# DISCREPANCY BETWEEN THE VECTOR DOMINANCE MODEL AND PION PRODUCTION BY POLARIZED PHOTONS* 

## by

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#### Abstract

The helicity density matrix for the process $\pi^{-} \mathrm{p} \rightarrow \pi^{+} \pi^{-} \mathrm{n}$ has been studied as a function of $\mathrm{M}_{\pi \pi^{*}}$. The ratio $\rho_{1-1}^{\text {hel }} / \rho_{11}^{\text {hel }}$ shows no variation with $\mathrm{M}_{\pi \pi}$ and it thus seems unlikely that the large discrepancy between the vector dominance prediction and the observed asymmetry of single pions produced by linearly polarized photons could be the result of density-matrix distortion caused by background interference.


[^0]The vector dominance model has been successfully used ${ }^{1}$ to relate single pion photoproduction

$$
\begin{align*}
& y \mathrm{p} \rightarrow \pi^{+} \mathrm{n}  \tag{1a}\\
& \gamma \mathrm{n} \rightarrow \pi^{-} \mathrm{p} \tag{1b}
\end{align*}
$$

to $\rho^{\circ}$ production by pions

$$
\begin{equation*}
\pi^{-} p \rightarrow \rho^{\circ} \mathrm{n} \tag{2}
\end{equation*}
$$

By taking the sum of the two photoproduction cross sections, the $\omega \rho$ interference terms drop out; since the $\omega$ term itself is expected to be only a few per cent, it is neglected, giving the $\rho$ dominance prediction

$$
\begin{equation*}
\frac{\sigma^{+}+\sigma^{-}}{2}=\frac{\pi \alpha}{\gamma_{\rho}^{2}}\left[\rho_{11} \frac{\mathrm{~d} \sigma}{\mathrm{dt}}\right]_{\pi-\rho \rightarrow \rho^{-} \mathrm{n}} \tag{3}
\end{equation*}
$$

where $\sigma^{+}$and $\sigma^{-}$are $\frac{\mathrm{d} \sigma}{\mathrm{dt}}$ for reactions (1a) and (1b), respectively.
Recently, experiments on reactions (1a) and (1b) ${ }^{3,4}$ have been performed using linearly polarized $\gamma$-rays produced coherently from crystals. These experiments yield the asymmetries

$$
\begin{equation*}
\mathrm{A}^{ \pm}=\frac{\sigma_{\perp}^{ \pm}-\sigma_{\| I}^{ \pm}}{\sigma_{\perp}^{ \pm}+\sigma_{\| I}^{ \pm}} \tag{4}
\end{equation*}
$$

where for example $\sigma_{1}^{+}$is the differential cross section for reaction (1a) with the electric vector of the photon perpendicular to the production plane. To eliminate the $\omega \rho$ interference terms the two cross sections are again summed together in an appropriate way giving the asymmetry

$$
\begin{equation*}
\mathrm{A}\left(\pi^{+}+\pi^{-}\right) \equiv \frac{\left(\sigma_{\perp}^{+}+\sigma_{1}^{-}\right)-\left(\sigma_{11}^{+}+\sigma_{11}^{-}\right)}{\left(\sigma_{\perp}^{+}+\sigma_{1}^{-}\right)+\left(\sigma_{11}^{+}+\sigma_{\|}^{-}\right)}=\frac{\mathrm{A}^{+}+\mathrm{RA}^{-}}{1+\mathrm{R}} \tag{5}
\end{equation*}
$$

where

$$
\begin{equation*}
\mathbf{R}=\sigma^{-} / \sigma^{+} . \tag{6}
\end{equation*}
$$

Values for $\mathrm{A}\left(\pi^{+}+\pi^{-}\right)$have been calculated from the experimental data ${ }^{2-5}$ at $\mathrm{k}=3.4 \mathrm{GeV}$ and are shown as crosses in the figure. Note that $\mathrm{R}, \mathrm{A}^{+}$and $\mathrm{A}^{-}$are themselves ratios and there should be no normalization problems in combining the results of several experiments.

