

Measurement of the Branching Fractions for Inclusive B^- and \bar{B}^0 Decays to Flavor-tagged D , D_s and A_c

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We report on the inclusive branching fractions of B^- and of \bar{B}^0 mesons decaying to $D^0 X$, $\bar{D}^0 X$, $D^+ X$, $D^- X$, $D_s^+ X$, $D_s^- X$, $\Lambda_c^+ X$, $\bar{\Lambda}_c^- X$, based on a sample of 88.9 million $B\bar{B}$ events recorded with the BABAR detector at the $\Upsilon(4S)$ resonance. Events are selected by completely reconstructing one B and searching for a reconstructed charmed particle in the rest of the event. We measure the number of charmed and of anti-charmed particles per B decay and derive the total charm yield per B^- decay, $n_c^- = 1.313 \pm 0.037 \pm 0.062_{-0.042}^{+0.063}$, and per \bar{B}^0 decay, $n_c^0 = 1.276 \pm 0.062 \pm 0.058_{-0.046}^{+0.066}$ where the first uncertainty is statistical, the second is systematic, and the third reflects the charm branching-fraction uncertainties.

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The dominant process for the decay of a b quark is $b \rightarrow cW^{*-}$ [1], resulting in a (flavor) correlated c quark and a virtual W . In the decay of the W , the production of a $\bar{u}d$ or a $\bar{c}s$ pair are both Cabibbo-allowed and should be equal, the latter being only suppressed by a phase-space factor. The first process dominates hadronic b decays, while the second can be easily distinguished as it will produce a (flavor) anti-correlated \bar{c} quark. Experimentally, correlated and anti-correlated charm production can be investigated through the measurement of the inclusive B -decay rates to flavor-tagged charmed mesons or baryons. Current measurements [2–4] of these rates have statistically limited precision and do not distinguish among the different B parent states.

Most of the charged and neutral D mesons produced in \bar{B} decays come from correlated production $\bar{B} \rightarrow DX$. However, a significant number of $\bar{B} \rightarrow \bar{D}X$ decays are expected through $b \rightarrow c\bar{c}s$ transitions, such as $\bar{B} \rightarrow D^{(*)}\bar{D}^{(*)}\bar{K}^{(*)}(n\pi)$. Although the branching fractions of the 3-body decays $\bar{B} \rightarrow D^{(*)}\bar{D}^{(*)}\bar{K}$ have been measured [5, 6], it is not clear whether they saturate $\bar{B} \rightarrow \bar{D}X$ transitions. It is therefore important to improve the precision on the branching fraction $\mathcal{B}(\bar{B} \rightarrow \bar{D}X)$.

By contrast, the anti-correlated D_s^- production $\bar{B} \rightarrow D_s^- D(n\pi)$ is expected to dominate \bar{B} decays to D_s mesons, since correlated production needs an extra $s\bar{s}$ pair created from the vacuum to give $\bar{B} \rightarrow D_s^+ K^-(n\pi)$. There is no prior published measurement of $\mathcal{B}(\bar{B} \rightarrow D_s^+ X)$.

All strangeless charmed baryons decay to Λ_c . Correlated Λ_c are produced in decays like $B^- \rightarrow \Lambda_c^+ \bar{p}\pi^-(\pi)$, while anti-correlated $\bar{\Lambda}_c^-$ should originate from $B^- \rightarrow \Xi_c \bar{\Lambda}_c^-(\pi)$. Another possibility is $B^- \rightarrow \Lambda_c^+ \bar{\Lambda}_c^- K^-$, the baryonic analogue of the $D\bar{D}K$ decay. The rates for Ξ_c production in B decays [7] are unknown, because there is no absolute measurement of Ξ_c decay branching fractions.

