

Measurement of $\sin 2\beta$ with *BABAR*

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We present updated results on time-dependent *CP*-violating asymmetries in neutral *B* decays to several *CP* eigenstates containing charmonium. In the Standard Model, the amplitude of these asymmetries is proportional to $\sin 2\beta$. We measure $\sin 2\beta = 0.741 \pm 0.067$ (stat) ± 0.034 (syst) from a data sample of about 88 million $\Upsilon(4S) \rightarrow B\bar{B}$ decays collected with the *BABAR* detector at the PEP-II asymmetric-energy *B* Factory at SLAC. We have also measured *CP*-violating asymmetries in open charm, and penguin modes sensitive to $\sin 2\beta$, which provide important consistency tests of the Standard Model.

1. INTRODUCTION

In the Standard Model of electroweak interactions, *CP* violation arises as a consequence of a complex phase in the three-generation Cabibbo-Kobayashi-Maskawa (CKM) quark-mixing matrix. Observations of *CP* violation in *B*⁰ decays to charmonium were reported last year by the *BABAR* [1] and Belle [2] collaborations. The PEP-II collider at SLAC has since delivered an additional 63 fb⁻¹, thereby approximately tripling the *BABAR* data sample near the $\Upsilon(4S)$ resonance. Here we report a more precise measurement [3], as well as updated and new results from non-charmonium *B*⁰ decays, using a sample of about 88 million $B\bar{B}$ decays. Changes in the analysis with respect to the previously published result [1] include the processing of all data with a uniform event reconstruction, a new flavor-tagging algorithm, and the addition of the decay mode $B^0 \rightarrow \eta_c K_s^0$. The *BABAR* detector and the measurement technique are described in detail in Refs. [4] and [5].

We reconstruct a sample of neutral *B* mesons (*B*_{*CP*}) decaying to *CP* eigenstates and examine each event for evidence that the recoiling *B* meson decayed as a *B*⁰ or \bar{B}^0 (flavor tag). The proper-time distribution for such events can be expressed in terms of a complex parameter λ that depends on both the *B*⁰- \bar{B}^0 oscillation amplitude and the amplitudes describing \bar{B}^0 and *B*⁰ decays to the

CP final state. The decay rate $f_{\pm}(f_{\pm})$, when the tagging meson is a *B*⁰ (\bar{B}^0), is given by

$$f_{\pm}(\Delta t) = \frac{e^{-|\Delta t|/\tau_{B^0}}}{4\tau_{B^0}} \left[1 \pm \frac{2\mathcal{I}m\lambda}{1+|\lambda|^2} \sin(\Delta m_d \Delta t) \mp \frac{1-|\lambda|^2}{1+|\lambda|^2} \cos(\Delta m_d \Delta t) \right], \quad (1)$$

where $\Delta t = t_{\text{rec}} - t_{\text{tag}}$ is the difference between the proper decay times of the reconstructed *B* meson (*B*_{rec}) and the tagging *B* meson (*B*_{tag}), τ_{B^0} is the *B*⁰ lifetime, and Δm_d is the *B*⁰- \bar{B}^0 oscillation frequency.

In the Standard Model, $\lambda = \eta_f e^{-2i\beta}$ for charmonium modes (e.g. $B^0 \rightarrow J/\psi K_s^0$), where η_f is the *CP* eigenvalue of the final state *f*, and $\beta \equiv \arg[-V_{cd}V_{cb}^*/V_{td}V_{tb}^*]$. The time-dependent *CP*-violating asymmetry then becomes

$$\begin{aligned} A_{CP}(\Delta t) &\equiv \frac{f_{+}(\Delta t) - f_{-}(\Delta t)}{f_{+}(\Delta t) + f_{-}(\Delta t)} \\ &= -\eta_f \sin 2\beta \sin(\Delta m_d \Delta t). \end{aligned} \quad (2)$$

2. TAGGING AND VERTEXING

A measurement of *A*_{*CP*} requires a determination of the experimental Δt resolution and the fraction *w* of events in which the tag assignment is incorrect. This mistag fraction reduces the observed *CP* asymmetry by a factor $1 - 2w$. Mistag fractions and Δt resolution functions are determined from a sample of neutral *B* mesons that

Table 1

Tagging efficiency, mistag fraction and quality factor for each category (measured from data).

