

Update of Time-Dependent CP Asymmetry Measurements in $b \rightarrow c\bar{c}s$ Decays

The *BABAR* Collaboration

August 13, 2008

Abstract

We present updated measurements of time-dependent CP asymmetries in fully reconstructed neutral B decays containing a charmonium meson. The measurements reported here use a data sample of $(465 \pm 5) \times 10^6$ $\Upsilon(4S) \rightarrow B\bar{B}$ decays collected with the *BABAR* detector at the PEP-II B factory. The time-dependent CP asymmetry parameters measured from $J/\psi K_S^0$, $J/\psi K_L^0$, $\psi(2S)K_S^0$, $\chi_{c1}K_S^0$, $\eta_c K_S^0$, and $J/\psi K^{*0}$ decays are

$$\begin{aligned} C_f &= 0.026 \pm 0.020(\text{stat}) \pm 0.016(\text{syst}), \\ S_f &= 0.691 \pm 0.029(\text{stat}) \pm 0.014(\text{syst}). \end{aligned}$$

Submitted to the 34th International Conference on High-Energy Physics, ICHEP 08,
29 July—5 August 2008, Philadelphia, Pennsylvania.

Stanford Linear Accelerator Center, Stanford University, Stanford, CA 94309

Work supported in part by Department of Energy contract DE-AC02-76SF00515.

The BABAR Collaboration,

B. Aubert, M. Bona, Y. Karyotakis, J. P. Lees, V. Poireau, E. Prencipe, X. Prudent, V. Tisserand
*Laboratoire de Physique des Particules, IN2P3/CNRS et Université de Savoie, F-74941 Annecy-Le-Vieux,
France*

J. Garra Tico, E. Grauges
Universitat de Barcelona, Facultat de Fisica, Departament ECM, E-08028 Barcelona, Spain

L. Lopez^{ab}, A. Palano^{ab}, M. Pappagallo^{ab}
INFN Sezione di Bari^a; Dipartimento di Fisica, Università di Bari^b, I-70126 Bari, Italy

G. Eigen, B. Stugu, L. Sun
University of Bergen, Institute of Physics, N-5007 Bergen, Norway

G. S. Abrams, M. Battaglia, D. N. Brown, R. N. Cahn, R. G. Jacobsen, L. T. Kerth, Yu. G. Kolomensky,
G. Lynch, I. L. Osipenkov, M. T. Ronan,¹ K. Tackmann, T. Tanabe
Lawrence Berkeley National Laboratory and University of California, Berkeley, California 94720, USA

C. M. Hawkes, N. Soni, A. T. Watson
University of Birmingham, Birmingham, B15 2TT, United Kingdom

H. Koch, T. Schroeder
Ruhr Universität Bochum, Institut für Experimentalphysik 1, D-44780 Bochum, Germany

D. Walker
University of Bristol, Bristol BS8 1TL, United Kingdom

D. J. Asgeirsson, B. G. Fulsom, C. Hearty, T. S. Mattison, J. A. McKenna
University of British Columbia, Vancouver, British Columbia, Canada V6T 1Z1

M. Barrett, A. Khan
Brunel University, Uxbridge, Middlesex UB8 3PH, United Kingdom

V. E. Blinov, A. D. Bukin, A. R. Buzykaev, V. P. Druzhinin, V. B. Golubev, A. P. Onuchin,
S. I. Serednyakov, Yu. I. Skovpen, E. P. Solodov, K. Yu. Todyshev
Budker Institute of Nuclear Physics, Novosibirsk 630090, Russia

M. Bondioli, S. Curry, I. Eschrich, D. Kirkby, A. J. Lankford, P. Lund, M. Mandelkern, E. C. Martin,
D. P. Stoker
University of California at Irvine, Irvine, California 92697, USA

S. Abachi, C. Buchanan
University of California at Los Angeles, Los Angeles, California 90024, USA

J. W. Gary, F. Liu, O. Long, B. C. Shen,¹ G. M. Vitug, Z. Yasin, L. Zhang
University of California at Riverside, Riverside, California 92521, USA

¹Deceased

V. Sharma

University of California at San Diego, La Jolla, California 92093, USA

C. Campagnari, T. M. Hong, D. Kovalskyi, M. A. Mazur, J. D. Richman

University of California at Santa Barbara, Santa Barbara, California 93106, USA

T. W. Beck, A. M. Eisner, C. J. Flacco, C. A. Heusch, J. Kroseberg, W. S. Lockman, A. J. Martinez,
T. Schalk, B. A. Schumm, A. Seiden, M. G. Wilson, L. O. Winstrom

University of California at Santa Cruz, Institute for Particle Physics, Santa Cruz, California 95064, USA

C. H. Cheng, D. A. Doll, B. Echenard, F. Fang, D. G. Hitlin, I. Narsky, T. Piatenko, F. C. Porter
California Institute of Technology, Pasadena, California 91125, USA

R. Andreassen, G. Mancinelli, B. T. Meadows, K. Mishra, M. D. Sokoloff

University of Cincinnati, Cincinnati, Ohio 45221, USA

P. C. Bloom, W. T. Ford, A. Gaz, J. F. Hirschauer, M. Nagel, U. Nauenberg, J. G. Smith, K. A. Ulmer,
S. R. Wagner

University of Colorado, Boulder, Colorado 80309, USA

R. Ayad,² A. Soffer,³ W. H. Toki, R. J. Wilson

Colorado State University, Fort Collins, Colorado 80523, USA

D. D. Altenburg, E. Feltresi, A. Hauke, H. Jasper, M. Karbach, J. Merkel, A. Petzold, B. Spaan, K. Wacker
Technische Universität Dortmund, Fakultät Physik, D-44221 Dortmund, Germany

M. J. Kobel, W. F. Mader, R. Nogowski, K. R. Schubert, R. Schwierz, A. Volk

Technische Universität Dresden, Institut für Kern- und Teilchenphysik, D-01062 Dresden, Germany

D. Bernard, G. R. Bonneaud, E. Latour, M. Verderi

Laboratoire Leprince-Ringuet, CNRS/IN2P3, Ecole Polytechnique, F-91128 Palaiseau, France

P. J. Clark, S. Playfer, J. E. Watson

University of Edinburgh, Edinburgh EH9 3JZ, United Kingdom

M. Andreotti^{ab}, D. Bettoni^a, C. Bozzi^a, R. Calabrese^{ab}, A. Cecchi^{ab}, G. Cibinetto^{ab}, P. Franchini^{ab},
E. Luppi^{ab}, M. Negrini^{ab}, A. Petrella^{ab}, L. Piemontese^a, V. Santoro^{ab}

INFN Sezione di Ferrara^a; Dipartimento di Fisica, Università di Ferrara^b, I-44100 Ferrara, Italy

R. Baldini-Ferroli, A. Calcaterra, R. de Sangro, G. Finocchiaro, S. Pacetti, P. Patteri, I. M. Peruzzi,⁴
M. Piccolo, M. Rama, A. Zallo

INFN Laboratori Nazionali di Frascati, I-00044 Frascati, Italy

A. Buzzo^a, R. Contri^{ab}, M. Lo Vetere^{ab}, M. M. Macri^a, M. R. Monge^{ab}, S. Passaggio^a, C. Patrignani^{ab},
E. Robutti^a, A. Santroni^{ab}, S. Tosi^{ab}

INFN Sezione di Genova^a; Dipartimento di Fisica, Università di Genova^b, I-16146 Genova, Italy