This analysis uses $\Upsilon(4S) \rightarrow B\bar{B}$ events in which either a B^+ or a B^0 meson (hereafter denoted B_{reco}) decays into a hadronic final state and is fully reconstructed. We then reconstruct D , D_s and Λ_c from the recoiling B^- (\bar{B}^0) meson and compare the flavor of the charm hadron with that of the B_{reco} , thus allowing separate measure-

ments of the B^- (\bar{B}^0) $\rightarrow D^0 X$, $D^+ X$, $D_s^+ X$, $\Lambda_c^+ X$ and B^- (\bar{B}^0) $\rightarrow \bar{D}^0 X$, $D^- X$, $D_s^- X$, $\bar{\Lambda}_c^- X$ branching fractions. We extract $\mathcal{B}(B^- \rightarrow \Lambda_c^+ \bar{\Lambda}_c^- K^-)$ from the missing-mass spectra of the $\Lambda_c^+ K^-$ or $\bar{\Lambda}_c^- K^-$ systems recoiling against the B_{reco} . We can then evaluate indirectly $\mathcal{B}(B^- \rightarrow \Xi_c X) = \mathcal{B}(B^- \rightarrow \bar{\Lambda}_c^- X) - \mathcal{B}(B^- \rightarrow \Lambda_c^+ \bar{\Lambda}_c^- K^-)$ and compute the average number of charm (anti-charm) particles per B^- decay, N_c^- ($N_{\bar{c}}^-$):

$$N_c^- = \sum_{X_c} \mathcal{B}(B^- \rightarrow X_c X), \quad (1)$$

$$N_{\bar{c}}^- = \sum_{\bar{X}_c} \mathcal{B}(B^- \rightarrow \bar{X}_c X), \quad (2)$$

where the sum is performed over $X_c = D^+$, D^0 , D_s^+ , Λ_c^+ , Ξ_c , $(c\bar{c})$ or $\bar{X}_c = D^-$, \bar{D}^0 , D_s^- , $\bar{\Lambda}_c^-$, $(c\bar{c})$, and $(c\bar{c})$ refers to all charmonium states collectively. We neglect $\bar{\Xi}_c$ production, as it requires both a $\bar{c}s$ and an $s\bar{s}$ pair in the decay to give $\bar{\Xi}_c \Omega_c$. We can sum N_c^- and $N_{\bar{c}}^-$ to obtain the average number of charm plus anti-charm quarks per B^- decay, $n_c^- = N_c^- + N_{\bar{c}}^-$ (and similarly for \bar{B}^0 decays). In addition to the theoretical interest [8–10], the fact that anti-correlated charmed particles are a background for many studies also motivates a more precise measurement of their production rates in B decays.

The measurements presented here are based on a sample of 88.9 million $B\bar{B}$ pairs (81.9 fb^{-1}) recorded at the $\Upsilon(4S)$ resonance with the BABAR detector at the PEP-II asymmetric-energy B -meson factory at SLAC. The BABAR detector is described in detail elsewhere [11]. Charged-particle trajectories are measured by a 5-layer double-sided silicon vertex tracker and a 40-layer drift chamber, both operating in a 1.5-T solenoidal magnetic field. Charged-particle identification is provided by the average energy loss (dE/dx) in the tracking devices and by an internally reflecting ring-imaging Cherenkov detector. Photons are detected by a CsI(Tl) electromagnetic calorimeter. We use Monte Carlo simulations of the BABAR detector based on GEANT4 [12] to optimize selection criteria and determine selection efficiencies.

We reconstruct B^+ and B^0 decays (B_{reco}) in the modes $B^+ \rightarrow \bar{D}^{(*)0}\pi^+$, $\bar{D}^{(*)0}\rho^+$, $\bar{D}^{(*)0}a_1^+$ and $B^0 \rightarrow D^{(*)-}\pi^+$, $D^{(*)-}\rho^+$, $D^{(*)-}a_1^+$. \bar{D}^0 candidates are reconstructed in

the $K^+\pi^-$, $K^+\pi^-\pi^0$, $K^+\pi^-\pi^+\pi^-$ and $K_S^0\pi^+\pi^-$ ($K_S^0 \rightarrow \pi^+\pi^-$) decay channels, while D^- are reconstructed in the $K^+\pi^-\pi^-$ and $K_S^0\pi^-\pi^-$ modes. D^* candidates are reconstructed in the $D^{*-} \rightarrow \bar{D}^0\pi^-$ and $\bar{D}^{*0} \rightarrow \bar{D}^0\pi^0$, $\bar{D}^0\gamma$ decay modes. The first kinematic variable used to identify fully reconstructed B decays is the beam-energy substituted mass, $m_{\text{ES}} = \sqrt{(s/2 + \mathbf{p}_i \cdot \mathbf{p}_B)^2 / E_i^2 - \mathbf{p}_B^2}$, where \mathbf{p}_B is the B_{reco} momentum and (E_i, \mathbf{p}_i) is the four-momentum of the initial e^+e^- system, both measured in the laboratory frame. The invariant mass of the initial e^+e^- system is \sqrt{s} . The second variable is $\Delta E = E_B^* - \sqrt{s}/2$, where E_B^* is the B_{reco} candidate energy in the center-of-mass frame. We require $|\Delta E| < n\sigma_{\Delta E}$ with $n = 2$ or 3 , depending on the decay mode, and using the measured resolution $\sigma_{\Delta E}$ for each decay mode.

