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Measurement of the $B^0 \to \phi K^{*0}$ Decay Amplitudes

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With a sample of about 227 million $B\overline{B}$ pairs recorded with the *BABAR* detector at the PEP-II storage ring we perform a full angular analysis of the decay $B^0 \to \phi K^{*0}(892)$. We measure the branching fraction to be $(9.2 \pm 0.9 \pm 0.5) \times 10^{-6}$ and determine the fractions of longitudinal and parity-odd transverse contributions as $f_L = 0.52 \pm 0.05 \pm 0.02$ and $f_\perp = 0.22 \pm 0.05 \pm 0.02$, respectively. The phases of the parity-even and parity-odd transverse amplitudes relative to the longitudinal amplitude are found to be $\phi_\parallel = 2.34^{+0.23}_{-0.20} \pm 0.05$ rad and $\phi_\perp = 2.47 \pm 0.25 \pm 0.05$ rad, respectively. We measure five *CP* asymmetries which provide important limits on *CP* violation originating from new physics. We also observe the decay $B^0 \to \phi K^{*0}(1430)$.

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The decay $B \to \phi K^*(892)$ is expected to have contributions from $b \to s$ loop transitions while the tree-level transition is suppressed in the Standard Model. Angular correlation measurements and asymmetries are particularly sensitive to amplitudes arising outside the Standard Model [1]. The first evidence for this decay was provided by the CLEO [2] and *BABAR* [3] experiments. The large fraction of transverse polarization observed by *BABAR* [4] and confirmed by BELLE [5] enables a full angular analysis described by ten parameters for contributing amplitudes and their relative phases.

The angular distribution of the $B \to \phi K^*$ decay products can be expressed as a function of $\mathcal{H}_i = \cos\theta_i$ and Φ, where θ_i is the angle between the direction of the K from the $K^* \to K\pi$ (θ_1) or $\phi \to K\overline{K}$ (θ_2) and the direction opposite the B in the vector resonance rest frame, and Φ is the angle between the two resonance decay planes. The differential decay width has three complex amplitudes A_λ corresponding to the vector meson helicity $\lambda = 0$ or ±1 [1, 6]. When the last two are expressed in terms of $A_\parallel = (A_{+1} + A_{-1})/\sqrt{2}$ and $A_\perp = (A_{+1} - A_{-1})/\sqrt{2}$ we have

$$\begin{aligned}\frac{8\pi}{9\Gamma}\,\frac{d^3\Gamma}{d\mathcal{H}_1 d\mathcal{H}_2 d\Phi} &= \frac{1}{|A_0|^2 + |A_\parallel|^2 + |A_\perp|^2} \times \{ \\ |A_0|^2\,\mathcal{H}_1^2\,\mathcal{H}_2^2 &+ \frac{1}{4}\,(|A_\parallel|^2 + |A_\perp|^2)\,(1-\mathcal{H}_1^2)\,(1-\mathcal{H}_2^2) \\ &+ \frac{1}{4}\,(|A_\parallel|^2 - |A_\perp|^2)\,(1-\mathcal{H}_1^2)\,(1-\mathcal{H}_2^2)\,\cos 2\Phi \\ &- \mathrm{Im}(A_\perp A_\parallel^*)\,(1-\mathcal{H}_1^2)\,(1-\mathcal{H}_2^2)\,\sin 2\Phi \\ &+ \sqrt{2}\,\mathrm{Re}(A_\parallel A_0^*)\,\mathcal{H}_1\,\mathcal{H}_2\,\sqrt{1-\mathcal{H}_1^2}\,\sqrt{1-\mathcal{H}_2^2}\,\cos\Phi \\ &- \sqrt{2}\,\mathrm{Im}(A_\perp A_0^*)\,\mathcal{H}_1\,\mathcal{H}_2\,\sqrt{1-\mathcal{H}_1^2}\,\sqrt{1-\mathcal{H}_2^2}\,\sin\Phi\ \}. \quad (1)\end{aligned}$$