²Now at Temple University, Philadelphia, Pennsylvania 19122, USA

³Now at Tel Aviv University, Tel Aviv, 69978, Israel

⁴Also with Università di Perugia, Dipartimento di Fisica, Perugia, Italy

K. S. Chaisanguanthum, M. Morii

Harvard University, Cambridge, Massachusetts 02138, USA

A. Adametz, J. Marks, S. Schenk, U. Uwer

Universität Heidelberg, Physikalisches Institut, Philosophenweg 12, D-69120 Heidelberg, Germany

V. Klose, H. M. Lacker

Humboldt-Universität zu Berlin, Institut für Physik, Newtonstr. 15, D-12489 Berlin, Germany

D. J. Bard, P. D. Dauncey, J. A. Nash, M. Tibbetts

Imperial College London, London, SW7 2AZ, United Kingdom

P. K. Behera, X. Chai, M. J. Charles, U. Mallik

University of Iowa, Iowa City, Iowa 52242, USA

J. Cochran, H. B. Crawley, L. Dong, W. T. Meyer, S. Prell, E. I. Rosenberg, A. E. Rubin

Iowa State University, Ames, Iowa 50011-3160, USA

Y. Y. Gao, A. V. Gritsan, Z. J. Guo, C. K. Lae

Johns Hopkins University, Baltimore, Maryland 21218, USA

N. Arnaud, J. Béquilleux, A. D’Orazio, M. Davier, J. Firmino da Costa, G. Grosdidier, A. Höcker,
V. Lepeltier, F. Le Diberder, A. M. Lutz, S. Pruvot, P. Roudeau, M. H. Schune, J. Serrano, V. Sordini,⁵
A. Stocchi, G. Wormser

*Laboratoire de l’Accélérateur Linéaire, IN2P3/CNRS et Université Paris-Sud 11, Centre Scientifique
d’Orsay, B. P. 34, F-91898 Orsay Cedex, France*

D. J. Lange, D. M. Wright

Lawrence Livermore National Laboratory, Livermore, California 94550, USA

I. Bingham, J. P. Burke, C. A. Chavez, J. R. Fry, E. Gabathuler, R. Gamet, D. E. Hutchcroft, D. J. Payne,
C. Touramanis

University of Liverpool, Liverpool L69 7ZE, United Kingdom

A. J. Bevan, C. K. Clarke, K. A. George, F. Di Lodovico, R. Sacco, M. Sigamani

Queen Mary, University of London, London, E1 4NS, United Kingdom

G. Cowan, H. U. Flaecher, D. A. Hopkins, S. Paramesvaran, F. Salvatore, A. C. Wren

*University of London, Royal Holloway and Bedford New College, Egham, Surrey TW20 0EX, United
Kingdom*

D. N. Brown, C. L. Davis

University of Louisville, Louisville, Kentucky 40292, USA

A. G. Denig M. Fritsch, W. Gradl, G. Schott

Johannes Gutenberg-Universität Mainz, Institut für Kernphysik, D-55099 Mainz, Germany

⁵Also with Università di Roma La Sapienza, I-00185 Roma, Italy

K. E. Alwyn, D. Bailey, R. J. Barlow, Y. M. Chia, C. L. Edgar, G. Jackson, G. D. Lafferty, T. J. West,
J. I. Yi

University of Manchester, Manchester M13 9PL, United Kingdom

J. Anderson, C. Chen, A. Jawahery, D. A. Roberts, G. Simi, J. M. Tuggle

University of Maryland, College Park, Maryland 20742, USA

C. Dallapiccola, X. Li, E. Salvati, S. Saremi

University of Massachusetts, Amherst, Massachusetts 01003, USA

R. Cowan, D. Dujmic, P. H. Fisher, G. Sciolla, M. Spitznagel, F. Taylor, R. K. Yamamoto, M. Zhao
*Massachusetts Institute of Technology, Laboratory for Nuclear Science, Cambridge, Massachusetts 02139,
USA*

P. M. Patel, S. H. Robertson

McGill University, Montréal, Québec, Canada H3A 2T8

A. Lazzaro^{ab}, V. Lombardo^a, F. Palombo^{ab}

INFN Sezione di Milano^a; Dipartimento di Fisica, Università di Milano^b, I-20133 Milano, Italy

J. M. Bauer, L. Cremaldi R. Godang,⁶ R. Kroeger, D. A. Sanders, D. J. Summers, H. W. Zhao

University of Mississippi, University, Mississippi 38677, USA

M. Simard, P. Taras, F. B. Viaud

Université de Montréal, Physique des Particules, Montréal, Québec, Canada H3C 3J7

H. Nicholson

Mount Holyoke College, South Hadley, Massachusetts 01075, USA

G. De Nardo^{ab}, L. Lista^a, D. Monorchio^{ab}, G. Onorato^{ab}, C. Sciacca^{ab}

*INFN Sezione di Napoli^a; Dipartimento di Scienze Fisiche, Università di Napoli Federico II^b, I-80126
Napoli, Italy*

G. Raven, H. L. Snoek

*NIKHEF, National Institute for Nuclear Physics and High Energy Physics, NL-1009 DB Amsterdam, The
Netherlands*

C. P. Jessop, K. J. Knoepfel, J. M. LoSecco, W. F. Wang

University of Notre Dame, Notre Dame, Indiana 46556, USA

G. Benelli, L. A. Corwin, K. Honscheid, H. Kagan, R. Kass, J. P. Morris, A. M. Rahimi,

J. J. Regensburger, S. J. Sekula, Q. K. Wong

Ohio State University, Columbus, Ohio 43210, USA

N. L. Blount, J. Brau, R. Frey, O. Igonkina, J. A. Kolb, M. Lu, R. Rahmat, N. B. Sinev, D. Strom,

J. Strube, E. Torrence

University of Oregon, Eugene, Oregon 97403, USA

⁶Now at University of South Alabama, Mobile, Alabama 36688, USA

G. Castelli^{ab}, N. Gagliardi^{ab}, M. Margoni^{ab}, M. Morandin^a, M. Posocco^a, M. Rotondo^a, F. Simonetto^{ab},
R. Stroili^{ab}, C. Voci^{ab}

INFN Sezione di Padova^a; Dipartimento di Fisica, Università di Padova^b, I-35131 Padova, Italy

P. del Amo Sanchez, E. Ben-Haim, H. Briand, G. Calderini, J. Chauveau, P. David, L. Del Buono,
O. Hamon, Ph. Leruste, J. Ocariz, A. Perez, J. Prendki, S. Sitt

*Laboratoire de Physique Nucléaire et de Hautes Energies, IN2P3/CNRS, Université Pierre et Marie
Curie-Paris6, Université Denis Diderot-Paris7, F-75252 Paris, France*

L. Gladney

University of Pennsylvania, Philadelphia, Pennsylvania 19104, USA

M. Biasini^{ab}, R. Covarelli^{ab}, E. Manoni^{ab},

INFN Sezione di Perugia^a; Dipartimento di Fisica, Università di Perugia^b, I-06100 Perugia, Italy

C. Angelini^{ab}, G. Batignani^{ab}, S. Bettarini^{ab}, M. Carpinelli^{ab,7}, A. Cervelli^{ab}, F. Forti^{ab}, M. A. Giorgi^{ab},
A. Lusiani^{ac}, G. Marchiori^{ab}, M. Morganti^{ab}, N. Neri^{ab}, E. Paoloni^{ab}, G. Rizzo^{ab}, J. J. Walsh^a

