background.

**Determination of \( \alpha \)**

Under the isospin SU(2) symmetry, the following relations hold [10]:

\[
\frac{1}{\sqrt{2}} A^{+-} = A^{+0} - A^{00}, \quad \frac{1}{\sqrt{2}} \tilde{A}^{+-} = \tilde{A}^{-0} - \tilde{A}^{00},
\]

where \( A^{ij} (\tilde{A}^{ij}) \) is the amplitude of \( B^0(\bar{B}^0) \rightarrow \rho^i\rho^j \) process. Neglecting possible electroweak penguins (EWP) amplitudes, which do not obey SU(2) isospin symmetry, \( A^{\pm 0} \) receives only tree amplitude contributions. The small \( A_{CP} \) value measured in \( B^+ \rightarrow \rho^+\rho^0 \) decay indicates that the contribution from EWPs is negligible, and the isospin analysis holds within an uncertainty of \( 1 - 2 \% \) [11]. Other possible isospin violation effects due to finite \( \rho \) width [12] are tested by varying the \( \pi \pi \) mass window. Such effects are below the current sensitivity. If \( A^{+0} \) and \( \tilde{A}^{-0} \) are aligned with a suitable choice of phases, the relations in Eq. (3) can be represented in the complex plane by two triangles, and the phase difference between \( A^{+-} \) and \( \tilde{A}^{+-} \) is \( 2\Delta \). Isospin relations similar to Eq. (3) hold separately for each polarization state. However, since \( f_L \sim 1 \), only the analysis of CP-even longitudinal decay is performed. Constraints on the CKM angle \( \alpha \) and on the penguin contribution \( \Delta \alpha \) are obtained from a confidence level (CL) scan. Assuming the isospin-triangle relations of Eq. (3), a \( \chi^2 \) for the five amplitudes \( A^{+0}, A^{-+}, A^{00}, \tilde{A}^{-+}, \tilde{A}^{00} \) is calculated from the measurements summarized in Table I, and minimized with respect to the parameters that don’t enter the scan. The 1 − CL values are then calculated from the probability of the minimized \( \chi^2 \). Results of such a scan are reported in Fig. 1. Since the measured BF of the several reactions satisfy \( B(B^+ \rightarrow \rho^+\rho^0) \approx B(\rho^+\rho^-) \gg B(B^0 \rightarrow \rho^0\rho^0) \), the
A partial wave analysis of the diffractively produced \( K\pi\pi \) to \( m \) used to determine the decay couplings and the mass poles of the K-matrix. Signal yields are extracted using a ML fit on \( K \) the negligible interference effects. In order to include interference effects in the fit, in a recent analysis by \( BABAR \) obtaining a 68% CL limit into a four-fold ambiguity, corresponding to peaks in the vicinity of 0°, 90° (two degenerate peaks), and 180°, \( BABAR \) obtains a 68% CL limit \(-1.8^\circ < \Delta\alpha < 6.7^\circ \). Taking only the solution consistent with the global CKM fits \[14\], \( \alpha \) is equal to \( 92.4^{+6.6}_{-6.5} \).

\[ \alpha \text{ FROM } B^0\rightarrow a_1(1260)^{\pm}\pi^\mp \]

A novel independent measurement of \( \alpha \) is performed by \( BABAR \) in the \( B\rightarrow a_1(1260)^{\pm}\pi^\mp \) decay. The TD decay rate of the \( B \) meson into the non-CP eigenstate \( a_1(1260)^{\mp}\pi^\mp \) is \[15\]

\[
f_q^{a_1}(\Delta t) \propto e^{-|\Delta t|/\tau} \left( 1 \pm q_{\text{CP}} \right) \left\{ 1 + q_{\text{eff}} \left[ (S \pm \Delta S) \sin(\Delta m_d\Delta t) + (C \pm \Delta C) \cos(\Delta m_d\Delta t) \right] \right\},
\]

and \( q_{\text{eff}} \) enters this equation via \( S \pm \Delta S = \sqrt{1 + (C \pm \Delta C)^2} \times \sin \left( 2\alpha_{\text{eff}} \pm \delta \right) \), where \( \delta \) is the strong phase between the tree amplitudes of \( B^0 \) decays to \( a_1^+ \pi^- \) and \( a_1^- \pi^+ \).

