Accelerator Physics Issues at LHC & Beyond Part II

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(A) Past and Future

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(B) The Large Hadron Collider (LHC)

- parameters, magnets, beam-beam effects, preinjectors, ion collisions,...
- A new phenomenon: Electron Cloud (today)

(C) LHC Upgrades, VLHC-I and II (today)

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Electron Cloud

Observed with the LHC beam:

1999 SPS, 2000 PS, 2000 PS-SPS transfer line

primary e- generated by photoemission or gas ionization; their number amplifies along a bunch train due to beam-induced multipacting

- (1) Build Up, Saturation, Decay
- (2) Wake Fields and Instabilities
- (3) Heat Load

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(4) LHC Approach



Intensity of 72-bunch LHC beam in SPS vs. time. Batch intensity (top) and bunch intensity for the first 4 bunches and last 4 bunches (where losses are visible after about 5 ms) of the batch (bottom). (Courtesy G. Arduini, 2001).

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Build Up, Saturation, Decay

 e^- production mechanisms:

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- residual gas ionization; typical rate $d^2 \lambda_e / (ds \ dt) \approx 5 \times 10^{11} e^- m^{-1} s^{-1}$
- synchrotron radiation and photo-emission; typical rate $d^2 \lambda_e / (ds \ dt) \approx 5 \times 10^{18} \text{ e}^{-} \text{ m}^{-1} \text{s}^{-1}$
- secondary emission: (1) true secondaries &
 (2) elastically reflected or rediffused; →
 exponential growth

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ENERGY DISTRIBUTION OF SECONDARY ELECTRON EMITTED BY COPPER



Normalized secondary electron energy distribution for conditioned copper, revealing three components: true secondaries $(E \ll E_p)$, elastically scattered $(E \approx E_p)$ and rediffused (in between). [N. Hilleret, 2001]

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Secondary emission yield for perpendicular incidence vs. primary electron energy with and w/o elastically scattered electrons. Parametrization based on measurements for LHC prototype chamber. [Ian Collins, 2000]

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Schematic of electron-cloud build up in the LHC beam pipe. [Courtesy Francesco Ruggiero]

Proper multipacting: $n_{\min} \equiv \frac{h_y^2}{N_b r_e L_{\text{sep}}} = 1$

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1		IZDIZD	DC	apa	TTA	DOD	
accelerator	PEP-II	KEKB	PS	SPS	LHC	PSR	SN
species	e^+	e^+	р	р	р	р	р
population N_b [10 ¹⁰]	10	3.3	10	10	10	5000	1000
spacing $L_{\rm sep}$ [m]	2.5	2.4	7.5	7.5	7.5	(108)	(248)
bunch length σ_z [m]	0.013	0.004	0.3	0.3	0.077	25	30
h. beam size σ_x [mm]	1.4	0.42	2.4	3	0.3	25	0.6
v. beam size σ_y [mm]	0.2	0.06	1.3	2.3	0.3	7.5	0.6
ch. $\frac{1}{2}$ size h_x [mm]	25	47	70	70	22	50	100
ch. $\frac{1}{2}$ size h_y [mm]	25	47	35	22.5	18	50	100
circumf. C [km]	2.2	3.0	0.63	6.9	27	0.09	0.22
beta function β	18	15	15	40	80	5	6
parameter n_{\min}	1	10	0.58	0.24	0.15	0.0002	0.000

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Simulation of Cloud Build Up (Schematic)



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indicators of e^- build up

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(1) nonlinear pressure rise → ρ_e
 (2) pick ups or dedicated e⁻ monitors → ρ_e
 (3) tune shift along the train → ρ_e
 (4) beam-size blow up along the train
 (5) luminosity drop

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example: magnitude of e^- cloud in the SPS (1) from pressure rise [O. Gröbner] : pressure balance reads $S_{\text{eff}}P/(k_BT) = Q$, where S_{eff} pumping speed in volume per meter per second, $Q = \alpha d\dot{\lambda}_e/ds$ total flux of molecules per unit length (α : desorption yield per electron) and $P = k_B T N/V$.

$$\frac{d\lambda_e}{ds} = \frac{T_{rev}}{\alpha k_B T} S_{\text{eff}} P$$

With P = 100 nTorr, $\alpha \approx 0.1$ and $S_{\text{eff}} \approx 20 \text{ l s}^{-1} \text{ m}^{-1}$:

$$\frac{d\lambda_e}{ds} \approx 10^{10} \frac{\text{electrons}}{\text{bunch} - \text{train meter}}$$

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(2) from damper pick-up measurements[W. Hoefle]:

a few 10^8 electrons per bunch passage are deposited on the pick-up; this amounts to $10^9 - 10^{10}$ per train, or, with an effective pick-up length of about 10 cm,

$$\frac{d\lambda_e}{ds} \approx 10^{10} \frac{\text{electrons}}{\text{bunch} - \text{train meter}}$$

The two estimates are consistent.