*INFN Sezione di Pisa^a; Dipartimento di Fisica, Università di Pisa^b; Scuola Normale Superiore di Pisa^c,
I-56127 Pisa, Italy*

D. Lopes Pegna, C. Lu, J. Olsen, A. J. S. Smith, A. V. Telnov

Princeton University, Princeton, New Jersey 08544, USA

F. Anulli^a, E. Baracchini^{ab}, G. Cavoto^a, D. del Re^{ab}, E. Di Marco^{ab}, R. Faccini^{ab}, F. Ferrarotto^a,
F. Ferroni^{ab}, M. Gaspero^{ab}, P. D. Jackson^a, L. Li Gioi^a, M. A. Mazzoni^a, S. Morganti^a, G. Piredda^a,
F. Polci^{ab}, F. Renga^{ab}, C. Voena^a

INFN Sezione di Roma^a; Dipartimento di Fisica, Università di Roma La Sapienza^b, I-00185 Roma, Italy

M. Ebert, T. Hartmann, H. Schröder, R. Waldi

Universität Rostock, D-18051 Rostock, Germany

T. Adye, B. Franek, E. O. Olaiya, F. F. Wilson

Rutherford Appleton Laboratory, Chilton, Didcot, Oxon, OX11 0QX, United Kingdom

S. Emery, M. Escalier, L. Esteve, S. F. Ganzhur, G. Hamel de Monchenault, W. Kozanecki, G. Vasseur,
Ch. Yèche, M. Zito

CEA, Irfu, SPP, Centre de Saclay, F-91191 Gif-sur-Yvette, France

X. R. Chen, H. Liu, W. Park, M. V. Purohit, R. M. White, J. R. Wilson

University of South Carolina, Columbia, South Carolina 29208, USA

M. T. Allen, D. Aston, R. Bartoldus, P. Bechtel, J. F. Benitez, R. Cenci, J. P. Coleman, M. R. Convery,
J. C. Dingfelder, J. Dorfan, G. P. Dubois-Felsmann, W. Dunwoodie, R. C. Field, A. M. Gabareen,
S. J. Gowdy, M. T. Graham, P. Grenier, C. Hast, W. R. Innes, J. Kaminski, M. H. Kelsey, H. Kim, P. Kim,
M. L. Kocian, D. W. G. S. Leith, S. Li, B. Lindquist, S. Luitz, V. Luth, H. L. Lynch, D. B. MacFarlane,
H. Marsiske, R. Messner, D. R. Muller, H. Neal, S. Nelson, C. P. O'Grady, I. Ofte, A. Perazzo, M. Perl,
B. N. Ratcliff, A. Roodman, A. A. Salnikov, R. H. Schindler, J. Schwiening, A. Snyder, D. Su,
M. K. Sullivan, K. Suzuki, S. K. Swain, J. M. Thompson, J. Va'vra, A. P. Wagner, M. Weaver, C. A. West,
W. J. Wisniewski, M. Wittgen, D. H. Wright, H. W. Wulsin, A. K. Yarritu, K. Yi, C. C. Young, V. Ziegler

Stanford Linear Accelerator Center, Stanford, California 94309, USA

⁷Also with Università di Sassari, Sassari, Italy

P. R. Burchat, A. J. Edwards, S. A. Majewski, T. S. Miyashita, B. A. Petersen, L. Wilden
Stanford University, Stanford, California 94305-4060, USA

S. Ahmed, M. S. Alam, J. A. Ernst, B. Pan, M. A. Saeed, S. B. Zain
State University of New York, Albany, New York 12222, USA

S. M. Spanier, B. J. Wogslund
University of Tennessee, Knoxville, Tennessee 37996, USA

R. Eckmann, J. L. Ritchie, A. M. Ruland, C. J. Schilling, R. F. Schwitters
University of Texas at Austin, Austin, Texas 78712, USA

B. W. Drummond, J. M. Izen, X. C. Lou
University of Texas at Dallas, Richardson, Texas 75083, USA

F. Bianchi^{ab}, D. Gamba^{ab}, M. Pelliccioni^{ab}
INFN Sezione di Torino^a; Dipartimento di Fisica Sperimentale, Università di Torino^b, I-10125 Torino, Italy

M. Bomben^{ab}, L. Bosisio^{ab}, C. Cartaro^{ab}, G. Della Ricca^{ab}, L. Lanceri^{ab}, L. Vitale^{ab}
INFN Sezione di Trieste^a; Dipartimento di Fisica, Università di Trieste^b, I-34127 Trieste, Italy

V. Azzolini, N. Lopez-March, F. Martinez-Vidal, D. A. Milanes, A. Oyanguren
IFIC, Universitat de Valencia-CSIC, E-46071 Valencia, Spain

J. Albert, Sw. Banerjee, B. Bhuyan, H. H. F. Choi, K. Hamano, R. Kowalewski, M. J. Lewczuk,
 I. M. Nugent, J. M. Roney, R. J. Sobie
University of Victoria, Victoria, British Columbia, Canada V8W 3P6

T. J. Gershon, P. F. Harrison, J. Ilic, T. E. Latham, G. B. Mohanty
Department of Physics, University of Warwick, Coventry CV4 7AL, United Kingdom

H. R. Band, X. Chen, S. Dasu, K. T. Flood, Y. Pan, M. Pierini, R. Prepost, C. O. Vuosalo, S. L. Wu
University of Wisconsin, Madison, Wisconsin 53706, USA

1 INTRODUCTION

The Standard Model (SM) of electroweak interactions describes CP violation as a consequence of an irreducible phase in the three-family Cabibbo-Kobayashi-Maskawa (CKM) quark-mixing matrix [1]. In the CKM framework, neutral B decays to CP eigenstates containing a charmonium and a $K^{(*)0}$ meson through tree-diagram dominated processes provide a direct measurement of $\sin 2\beta$ [2], where the angle β is defined in terms of the CKM matrix elements V_{ij} for quarks i, j as $\arg[-(V_{cd}V_{cb}^*)/(V_{td}V_{tb}^*)]$.

We identify (tag) the initial flavor of the reconstructed B candidate, B_{rec} , using information from the other B meson, B_{tag} , in the event. The decay rate g_+ (g_-) for a neutral B meson decaying to a CP eigenstate accompanied by a B^0 (\bar{B}^0) tag can be expressed as

$$g_{\pm}(\Delta t) = \frac{e^{-|\Delta t|/\tau_{B^0}}}{4\tau_{B^0}} \left\{ (1 \mp \Delta w) \pm (1 - 2w) \times \left[-\eta_f S \sin(\Delta m_d \Delta t) - C \cos(\Delta m_d \Delta t) \right] \right\} \quad (1)$$

where

$$S = -\eta_f \frac{2\text{Im}\lambda}{1 + |\lambda|^2},$$

$$C = \frac{1 - |\lambda|^2}{1 + |\lambda|^2},$$

the CP eigenvalue $\eta_f = +1$ (-1) for a CP even (odd) final state, $\Delta t \equiv t_{\text{rec}} - t_{\text{tag}}$ is the difference between the proper decay times of B_{rec} and B_{tag} , τ_{B^0} is the neutral B lifetime, and Δm_d is the mass difference of the B meson mass eigenstates determined from B^0 - \bar{B}^0 oscillations [3]. We assume that the corresponding decay-width difference $\Delta\Gamma_d$ is zero. Here, $\lambda = (q/p)(\bar{A}/A)$ [4], where q and p are complex constants that relate the B -meson flavor eigenstates to the mass eigenstates, and \bar{A}/A is the ratio of amplitudes of the decay without mixing of a \bar{B}^0 or B^0 to the final state under study. The average mistag probability w describes the effect of incorrect tags, and Δw is the difference between the mistag probabilities for B^0 and \bar{B}^0 mesons. The sine term in Eq. 1 results from the interference between direct decay and decay after $B^0 - \bar{B}^0$ oscillation. A non-zero cosine term arises from the interference between decay amplitudes with different weak and strong phases (direct CP violation $|\bar{A}/A| \neq 1$) or from CP violation in $B^0 - \bar{B}^0$ mixing ($|q/p| \neq 1$). In the SM, CP violation in mixing and direct CP violation in $b \rightarrow c\bar{c}s$ decays are both negligible [4]. Under these assumptions, $\lambda = \eta_f e^{-2i\beta}$, and $C = 0$. Thus, the time-dependent CP -violating asymmetry is