**Experimental Inputs**

Since an isospin SU(2) analysis of the \( B^0\rightarrow a_1(1260)^{\pm}\pi^\mp \) decay is not feasible \[16\], a flavor-SU(3) based approach is used to extract the information on \( \Delta\alpha \). The experimental inputs needed for the SU(3) analysis are: TD parameters and BF of \( B^0\rightarrow a_1(1260)^{\pm}\pi^\mp \), BFs of \( B\rightarrow a_1(1260)^{\pm}K \), and BF of \( B\rightarrow K_{1A}\pi \). The BFs and TD parameters of \( B\rightarrow a_1(1260)\pi \) and \( B\rightarrow a_1(1260)K \) decays have been measured in the last few years \[17, 18\].

The \( B\rightarrow K_{1A}\pi \) BF can be extracted using the combined branching fraction of \( B \) decays to \( K_1(1400)\pi \) and \( K_1(1270)\pi \), and the relative magnitude and phase of \( B\rightarrow K_1(1270)\pi \) and \( B\rightarrow K_1(1400)\pi \) amplitudes. \( K_1(1270) \) and \( K_1(1400) \) are both axial vector mesons, they have overlapping mass distributions and both decays to \( K\pi\pi \), hence they undergo not-negligible interferenece effects. In order to include interference effects in the fit, in a recent analysis by \( BABAR \) \[19\] the \( K_1 \) system is parameterized in terms of a two-resonance, six-channel \( K \)-matrix model \[20\] in the \( P \)-vector approach \[21\]. A partial wave analysis of the diffractively produced \( K\pi\pi \) system performed by the ACCMOR collaboration \[20\] is used to determine the decay couplings and the mass poles of the \( K \)-matrix. Signal yields are extracted using a ML fit to \( m_{\text{BS}}, \Delta E, \text{and } S\mathcal{C} \). The invariant mass of the resonant \( K\pi\pi \) system is sensitive to the production parameters of the \( K_1 \) system. The combined signal branching fractions are \( B(B^0\rightarrow K_1(1270)^{\mp}\pi^\mp + K_1(1400)^{\mp}\pi^\mp) = 31^{+7}_{-5} \times 10^{-6} \).
and \( B(B^+ \to K_1(1270)^0\pi^+ + K_1(1400)^0\pi^+) = 29^{+20}_{-17} \times 10^{-6} \), where the error includes both statistical and systematic contributions. The information about the fraction and phase of the two resonances is used to calculate the contribution of the \( K_{1A} \) meson which belongs to the same SU(3) octet as the \( a_1 \) meson. The results are \( B(B^0 \to K_{1A}^0\pi^-) = 14^{+9}_{-10} \times 10^{-6} \) and \( B(B^+ \to K_{1A}^+\pi^+) < 36 \times 10^{-6} \), where the latter upper limit is evaluated at the 90% confidence level [19].

**Determination of \( \alpha \)**

Using a flavor-SU(3) based approach, the size of penguin amplitudes contributing to the decay is related to the branching fractions of the \( \Delta S = 1 \) partners of the \( B^0 \to a_1^{\pm}\pi^\mp \) decays: \( B \to a_1K \) and \( B \to K_{1A}\pi \) [15, 22]. Nonfactorizable contributions to transition amplitudes from exchange and weak annihilation diagrams are neglected [15, 22]. \( \hat{\delta} \) is assumed to be negligible. \( a_1 \) and \( K_{1A} \) form factors, that are needed in the analysis, are obtained from the study of \( \tau \) decays [23]. A Monte Carlo method is used to derive the 68% and 90% CL upper limits for \( \Delta \alpha \) [19]. The result is \( |\Delta \alpha| < 11^\circ \) (13\(^\circ\)) at the 68% (90%) CL. Combining this bound with the results from \( B^0 \to a_1(1260)^{\mp}\pi^{\mp} \) TD analysis [18], the final result is \( \alpha = (79 \pm 7 \pm 11)^\circ \) for the solution compatible with the CKM global fits, where the first error is statistical and systematic combined and the second is due to penguin pollution.

**MEASURING \( \gamma \)**

The CKM angle \( \gamma \) is measured by exploiting the interference between the \( b \to c\bar{u}s \) and \( b \to u\bar{c}s \) tree amplitudes in \( CP \)-violating \( B \to D^{(*)}K^{(*)} \) decays. Such amplitudes also depend on the magnitude ratio \( r_B \equiv |A(b \to ac)/A(b \to u\bar{c})| \), and the relative strong phase \( \delta_B \) between the CKM favored and suppressed amplitudes. These hadronic parameters depend on the \( B \) decay and are extracted from data. In the following we report the results of full-luminosity updates of three different analysis [24–26], based on the GGSZ [27], ADS [28], and GLW [29] methods, respectively.