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Sum and difference signal on damper pick-up during the passage of an LHC batch in the SPS $(1\mu s/div)$. (Courtesy W. Hofle, 2001).

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Simulated electron-cloud build up for an SPS dipole chamber, with and without elastic electron reflection. Saturation at $\lambda_{e,\text{sat}} \sim N_b/L_{\text{sep}} \approx 1.3 \times 10^{10} \text{ m}^{-1} \rightarrow \text{`neutralization' density}$ $\rho_{\text{sat}} \approx N_b/(\pi h_x h_y L_{\text{sep}}) \approx 3 \times 10^{12} \text{ m}^{-3}.$

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Fourier spectrum of the vertical oscillations of the LHC-beam bunch centroids as a function of bunch number

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Wake Fields and Instability due to Electron Cloud

• Multi-Bunch Instability

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- Coherent and Incoherent Tune Shift, etc.
- Single-Bunch Instability
 strong head-tail (TMCI), regular head-tail,
 transverse & longitudinal wakes,
 potential-well distortion

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Snapshots of the horizontal and vertical electron phase space (top) and their projections onto the position axes (bottom). (Courtesy G. Rumolo, 2001).

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adiabatic trapping (B. Richter, SLAC, March 2000)

WKB approximation \rightarrow adiabaticity condition

$$A \equiv \sigma_z \omega_{e,y} \sqrt{8e} / c \gg 1$$

where e = 2.718...

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 $A \approx 10$ for KEKB, PEP-II, PS, SPS, LHC!

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Single-Bunch Instability



Simulated bunch shape after 0, 250 and 500 turns (centroid and rms beam size shown) in the CERN SPS with an e⁻ cloud density of $\rho_e = 10^{12}$ m⁻³, without (left) and with (right) proton space charge (Courtesy G. Rumolo).

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Beam size evolution for an SPS bunch interacting with an electron cloud as predicted by different simulation approaches. (Courtesy G Rimolo, 2001).

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Wake force W_1 induced by an electron cloud; each line represents a different cloud size. Left: KEKB; right: SPS. [K. Ohmi et al., HEACC'01]. (Courtesy K. Ohmi, 2001).

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Wake force in V/m/C computed by displacing slice 1 and 40 (out of 100) of a Gaussian bunch (Courtesy G. Rumolo, 2001).

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TMCI calculation: betatron side band frequencies $(\omega - \omega_{\beta})/\omega_s$ vs. $cR_s/Q \propto \rho_e$ for KEKB LER. [K. Ohmi et al., HEACC'01]. (Courtesy K. Ohmi, 2001).

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estimated TMCI thresholds

accelerator	PEP-II	KEKB	PS	SPS	LHC	PSR	SNS
e^- osc./bunch	0.8	1.0	1	0.75	3	34	970
$n_{ m osc} \equiv \omega_e \sigma_z / (\pi c)$							
TMCI threshold	1	0.5	5	0.25	3	(0.6)	(0.5)
$ ho_e \ [10^{12} \ { m m}^{-3}]$							
density ratio	19	4	0.35	11	4	(92)	(27)
$ ho_{e,\mathrm{sat}}/ ho_{e,\mathrm{thresh}}$							

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Detail of the sum (top) and delta (bottom) signals at the SPS provided by the wide-band transverse pick-up in the vertical plane. Head-tail motion inside the bunches is visible. (Courtesy G. Arduini, 2001). Wake period determined from measured head-tail motion: $\lambda_{e^-,\text{wake}} \approx \sigma_z$! (K. Cornelis).

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y_{centroid} (mm)



Simulated centroid motion and vertical beam size with zero and positive chromaticity in the SPS ($\xi_y = 0.2$). Machine broadband impedance is also included. (Courtesy G. Rumolo, 2001).