Category	ε (%)	w (%)	Q (%)
Lepton	9.1 ± 0.2	3.3 ± 0.6	7.9 ± 0.3
Kaon I	16.7 ± 0.2	10.0 ± 0.7	10.7 ± 0.4
Kaon II	19.8 ± 0.3	20.9 ± 0.8	6.7 ± 0.4
Inclusive	20.0 ± 0.3	31.5 ± 0.9	2.7 ± 0.3
All	65.6 ± 0.5		28.1 ± 0.7

decay to flavor eigenstates (B_{flav}) consisting of the channels $D^{(*)-}h^+(h^+ = \pi^+, \rho^+, \text{ and } a_1^+)$ ¹ and $J/\psi K^{*0}(K^{*0} \rightarrow K^+\pi^-)$.

We use multivariate algorithms to identify signatures of B decays that indicate the flavor of B_{tag} . Based on the output of a multi-level neural network, we assign each event to one of four hierarchical, mutually exclusive tagging categories. The **Lepton** category contains events with an identified lepton. Events with kaon and soft pion candidates (but not identified leptons) are assigned to the **Kaon I** or **Kaon II** category. All other events are assigned to the **Inclusive** category or excluded from further analysis based on the estimated mistag probability. The tagging efficiencies ε_i for the four tagging categories are measured from data and are summarized in Table 1. The figure of merit for tagging is the effective tagging efficiency $Q \equiv \sum_i \varepsilon_i (1 - 2w_i)^2$. This algorithm improves Q by about 7% (relative) over the algorithm used in Ref. [5].

The time interval Δt between the two B decays is calculated from the measured separation Δz between the decay vertices of B_{rec} and B_{tag} along the collision (z) axis. We employ constraints from the beam spot location and the B_{rec} momentum to improve vertex reconstruction efficiency and resolution. The fraction of events satisfying the vertex quality requirements is 95%. The r.m.s. Δt resolution for practically all accepted events is 1.1 ps.

¹Charge conjugation is implied throughout this paper, unless explicitly stated.

3. CHARMONIUM MODES

3.1. Event sample

We reconstruct B^0 decays to the following final states containing charmonium: $\eta_f = -1$ modes $J/\psi K_S^0$, $\psi(2S)K_S^0$, $\chi_{c1}K_S^0$, $\eta_c K_S^0$; $\eta_f = +1$ mode $J/\psi K_L^0$; and $J/\psi K^{*0}(K^{*0} \rightarrow K_S^0 \pi^0)$. This last mode, $J/\psi K^{*0}$, is not a pure CP eigenstate, but is a combination of even ($L=0, 2$) and odd ($L=1$) orbital angular momenta final states. In this analysis we do not separate the angular components, so the measured CP asymmetry is reduced by a factor $1 - 2R_\perp$, where R_\perp is the fraction of the $L=1$ component. We have measured $R_\perp = (16.0 \pm 3.5)\%$ [6], which gives an effective $\eta_f = 0.65 \pm 0.07$ for the $J/\psi K^{*0}$ mode after acceptance corrections.

The event selection criteria for the B_{CP} sample are described in Ref. [3]. The selection includes two main discriminating variables constructed from the measured energy and momentum of the final state: the difference ΔE between the decay energy and the beam energy in the center-of-mass frame, and the beam-energy substituted mass $m_{\text{ES}} = \sqrt{(E_{\text{beam}}^{\text{cm}})^2 - (p_B^{\text{cm}})^2}$. Since it is not well measured, we determine the K_L^0 momentum in the $J/\psi K_L^0$ mode by applying one of the kinematic constraints (effectively m_{ES}), leaving only one (ΔE) of the two discriminating variables for the selection of this mode. Figure 1 shows the m_{ES} distribution for modes containing a K_S^0 or K^{*0} and the ΔE distribution for the $J/\psi K_L^0$ candidates.

To determine numbers of events and purities, a signal region 5.270 (5.273) $< m_{\text{ES}} < 5.290$ (5.288) GeV/c^2 is used for modes containing K_S^0 (K^{*0}). In the $J/\psi K_L^0$ mode, the ΔE resolution the signal region is defined by $|\Delta E| < 10$ MeV. The signal region contains 2641 events that satisfy the tagging and vertexing requirements. In Table 2 we list the number of events and the signal purity for the tagged B_{CP} candidates.