**GGSZ method:** \( B^\pm \to D^{(*)}K^{(*)}\pm, D \to K_{S}^0h^+h^- \)

In the GGSZ method [27], the information on \( \gamma \) is extracted from the Dalitz-plot distribution of the \( D \) daughters. The variables sensitive to \( CP \) violation are \( x_\pm \equiv r_B \cos(\delta_B \pm \gamma) \) and \( y_\pm \equiv r_B \sin(\delta_B \pm \gamma) \).

\( B^\pm \to DK^\pm, D^*K^\pm \) (\( D \to D\gamma \) and \( D\sigma^0 \)), and \( DK^{*\pm} \) (\( K^{*\pm} \to K_{S}^0\pi^\pm \)) decays are studied, with \( D \to K_{S}^0h^+h^- (h = \pi, K) \) [24]. Signal yields are extracted using a ML fit to \( m_{ES}, \Delta E, \) and \( \mathcal{S}C \). A fit to the Dalitz-plot distribution of the \( D \) daughters is used to determine 2D confidence regions for \( x_\pm \) and \( y_\pm \), which are shown in Fig. 2. The Dalitz plot model for \( D^0 \) and \( \bar{D}^0 \) decay is studied using the large (\( \approx 6.2 \times 10^5 \)) and very pure \( D^\pm \to D\sigma \) control sample [30]. The Dalitz model includes a non-resonant part and several intermediate \( K_{S}^0h \) or \( h^+h^- \) quasi-two-body decays. The fitted signal yields are about 1000, 500 and 200 events for \( B \to DK, B \to D^*K, \) and \( B \to DK^* \), respectively. The fitted \( (x_\pm, y_\pm) \) values are reported in Table II.

<table>
<thead>
<tr>
<th></th>
<th>( B^\pm \to DK^\pm )</th>
<th>( B^\pm \to D^*K^\pm )</th>
<th>( B^\pm \to DK^{*\pm} )</th>
</tr>
</thead>
<tbody>
<tr>
<td>( x_+ )</td>
<td>(-0.103 \pm 0.037 \pm 0.006 \pm 0.007 )</td>
<td>(0.147 \pm 0.054 \pm 0.017 \pm 0.003 )</td>
<td>(-0.161 \pm 0.083 \pm 0.029 \pm 0.006 )</td>
</tr>
<tr>
<td>( y_+ )</td>
<td>(-0.021 \pm 0.048 \pm 0.004 \pm 0.009 )</td>
<td>(-0.032 \pm 0.077 \pm 0.008 \pm 0.006 )</td>
<td>(0.045 \pm 0.106 \pm 0.036 \pm 0.008 )</td>
</tr>
<tr>
<td>( x_- )</td>
<td>(0.060 \pm 0.039 \pm 0.007 \pm 0.006 )</td>
<td>(-0.104 \pm 0.051 \pm 0.019 \pm 0.002 )</td>
<td>(0.075 \pm 0.096 \pm 0.029 \pm 0.007 )</td>
</tr>
<tr>
<td>( y_- )</td>
<td>(0.062 \pm 0.045 \pm 0.004 \pm 0.006 )</td>
<td>(0.052 \pm 0.064 \pm 0.009 \pm 0.007 )</td>
<td>(0.127 \pm 0.095 \pm 0.027 \pm 0.006 )</td>
</tr>
<tr>
<td>( r_B )</td>
<td>(0.096 \pm 0.029 )</td>
<td>(0.134 \pm 0.034 )</td>
<td>(0.149 \pm 0.066 )</td>
</tr>
<tr>
<td>( \delta_B )</td>
<td>((119^{+15}_{-13})^\circ )</td>
<td>((-82 \pm 21)^\circ )</td>
<td>((111 \pm 32)^\circ )</td>
</tr>
</tbody>
</table>

**TABLE II: GGSZ analysis results: \( x_\pm, y_\pm, r_B, \) and \( \delta_B \).** For \( x_\pm \) and \( y_\pm \) the errors are statistical, systematic and Dalitz-model, respectively. For \( r_B \) and \( \delta_B \) the error is statistical and systematic combined. For \( B \to DK^{*\pm} \) decay we report the value of \( kr_B \), where \( k = 0.9 \pm 0.1 \) takes into account the \( K^* \) finite width.