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Electron-Cloud Heat Load



Energy distribution of e^{-s} incident on LHC chamber wall for a chamber radius r = 158 mm (left) and 29 mm (right) (G. Rumolo).

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Snapshot of transverse e⁻ distribution in an LHC dipole chamber (F.Z., 1997). Parameters: $\delta_{\text{max}} = 1.3$, $\epsilon_{\text{max}} = 450$ eV, R = 0.1, and $Y^* = 0.025$.

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Average arc heat load and cooling capacity as a function of bunch population N_b , for various δ_{max} . Other parameters are $\epsilon_{\text{max}} = 240$ eV, R = 5%, Y = 5%, and elastic electron reflection is included.

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LHC Recipe

- in arc dipoles: use sawtooth chamber to reduce photon reflections
- coat all warm sections with getter material TiZr (low secondary emission yield)
- rely on surface scrubbing during the commissioning to reduce the maximum secondary emission yield to a value of 1.1

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00101/05-490 Zahnhöhe 33.6 dump HGBN-TO Zahnhöhe 38.6 µm The second Material ALCON ELECTRA Algemenes. CERN - Profilerungsversuch Stahl - Cu Band Reflectionswinkel 100. Albinitial Inkel 99.07 500 200 µm 200 µm 906P0471 200 Probe vom 17.06.98 Links 200 Probe vom 17.05,99 Links 900P0400 00 Probe vom 17.08.99 Links Officered State P00361 Metallographie 12:06.1999

Sawtooth chamber protoype; the sawtooth reduces the photon reflectivity R to 1.3% [co-laminated Cu: $R \approx 80\%$]. (Ian Collins).

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Comparison of dose dependence of the Secondary Emission Yield as maeasured at CERN and SLAC (N. Hilleret et al., 2001).

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Newly Installed SPS e^- Cloud Detectors

- Pick-ups for e⁻ characteristics
 - e⁻ cloud build up, e⁻ energy distribution, triggering on the batch
- Behavior of e⁻ in a dipole magnetic field
 - 'strip detector'
 - 'triangle detector'

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- Scrubbing effect by in-situ measurement of secondary emission yield
- Ion detectors to exclude ion-stimulated desorption
- WAM_PAC Cu calorimeter to directly measure heat load from e^{-1} cloud

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Cu Calorimeter



SSWG 17/7/2001

WAMPAC in BA4



SSWG 17/7/2001




Triangle detector



SPS Machine as a vacuum test bench for the electron cloud studies with LHC type beams

SSWG prepared by J.M. JIMENEZ CERN Division LHC - Vacuum Group



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Strip detector



SSWG prepared by J.M. JIMENEZ CERN Division LHC - Vacuum Group



Stripe signal for 8.5 mm orbit bump with 2 s duration. The stripe clearly follows the beam. (G. Arduini, et al., 2001).

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Electron-stripe intensity in a 2-T dipole field vs. bunch population. Threshold at 2×10^{10} protons per bunch.

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e⁻ cloud behaviour in a dipole magnetic field - Magnetic field passing through 0 Gauss

e⁻ stripe vs. magnetic field; signal disappears for $|B| \leq 15$ G.

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At higher intensity $N_b \approx 6 \times 10^{10}$ two stripes are observed! Spacing consistent with simulation.

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(C) Beyond LHC: LHC-II and VLHC

- higher luminosity and/or energy
- more bunches?
- crossing angle & crabbing
- magnets (stronger and/or cheaper)
- synchrotron radiation
- emittance control
- \bullet collective effects & electron cloud
- IP debris, quench limits, & safe beam abort
- quasi-continuous beams?

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Luminosity for Bunched Beams

$$L = \frac{N_b^2 f_{\rm rep} H_D}{4\pi \sigma_x \sigma_y} \ \eta_L$$

For a horizontal crossing, η_L is

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$$\eta_L = \frac{2}{\sigma_z \sqrt{\pi}} \int_0^\infty \frac{\exp\left(-\left(\frac{z}{\sigma_z}\right)^2 \left\{1 + \frac{\theta_c^2}{4\theta_d^2} \left[\frac{1}{1 + (z/\beta_x^*)^2}\right]^2\right\}\right)}{\sqrt{(1 + (z/\beta_x^*)^2) (1 + (z/\beta_y^*)^2)}}$$

where $\theta_d = \sigma_x / \sigma_z$, θ_c the full crossing angle.