In this analysis, we measure the branching fraction, obtained from the number of reconstructed signal events n_{sig}, the polarization parameters $f_L = |A_0|^2/\Sigma|A_\lambda|^2$, $f_\perp = |A_\perp|^2/\Sigma|A_\lambda|^2$, and the relative phases $\phi_\parallel = \arg(A_\parallel/A_0)$, $\phi_\perp = \arg(A_\perp/A_0)$. We allow for *CP*-violating differences between the $\overline{B}^0$ $(Q=+1)$ and B^0 $(Q=-1)$ decay amplitudes ($\overline{A}_\lambda$ and A_λ), where the flavor sign Q is determined in the self-tagging final state with a $\overline{K}^*$ or K^*:

$$\begin{aligned} n_{\mathrm{sig}}^Q &= n_{\mathrm{sig}}\,(1 + Q\,\mathcal{A}_{CP})/2; \quad (2) \\ f_L^Q = f_L\,(1 + Q\,\mathcal{A}_{CP}^0);& \quad f_\perp^Q = f_\perp\,(1 + Q\,\mathcal{A}_{CP}^\perp); \\ \phi_\parallel^Q = \phi_\parallel + Q\,\Delta\phi_\parallel;& \quad \phi_\perp^Q = \phi_\perp + \frac{\pi}{2} + Q\,(\Delta\phi_\perp + \frac{\pi}{2}).\end{aligned}$$

If one loop diagram dominates the decay amplitude, the three direct *CP* asymmetries $\mathcal{A}_{CP}$, $\mathcal{A}_{CP}^0$, and $\mathcal{A}_{CP}^\perp$, and the two weak-phase differences $\Delta\phi_\parallel$ and $\Delta\phi_\perp$ are expected to be negligible. From the above parameters one can derive vector triple-product asymmetries $\mathcal{A}_T^\parallel$ and $\mathcal{A}_T^0$ as discussed in Ref. [1]:

$$\mathcal{A}_T^{\parallel,0} = \frac{1}{2}\left(\frac{\mathrm{Im}(A_\perp A_{\parallel,0}^*)}{\Sigma|A_\lambda|^2} + \frac{\mathrm{Im}(\overline{A}_\perp \overline{A}_{\parallel,0}^*)}{\Sigma|\overline{A}_\lambda|^2}\right). \quad (3)$$

We use data collected with the *BABAR* detector [7] at the PEP-II asymmetric-energy e^+e^- collider [8] operated at the center-of-mass (CM) energy of the $\Upsilon(4S)$ resonance ($\sqrt{s} = 10.58$ GeV). These data represent an integrated luminosity of about 205 fb^{-1}, corresponding to 226.6 ± 2.5 million $B\overline{B}$ pairs.

Charged-particle momenta are measured in a tracking system consisting of a five-layer double-sided silicon vertex tracker and a 40-layer central drift chamber, both immersed in a 1.5-T solenoidal magnetic field. Charged-particle identification is provided by measurements of the energy loss ($\mathrm{d}E/\mathrm{d}x$) in the tracking devices and by a ring-imaging Cherenkov detector.

We fully reconstruct $\overset{(-)}{B}{}^0 \to \phi \overset{(-)}{K}{}^{*0}$ candidates from their decay products $\phi \to K^+K^-$ and $\overset{(-)}{K}{}^{*0} \to K^\pm\pi^\mp$ as discussed in Ref. [4]. Charged track candidates are required to originate from a single vertex near the interaction point. We identify B meson candidates kinematically using the beam-energy-substituted mass $m_{\mathrm{ES}} = [(s/2 + \mathbf{p}_i \cdot \mathbf{p}_B)^2/E_i^2 - \mathbf{p}_B^2]^{1/2}$ and the energy difference $\Delta E = (E_iE_B - \mathbf{p}_i\cdot\mathbf{p}_B - s/2)/\sqrt{s}$, where $(E_i, \mathbf{p}_i)$ is the initial state four-momentum obtained from the beam momenta, and $(E_B, \mathbf{p}_B)$ is the four-momentum of the reconstructed B candidate. The requirements on the K^* and ϕ invariant masses are $0.75 < m_{K\pi} < 1.05$ and $0.99 < m_{K\overline{K}} < 1.05$ (GeV). We move the selection window to $1.13 < m_{K\pi} < 1.73$ (GeV) in the study of the higher-mass K^* resonances.