In the m_{ES} spectra (Fig. 1), we define a signal region with $5.274 < m_{\text{ES}} < 5.290$ GeV/c^2 and a background control region with $5.220 < m_{\text{ES}} < 5.260$ GeV/c^2 . For each of the B -decay modes, the combinatorial background in the signal region is derived from a fit to the m_{ES} distribution that uses an empirical phase-space threshold function [13] for the background, together with a signal function [14] peaked at the B meson mass. The numbers of reconstructed B^+ and B^0 candidates, $N_{B^+} = 85840 \pm 1910$ (syst.) and $N_{B^0} = 48322 \pm 590$ (syst.), are then obtained by subtracting this background from the total number of events found in the signal region. These measured B meson yields provide the normalization of all branching fraction measurements reported below. The systematic uncertainties quoted above are computed by varying the boundaries of the signal and background regions, and by comparing the shapes of the threshold function [13] in the data and in the simulation.

The contamination of B^0 events in the B^+ signal induces a background which peaks near the B mass. From the Monte Carlo simulation, the fraction of B^0 events in the reconstructed B^+ signal sample is found to be $c_0 = 0.034$, and the fraction of B^+ events in the reconstructed B^0 signal sample to be $c_+ = 0.019$. A 100 % systematic uncertainty is conservatively assigned to these numbers but they will have a small effect on the final results.

We now turn to the analysis of inclusive D , D_s and A_c production in the decays of the \bar{B} 's that recoil against the reconstructed B . Charmed particles X_c (correlated production) are distinguished from anti-charmed particles \bar{X}_c (anti-correlated production). They are reconstructed from charged tracks that do not belong to the B_{reco} . The decay modes considered are listed in Table I.

For charged B decays, Fig. 2 shows the D , D_s , and A_c mass spectra of correlated and anti-correlated candidates recoiling against B 's reconstructed in the m_{ES} signal region, for some selected decay modes. These spectra are fitted with the sum of a Gaussian signal and a linear background (including a satellite peak for some

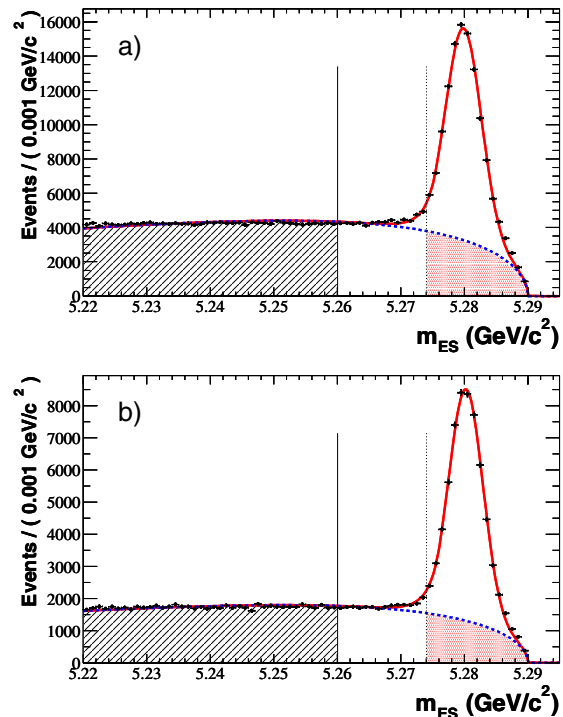


FIG. 1: m_{ES} spectra of reconstructed (a) B^+ and (b) B^0 candidates. The full vertical line shows the upper limit of the background control region (hatched), the dotted vertical line the lower limit of the B signal region. The crossed area shows the background under the B signal. The solid curve is the sum of the fitted signal and background, the dashed curve is the background component only.