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Beam-Beam Tune Shifts for Bunched Beams

Assuming $\beta_x^* \approx \beta_y^*$ and $\epsilon_x \approx \epsilon_y$ the beam-beam tune shifts for a particle at the center of the bunch are

$$\begin{split} \Delta Q_x &= -\frac{N_b r_p}{2\pi\gamma} \frac{1}{\sqrt{2\pi\sigma_z}} \int_{\infty}^{\infty} \left(\beta^* + \frac{s^2}{\beta^*}\right) \left[\left(\frac{1}{(\beta^* + s^2/\beta^*)\epsilon} + \frac{1}{\theta_c^2 s^2}\right) \right. \\ &\left. \exp\left(-\frac{\theta_c^2 s^2}{2\left(\beta^* + s^2/\beta^*\right)\epsilon}\right) - \frac{1}{\theta_c^2 s^2} \right] \exp\left(-\frac{s^2}{2\sigma_z^2}\right) \, ds \\ \Delta Q_y &= -\frac{N_b r_p}{2\pi\gamma} \frac{1}{\sqrt{2\pi\sigma_z}} \int_{\infty}^{\infty} \left(\beta^* + \frac{s^2}{\beta^*}\right) \left[\frac{1}{\theta_c^2 s^2} \left(1 - \frac{\theta_c^2 s^2}{2(\beta^* + s^2/\beta^*)\epsilon}\right)\right) \right] \exp\left(-\frac{s^2}{2\sigma_z^2}\right) \, ds \,, \end{split}$$

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Luminosity (left) and total beam-beam tune shift (right) vs. crossing angle; parameters: $N_b = 1.7 \times 10^{11}$, $\beta^* = 0.25$ m, $\sigma_z = 7.7$ cm, $n_b = 2800$, $\gamma \epsilon_{\perp} = 3.75$ µm.

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Applying a deflection of opposite sign to the head and tail of each bunch, luminosity loiss due to the crossing angle is avoided.

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Crab Cavities cont'd

Distance between last quadrupole and IP about 20 m. Outer quadrupole radius 25 cm. Two separate final quadrupoles require $\theta_c \geq 25$ mrad. Transverse crab deflecting voltage:

$$V_{\perp} = \frac{cE \tan \theta_x/2}{e\omega_{rf} \sqrt{\beta_x^* \beta_{\rm crab}}}$$

variable	symbol	KEKB HER	LHC
beam energy	E	$8.0 \mathrm{GeV}$	$7 { m TeV}$
RF frequency	$f_{ m rf}$	$508.9 \mathrm{~MHz}$	$1.3~\mathrm{GHz}$
half crossing angle	$ heta_c/2$	11 mrad	$12.5 \mathrm{mrad}$
IP beta function	eta_x^*	$0.33 \mathrm{\ m}$	$0.25 \mathrm{~m}$
cavity beta function	eta_x	100 m	2000 m
required kick voltage	V_{\perp}	$1.44 \mathrm{MV}$	$144 \mathrm{MV}$

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phase diagram ofNb₃Sn





 $\epsilon = 0$ after cool-down to 4.2 K

Stronger Magnets? Nb_3Sn instead of NbTi

year	group	type	field/gradient
1982	CERN	quad	$71 \mathrm{T/m}$
1983	CERN/Saclay	dipole	$5.3~\mathrm{T}$
1985	LBL	dipole D10	8 T
1986	KEK	dipole	$4.5 \mathrm{~T}$
1988	BNL	dipole	$7.6 \mathrm{~T}$
1991	CERN-ELIN	dipole	$9.5~\mathrm{T}$
1995	LBNL	hybrid dipole D19H	$8.5~\mathrm{T}$
1995	UT-CERN	dipole MSUT	11.2 T
1996	LBNL	dipole D20	13.3 T
2001	LBNL	common coil dipole	14.4 T

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LHC-II Parameters

parameter	LHC	LHC-II
beam energy E [TeV]	7	14
dipole field B [T]	8.39	16.8
total energy/beam $[MJ]$	334	1130
number of bunches n_b	2800	5600
bunch population N_b [10 ¹¹]	1.05	1.05
rms IP beam size $\sigma_{x,y}^*$ [µm]	15.9	7.4
rms IP div. $\sigma^*_{x',y'}$ [µrad]	31.7	34
IP beta $\beta_{x,y}^*$ [m]	0.5	0.22
beam-beam tune shift / IP $\xi_{x,y}$	0.0034	0.005
crossing angle $\theta_c \ [\mu rad]$	300	300
rms bunch length σ_z [cm]	7.7	4.0