To reject the dominant quark-antiquark continuum background, we require $|\cos\theta_T| < 0.8$, where θ_T is the angle between the B-candidate thrust axis and that of the rest of the tracks and neutral clusters in the event, calculated in the CM frame. We also construct a Fisher discriminant, $\mathcal{F}$, further discriminating between signal and background, that combines the following variables: the polar angles of the B-momentum vector and the B-candidate thrust axis with respect to the beam axis in the CM frame, and the two Legendre moments L_0 and L_2 of the energy flow around the B-candidate thrust axis [9].

Contamination from other B decays is small (about 2% of the total background) according to Monte Carlo (MC) simulation [10] and is taken into account in the fit described below. We remove signal candidates that have decay products with invariant mass within 12 MeV of the nominal mass values for $D_s^\pm$ or $D^\pm \to \phi\pi^\pm$.

We use an unbinned, extended maximum-likelihood fit to extract simultaneously the signal yield and angular distributions from a sample of selected events. There are several event categories j: signal, continuum $q\overline{q}$, combinatoric $B\overline{B}$ background, $B \to \phi K\pi$ with a non-resonant S-wave $K^\pm\pi^\mp$ contribution, and $B \to f_0(980)K^*$ with a broad S-wave K^+K^- contribution. The likelihood for each candidate i is defined as $\mathcal{L}_i = \sum_{j,k} n_j^k\, \mathcal{P}_j^k(\vec{x}_i;\vec{\alpha};\vec{\beta})$, where each of the $\mathcal{P}_j^k(\vec{x}_i;\vec{\alpha};\vec{\beta})$ is the probability density function (PDF) for variables $\vec{x}_i = \{m_{\rm ES},\ \Delta E,\ \mathcal{F},\ m_{K\pi},\ m_{K\overline{K}},\ \mathcal{H}_1,\ \mathcal{H}_2,\ \Phi,\ Q\}$. The flavor index k corresponds to the measured value of Q, that is $\mathcal{P}_j^k \equiv \mathcal{P}_j \times \delta_{kQ}$. The n_j^k is the number of events with the flavor k in the category j.

The PDF $\mathcal{P}_j^k(\vec{x}_i;\vec{\alpha};\vec{\beta})$ for a given candidate i is the product of the PDFs for each of the variables and a joint PDF for the helicity angles and resonance masses as discussed below. The signal angular distributions are parameterized with the set $\vec{\alpha} = \{f_L,\ f_\perp,\ \phi_\parallel,\ \phi_\perp,\ \mathcal{A}_{CP}^0,\ \mathcal{A}_{CP}^\perp,\ \Delta\phi_\parallel,\ \Delta\phi_\perp\}$ which are left free to vary in the fit. The other PDF parameters $\vec{\beta}$ are extracted from MC simulation and data in $m_{\rm ES}$ and ΔE sidebands and are fixed in the fit. The MC resolutions are adjusted by comparing data and MC in calibration channels with similar kinematics and topology, such as $B^0 \to D^-\pi^+$ with $D^- \to K^+\pi^-\pi^-$. The PDF parameterization for each event candidate accounts for the loss of acceptance near $\mathcal{H}_1 = 0.8$ due to the $D_s^\pm$ and $D^\pm$ rejection requirements.

We use a three-dimensional description for the helicity part of the signal PDF, using the ideal angular distribution from Eq. (1) multiplied by an acceptance function $\mathcal{G}(\mathcal{H}_1,\mathcal{H}_2,\Phi)$ parameterized with empirical polynomial functions. The detector acceptance effects are found to be uniform in Φ, and we factor the $\mathcal{H}_1$ and $\mathcal{H}_2$ dependence as $\mathcal{G} \equiv \mathcal{G}_1(\mathcal{H}_1) \times \mathcal{G}_2(\mathcal{H}_2)$. We use two Gaussian functions for the parameterization of the signal PDFs for ΔE, $m_{\rm ES}$, and $\mathcal{F}$. A relativistic P-wave Breit-Wigner distribution, convoluted with a Gaussian resolution function, is used for the resonance masses.