$$A_{CP}(\Delta t) \equiv \frac{g_+(\Delta t) - g_-(\Delta t)}{g_+(\Delta t) + g_-(\Delta t)} \quad (2)$$

$$= -(1 - 2w)\eta_f S \sin(\Delta m_d \Delta t),$$

and $S = \sin 2\beta$. If we relax the assumption that $C = 0$, then $S = \sqrt{1 - C^2} \sin 2\beta$.

Previous *BABAR* measurements have reported time-dependent CP asymmetries in terms of the parameters $\sin 2\beta$ and $|\lambda|$. In this paper we report results in terms of $C_f = \eta_f C$ and $S_f = \eta_f S$ to be consistent with other time-dependent CP asymmetry measurements. We reconstruct B^0 decays to the final states $J/\psi K_S^0$, $J/\psi K_L^0$, $\psi(2S)K_S^0$, $\chi_{c1}K_S^0$, $\eta_c K_S^0$, and $J/\psi K^{*0}$ ($K^{*0} \rightarrow K_S^0 \pi^0$) [5]. The $J/\psi K_L^0$ final state is CP even, and the $J/\psi K^{*0}$ final state is an admixture of CP even and CP odd amplitudes. Ignoring the angular information in $J/\psi K^{*0}$ results in a dilution of the measured CP asymmetry by a factor $1 - 2R_{\perp}$, where R_{\perp} is the fraction of the $L=1$ contribution. In Ref. [6] we

have measured $R_{\perp} = 0.233 \pm 0.010$ (stat) ± 0.005 (syst), which gives an effective $\eta_f = 0.504 \pm 0.033$ for $f = J/\psi K^{*0}$, after acceptance corrections. In addition to measuring a combined S_f and C_f for the CP modes described above, we measure S_f and C_f for each mode individually, for the $J/\psi K_S^0$ mode where we split this into samples with $K_S^0 \rightarrow \pi^+\pi^-$ and $\pi^0\pi^0$, and for the channel $J/\psi K^0$ (combining the K_S^0 and K_L^0 final states). Since our last published result [7], we have added $82 \times 10^6 B\bar{B}$ decays and applied improved event reconstruction algorithms to the entire dataset.

2 THE DATASET AND BABAR DETECTOR

The results presented in this paper are based on data collected with the BABAR detector at the PEP-II asymmetric energy e^+e^- storage rings [8] operating at the Stanford Linear Accelerator Center. At PEP-II, 9.0 GeV electrons and 3.1 GeV positrons collide at a center-of-mass energy of 10.58 GeV which corresponds to the $\Upsilon(4S)$ resonance. The asymmetric energies result in a boost from the laboratory to the center-of-mass (CM) frame of $\beta\gamma \approx 0.56$. The dataset analyzed has an integrated luminosity of 425.7 fb^{-1} corresponding to $(465 \pm 5) \times 10^6 B\bar{B}$ pairs recorded at the $\Upsilon(4S)$ resonance (on-peak).

The BABAR detector is described in detail elsewhere [9]. Surrounding the interaction point is a five-layer double-sided silicon vertex tracker (SVT) which measures the impact parameters of charged particle tracks in both the plane transverse to, and along the beam direction. A 40-layer drift chamber (DCH) surrounds the SVT and provides measurements of the momenta for charged particles. Charged hadron identification is achieved through measurements of particle energy-loss in the tracking system and the Cherenkov angle obtained from a detector of internally reflected Cherenkov light. A CsI(Tl) electromagnetic calorimeter (EMC) provides photon detection, electron identification, and π^0 reconstruction. The aforementioned components are surrounded by a solenoid magnet, that provides a 1.5 T magnetic field. Finally, the flux return of the magnet is instrumented in order to allow discrimination of muons from pions. For the most recent 211.7 fb^{-1} of data, a portion of the resistive plate chambers constituting the muon system has been replaced by limited streamer tubes [10, 11].

We use a right-handed coordinate system with the z axis along the electron beam direction and the y axis upward, with the origin at the nominal beam interaction point. Unless otherwise stated, kinematic quantities are calculated in the laboratory rest frame. The other reference frame which we commonly use is the CM frame of the colliding electrons and positrons.

We use Monte Carlo (MC) simulated events generated using the GEANT4 [12] and EvtGen [13] based BABAR simulation.

3 RECONSTRUCTION OF B CANDIDATES

We select two samples of events in order to measure the time-dependent CP asymmetry parameters S_f and C_f . These are a sample of fully reconstructed B meson decays to flavor eigenstates (B_{flav}) and a sample of signal events used in the extraction of the CP parameters (B_{CP}). The B_{flav} sample consists of B^0 decays to $D^{(*)-}(\pi^+, \rho^+, a_1^+)$ and $J/\psi K^{*0}$ (where $K^{*0} \rightarrow K^+\pi^-$) final states. We use the B_{flav} sample of events to determine dilution parameters (mistag probabilities). The B_{CP} sample of events consists of B^0 decays to $J/\psi K_S^0$, $J/\psi K_L^0$, $\psi(2S)K_S^0$, $\eta_c K_S^0$, $\chi_{c1} K_S^0$, and $J/\psi K^{*0}$. We assume the interference between the CP side and the tag side reconstruction is negligible and therefore the dilution parameters are assumed to be the same for the B_{flav} and B_{CP} samples. We

also select a sample of fully reconstructed charged B meson decays to $J/\psi K^+$, $\psi(2S)K^+$, $\chi_{c1}K^+$, $\eta_c K^+$, and $J/\psi K^{*+}$ (where $K^{*+} \rightarrow K^+\pi^0$) final states to use as a control sample.

The event selection is unchanged from that described in Ref [7]. Events that pass the selection requirements are refined using kinematic variables. The $J/\psi K_L^0$ mode has the requirement that the difference ΔE between the candidate CM energy and the beam energy E_{beam}^* in the CM satisfies $|\Delta E| < 80 \text{ MeV}$. We require that the beam-energy substituted mass $m_{ES} = \sqrt{(E_{\text{beam}}^*)^2 - (p_B^*)^2}$ is greater than $5.2 \text{ GeV}/c^2$ for all other categories of events, where p_B^* is the B momentum in the e^+e^- CM frame.

We calculate the proper time difference Δt between the two B decays from the measured separation Δz between the decay vertices of B_{rec} and B_{tag} along the collision (z) axis [14]. The z position of the B_{rec} vertex is determined from the charged daughter tracks. The B_{tag} decay vertex is determined by fitting tracks not belonging to the B_{rec} candidate to a common vertex, and including constraints from the beam spot location and the B_{rec} momentum [14]. Events are accepted if the calculated Δt uncertainty is less than 2.5 ps and $|\Delta t|$ is less than 20 ps . The fraction of signal MC events satisfying these requirements is 95% .