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LHC-II Parameters (cont'd)

bunch spacing $L_{\rm sep}$ [m]	7.48	3.74
SR power $P_{\rm SR}$ [kW]	3.6	114
SR dipole heat load dP/ds [W/m]	0.2	6.6
rms transv. emittance $\gamma \epsilon_{x,y}$ [µm]	3.75	$3.75 \rightarrow 1.0$
eq. horiz. emittance $\gamma \epsilon_x^{eq}$ [µm]	2.03^{\star}	1.07^{\star}
longit. emittance ϵ_L (σ) [eVs]	0.2	0.15^{\star}
damp. time $\tau_{x,\text{SR}}$ [hr]	52	6.5
IBS growth time $\tau_{x,\text{IBS}}$ [hr]	142	345 (in.)
events per crossing	18	90
peak luminosity $L [10^{34} \text{ cm}^{-2} \text{s}^{-1}]$	1.0	10.
lum. lifetime τ [hr]	10	3.2

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Emittance Evolution

synchrotron radiation amplitude damping time

$$\tau_z J_z = \left(\frac{3(m_p c^2)^3}{e^2 c^3 r_p Z^2}\right) \frac{1}{B^2 E} \left(\frac{C}{2\pi\rho}\right) \approx \frac{16644 \text{hr}}{E[\text{TeV}]B[\text{T}]^2} \left(\frac{C}{2\pi\rho}\right) \frac{A^4}{Z^4}$$

damping decrement (for 2 IPs)

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$$\delta = \frac{T_0}{n_{\rm IP}\tau_{x,y}} \approx 5.7 \times 10^{-13} \ E[{\rm TeV}]^2 B[{\rm T}] \frac{Z^3}{A^4}$$

does this affect the maximum beam-beam tune shift? maximum ξ : measurements and simulations fitted by $\xi_{\rm max} \propto 0.009 + 0.021 \ (\delta/10^{-4})^{0.5}$

[E. Keil & R. Talman, 1983; S. Peggs, 1999; R. Assmann et al., 2000]

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Tune shift parameter vs. damping decrement. [LEP data courtesy of R. Assmann; not beam-beam limited]

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more important consequence of synchrotron radiation: shrinkage of emittance during the store situation different from e^- storage rings; $\tau_{SR} \sim$ hours SR equilibrium emittance:

$$\epsilon_{x,N}^{\mathrm{SR}} \approx \frac{55}{32\sqrt{3}} \frac{\dot{\lambda}_A}{J_x} \left(\frac{\gamma^3}{Q_\beta^3}\right) \left(\frac{C}{2\pi\rho}\right)^3$$

for LHC-II and HF VLHC 2–3 orders of magnitude below desired design emittance!

 \rightarrow large beam-beam tune shifts, halo, background,...? (J. Gareyte)

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equilibrium emittance determined by balance of radiation damping and **intrabeam scattering**

$$\frac{1}{\tau_{\mathbf{x},\text{IBS}}} \approx \frac{\mathbf{cr_p^2 N_b L_c}}{16 \mathbf{Q} \epsilon_{\mathbf{x},\mathbf{N}}^2 \sqrt{\kappa} \sqrt{\kappa + 1} \gamma \sigma_{\mathbf{z}} \sigma_{\delta}} \quad [\text{J. Wei}]$$

where $L_c \approx 20$. Asymptotically, for $\gamma \gg Q_\beta$: $1/\tau_{\delta,\text{IBS}} \approx 1/\tau_{x,\text{IBS}}$ and $\sigma_{\delta} \approx Q_{\beta}^{3/2} \sqrt{\epsilon_x/\rho}$ IBS equilibrium emittance:

$$\epsilon_{\mathbf{x},\mathbf{N}}^{\mathrm{IBS}} = \frac{\rho^{5/6} \mathbf{N}_{\mathbf{b}}^{1/3}}{\mathbf{Q}_{\beta} \gamma^{7/6}} \left(\frac{\mathbf{Z} \mathbf{f}_{\mathrm{rf}} \mathbf{e} \mathbf{V}_{\mathrm{rf}}}{\mathbf{c} \mathbf{E} \kappa (\kappa + 1)} \right)^{1/6} \left(\frac{\mathbf{C}}{2\pi \rho} \right)^{1/6} \left(\frac{\mathbf{3} \mathbf{r}_{\mathbf{p}} \mathbf{L}_{\mathbf{c}}}{\mathbf{16}} \right)^{1/3}$$

$$f_{\mathrm{rf}}: \text{ rf frequency; } V_{\mathrm{rf}}: \text{ total rf voltage}$$