A frequentist approach is used to obtain \( r_B \), \( \delta_B \), and \( \gamma \) for each decay mode. Results of this analysis are reported in Fig. 3. The values of \( r_B \) and \( \delta_B \) for each \( B \) decay mode are reported in Table II. \( r_B \) is found to be \( \approx 0.1 \), as expected by the theory. CKM angle \( \gamma \) is found to be equal to \( (68 \pm 14 \pm 4 \pm 3)^\circ \) (mod \( 180^\circ \)), where the three uncertainties are statistical, systematic and Dalitz-model, respectively. The distance \( d \) between \( (x_+, y_+) \) and \( (x_-, y_-) \) is sensitive to direct \( CP \)-violation (\( d = 0 \) in case of no \( CPV \)). Results of the analysis indicate a 3.5\( \sigma \) evidence of direct \( CPV \).
sign kaons are produced through a CKM favored (suppressed) decay. Events that have same (right) sign kaons are produced in CKM favored decays. Events that have opposite (wrong) sign kaons are produced in CKM suppressed decays. The ADS method \[28\] exploit the interference between these decay chains. Since the total suppression factor is equal for the two decay chains, interfering amplitudes have similar magnitude, thus large interference effects are expected. On the other side, the large suppression implies a \(B^F \approx O(10^{-7})\). Defining the wrong-to-right sign decay amplitude ratio \(R^\pm \equiv \frac{\Gamma([K^+\pi^0]DK^*)}{\Gamma([K^+\pi^-]DK^*)}\), the following definitions hold

\[
R_{ADS} = \frac{1}{2}(R^+ + R^-) = r_B^2 + r_D^2 \pm 2r_Br_D\cos(\delta_B + \delta_D)\cos\gamma, \tag{4}
\]

\[
A_{ADS} = \frac{R^- - R^+}{R^+ + R^-} = 2r_Br_D\sin\gamma\sin(\delta_B + \delta_D), \tag{5}
\]

where \(r_D\) and \(\delta_D\) are the ratio and the relative phase between the CKM suppressed and the CKM favored amplitude for \(D\) decay, and the \(+(-)\) sign in Eq. (4) is used for \(D\) and \(D^* \rightarrow D\pi^0\) \((D^* \rightarrow D\gamma)\) decays.

To enhance signal purity, a tight \(\Delta E\) cut is applied. Specialized selection criteria are applied to suppress \(K - \pi\) misidentification and \(B\rightarrow D\pi\), \(D\rightarrow K^+K^-\) decays, that are the main sources of \(m_{BG}\)-peaking background. The signal yields, \(R_{ADS}\), and \(A_{ADS}\) are determined from fits to \(m_{BG}\) and \(SC\). In Table III we report the measured values for the signal yield, \(R_{ADS}\), \(A_{ADS}\), and the signal statistical significance (including systematic uncertainties), for each decay mode. A frequentist approach is used to extract the unknowns of Eq. (4)–(5) from the measured
observables. The values of $\delta_D$ and $r_D$ are fixed to those reported in [31]. Results are reported in Fig. 4. Despite a poor sensitivity to $\gamma$, which is bound to be $54^\circ < \gamma < 83^\circ$, a good determination of $r_B$ is obtained: $r_B^{DK^\pm} = 0.095^{+0.051}_{-0.041}$, $r_B^{D^*K^\pm} = 0.096^{+0.045}_{-0.051}$.

<table>
<thead>
<tr>
<th>$B^+\rightarrow DK^+$</th>
<th>$B^+\rightarrow D_{DS}^0 K^+$</th>
<th>$B^+\rightarrow D_{D_s} K^+$</th>
</tr>
</thead>
<tbody>
<tr>
<td>Wrong-sign Signal Yield</td>
<td>19.4 ± 9.6</td>
<td>10.3 ± 5.5</td>
</tr>
<tr>
<td>Stat. Sign. ($\sigma$)</td>
<td>2.1</td>
<td>2.2</td>
</tr>
<tr>
<td>$R_{ADS}$ (10^{-2})</td>
<td>1.1 ± 0.6 ± 0.2</td>
<td>1.8 ± 0.9 ± 0.4</td>
</tr>
<tr>
<td>$A_{ADS}$</td>
<td>$-0.86 \pm 0.47^{+0.12}_{-0.10}$</td>
<td>$+0.77 \pm 0.35 \pm 0.12$</td>
</tr>
</tbody>
</table>