4 B MESON FLAVOR TAGGING

A key ingredient in the measurement of time-dependent CP asymmetries is to determine whether at the time of the B_{tag} decay the B_{rec} was a B^0 or a \bar{B}^0 . This ‘flavor tagging’ is achieved with the analysis of the decay products of the recoiling B meson B_{tag} . The overwhelming majority of B mesons decay to a final state that is flavor-specific, i.e. only accessible from either a B^0 or a \bar{B}^0 , but not from both. The purpose of the flavor tagging algorithm is to determine the flavor of B_{tag} with the highest possible efficiency ϵ_{tag} and lowest possible probability w of assigning the wrong flavor to B_{tag} . It is not necessary to fully reconstruct B_{tag} in order to estimate its flavor. The figure of merit for the performance of the tagging algorithm is the effective tagging efficiency

$$Q = \epsilon_{\text{tag}}(1 - 2w)^2, \quad (3)$$

which is related to the statistical uncertainty σ_S and σ_C in the coefficients S and C through

$$\sigma_{S,C} \propto \frac{1}{\sqrt{Q}}. \quad (4)$$

We use a neural network based technique [14, 7] that isolates primary leptons, kaons and pions from B decays to final states containing D^* mesons, and high momentum charged particles from B decays, to determine the flavor of the B_{tag} . The output of this algorithm is divided into seven mutually-exclusive categories. These are (in order of decreasing signal purity) **Lepton**, **Kaon I**, **Kaon II**, **Kaon-Pion**, **Pion**, **Other** and **Untagged**. The performance of this algorithm is determined using the B_{flav} sample. The **Untagged** category of events contain no flavor information, so carry no weight in the time-dependent analysis, and are not used here.

5 LIKELIHOOD FIT METHOD

We determine the composition of our final sample by performing simultaneous fits to the m_{ES} distributions for the full B_{CP} and B_{flav} samples, except for the $J/\psi K_L^0$ sample for which we fit the ΔE distribution.

We define a signal region $5.27 < m_{ES} < 5.29 \text{ GeV}/c^2$ ($|\Delta E| < 10 \text{ MeV}$ for $J/\psi K_L^0$), which contains 15481 CP candidate events that satisfy the tagging and vertexing requirements (see Table 1). For all modes except $\eta_c K_S^0$ and $J/\psi K_L^0$, we use simulated events to estimate the fractions of events that peak in the m_{ES} signal region due to cross-feed from other decay modes (peaking background). For the $\eta_c K_S^0$ mode, the cross-feed fraction is determined from a fit to the $m_{KK\pi}$ and m_{ES} distributions in data. For the $J/\psi K_L^0$ decay mode, the sample composition, effective η_f , and ΔE distribution of the individual background sources are determined either from simulation (for $B \rightarrow J/\psi X$) or from the $m_{\ell+\ell^-}$ sidebands in data (for non- J/ψ background). Figure 1 shows the distributions of m_{ES} obtained for the B_{CP} and B_{flav} events, and ΔE obtained for the $J/\psi K_L^0$ events.

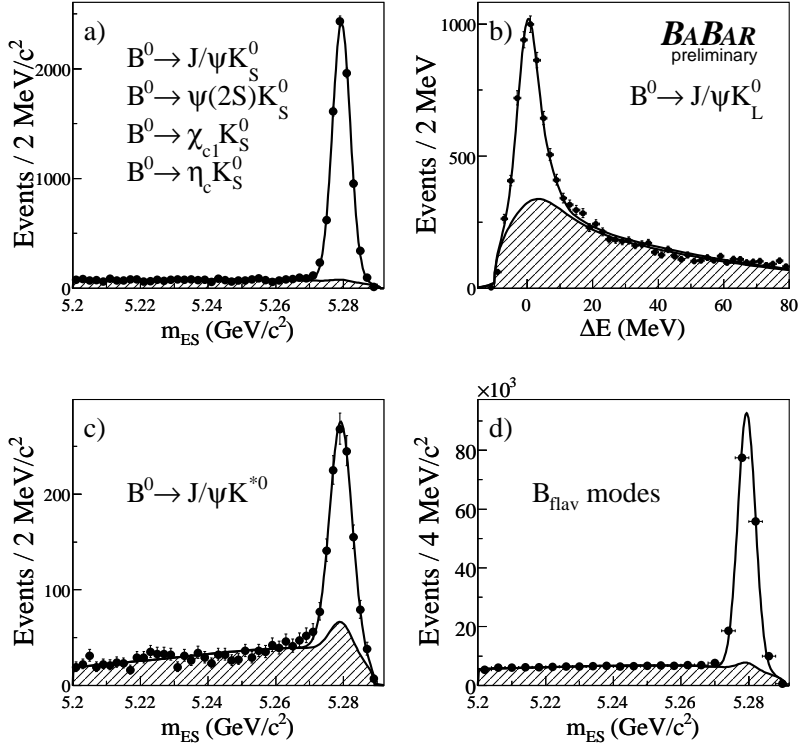


Figure 1: Distributions for B_{CP} and B_{flav} candidates satisfying the tagging and vertexing requirements: a) m_{ES} for the final states $J/\psi K_S^0$, $\psi(2S) K_S^0$, $\chi_{c1} K_S^0$, and $\eta_c K_S^0$, b) ΔE for the final state $J/\psi K_L^0$, c) m_{ES} for $J/\psi K^{*0}$ ($K^{*0} \rightarrow K_S^0 \pi^0$), and d) m_{ES} for the B_{flav} sample. In each plot, the shaded region is the estimated background contribution.

We determine S_f and C_f from a simultaneous maximum likelihood fit to the Δt distribution of the tagged B_{CP} and B_{flav} samples. The Δt distributions of the B_{CP} sample are modeled by Eq. 1. Those of the B_{flav} sample evolve according to Eq. 1 with $S_f = C_f = 0$. The observed amplitudes for the CP asymmetry in the B_{CP} sample and for flavor oscillation in the B_{flav} sample are reduced by the same factor, $1 - 2w$, due to flavor mistags. The Δt distributions for the signal are convolved with a resolution function common to both the B_{flav} and B_{CP} samples, modeled by the sum of three Gaussian functions [14]. The combinatorial background is incorporated with an empirical

description of its Δt spectra, containing zero and non-zero lifetime components convolved with a resolution function [14] distinct from that of the signal. The peaking background is assigned the same Δt distribution as the signal but with $S_f = C_f = 0$, and uses the same Δt resolution function as the signal. As the non-zero lifetime component of the combinatorial background contains both events that are mixed and un-mixed, we allow the value of Δm_d for this component to float in the fit.

In addition to S_f and C_f , there are 69 free parameters in the CP fit. For the signal, these are

- 7 parameters for the Δt resolution,
- 12 parameters for the average mistag fractions w and the differences Δw between B^0 and \bar{B}^0 mistag fractions for each tagging category,
- 7 parameters for the difference between B^0 and \bar{B}^0 reconstruction and tagging efficiencies.

The background is described by

- 24 mistag fraction parameters,
- 3 parameters for the Δt resolution,
- 4 parameters for the B_{flav} time dependence,
- 8 parameters for possible CP violation in the background, including the apparent CP asymmetry of non-peaking events in each tagging category,
- 1 parameter for possible direct CP violation in the $\chi_{c1} K_S^0$ background to $J/\psi K^{*0}$, and
- 3 parameters for possible direct CP violation in the $J/\psi K_L^0$ mode, coming from $J/\psi K_S^0$, $J/\psi K^{*0}$, and the remaining J/ψ backgrounds.