 $\epsilon_y = \kappa \epsilon_x$ due to coupling and spurious D; assume $\kappa = 1$ for LHC-II

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Evolution of transverse emittance vs. time in LHC-II.

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Evolution of **beam current** during a store in LHC-II.

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Evolution of **beam-beam tune shift** vs. time in LHC-II.

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Evolution of **luminosity** during a store in LHC-II.

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Collective Effects

 loss of Landau damping for higher-order longitudinal modes (F. Ruggiero, J. Rogers):

$$\sigma_s \ge \frac{C}{2\pi} \left[\frac{\pi^3 N_b f_{\text{rev}} e}{6 h_{\text{rf}}^3 V_{\text{rf}}} \operatorname{Im} \left(\frac{Z_L}{n} \right)_{\text{eff}} \right]^{1/5}$$

- longitudinal microwave instability
- transverse coupled-bunch resistive-wall instability
- electron cloud

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simulation with SR damping, IBS, particle consumption



Evolution of **rms bunch length** during a store in LHC-II, and instability thresholds for $\text{Im}(Z_L/n)_{\text{eff}} \approx 0.1 \Omega$ (LHC).

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simulation with SR damping, IBS, particle consumption



Evolution of **rms bunch length** during a store in LHC-II, when after 3 hours **noise** maintains $\epsilon_L \ge 0.104$ eVs.

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Total Beam Current

synchrotron radiation power

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$$P_{\rm SR} = \frac{C_{\gamma} E^4 N_b n_b c}{C \rho} = U_0 f_{\rm rev} n_b N_b$$

using L and ξ this can be rewritten as

$$P_{\rm SR} = \left(\frac{8\pi r_p^{3/2}}{\sqrt{3cE_A}}\right) \frac{\kappa}{1+\kappa^2} \frac{E^{3/2}L\beta_x^*}{\xi\sqrt{J_z\tau_z}} \sqrt{\frac{C}{2\pi\rho}}$$

scaling:

 $B = \text{constant} \rightarrow J_z \tau_z \propto 1/E \text{ and } P_{\text{SR}} \propto E^2 L$ $B \propto E^{1/2} \rightarrow J_z \tau_z \propto 1/E^2 \text{ and } P_{\text{SR}} \propto E^{5/2} L.$

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Wall plug power density vs. SR load for different solutions: cold BS, warm BS/shield & photon stop (P. Bauer et al.).

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Sketch of the proposed VLHC-II photon stop (P. Bauer et al., PAC2001).

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Electron Cloud Heat Load for Shorter Bunch Spacing (LHC Luminosity Upgrade)



Average arc heat load as a function of bunch population for bunch spacings of 12.5 ns, 15 ns, and 25 ns, and a maximum secondary emission yield $\delta_{\text{max}} = 1.1$. Elastically reflected electrons are included.

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Average arc heat load as a function of bunch spacing, for $\delta_{\text{max}} = 1.1$ and various bunch populations.

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Fermilab-TM-2149 June 11, 2001

www.vlhc.org

Design Study for a Staged Very Large Hadron Collider



Snowmass



Stage 2 VLHC Tunnel




Transmission Line Magnet



- ✤ 2-in-1 warm iron
- Superferric: 2T bend field
- 100kA Transmission Line
- alternating gradient (no quadrupoles needed)
- ✤ 65m Length
- Self-contained including Cryogenic System and Electronics Cabling
- Warm Vacuum System

Snowmass

Operating Margin &rified





17 meter 100 kA test loop

Seven Designs Tested -0.8Kmargin at design current of 87.5 kA -25kA margin at nominal peak temperature of 6.0 K (similar margins for three variants used in Design Report) Very Large Hadron Collider

VLHC DESIGN STUDY SITE LAYOUT



'Continous Beams' or Super Bunches

- ISR was extremely successful
- continuous beams abandoned due to scarcity of antiprotons, no longer a problem
- no PACMAN bunches!
- no electron cloud!
- use induction acceleration modules, 25 kV/m, to generate long bunches bounded by barrier buckets (K. Takayama)
- stochastic cooling
- higher current
- route to high luminosity

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exciting new development

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Schematic of Super Bunches in a High-Luminosity Collider (K. Takayama et al.)