**TABLE III:** Results for $B\rightarrow DK$ ADS analysis.

**FIG. 4:** 1-CL as a function of $\gamma$ (left) and $r_B$ (right) from the $B\rightarrow D^{(*)}K$ ADS study.

**GLW method:** $B^\pm\rightarrow DK^{(*)\pm}$, $D\rightarrow f_{(CP)}$

In the GLW method [29], the $b\rightarrow c\bar{u}s$ and $b\rightarrow u\bar{c}s$ amplitude interference is studied via $D$ meson decays into $CP$ eigenstates. The decay amplitudes are used to build the following quantities

$$ R_{CP\pm} = \frac{\Gamma(B^\pm\rightarrow D^{0}_{CP} K^-) + \Gamma(B^\pm\rightarrow D^{0}_{CP} K^-)}{\Gamma(B^\pm\rightarrow D^{0}_{CP} K^-) + \Gamma(B^\pm\rightarrow D^{0}_{CP} K^-) + 1 + \lambda_{CP} r_B \cos \gamma \cos \delta_B,} $$

$$ A_{CP\pm} = \frac{\Gamma(B^\pm\rightarrow D^{0}_{CP} K^-) - \Gamma(B^\pm\rightarrow D^{0}_{CP} K^-)}{\Gamma(B^\pm\rightarrow D^{0}_{CP} K^-) - \Gamma(B^\pm\rightarrow D^{0}_{CP} K^-) + 2 \lambda_{CP} r_B \sin \gamma \sin \delta_B.} $$

where $D^{0}_{CP} (D^0)$ indicates a $D$ decay into a $CP$ (flavor) eigenstate, and $\lambda_{CP}$ the $CP$-eigenvalue of the final state. $B^\pm\rightarrow DK^\pm$ decays are reconstructed, with $D$ mesons decaying to $CP$-even ($K^+K^-, \pi^+\pi^-$), $CP$-odd ($K^0\pi^0, K^0\phi, K^0\omega$), and flavor ($K^-\pi^+$) eigenstates [26]. The signal yields are measured, and the partial decay rates determined with a ML fit to $m_{ES}, \Delta E$ and $SC$. The fitted signal yield is about 500 events for both $CP$-even and $CP$-odd final states. $A_{CP\pm}$ ($A_{CP\mp}$) and $R_{CP\pm}$ ($R_{CP\mp}$) are extracted from data and are equal to $0.25 \pm 0.06 \pm 0.02$ ($-0.09 \pm 0.07 \pm 0.02$) and $1.18 \pm 0.09 \pm 0.05$ ($1.07 \pm 0.08 \pm 0.04$), respectively. The four observables of Eq. (6)–(7) are used to obtain $\gamma$, $r_B$ and $\delta_B$, using a frequentist approach. The results are $0.24 < r_B < 0.45$ ($0.06 < r_B < 0.51$) and, modulo $180^\circ$, $11.3^\circ < \gamma < 22.7^\circ$ or $80.9^\circ < \gamma < 99.1^\circ$ or $157.3^\circ < \gamma < 168.7^\circ$ ($7.0^\circ < \gamma < 173.0^\circ$) at the 68% (95%) CL, and are shown in Fig. 5.

In order to compare these results with those from GGSZ method [24], $x_\pm = \frac{1}{4} |R_{CP\pm}(1 \mp A_{CP\pm}) - R_{CP\mp}(1 \mp A_{CP\mp})|$ is computed: $x_+ = -0.057 \pm 0.039 \pm 0.015$ and $x_- = 0.132 \pm 0.042 \pm 0.018$. Data from the $D\rightarrow K^0\phi, \phi\rightarrow K^+K^-$ decay are not used to determine $x_\pm$, since they are already used in the GGSZ analysis [24].
CONCLUSION

We have reported results for $\alpha$ and $\gamma$ measurement at BABAR. $\alpha$ is measured in $B \to \rho \rho$ decay at 6° level. A novel independent measurement of $\alpha$ in $B \to a_1(1260)\pi$ decay is performed. The measurement of angle $\gamma$ using the full BABAR dataset is performed using three different techniques. Results of the different techniques are consistent inside the experimental uncertainties. The angle $\gamma$ is measured with a precision of about 15°. The hadronic parameter $r_B$ is found to be $\approx 0.1$, as expected by the theory. Finally, evidences of direct CPV at 3.5$\sigma$ level are reported in $B \to D(\ast)K$ decays.

Acknowledgments

I would like to thank all my BABAR collaborators and especially F. Palombo, A. Gaz, and V. Poiréau for their help in preparing this paper.

[6] Charge conjugation is implied through the paper, unless otherwise specified.