

SLAC-PUB-1055
(TH) (EXP)
June 1972

Scaling and Locality with No Pomeron in Electroproduction*

Ashok suri
Division of Natural Sciences
University of California, Santa Cruz
Santa Cruz, California 95060

* Work supported by the U. S. Atomic Energy Commission.

(Submitted for publication)

ABSTRACT

A recent sum rule of Brandt and Ng implies the existence of a pomeron contribution in electroproduction. This is shown to be based on an unrealistic assumption. Therefore the conclusions drawn from it are physically unfounded. In particular, electroproduction need not have a pomeron but still scale in a causal model.

$$\int_0^1 \bar{F}(\omega) / \omega \, d\omega = \pi \neq 0.$$

This shows that the Brandt-Ng sum rule is invalid and there is, therefore, no need for a pomeron.

In configuration space, our example corresponds to

$$vW_2(\kappa, \nu) = \kappa v \tilde{C}(\kappa, \nu),$$

where the fourier transform of \tilde{C} gives (in the frame, $P=0$)

$$C(y^2, y_0) = -8\varepsilon(y_0)\theta(y^2) \left\{ \sin y_0/y_0^3 - 3 \cos y_0/y_0^4 + 3 \sin y_0/y_0^3 \int_0^\infty da \, a(1-a)e^{-a} J_1(\sqrt{ay^2})/\sqrt{ay^2} \right\}.$$

This clearly has the proper light cone singularity for scaling and corresponds to short-range terms discussed by suri and Yennie⁴. These short-range terms are essential even in the presence of Regge behavior to maintain proper support restrictions.⁴

To sum up, our counter example shows that the Brandt-Ng sum rule is invalid. The trouble with the sum rule arises not from the general physical principles incorporated in (1) and (2), but from the specific assumption (3), which is incorrect. Therefore, the results based on the validity of this sum rule need not be correct. In particular electroproduction may have no contribution from the pomeron but still scale in a causal model.

ACKNOWLEDGEMENTS

We wish to thank S. J. Brodsky, J. F. Gunion, C. H. Llewellyn-Smith and J. D. Bjorken for useful comments. Thanks are due to S. D. Drell for the hospitality at SLAC.

APPENDIX

Using the DGS representation, it is easy to show ^{2, 6} that

$$\nu W_2 \equiv \bar{W}(\kappa, \nu) = \kappa \nu \int_{-1}^1 db \int_0^\infty da \sigma(a, b) \delta(\kappa + 2b\nu - a) \epsilon(\nu + b).$$

Scaling then requires that

$$\int_0^\infty da \sigma(a, b) = 0, \quad F_2(\omega) = -\omega/2 \int_0^\infty da a \left\{ \frac{\partial}{\partial \omega} \sigma(a, \omega) \right\}.$$

From this it follows that

$$\int_0^1 d\omega/\omega F_2(\omega) = -\frac{1}{2} \int_0^\infty da a \{ \sigma(a, 1) - \sigma(a, 0) \}.$$

Mass spectrum condition and smoothness requires $\sigma(a, 1) = 0$. In the non-pomeron case (or for the non-pomeron part), Brandt and Ng incorrectly assume that $F(\omega) \rightarrow 0$ as $\omega \rightarrow 0$ requires $\bar{\sigma}(a, 0) = 0$ and they get the sum rule:

$$\int_0^1 d\omega/\omega \bar{F}_2(\omega) = 0,$$

where they argue that $\bar{F}_2(\omega) = F_2(\omega) - F_2(0)$.

REFERENCES

1. Richard A. Brandt and Wing-Chiu Ng, Locality in Electroproduction, NYU Technical Report No. 13/72 (May 1972), New York University, New York, N.Y. 10003. Our notation follows their paper.
2. S. Deser, W. Gilbert and E.G.G. Sudarshan, Phys. Rev. 115, 731 (1959).
3. Ashok suri, Phys. Rev. D4, 570 (1971) and references therein.
4. Ashok suri and Donald R. Yennie, Ann. of Phy. 71, (July, 1972) (to be published).
5. S.D. Drell and Tung-Mow Yan, Ann. of Phy. 66, 578 (1971).
6. R.A. Brandt, Phy. Rev. D1, 2808 (1970).