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Continuous Beams – Luminosity

$$L = \frac{c\lambda_1\lambda_2 l_{\text{det}}}{4\pi\sigma_0^2} \ K\left(\frac{l}{2\beta^*}, \frac{\beta^*\theta_c}{\sigma_0}\right)$$

where

$$K(\xi,\eta) = \frac{1}{\xi} \int_{-\xi}^{\xi} \frac{1}{1+u^2} \exp\left[-\frac{\eta^2}{4} \frac{u^2}{1+u^2}\right] du$$

The integral $K(\xi, \eta)$ is defined such that $K(\xi, \eta) \to 2$ for $\xi, \eta \to 0$ (E. Keil, et al., 1972/73).

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Continuous Beams – Tune Shift

For horizontal crossing, the beam-beam tune shifts are

$$\Delta Q_x = \frac{2\lambda r_p l}{4\pi\gamma\epsilon_{\perp}} I_x \left(\frac{l}{2\beta^*}, \frac{\beta^*\theta_c}{\sigma_0}\right)$$
$$\Delta Q_y = \frac{2\lambda r_p l}{4\pi\gamma\epsilon_{\perp}} I_y \left(\frac{l}{2\beta^*}, \frac{\beta^*\theta_c}{\sigma_0}\right)$$

where

$$I_{x}(\xi,\eta) = \frac{1}{\xi\eta^{2}} \int_{-\xi}^{+\xi} (1+u^{2}) \left[\left(u^{-2} + \frac{\eta^{2}}{1+u^{2}} \right) \exp\left(-\frac{\eta^{2}}{2} \frac{u^{2}}{1+u^{2}} \right) -u^{-2} \right] du$$

$$I_{y}(\xi,\eta) = \frac{1}{\xi\eta^{2}} \int_{-\xi}^{+\xi} (1+u^{-2}) \left[1 - \exp\left(-\frac{\eta^{2}}{2} \frac{u^{2}}{1+u^{2}} \right) \right] du$$

and the interaction happens between -l/2 and l/2. The integrals $I_{x,y}(\xi,\eta)$ are defined such that $I_{x,y}(\xi,\eta) \to 1$ for $\eta \to 0$ and all ξ .

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Luminosity (left) and beam-beam tune shifts (right) as a function of crossing angle, for a continuous beam with a line density $\lambda = 8.8 \times 10^{11} \text{ m}^{-1}$ (40 A current), $\beta^* = 0.25 \text{ m}$, $l_{\text{det}} = 1 \text{ m}$, l = 20 m, and $\gamma \epsilon_{\perp} = 3.75 \ \mu \text{m}$.

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Optimization of Continuous Beam Parameters – Length & Number & Charge of Super-Bunches? Ongoing Study at CERN

- maximum luminosity
- maximum beam-beam tune shift
- acceptable heat load
- timing constraints by (induction) rf system
- injectors and filling time
- beam abort system

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(D) Conclusions

- hadron colliders have performed exceedingly well in the past
- the LHC will break new territory:
 - highest energy (14 TeV) and highest luminosity (10^{35} cm⁻²s⁻¹) ever
 - long-range collisions
 - strong-strong collisions
 - electron cloud

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- radiation damping stronger than IBS

- beyond LHC: LHC upgrades and various stages of VLHC, Eloisatron, ...
 - higher fields or larger circumference (\rightarrow peculiar collective effects)
 - more synchrotron radiation; possibly more electron cloud
 - new exciting development:
 'quasi-continuous beams' (closing the circle to the ISR)

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Thanks

G. Arduini, R. Assmann, P. Bagley, P. Bauer, F. Bordry,
L. Bottura, D. Brandt, O. Brüning, I. Collins, K. Cornelis,
A. Faus-Golfe, W. Fischer, J. Gareyte, O. Gröbner,
H. Grote, G. Guignard, W. Herr, J.B. Jeanneret,
J.M. Jimenez, C. Johnstone, J. Jowett, E. Keil,
J.-P. Koutchouk, K.-H. Mess, K. Ohmi, S. Peggs, F. Pilat,
L. Rossi, F. Ruggiero, G. Rumolo, F. Schmidt,
R. Schmidt, E. Shaposhnikova, V. Shiltsev, M. Syphers,
T. Taylor, R. Thomas, A. Verdier, L. Vos, J. Wei,...