TABLE I: Summary of the $B^0 \to \phi K^{*0}(892)$ fit results. We show results for the ten primary signal fit parameters defined in Eq. (2) and the derived parameters: reconstruction efficiency ϵ which depends on decay polarization, branching fraction $\mathcal{B}$, and triple-product asymmetries from Eq. (3). All results include systematic errors, which are quoted following the statistical errors. For the dominant correlations we give the coefficients in the last column.

Fit parameter	Fit result	Correlation
$n_{\rm sig}$ (events)	$201 \pm 20 \pm 6$	
f_L	$0.52 \pm 0.05 \pm 0.02$	$\}\ -46\%$
$f_\perp$	$0.22 \pm 0.05 \pm 0.02$	
$\phi_\parallel$ (rad)	$2.34^{+0.23}_{-0.20} \pm 0.05$	$\}\ +70\%$
$\phi_\perp$ (rad)	$2.47 \pm 0.25 \pm 0.05$	
$\mathcal{A}_{CP}$	$-0.01 \pm 0.09 \pm 0.02$	
$\mathcal{A}_{CP}^0$	$-0.06 \pm 0.10 \pm 0.01$	$\}\ -45\%$
$\mathcal{A}_{CP}^\perp$	$-0.10 \pm 0.24 \pm 0.05$	
$\Delta\phi_\parallel$ (rad)	$0.27^{+0.20}_{-0.23} \pm 0.05$	$\}\ +70\%$
$\Delta\phi_\perp$ (rad)	$0.36 \pm 0.25 \pm 0.05$	
ϵ (%)	9.7 ± 0.5	
$\mathcal{B}$	$(9.2 \pm 0.9 \pm 0.5) \times 10^{-6}$	
$\mathcal{A}_T^\parallel$	$-0.02 \pm 0.04 \pm 0.01$	
$\mathcal{A}_T^0$	$+0.11 \pm 0.05 \pm 0.01$	

Parameterization of the non-resonant B-decay contributions is identical to that of the signal for $m_{\rm ES}$, ΔE, and $\mathcal{F}$, but is different for the angular and invariant mass distributions. In particular, a broad invariant mass distribution accounts for all potential S-wave contributions leaking into the mass selection window. For the combinatorial background, we use polynomials, except for $m_{\rm ES}$ and $\mathcal{F}$ distributions which are parameterized by an empirical phase-space function and by the two Gaussian functions, respectively. Resonance production occurs in the background and this is taken into account in the PDF. The background $\mathcal{H}_i$ distribution is separated into contributions from combinatorial background and from real vector mesons.

We allow for multiple candidates in a given event by assigning to each a weight of $1/N_i$, where N_i is the number of candidates in the same event. The average number of candidates per event is 1.04. The extended likelihood for a sample of $N_{\rm cand}$ candidates is

$$\mathcal{L} = \exp\left(-\sum_j n_j\right) \prod_{i=1}^{N_{\rm cand}} \exp\left(\frac{\ln \mathcal{L}_i}{N_i}\right). \qquad (4)$$

The event yields n_j, asymmetries $\mathcal{A}_j$, and the signal po-

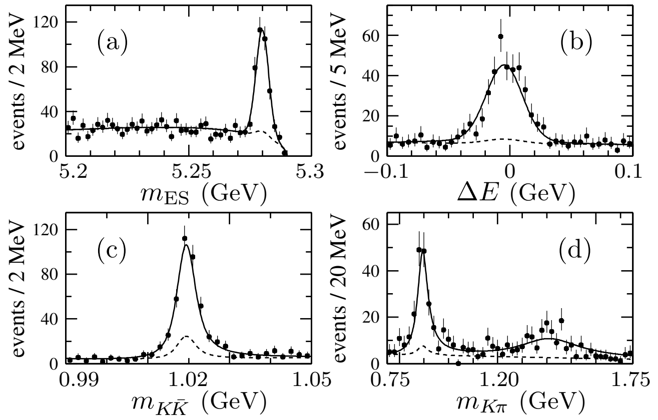

FIG. 1: Projections onto the variables m_{ES} (a), ΔE (b), $m_{K\overline{K}}$ (c), and $m_{K\pi}$ (d) for the signal $B^0 \to \phi K^{*0}(892)$ and $\phi K^{*0}(1430)$ candidates combined.