The quantity $\mathrm{A}\left(\pi^{+}+\pi^{-}\right)$can be compared with data from reaction (2) using the $\rho$ dominance model; as before, the few per cent $\omega$ term is neglected. While the probability for the $\rho^{\circ}$ in reaction (2) to have helicity 1 is just $\rho_{11}^{\text {hel }}$, the probability for it to have linear polarization normal to its direction of motion and parallel or perpendicular to the reaction plane is given by ${ }^{6}$

$$
\begin{align*}
& P_{1}=\rho_{11}^{\mathrm{hel}}+\rho_{1-1}^{\mathrm{hel}},  \tag{7}\\
& \mathrm{P}_{\mathrm{II}}=\rho_{11}^{\mathrm{hel}}-\rho_{1-1}^{\mathrm{hel}} .
\end{align*}
$$

The $\rho$ dominance model then predicts

$$
\begin{align*}
& \sigma_{1}^{+}+\overline{\sigma_{1}}=\frac{\pi \alpha}{\gamma_{\rho}^{2}}\left(\begin{array}{c}
\text { hel }
\end{array} \rho_{11}^{\mathrm{hel}}+\rho_{1-1}\right) \frac{\mathrm{d} \sigma}{\mathrm{dt}}  \tag{8}\\
& \sigma_{11}^{+}+\overline{\sigma_{11}}=\frac{\pi \alpha}{\gamma_{\rho}^{2}}\left(\rho_{11}^{\mathrm{hel}}-\rho_{1-1}^{\mathrm{hel}}\right) \frac{\mathrm{d} \sigma}{\mathrm{dt}}
\end{align*}
$$

from which it follows

$$
\begin{equation*}
\mathrm{A}\left(\pi^{+}+\pi^{-}\right)=\frac{\rho_{1-1}^{\mathrm{hel}}}{\rho_{11}^{\mathrm{hel}}} \tag{9}
\end{equation*}
$$

The ratio of the density matrix elements has been plotted with solid points in the figure. The density matrices were obtained by direct fits in the helicity frame
to the angular distributions of the $4 \mathrm{GeV} / \mathrm{c}$ data of the Notre Dame-Purdue-SLAC Collaboration. ${ }^{7}$ As can be seen from the figure, there is a large discrepancy, ${ }^{8}$ especially at $|t|=0.2$ and $0.4(\mathrm{GeV} / \mathrm{c})^{2}$ where $\rho_{1-1}^{\text {hel }} / \rho_{11}^{\text {he }}$ is negative while $\mathrm{A}\left(\pi^{+}+\pi^{-}\right)$is near 0.5 . This discrepancy has been pointed out by the DESY group. ${ }^{3}$ Unfortunately, there is a non- $\rho$ background of perhaps $20 \%$ in the $\rho$ region (taken as $700 \leq \mathrm{M}_{\pi \pi} \leq 850 \mathrm{MeV}$ in the fits). To test whether the discrepancy could be due to interference of the $\rho$ with the background, we have studied the density matrix as a function of the effective $\pi \pi$ mass from 500 to 1050 MeV . One might expect that any appreciable distortion of the density matrix by backgrounds would also lead to a rapid variation of the matrix elements as $M_{\pi \pi}$ passes through $M_{\rho}$. In each $t$ interval the ratio $\rho_{1-1}^{\text {hel }} / \rho_{11}^{\text {hel }}$ shows no trend with $M_{\pi \pi}$, however, all mass intervals giving a ratio consistent with that obtained for the $\rho$ region. Values of $\rho_{1-1}^{\text {hel }} / \rho_{11}^{\text {hel }}$ obtained from fits to events with $M_{\pi \pi}$ between 575 and 675 MeV plus those between 875 and 975 MeV are shown with open circles in the figure; in this background region there are roughly equal numbers of $\rho$ and background events. The lack of variation with $M_{\pi \pi}$ would seem to climinate backgroundinterference effects as a possible source of the vector dominance discrepancy.