channels [15]). The shaded areas correspond to well reconstructed D , D_s or A_c from the combinatorial B_{reco} background. They are obtained from data in the m_{ES} background control region, normalized to the number of combinatorial background events expected under the B_{reco} peak. The background-subtracted reconstructed signal yields are listed in Table I. The reconstruction efficiencies for each charmed (anti-charmed) final state $X_c \rightarrow f$ ($\bar{X}_c \rightarrow \bar{f}$) are computed from the simulation as a function of the charmed-particle momentum in the B^- center-of-mass frame, and are applied event-by-event to obtain the efficiency-corrected charm signal yields $N(X_c \rightarrow f)$ ($N(\bar{X}_c \rightarrow \bar{f})$). The final branching fractions are computed from these yields, the number of B_{reco} , and the intermediate branching fractions $\mathcal{B}(X_c \rightarrow f)$ taken from [16]. They are given by

$$\mathcal{B}(B^- \rightarrow X_c X) = \frac{N(X_c \rightarrow f)}{N_{B^+} \times \mathcal{B}(X_c \rightarrow f)} - c_0 \mathcal{B}_0. \quad (3)$$

Here the raw branching fraction for $B^- \rightarrow X_c X$ is modified by a small corrective term, $c_0 \mathcal{B}_0$, that accounts for the B^0 contamination in the reconstructed B^+ sample. The factor \mathcal{B}_0 depends on the measured $\bar{B}^0 \rightarrow X_c X$ and

TABLE I: Charmed-particle signal yields and B branching fractions per decay mode. The first uncertainty is statistical, the second is systematic (but does not include the charm branching fraction uncertainties).

X_c decay mode	$B^- \rightarrow X_c X$		$B^- \rightarrow \bar{X}_c X$		$\bar{B}^0 \rightarrow X_c X$		$\bar{B}^0 \rightarrow \bar{X}_c X$	
	yield	$\mathcal{B}(\%)$	yield	$\mathcal{B}(\%)$	yield	$\mathcal{B}(\%)$	yield	$\mathcal{B}(\%)$
$D^0 \rightarrow K^- \pi^+$	1273 \pm 42	79.2 \pm 2.6 \pm 3.9	160 \pm 16	9.3 \pm 1.0 \pm 0.5	397 \pm 24	50.3 \pm 3.4 \pm 2.4	139 \pm 14	7.3 \pm 2.2 \pm 0.5
$\rightarrow K^- \pi^+ \pi^- \pi^+$	998 \pm 65	80.6 \pm 5.3 \pm 7.5	173 \pm 30	13.4 \pm 2.4 \pm 1.3	332 \pm 36	56.2 \pm 6.8 \pm 5.4	83 \pm 23	1.8 \pm 4.4 \pm 0.5
$D^+ \rightarrow K^- \pi^+ \pi^+$	262 \pm 29	9.8 \pm 1.2 \pm 1.2	98 \pm 20	3.8 \pm 0.9 \pm 0.4	452 \pm 31	39.7 \pm 3.0 \pm 2.8	125 \pm 18	2.3 \pm 1.8 \pm 0.3
$D_s^+ \rightarrow \phi \pi^+$	11 \pm 5	2.2 \pm 1.1 \pm 0.3	82 \pm 11	16.5 \pm 2.3 \pm 1.7	24 \pm 6	8.3 \pm 2.8 \pm 0.8	28 \pm 6	9.9 \pm 2.9 \pm 1.0
$\rightarrow \bar{K}^{*0} K^+$	0 \pm 3	0.0 \pm 1.1 \pm 0.2	55 \pm 11	18.0 \pm 3.5 \pm 1.7	3 \pm 4	0.0 \pm 2.8 \pm 0.1	14 \pm 5	9.9 \pm 4.1 \pm 1.2
$\rightarrow K_s^0 K^+$	0 \pm 3	0.0 \pm 0.9 \pm 0.2	31 \pm 9	9.2 \pm 2.7 \pm 0.8	12 \pm 5	5.0 \pm 3.4 \pm 0.4	23 \pm 6	13.3 \pm 4.3 \pm 1.0
$\Lambda_c^+ \rightarrow p K^- \pi^+$	41 \pm 9	3.5 \pm 0.8 \pm 0.3	33 \pm 9	2.9 \pm 0.8 \pm 0.3	28 \pm 8	4.9 \pm 1.7 \pm 0.4	16 \pm 6	2.0 \pm 1.2 \pm 0.2