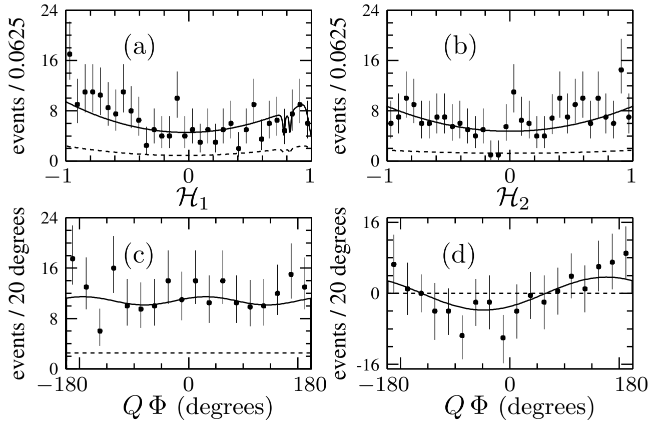

FIG. 2: Projections onto the variables $\mathcal{H}_1$ (a), $\mathcal{H}_2$ (b), $Q\,\Phi$ (c), and the differences between the $Q\,\Phi$ projections for events with $\mathcal{H}_1\,\mathcal{H}_2 > 0$ and with $\mathcal{H}_1\,\mathcal{H}_2 < 0$ (d) for the signal $B^0 \to \phi K^{*0}(892)$ candidates.

larization parameters $\vec{\alpha}$ are obtained by maximizing $\mathcal{L}$.

The results of our maximum likelihood fit to the sample of $B^0 \to \phi K^{*0}(892)$ candidates are summarized in Table I. We also repeat the fit with the requirement $1.13 < m_{K\pi} < 1.73$ (GeV) and without the angular information. We observe 181 ± 17 events (statistical errors only) of the decays $B^0 \to \phi K^{*0}(1430)$ with statistical significance greater than 10σ. In Fig. 1–3 we show projections onto the variables, where data distributions are shown with a requirement on the signal-to-background probability ratio $\mathcal{P}_{\mathrm{sig}}/\mathcal{P}_{\mathrm{bkg}}$ calculated with the plotted variable excluded. The solid (dashed) lines show the signal-plus-background (background) PDF projections.

In the analysis of the decay $B^0 \to \phi K^{*0}(892)$ for any given set of values $(\phi_{\parallel}, \phi_{\perp}, \Delta\phi_{\parallel}, \Delta\phi_{\perp})$ simple transformations of the angles, for example $(-\phi_{\parallel}, \pi - \phi_{\perp}, -\Delta\phi_{\parallel}, -\Delta\phi_{\perp})$, give rise to other sets of values which satisfy Eq. (1) in an identical manner. To resolve this ambiguity, the set of values lying closest to the theoretical

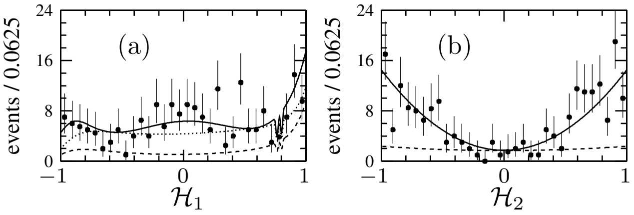

FIG. 3: Projections onto the variables $\mathcal{H}_1$ (a) and $\mathcal{H}_2$ (b) for the signal $B^0 \to \phi K^{*0}(1430)$ candidates. The difference between the solid and dotted lines in (a) shows the contribution of the tensor state to the angular distribution.

expectation $(\pi, \pi, 0, 0)$ [1, 6, 11] is chosen. In Fig. 4 we show likelihood function contour plots.