At high energies $\sigma_{\perp}$ and $\sigma_{\|}$correspond to $P(-1)^{J}=+1$ and -1 (natural and unnatural parity) exchange, respectively. ${ }^{9}$ The data thus indicate that while photoproduction procceds principally via natural-parity exchange ( $\sim 75 \%$ of the time), the raction $\pi^{-} p \rightarrow \rho^{\circ} n$ has roughly equal contributions from the two spin-parity sequences, depending somewhat on $t$.

Taking the discrepancy at face value, the relative strength of the $\rho$-dominance terms in $\pi^{ \pm}$photoproduction can be estimated by assuming the photoproduction cross section to bc the sum of two contributions, the first being the $\rho$-dominance contribution, with the ratio of natural to unnatural parity exchanges determined by
the model as

$$
\begin{equation*}
\frac{\left(\sigma_{1}^{+}+\sigma_{1}^{-}\right)_{\rho}}{\left(\sigma_{11}^{+}+\sigma_{11}^{-}\right)_{\rho}}=\frac{1+\rho_{1-1} / \rho_{11}}{1-\rho_{1-1} / \rho_{11}} \tag{10}
\end{equation*}
$$

and the second contribution being an ad hoc term contributing only to $\sigma_{\perp}^{+}+\sigma_{\perp}^{-}$. These assumptions give an upper limit to the fraction of pions photoproduced via the $\rho$-dominance terms:

$$
\begin{equation*}
\mathrm{F}=\frac{1-\mathrm{A}\left(\pi^{+}+\pi^{-}\right)}{1-\rho_{1-1} / \rho_{11}} \tag{11}
\end{equation*}
$$

Using the data shown in the figure, this upper limit is $(30 \pm 6) \%,(39 \pm 7) \%$ and $(60 \pm 21) \%$ at $-\mathrm{t}=0.2,0.4$, and $0.6 \mathrm{GeV}^{2}$ respectively. Thus, the present data imply that the original success of the $\rho$-dominance model in fitting the photoproduction differential cross section ${ }^{1}$ must be regarded as fortuitous and that the unknown second term must have a $t$-dependence similar to the $\rho$-dominance term; furthermore, the value of $\gamma_{\rho}^{-2}$ must be roughly $1 / 2$ (or less) that used previously. ${ }^{10}$

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6. Notre Dame-Purdue-SLAC Collaboration, "Compilation of $\pi^{-} \mathrm{p}$ Data at $4 \mathrm{GeV} / \mathrm{c}$ ", to be published in the Physical Review.
7. In this momentum transfer region $R$ is considerably less than one and it is the $\pi^{+}$asymmetry which dominates $\mathrm{A}\left(\pi^{+}+\pi^{-}\right)$; even if one were to ignore the experimental values and arbitrarily set $\mathrm{A}^{-}$equal to the limit -1 , there would still be a large discrepancy at $|\mathrm{t}|=0.2$ and $0.4(\mathrm{GeV} / \mathrm{c})^{2}$ where $\mathrm{A}\left(\pi^{+}+\pi^{-}\right)$ calculated in this way would be $0.19 \pm 0.08$ and $0.23 \pm 0.09$, respectively.
8. P. Stichel, Zeit. für Physik 180, 170 (1964);
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9. This would imply $\gamma_{\rho}^{2} / 4 \pi \geq 1$, in agreement with the preliminary values reported to the Vienna Conference by the Cornell and SLAC $\rho$-photoproduction groups.

## FIGURE CAPTION

Figure 1 - Comparison of experimental values of $\mathrm{A}\left(\pi^{+}+\pi^{-}\right)$, the photoproduction asymmetry, at 3.4 GeV (crosses), with the ratio of helicity-frame density matrix elements $\rho_{1-1} / \rho_{11}$ for the process $\pi^{-} \rho \rightarrow \pi^{+} \pi^{-}$n at $4 \mathrm{GeV} / \mathrm{c}$. The closed circles are the ratio of density matrix elements calculated from events in the $\rho$ region ( $\mathrm{M}_{\pi \pi}$ between 700 and 850 MeV ) while the open circles were calculated with events in the background regions, 575 to 675 MeV and 875 to 975 MeV .


Fig. 1


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