$B^0 \rightarrow X_c X$ branching fractions, and on the $B^0 - \bar{B}^0$ mixing parameter χ_d [16]. It ranges from less than 3% for Λ_c to as much as 50% for correlated D^0 and D^+ . Doubly Cabibbo-suppressed D^0 decays are also taken into account. The branching fractions and their errors are given in Table I. The statistical and systematic uncertainties are computed separately for each channel. For example, the 3.9% absolute systematic uncertainty on $\mathcal{B}(B^- \rightarrow D^0(K^- \pi^+)X)$ reflects the quadratic sum of 1.8% attributed to N_{B^+} , 1.3% to the error on the rate of true D 's in the B combinatorial background, 0.8% to the Monte Carlo statistics, 1.2% to the track-finding efficiency, 2.5% to the particle identification, 1.2% to c_0 , and 0.1% to \mathcal{B}_0 . We combine the results from the different D^0 and D_s decay modes to extract the final branching fractions listed in Table II.

TABLE II: Combined B^- branching fractions. The first uncertainty is statistical, the second is systematic, and the third reflects charm branching-fraction uncertainties [16].

X_c	correlated	anti-correlated
	$\mathcal{B}(B^- \rightarrow X_c X)(\%)$	$\mathcal{B}(B^- \rightarrow \bar{X}_c X)(\%)$
D^0	79.3 \pm 2.5 \pm 4.0 $^{+2.0}_{-1.9}$	9.8 \pm 0.9 \pm 0.5 $^{+0.3}_{-0.3}$
D^+	9.8 \pm 1.2 \pm 1.2 $^{+0.8}_{-0.7}$	3.8 \pm 0.9 \pm 0.4 $^{+0.3}_{-0.3}$
D_s^+	0.5 \pm 0.6 \pm 0.2 $^{+0.2}_{-0.1}$ < 2.2 at 90% CL	14.3 \pm 1.6 \pm 1.5 $^{+4.9}_{-3.0}$
Λ_c^+	3.5 \pm 0.8 \pm 0.3 $^{+1.3}_{-0.8}$	2.9 \pm 0.8 \pm 0.3 $^{+1.1}_{-0.6}$

To extract N_c from these numbers, we need to evaluate the contribution of $B^- \rightarrow \Lambda_c^+ \bar{\Lambda}_c^- K^-$. Combining the four-momenta of the recoiling B^- , of a K^- and of the reconstructed Λ_c^+ or $\bar{\Lambda}_c^-$ candidate, we compute the missing mass: the absence of signal at the Λ_c mass excludes a significant contribution of this process. We therefore take $\mathcal{B}(B^- \rightarrow \Xi_c X) = \mathcal{B}(B^- \rightarrow \bar{\Lambda}_c^- X)$ in the computation of N_c . Using Eqs. 1 and 2 and taking $\mathcal{B}(B^- \rightarrow (c\bar{c})X) =$

$(2.3 \pm 0.3)\%$ [17] [18], one obtains:

$$\begin{aligned} N_c^- &= 0.983 \pm 0.030 \pm 0.046_{-0.023}^{+0.028}, \\ N_{\bar{c}}^- &= 0.330 \pm 0.022 \pm 0.020_{-0.031}^{+0.051}, \\ n_c^- &= 1.313 \pm 0.037 \pm 0.062_{-0.042}^{+0.063}. \end{aligned}$$

The reconstruction of D , D_s and Λ_c from \bar{B}^0 decays is performed in the same way as that in the B^- analysis. The corresponding yields are listed in Table I. We then compute for each decay channel $X_c \rightarrow f$ the efficiency-corrected signal yields $N(X_c \rightarrow f)$ ($N(\bar{X}_c \rightarrow \bar{f})$) and define the raw branching fractions \mathcal{B}_c and $\bar{\mathcal{B}}_c$ as

$$\mathcal{B}_c = N(X_c \rightarrow f) / [N_{B^0} \times \mathcal{B}(X_c \rightarrow f)] \quad (4)$$

$$\bar{\mathcal{B}}_c = N(\bar{X}_c \rightarrow \bar{f}) / [N_{B^0} \times \mathcal{B}(X_c \rightarrow f)]. \quad (5)$$