We find the decay $B^0 \to \phi K^{*0}(1430)$ to be predominantly longitudinally polarized based on the $\mathcal{H}_2$ angular distribution in Fig. 3 (b). The width [12] and the angular distribution of the $K^{*0}(1430)$ resonance structure are not consistent with the pure $K_2^{*0}(1430)$ tensor state at more than 10σ. However, the angular distribution provides evidence (with statistical significance of 3.2σ) of the longitudinally polarized tensor $K_2^{*0}(1430)$ contribution in addition to the scalar $K_0^{*0}(1430)$, see Fig. 3 (a).

Our $B^0 \to \phi K^{*0}(892)$ fit is performed with the $B \to f_0 K^*$ and $B \to \phi K\pi$ contributions unconstrained. We obtain the event yields 25 ± 10 and 11 ± 15, respectively. The systematic uncertainties due to interference are estimated using generated samples with conservative assumptions about the S-wave intensity and the interference phase. Additional systematic uncertainty originating from B background is taken as the difference between the fit results with the combinatoric $B\overline{B}$ background component fixed to zero and fixed to the expectation from MC.

We vary the PDF parameters within their respective uncertainties, and derive the associated systematic errors. The biases from the finite resolution of the helicity angle measurement and the dilution due to the presence of fake combinations are estimated with MC simulation.

The systematic errors in efficiencies are dominated by those in track finding and particle identification. Other systematic effects arise from event-selection criteria, ϕ and K^{*0} branching fractions, MC statistics, and number of B mesons. We calculate the efficiencies using the measured polarization and assign a systematic error corresponding to the total polarization uncertainty. We find the uncertainty in the charge asymmetry due to the track reconstruction and identification to be less than 0.02.

In summary, we have performed a full angular analysis and searched for *CP* violation in the angular distribution with the decays $\overset{(-)}{B}{}^0 \to \phi \overset{(-)}{K}{}^{*0}(892)$. Our results are summarized in Table I. We observe, with more than 5σ significance, non-zero contributions from all of the three amplitudes $|A_0|$, $|A_{\perp}|$, and $|A_{\parallel}|$, see Fig. 4 (a). We also find 3σ evidence for non-zero final-state-interaction

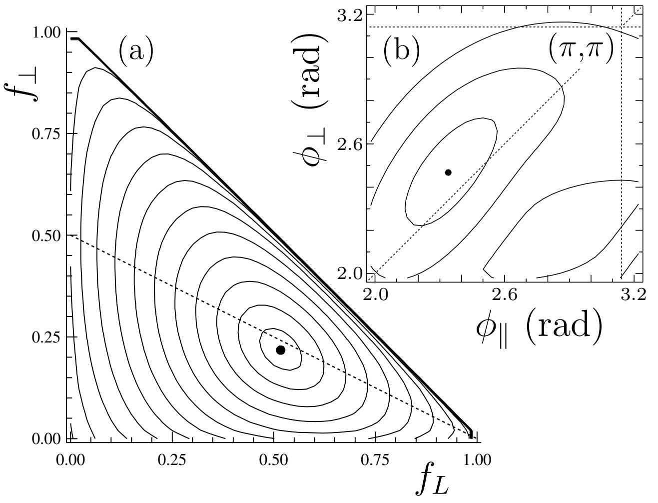

FIG. 4: Likelihood function contours with 1σ intervals for polarization (a) and phase (b) measurements in the $B^0 \to \phi K^{*0}(892)$ analysis. The fit results are shown with dots. Diagonal dashed lines $f_\perp = (1 - f_L)/2$ and $\phi_\perp = \phi_\parallel$ correspond to $|A_{+1}| \gg |A_{-1}|$. In (b) the (π, π) point is indicated by the crossed dashed lines.

phases, see Fig. 4 (b). These results supersede our earlier measurements in this channel [3, 4]. We also observe the decay $B^0 \to \phi K^{*0}(1430)$.

For B decays to light charmless particles we expect the hierarchy of decay amplitudes to be $|A_0| \gg |A_{+1}| \gg |A_{-1}|$ under the assumption of pure loop diagram contribution, which is analogous to the discussion in Ref. [11]. Our measurements with the decay $B^0 \to \phi K^{*0}(892)$ do not agree with the first inequality but agree with the previous measurements in Ref. [4, 5]. This suggests other contributions to the decay amplitude, previously neglected, either within or beyond the Standard Model [1, 13].

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