The effective $|\lambda|$ of the non- J/ψ background is fixed from a fit to the J/ψ -candidate sidebands in $J/\psi K_L^0$. We fix $\tau_{B^0} = 1.530$ ps and $\Delta m_d = 0.507$ ps [3]. The determination of the mistag fractions and Δt resolution function parameters for the signal is dominated by the B_{flav} sample, which is about 10 times more abundant than the CP sample.

6 LIKELIHOOD FIT VALIDATION

Before fitting the data in order to extract CP asymmetry parameters, we validate the integrity of the likelihood. We perform three sets of tests in order to validate the fit. The first of these tests consists of generating ensembles of simulated experiments from the PDFs and fitting each simulated experiment. The distribution of fitted S_f and C_f parameters are required to be unbiased, and we verify that the uncertainties are extracted correctly from the fit by requiring that the distribution of the pull \mathcal{P} on a parameter \mathcal{O} , given by $\mathcal{P} = (\mathcal{O}_{\text{fit}} - \mathcal{O}_{\text{gen}})/\sigma(\mathcal{O}_{\text{fit}})$, is a Gaussian centered about zero with a width of one. The quantity \mathcal{O}_{fit} is the fitted value, with a fitted error of $\sigma(\mathcal{O}_{\text{fit}})$, and \mathcal{O}_{gen} is the generated value.

The second test involves fitting simulated CP events with the full *BABAR* detector simulation. We require that the \mathcal{P} distributions for these signal-only simulated experiments are centered about zero with a width of one. We assign a systematic uncertainty corresponding to any deviations

and the statistical uncertainties of the mean values of the fitted S_f and C_f distributions from the generated values.

The third test on our ability to extract S_f and C_f correctly is to perform null tests on control samples of neutral and charged B events where S_f and C_f should equal zero. We use charged B decays to $J/\psi K^\pm$, $\psi(2S)K^\pm$, $\chi_{c1}K^\pm$, $J/\psi K^{*\pm}$ with $K^{*\pm} \rightarrow K^\pm\pi^0$ and $K_S^0\pi^\pm$, and neutral B_{flav} decays for this purpose. The parameters S_f and C_f are zero for these modes within the SM.

7 RESULTS

The fit to the B_{CP} and B_{flav} samples yields $S_f = 0.691 \pm 0.029$ and $C_f = 0.026 \pm 0.020$, where the errors are statistical only. The correlation between these two parameters is $+0.3\%$. We also perform a separate fit in which we allow different S_f and C_f values for each charmonium decay mode, a fit to the $J/\psi K_S^0 (\pi^+\pi^- + \pi^0\pi^0)$ mode, and a fit to the $J/\psi K^0 (K_S^0 + K_L^0)$ sample. We split the data sample by run period and by tagging category. We perform the CP measurements on control samples with no expected CP asymmetry. The results of these fits are summarized in Table 1. Figure 2 shows the Δt distributions and asymmetries in yields between events with B^0 tags and \bar{B}^0 tags for the $\eta_f = -1$ and $\eta_f = +1$ samples as a function of Δt , overlaid with the projection of the likelihood fit result. Figure 3 shows the Δt distributions and asymmetry for $J/\psi K_S^0$ events only. We also performed the CP fit using the $\sin 2\beta$ and $|\lambda|$ parameters, which yields $\sin 2\beta = 0.691 \pm 0.029$ and $|\lambda| = 0.974 \pm 0.020$.

The dominant systematic errors on S_f are due to limited knowledge of various background properties, including possible differences between the B_{flav} and B_{CP} tagging performances, to the description of the Δt resolution functions, uncertainties in $J/\psi K_L^0$ -specific backgrounds and in the amounts of peaking backgrounds and their CP asymmetries, and to the uncertainties in the values of the physics parameters $\Delta m_d, \tau_B, \Delta\Gamma_d/\Gamma_d$. The only sizable systematic uncertainties on C_f are due to the CP content of the peaking backgrounds and to the possible interference between the suppressed $\bar{b} \rightarrow \bar{u}c\bar{d}$ amplitude with the favored $b \rightarrow c\bar{u}d$ amplitude for some tag-side B decays [18]. The total systematic error on S_f (C_f) is 0.014 (0.016). The main systematic uncertainties on both S_f and C_f for the full sample, for the seven individual modes, and for the fits to the $J/\psi K^0$ and $J/\psi K_S^0$ samples are summarized in Tables 2 and 3.

8 CONCLUSIONS

We report improved measurements of the time-dependent CP asymmetry parameters that supersede our previous results [7]. These measurements are given in terms of C_f and S_f for the first time with our data sample. We measure

$$\begin{aligned} C_f &= 0.026 \pm 0.020(\text{stat}) \pm 0.016(\text{syst}), \\ S_f &= 0.691 \pm 0.029(\text{stat}) \pm 0.014(\text{syst}), \end{aligned}$$

providing a model independent constraint on the position of the apex of the Unitarity Triangle [15]. Our measurements agree with previous published results [7, 19] and with the theoretical estimates of the magnitudes of CKM matrix elements within the context of the SM [20]. We also report measurements of C_f and S_f for each of the decay modes within our CP sample and of the $J/\psi K^0 (K_S^0 + K_L^0)$ sample.

Table 1: Number of events N_{tag} and signal purity P in the signal region after tagging and vertexing requirements, and results of fitting for CP asymmetries in the B_{CP} sample and various subsamples. In addition, fit results for the B_{flav} and B^+ control samples demonstrate that no artificial CP asymmetry is found where we expect no CP violation ($S_f = 0, C_f = 0$). Errors are statistical only.

Sample	N_{tag}	$P(\%)$	S_f	C_f
Full CP sample	15481	76	0.691 ± 0.029	0.026 ± 0.020
$J/\psi K_S^0(\pi^+\pi^-)$	5426	96	0.666 ± 0.039	0.019 ± 0.028
$J/\psi K_S^0(\pi^0\pi^0)$	1324	87	0.629 ± 0.092	0.093 ± 0.063
$\psi(2S)K_S^0$	861	87	0.905 ± 0.101	0.092 ± 0.077
$\chi_{c1}K_S^0$	385	88	0.619 ± 0.161	0.133 ± 0.109
$\eta_c K_S^0$	381	79	0.930 ± 0.160	0.082 ± 0.125
$J/\psi K_L^0$	5813	56	0.698 ± 0.062	-0.030 ± 0.050
$J/\psi K^{*0}$	1291	67	0.608 ± 0.241	0.028 ± 0.084
$J/\psi K^0$	12563	77	0.670 ± 0.031	0.019 ± 0.023
$J/\psi K_S^0$	6750	95	0.660 ± 0.036	0.029 ± 0.026
$\eta_f = -1$	8377	93	0.688 ± 0.032	0.041 ± 0.023
1999-2002 data	3079	78	0.736 ± 0.061	0.013 ± 0.045
2003-2004 data	4916	77	0.721 ± 0.050	0.047 ± 0.037
2005-2006 data	4721	76	0.634 ± 0.051	0.046 ± 0.035
2007 data	2765	75	0.666 ± 0.071	-0.017 ± 0.049
Lepton	1740	83	0.734 ± 0.052	0.079 ± 0.038
Kaon I	2187	78	0.617 ± 0.054	-0.045 ± 0.039
Kaon II	3630	76	0.695 ± 0.057	0.073 ± 0.039
Kaon-Pion	2882	74	0.746 ± 0.087	0.006 ± 0.061
Pion	3053	76	0.726 ± 0.135	0.018 ± 0.092
Other	1989	74	0.767 ± 0.349	-0.168 ± 0.238
B_{flav} sample	166276	83	0.021 ± 0.009	0.012 ± 0.006
B^+ sample	36082	94	0.021 ± 0.015	0.013 ± 0.011