After correcting these numbers for $B^0 \bar{B}^0$ mixing, we obtain the final branching fraction for $\bar{B}^0 \rightarrow X_c X$:

$$\mathcal{B}(\bar{B}^0 \rightarrow X_c X) = \frac{\mathcal{B}_c - \chi_d(\mathcal{B}_c + \bar{\mathcal{B}}_c) - c_+ \mathcal{B}_+}{1 - 2\chi_d}, \quad (6)$$

where $\chi_d = 0.181 \pm 0.004$ is the $B^0 - \bar{B}^0$ mixing parameter [16]. The correcting factor \mathcal{B}_+ accounts for B^+ contamination in the B^0 sample and depends on $\mathcal{B}(B^- \rightarrow X_c X)$ and $\mathcal{B}(B^+ \rightarrow X_c X)$. The results are given in Table I. Combining the different D^0 or D_s modes, we obtain the final branching fractions listed in Table III.

To compute N_c , we neglect $\bar{B}^0 \rightarrow \Lambda_c^+ \bar{\Lambda}_c^- K^0$ production and assume that $\mathcal{B}(\bar{B}^0 \rightarrow \Xi_c X) = \mathcal{B}(\bar{B}^0 \rightarrow \bar{\Lambda}_c^- X)$. Substituting \bar{B}^0 for B^- in Eqs. 1 and 2 and taking $\mathcal{B}(B^0 \rightarrow (c\bar{c})X) = (2.3 \pm 0.3)\%$ [17] [18], we obtain:

$$\begin{aligned} N_c^0 &= 1.039 \pm 0.051 \pm 0.049_{-0.031}^{+0.039}, \\ N_{\bar{c}}^0 &= 0.237 \pm 0.036 \pm 0.012_{-0.024}^{+0.039}, \\ n_c^0 &= 1.276 \pm 0.062 \pm 0.058_{-0.046}^{+0.066}. \end{aligned}$$

We also compute the fraction of anti-correlated charm production in B decays, $w(\bar{X}_c) = \mathcal{B}(\bar{B} \rightarrow \bar{X}_c X) / (\mathcal{B}(\bar{B} \rightarrow X_c X) + \mathcal{B}(\bar{B} \rightarrow \bar{X}_c X))$. Here, many systematic uncertainties cancel (tracking, K identification,