3.2. Measurement of $\sin 2\beta$

We determine $\sin 2\beta$ with a simultaneous unbinned maximum likelihood fit to the Δt distributions of the tagged B_{CP} and B_{flav} samples. In

this fit, the Δt distributions of the B_{CP} sample are described by Eq. 1 with $|\lambda| = 1$, i.e. with the assumption of no direct CP violation. The Δt distributions of the B_{flav} sample evolve according to the known frequency for flavor oscillation in B^0 mesons. The observed amplitudes for the CP asymmetry and flavor oscillation are reduced by the same factor, $1 - 2w$, due to flavor mistags.

There are 34 free parameters in the fit: $\sin 2\beta$ (1), the average mistag fractions w and the differences Δw between B^0 and \bar{B}^0 mistag fractions for each tagging category (8), parameters for the signal Δt resolution (8), and parameters for background time dependence (6), Δt resolution (3), and mistag fractions (8). We fix $\tau_{B^0} = 1.542$ ps and $\Delta m_d = 0.489 \text{ ps}^{-1}$ [7]. The fit to the B_{CP} and B_{flav} samples yields

$$\sin 2\beta = 0.741 \pm 0.067 \text{ (stat)} \pm 0.034 \text{ (syst)}.$$

Figure 2 shows the Δt distributions and the raw asymmetry $(N_{B^0} - N_{\bar{B}^0}) / (N_{B^0} + N_{\bar{B}^0})$ versus Δt for the $\eta_f = -1$ and $\eta_f = +1$, overlaid with the projection of the likelihood fit result.

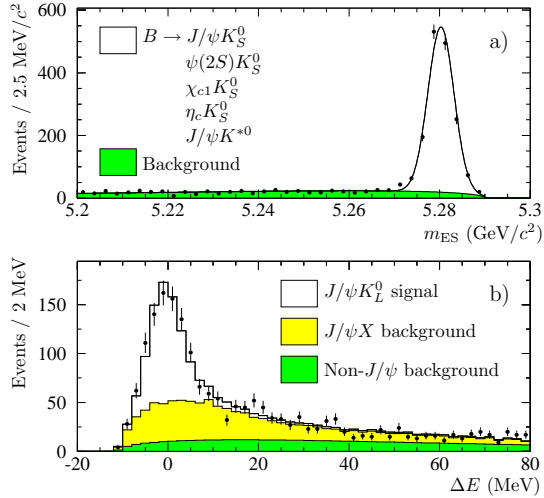


Figure 1. Distributions for B_{CP} candidates satisfying the tagging and vertexing requirements: a) all final states except $J/\psi K_L^0$, and b) $J/\psi K_L^0$.

Table 2

Number of events in the signal region after tagging and vertexing requirements, signal purity, and results of the CP asymmetry fits.

Sample	N_{tag}	$P(\%)$	$\sin 2\beta$
$(J/\psi, \psi(2S), \chi_{c1}, \eta_c)K_S^0$	1506	94	0.76 ± 0.07
$J/\psi K_L^0$ ($\eta_f = +1$)	988	55	0.72 ± 0.16
$J/\psi K^{*0}$ ($K^{*0} \rightarrow K_S^0 \pi^0$)	147	81	0.22 ± 0.52
Full CP sample	2641	78	0.74 ± 0.07
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$(J/\psi, \psi(2S), \chi_{c1}, \eta_c)K_S^0$ only ($\eta_f = -1$)			
$J/\psi (K_S^0 \rightarrow \pi^+ \pi^-)$	974	97	0.82 ± 0.08
$J/\psi (K_S^0 \rightarrow \pi^0 \pi^0)$	170	89	0.39 ± 0.24
$\psi(2S) (K_S^0 \rightarrow \pi^+ \pi^-)$	150	97	0.69 ± 0.24
$\chi_{c1} K_S^0$	80	95	1.01 ± 0.40
$\eta_c K_S^0$	132	73	0.59 ± 0.32
Lepton category	220	98	0.79 ± 0.11
Kaon I category	400	93	0.78 ± 0.12
Kaon II category	444	93	0.73 ± 0.17
Inclusive category	442	92	0.45 ± 0.28
B^0 tags	740	94	0.76 ± 0.10
\bar{B}^0 tags	766	93	0.75 ± 0.10
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B_{flav} sample	25375	85	0.02 ± 0.02
B^+ sample	22160	89	0.02 ± 0.02

The dominant sources of systematic error are the uncertainties in the background content, the assumed parameterization of the Δt resolution function, and possible differences between the B_{flav} and B_{CP} mistag fractions. Most systematic errors are determined with data and will continue to decrease with additional statistics.