9 ACKNOWLEDGMENTS

We are grateful for the extraordinary contributions of our PEP-II colleagues in achieving the excellent luminosity and machine conditions that have made this work possible. The success of this project also relies critically on the expertise and dedication of the computing organizations that support *BABAR*. The collaborating institutions wish to thank SLAC for its support and the kind hospitality extended to them. This work is supported by the US Department of Energy and National Science Foundation, the Natural Sciences and Engineering Research Council (Canada), the Commissariat à l’Energie Atomique and Institut National de Physique Nucléaire et de Physique des Particules (France), the Bundesministerium für Bildung und Forschung and Deutsche Forschungsgemeinschaft (Germany), the Istituto Nazionale di Fisica Nucleare (Italy), the Foundation for Fundamental Research on Matter (The Netherlands), the Research Council of Norway, the Ministry of Education and Science of the Russian Federation, Ministerio de Educación y Ciencia (Spain), and the Science and Technology Facilities Council (United Kingdom). Individuals have received

Table 2: Main systematic uncertainties on S_f and C_f for the full CP sample, and for the $J/\psi K^0$, $J/\psi K_S^0$, and $J/\psi K_L^0$ samples. For each source of systematic uncertainty, the first line gives the error on S_f and the second line the error on C_f . The total systematic error (last row) also includes smaller effects not explicitly mentioned in the table.

Source/sample		Full	$J/\psi K^0$	$J/\psi K_S^0$	$J/\psi K_L^0$
Beamspot	S_f	0.0013	0.0021	0.0027	0.0000
	C_f	0.0006	0.0010	0.0021	0.0001
Mistag differences	S_f	0.0077	0.0057	0.0059	0.0083
	C_f	0.0047	0.0069	0.0053	0.0052
Δt resolution	S_f	0.0067	0.0068	0.0069	0.0071
	C_f	0.0027	0.0029	0.0034	0.0070
$J/\psi K_L^0$ background	S_f	0.0057	0.0063	0.0000	0.0271
	C_f	0.0007	0.0008	0.0000	0.0036
Background fraction and CP content	S_f	0.0046	0.0034	0.0036	0.0044
	C_f	0.0029	0.0021	0.0009	0.0107
m_{ES} parameterization	S_f	0.0022	0.0020	0.0026	0.0006
	C_f	0.0004	0.0005	0.0008	0.0002
$\Delta m_d, \tau_B, \Delta\Gamma_d/\Gamma_d$	S_f	0.0030	0.0033	0.0036	0.0040
	C_f	0.0013	0.0012	0.0011	0.0013
Tag-side interference	S_f	0.0014	0.0014	0.0014	0.0014
	C_f	0.0143	0.0143	0.0143	0.0143
Fit bias (MC statistics)	S_f	0.0023	0.0044	0.0041	0.0063
	C_f	0.0026	0.0044	0.0041	0.0060
Total	S_f	0.0135	0.0131	0.0119	0.0311
	C_f	0.0164	0.0187	0.0167	0.0270

Table 3: Main systematic uncertainties on S_f and C_f for the $J/\psi K_S^0(\pi^+\pi^-)$, $J/\psi K_S^0(\pi^0\pi^0)$, $\psi(2S)K_S^0$, $\chi_{c1}K_S^0$, $\eta_c K_S^0$, and $J/\psi K^{*0}(K^{*0} \rightarrow K_S^0\pi^0)$ decay modes. For each source of systematic uncertainty, the first line gives the error on S_f and the second line the error on C_f . The total systematic error (last row) also includes smaller effects not explicitly mentioned in the table.

Source/sample		$J/\psi K_S^0(\pi^+\pi^-)$	$J/\psi K_S^0(\pi^0\pi^0)$	$\psi(2S)K_S^0$	$\chi_{c1}K_S^0$	$\eta_c K_S^0$	$J/\psi K^{*0}$
Beamspot	S_f	0.0027	0.0020	0.0078	0.0284	0.0010	0.0058
	C_f	0.0017	0.0032	0.0084	0.0115	0.0001	0.0001
Mistag differences	S_f	0.0075	0.0074	0.0089	0.0065	0.0064	0.0117
	C_f	0.0039	0.0046	0.0052	0.0067	0.0047	0.0019
Δt resolution	S_f	0.0072	0.0074	0.0072	0.0099	0.0163	0.0259
	C_f	0.0030	0.0043	0.0070	0.0039	0.0036	0.0062
$J/\psi K_L^0$ background	S_f	0.0001	0.0000	0.0001	0.0000	0.0001	0.0001
	C_f	0.0000	0.0000	0.0000	0.0000	0.0000	0.0000
Background fraction and CP content	S_f	0.0032	0.0073	0.0156	0.0174	0.0506	0.0564
	C_f	0.0012	0.0034	0.0056	0.0098	0.0187	0.0256
m_{ES} parameterization	S_f	0.0021	0.0089	0.0238	0.0061	0.0023	0.0372
	C_f	0.0007	0.0063	0.0008	0.0017	0.0005	0.0080
$\Delta m_d, \tau_B, \Delta\Gamma_d/\Gamma_d$	S_f	0.0031	0.0073	0.0157	0.0025	0.0158	0.0140
	C_f	0.0014	0.0013	0.0010	0.0009	0.0020	0.0013
Tag-side interference	S_f	0.0014	0.0014	0.0014	0.0014	0.0014	0.0014
	C_f	0.0143	0.0143	0.0143	0.0143	0.0143	0.0143
Fit bias (MC statistics)	S_f	0.0048	0.0040	0.0079	0.0072	0.0073	0.0271
	C_f	0.0042	0.0030	0.0019	0.0042	0.0070	0.0389
Total	S_f	0.0129	0.0179	0.0365	0.0398	0.0566	0.0876
	C_f	0.0160	0.0187	0.0209	0.0257	0.0271	0.0540