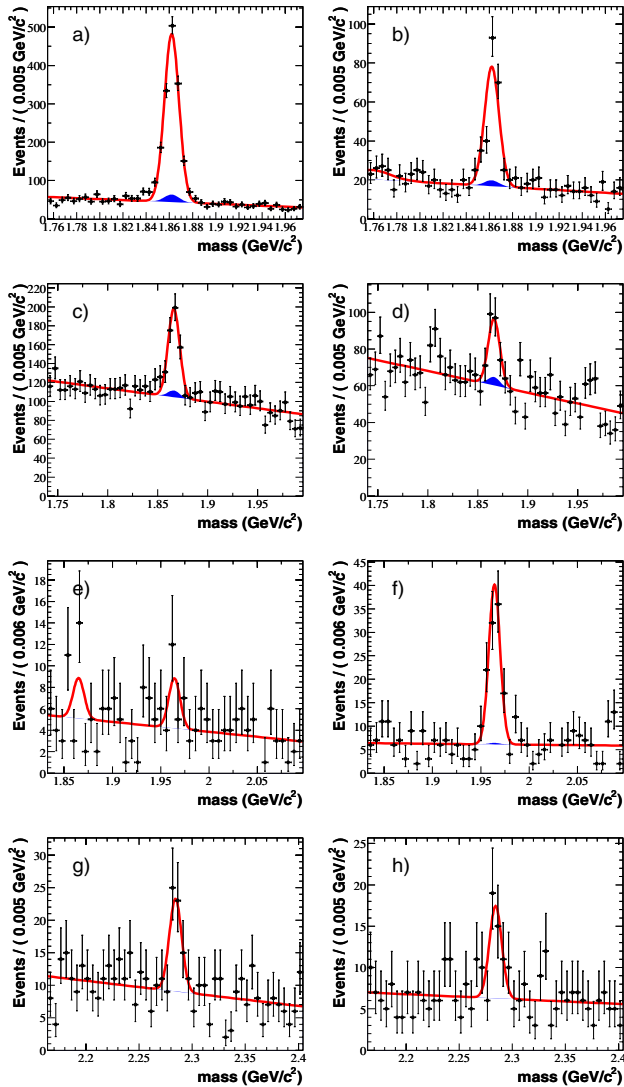


FIG. 2: Correlated (left) and anti-correlated (right) charmed particle mass spectra in the recoil of B^+ events, for (a),(b) $D^0 \rightarrow K^-\pi^+$; (c),(d) $D^+ \rightarrow K^-\pi^+\pi^+$; (e),(f) $D_s^+ \rightarrow \phi\pi^+$; and (g),(h) $A_c^+ \rightarrow pK^-\pi^+$. The solid curve is the sum of a Gaussian signal and of a linear background plus independent satellite contributions [15]. The shaded areas show the contribution of well reconstructed D , D_s or A_c in the B^+ combinatorial background.

D branching fractions, B counting). The results are given in Table IV. We obtain an upper limit on the correlated D_s^+ fraction in B^- decays: $\mathcal{B}(B^- \rightarrow D_s^+ X)/\mathcal{B}(B^- \rightarrow D_s^\pm X) < 0.126$ at 90% CL.

In conclusion, we have measured for the first time the branching fractions for inclusive decays of B mesons to flavor-tagged D , D_s and A_c , separately for B^- and \bar{B}^0 . We observe significant production of anti-correlated D^0 and D^+ mesons in B decays (Table IV), with the branching fractions detailed in Tables II and III. The correlated

TABLE III: Combined \bar{B}^0 branching fractions. The first uncertainty is statistical, the second is systematic, and the third reflects charm branching-fraction uncertainties [16].

X_c	correlated $\mathcal{B}(\bar{B}^0 \rightarrow X_c X)(\%)$	anti-correlated $\mathcal{B}(\bar{B}^0 \rightarrow \bar{X}_c X)(\%)$
D^0	$51.1 \pm 3.1 \pm 2.5^{+1.3}_{-1.3}$	$6.3 \pm 1.9 \pm 0.5^{+0.2}_{-0.2}$
D^+	$39.7 \pm 3.0 \pm 2.8^{+2.8}_{-2.5}$	$2.3 \pm 1.8 \pm 0.3^{+0.2}_{-0.2}$
		< 5.1 at 90% CL
D_s^+	$3.9 \pm 1.7 \pm 0.4^{+1.3}_{-0.8}$	$10.9 \pm 2.1 \pm 0.8^{+3.8}_{-2.3}$
	< 8.7 at 90% CL	
A_c^+	$4.9 \pm 1.7 \pm 0.4^{+1.8}_{-1.0}$	$2.0 \pm 1.2 \pm 0.2^{+0.7}_{-0.4}$
		< 3.8 at 90% CL

TABLE IV: Fraction w of anti-correlated charm.

Mode	B^- decays	\bar{B}^0 decays
$\bar{D}^0 X$	$0.110 \pm 0.010 \pm 0.003$	$0.110 \pm 0.031 \pm 0.008$
$D^- X$	$0.278 \pm 0.052 \pm 0.009$	$0.055 \pm 0.040 \pm 0.006$
$D_s^- X$	$0.966 \pm 0.039 \pm 0.012$	$0.733 \pm 0.092 \pm 0.010$
$\bar{A}_c^- X$	$0.452 \pm 0.090 \pm 0.003$	$0.286 \pm 0.142 \pm 0.007$

D_s production in B^- decays is measured to be small.

As expected, the sum of all correlated charm branching fractions, N_c , is compatible with 1, for charged as well as for neutral B 's. The numbers of charmed particles per B^- decay ($n_c^- = 1.313 \pm 0.037 \pm 0.062^{+0.063}_{-0.042}$) and per \bar{B}^0 decay ($n_c^0 = 1.276 \pm 0.062 \pm 0.058^{+0.066}_{-0.046}$) are consistent with previous measurements [2, 17, 19] and with theoretical expectations [8–10].

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[§] Deceased

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