The large B_{CP} sample allows a number of consistency checks, including separation of the data by decay mode, tagging category, and B_{tag} flavor. The results of fits to these subsamples are found to be statistically consistent. Fits to the control samples of non- CP decay modes (the B_{flav} sample and B^+ mesons decaying to the final states $J/\psi K^{(*)+}$, $\psi(2S)K^+$, $\chi_{c1}K^+$, $\eta_c K^+$, and $\bar{D}^{(*)0} \pi^+$) indicate no statistically significant asymmetry, as expected. The results of these checks are given in Table 2.

This measurement of $\sin 2\beta$ improves upon the precision of each of the previous measurements [1,2] by a factor of two. The measured value is in good agreement with the expectation implied by the measurements and theoretical esti-

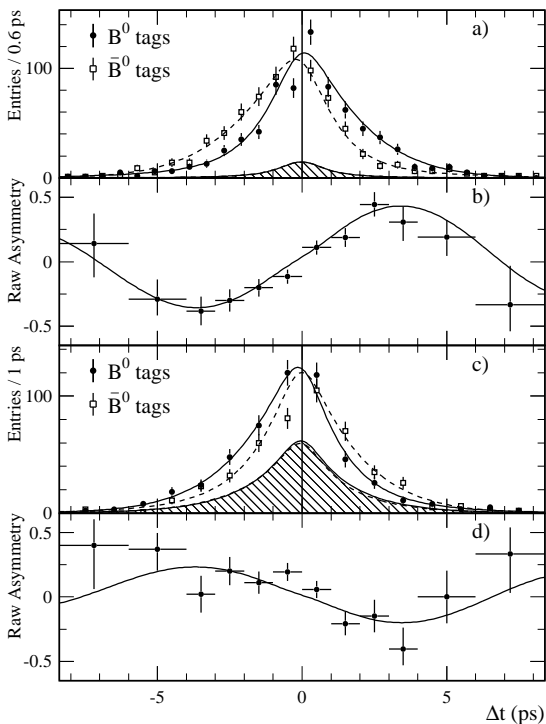


Figure 2. Number of tagged events, raw asymmetry, and fit projections as functions of Δt for $\eta_f = -1$ (a & b) and $\eta_f = +1$ (c & d) candidates. The shaded regions represent the background contributions.

mates of the magnitudes of CKM matrix elements in the context of the Standard Model, and provides a precise and model-independent constraint on the position of the apex of the Unitarity Triangle.

3.3. Direct CP violation

We also measure the parameter $|\lambda|$ in Eq. 1 from a fit to the $\eta_f = -1$ sample, which has high purity and requires minimal assumptions on the effect of backgrounds. This parameter is sensitive to the difference in the number of B^0 - and \bar{B}^0 -tagged events. In order to account for differences in reconstruction and tagging efficiencies for B^0 and \bar{B}^0 mesons, we incorporate five ad-

ditional free parameters in this fit. We obtain $|\lambda| = 0.948 \pm 0.051$ (stat) ± 0.030 (syst). We simultaneously fit the coefficient of the sine term in Eq. 1 and find it to be 0.759 ± 0.074 (stat). These results are consistent with the no direct CP violation hypothesis ($|\lambda| = 1$) and with our $\sin 2\beta$ measurement.

4. OPEN CHARM MODE

Within the context of the Standard Model, the open charm decay $B^0 \rightarrow D^{*+}D^{*-}$ is expected to have a time-dependent CP -violating asymmetry [8]. Up to corrections due to theoretically uncertain penguin diagram contributions [9], the Cabibbo-suppressed tree diagram produces an asymmetry sensitive to $\sin 2\beta$. Factorization models predict that the penguin-induced correction is small [10], so comparisons with the CP asymmetry in charmonium modes provides a valuable test of these models and a consistency check on the Standard Model.