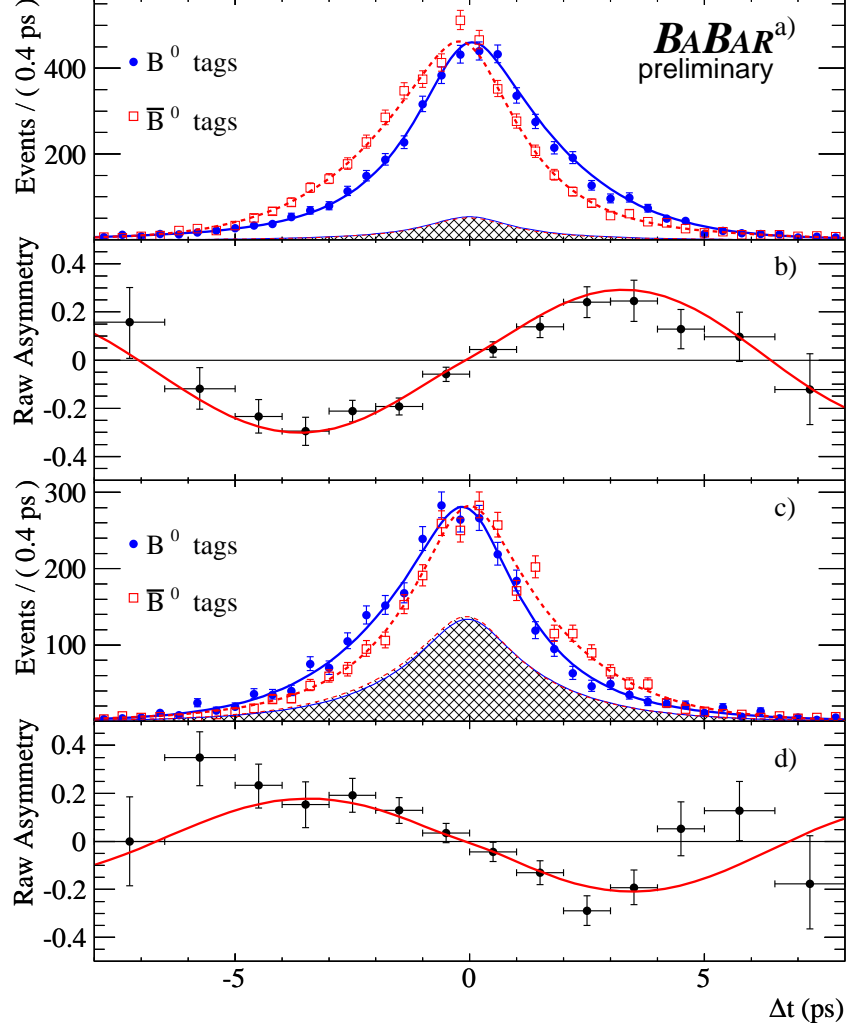


Figure 2: a) Number of $\eta_f = -1$ candidates ($J/\psi K_S^0$, $\psi(2S)K_S^0$, $\chi_{c1}K_S^0$, and $\eta_c K_S^0$) in the signal region with a B^0 tag (N_{B^0}) and with a \bar{B}^0 tag ($N_{\bar{B}^0}$), and b) the raw asymmetry, $(N_{B^0} - N_{\bar{B}^0})/(N_{B^0} + N_{\bar{B}^0})$, as functions of Δt ; c) and d) are the corresponding distributions for the $\eta_f = +1$ mode $J/\psi K_L^0$. The solid (dashed) curves represent the fit projections in Δt for B^0 (\bar{B}^0) tags. The shaded regions represent the estimated background contributions.

support from the Marie-Curie IEF program (European Union) and the A. P. Sloan Foundation.

References

- [1] N. Cabibbo, Phys. Rev. Lett. **10**, 531 (1963); M. Kobayashi and T. Maskawa, Prog. Th. Phys. **49**, 652 (1973).
- [2] A.B. Carter and A.I. Sanda, Phys. Rev. D **23**, 1567 (1981); I.I. Bigi and A.I. Sanda, Nucl. Phys. **193**, 85 (1981).

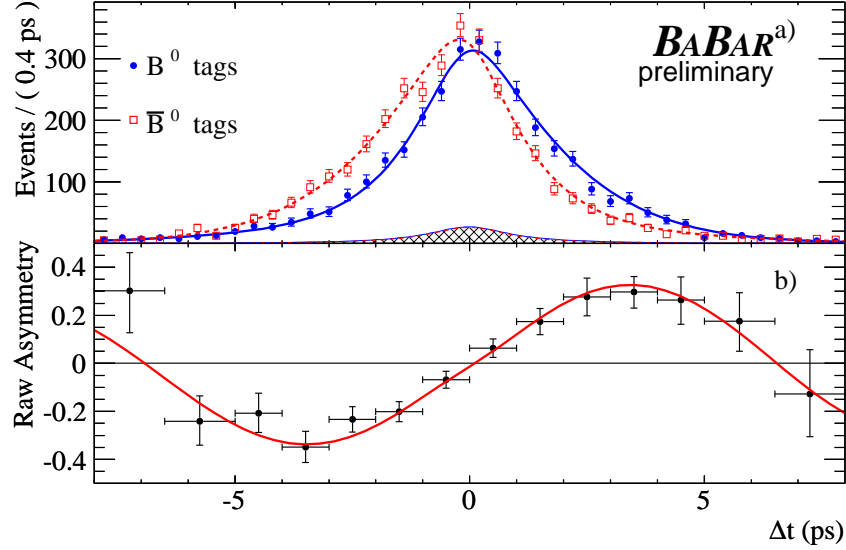


Figure 3: a) Number of $J/\psi K_S^0$ candidates in the signal region with a B^0 tag (N_{B^0}) and with a \bar{B}^0 tag ($N_{\bar{B}^0}$), and b) the raw asymmetry, $(N_{B^0} - N_{\bar{B}^0}) / (N_{B^0} + N_{\bar{B}^0})$, as functions of Δt . The solid (dashed) curves represent the fit projections in Δt for B^0 (\bar{B}^0) tags. The shaded regions represent the estimated background contributions.

- [3] W.-M. Yao *et al.*, (Particle Data Group), J. Phys. **G33**, 1 (2006).
- [4] See, for example, D. Kirkby and Y. Nir, pages 146-155 in Ref. [3].
- [5] Charge conjugation is implied throughout this paper.
- [6] B. Aubert *et al.*, (BABAR Collaboration), Phys. Rev. D **76**, 031102(R) (2007).
- [7] B. Aubert *et al.*, (BABAR Collaboration), Phys. Rev. Lett. **99**, 171803 (2007).
- [8] PEP-II Conceptual Design Report, SLAC-R-418 (1993).
- [9] BABAR Collaboration, B. Aubert *et al.*, Nucl. Instrum. Methods Phys. Res., Sect. A **479**, 1 (2002).
- [10] G. Benelli and K. Honscheid and E. A. Lewis and J. J. Regensburger and D. S. Smith, Nuclear Science Symposium Conference Record, 2005 IEEE, **2**, 1145 (2005).
- [11] M. R. Convery and P. C. Kim and H. P. Paar and C. H. Rogers and R. H. Schindler and S. K. Swain and C. C. Young, Nucl. Instrum. Methods Phys. Res., Sect. A **556**, 134 (2006).
- [12] S. Agostinelli *et al.*, (GEANT4 Collaboration), Nucl. Instrum. Methods Phys. Res., Sect. A **506**, 250 (2003).
- [13] D. Lange *et al.*, Nucl. Instrum. Methods Phys. Res., Sect. A **462**, 152 (2001).
- [14] B. Aubert *et al.*, (BABAR Collaboration), Phys. Rev. D **66**, 032003 (2002).

- [15] See, for example, A. Ceccucci, Z. Ligetti, and Y. Sakai, pages 138-145 in Ref. [3]
- [16] R.Itoh, *et al.*, (The Belle Collaboration), Phys. Rev. Lett. **95** 091601 (2005); B. Aubert *et al.*. (BABAR Collaboration), Phys. Rev. D **71** 032005 (2005).
- [17] P.Krokovny *et al.*, (The Belle Collaboration), Phys. Rev. Lett. **97** 081801 (2006); B. Aubert *et al.*, (BABAR Collaboration), Phys. Rev. D **99** 231802 (2007).
- [18] O. Long, M. Baak, R. N. Cahn, and D. Kirkby, Phys. Rev. D **68**, 034010 (2003).
- [19] K. F. Chen, *et al.*, (Belle Collaboration), Phys. Rev. Lett. **98**, 031802 (2007).
- [20] M. Ciuchini *et al.*, Z. Phys. C **68**, 239 (1995).