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Web addresses

- LHC http://lhc.web.cern.ch/lhc/
- LHC beam-beam effects http://wwwslap.cern.ch/collective/zwe/lhcbb/Welcome.html
- LHC electron cloud http://wwwslap.cern.ch/collective/electroncloud/electron-cloud.html
- Accelerator Physics Group of the CERN SL (SPC+LHC) Division http://wwwslap.cern.ch/
- VLHC http://vlhc.org/

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Extended Parameter Set for pp or $p\bar{p}$ Colliders x

acc.	$\mathrm{Sp} \bar{\mathrm{p}} \mathrm{S}$	TeV2a	LHC	LHC-II	VLHC-I	VLHC-II
$E [{\rm TeV}]$	0.32	0.98	7	14	20	87.5
B [T]	1.4	4.34	8.4	16.8	2	9.8
$\frac{\text{energy}}{\text{beam}}$ [MJ]	0.05	1	334	1320	3328	4200
$C \; [\mathrm{km}]$	6.9	6.28	26.7	26.7	233	233
n_b	6	36	2800	5600	40000	40000
$N_b \ [10^{11}]$	1.7~(p)	2.7~(p)	1.05	1.05	0.26	0.075
	$0.8~(ar{p})$	$\sim 1.0~(\bar{p})$				
$\hat{L} \left[\frac{1}{10^{34} \mathrm{cm}^{-2} \mathrm{s}^{-1}} \right]$	0.0006	~ 0.02	1.00	10.	1.0	2.0
$\sigma^*_{x,y}~[\mathrm{\mu m}]$	80, 40	32	15.9	7.4^{\star}	4.6	$3.4 \rightarrow 0.79$
$\sigma^*_{x',y'}$ [µrad]	136, 272	91	31.7	34^{\star}	15	$5 \rightarrow 1$

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acc.	$\mathrm{Sp}\bar{\mathrm{p}}\mathrm{S}$	TeV2a	LHC	LHC-II	VLHC-I	VLHC-
$\beta^*_{x,y}$ [m]	0.6,0.15	0.35	0.5	0.22	0.3	0.71
no. of IPs	3	2	2(4)	2(4)	2	2
bb ts./IP $\xi_{x,y}$	0.005	0.01	0.0034	$0.003 {\leftrightarrow} 0.005$	0.002	$\rightarrow 0.008$
$\theta_c \; [\mu \mathrm{rad}]$	0	0	300	300	153	10
$\sigma_z [{ m cm}]$	30	37	7.7	4.0	3.0	$\rightarrow 1.5$
$L_{\rm sep}$ [m]	1150	119	7.48	3.74	5.645	5.645
$P_{\rm SR}$ [kW]		$< 10^{-3}$	3.6	114	7	1095
dP/ds [W/m]		$\ll 10^{-3}$	0.2	6.6	0.03	4.7
$ au_{\mathrm{IBS}} \ [\mathrm{hr}]$	10	50(?)	142	345 (in.)	400	$4000 \rightarrow 1$
$ au_{y,\mathrm{SR}} \; [\mathrm{hr}]$		1200	52	6.5	200	2
d.d./IP $\delta~[10^{-10}]$		0.025	2.5	20	5	400
events/cross.		~ 6	18	90	21	54

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acc.	$\mathrm{Sp}\bar{\mathrm{p}}\mathrm{S}$	TeV2a	LHC	LHC-II	VLHC-I	VLHC-II
lum. lifet. τ_L [hr]	9	9	10	3.2	24	8
tune Q_{eta}	26	~ 20	63	63	220	220
$\gamma\epsilon_{x,y}~[\mu{ m m}]$	3.75	~ 3	3.75	$3.75 \rightarrow 1.0$	1.5	$1.6 \rightarrow 0.04$
$\gamma \epsilon_x^{eq} [\mu { m m}]$		$\sim 10^{\star}$	2.03	1.07	1.0	0.06
$\epsilon_L (\sigma) [eVs]$	0.11	0.11	0.2	$\rightarrow 0.15$	0.4	$0.4 \rightarrow 0.1$

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