The selection of $B^0 \rightarrow D^{*+}D^{*-}$ candidates is described in Ref [11] and results in a sample of 102 tagged events with a purity of 82%. The $B^0 \rightarrow D^{*+}D^{*-}$ mode consists of three partial wave contributions with two different CP parities: even (S - and D -waves) and odd (P -wave). Since CP -even and odd components can be separated by a one-dimensional angular analysis in the transversity basis [12], we parameterize the tagged B^0 decay rate in terms of the transversity angle and fit for the complex CP -even (odd) parameters $\lambda_+(\lambda_-)$, and R_\perp the fraction of signal events with CP -odd parity.

From a time-integrated fit and corrected for acceptance, we measure $R_\perp = 0.07 \pm 0.06$ (stat) ± 0.03 (syst). Since the CP -odd fraction is small, we fix the CP -odd parameters in the time-dependent fit to those measured from the charmonium analysis, assuming the penguin diagram contribution is negligible, i.e. $|\lambda_\perp| = 1$ and $\text{Im}(\lambda_\perp) = -0.741$. With a suitably large statistical sample, one can fit the CP -odd and even parameters simultaneously.

Performing a time-dependent fit including both cosine and sine terms in Eq. 1 yields the prelimi-

nary result [11]

$$\begin{aligned}\text{Im}(\lambda_+) &= 0.31 \pm 0.43(\text{stat}) \pm 0.13(\text{syst}) \\ |\lambda_+| &= 0.98 \pm 0.25(\text{stat}) \pm 0.09(\text{syst}).\end{aligned}$$

If the penguin contribution can be neglected, then $|\lambda_+| = 1$ and $\text{Im}(\lambda_+) = -\sin 2\beta$. By making this assumption and refitting, we find that the change in the likelihood corresponds to a 2.7 standard deviation from the result of the fit in charmonium ($\sin 2\beta = 0.741$).

5. PENGUIN MODES

Since the tree diagram is Cabibbo-suppressed, the decay $B^0 \rightarrow J/\psi \pi^0$ ($\eta_f = +1$) has comparable tree and penguin contributions. The tree diagram has the same weak phase as charmonium modes containing neutral kaons (e.g. $B^0 \rightarrow J/\psi K_S^0$); however, a portion of the penguin diagram contributes a different weak phase. We therefore perform a time-dependent CP -violating asymmetry fit for the coefficients ($S_{J/\psi \pi^0}$ and $C_{J/\psi \pi^0}$) of the sine and cosine terms in Eq. 1.

We select 49 tagged events with a purity of 59% and measure the preliminary result [13]

$$\begin{aligned}C_{J/\psi \pi^0} &= 0.38 \pm 0.41 (\text{stat}) \pm 0.09 (\text{syst}), \\ S_{J/\psi \pi^0} &= 0.05 \pm 0.49 (\text{stat}) \pm 0.16 (\text{syst}).\end{aligned}$$

In the absence of the penguin contribution, we expect $C_{J/\psi \pi^0} = 0$ and $S_{J/\psi \pi^0} = -\sin 2\beta$.

The charmless decay $B^0 \rightarrow \phi K_S^0$ ($\eta_f = -1$) is dominated by gluonic penguin diagrams and within the Standard Model should have the same weak phase as the charmonium modes. The predicted deviation of the effective $\sin 2\beta$ measured in this mode, compared to the Standard Model parameter, is smaller than 4% [14]. A measurement of $\sin 2\beta$ in ϕK_S^0 therefore probes new physics participating in penguin loops.

From the analysis described in Ref. [15], we find 66 tagged candidates with a purity of 50%. Setting $|\lambda| = 1$ and fitting only the coefficient $S_{\phi K_S^0}$ of the sine term in Eq. 1, we determine $S_{\phi K_S^0} = -0.19_{-0.50}^{+0.52}(\text{stat}) \pm 0.09(\text{syst})$. In the absence of new physics, $S_{\phi K_S^0} = \sin 2\beta$ at the few percent level.

In these first measurements from *BABAR* on the CP -violating asymmetry in penguin modes sensitive to $\sin 2\beta$, the statistical precision of the results do not yet provide a meaningful test of the Standard Model; however, high statistics in the future provide a good opportunity to challenge the theory.

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