Rare B Decays from BABAR

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- Introduction: what's "rare"?
- What can we learn?
- Electroweak penguin mediated decays massive propagators, $|V_{t(s,d)}|, \ldots$
- Hadronic two-body decays \rightarrow CKM angles (*CP* violating), matrix elements, ...
- Conclusions

Quark couplings in the standard model

Weak isospin doublet members (d', s', b') are mixed flavor states:

$$\begin{pmatrix} d'\\s'\\b' \end{pmatrix} = \begin{pmatrix} V_{ud} & V_{us} & V_{ub}\\V_{cd} & V_{cs} & V_{cb}\\V_{td} & V_{ts} & V_{tb} \end{pmatrix} \begin{pmatrix} d\\s\\b \end{pmatrix}$$

V is the (unitary) Cabibbo-Kobayashi-Maskawa (CKM) mixing matrix. Unitarity \Rightarrow 4 free parameters, including 1 phase:

$$V \simeq \begin{pmatrix} 1 - \frac{1}{2}\lambda^2 & \lambda & A\lambda^3(\rho - i\eta) \\ -\lambda & 1 - \frac{1}{2}\lambda^2 - iA^2\lambda^4\eta & A\lambda^2 \\ A\lambda^3(1 - \rho - i\eta) & -A\lambda^2 & 1 \end{pmatrix}$$
$$\lambda \simeq \sin \theta_c \simeq 0.22$$
$$A \sim 1$$

Unitarity Triangle

 $V^{\dagger}V = 1: 1^{\text{st}} \text{ column} \times (3^{\text{rd}} \text{ column})^* \rightarrow$ $V_{ud}V_{ub}^* + V_{td}V_{tb}^* + V_{cd}V_{cb}^* = 0 \qquad \frac{V_{ub}^*}{V_{cd}V_{cb}^*} + \frac{V_{td}}{V_{cd}V_{cb}^*} + 1 = 0$ $(V_{tb} \simeq V_{ud} \simeq 1)_{(\rho, \eta)}$ or $-[\rho + i\eta] - [1 - (\rho + i\eta)] + 1 = 0$ α $\rho + ir$ $1 - (\rho + i\eta)$ β γ (0, 0)1 (1, 0)

Major B decay modes are

What is "rare"

Here we mean modes with neither open nor hidden charm:

 $b \rightarrow u$ (CKM suppressed), $|V_{ub}/V_{cb}| \sim 0.06 - 0.10$ $b \rightarrow s, d$ (induced FCNC, penguin loops)



Why focus on RARE B decays?

- Heavy quark systems are relatively amenable to theoretical treatment $(m_b \gg \Lambda_{QCD}).$
- Need to flesh out S. M. with measurements (currently imprecise) of flavor mixing.
- CKM suppression of $b \rightarrow u$ tree diagrams ...





- Particles too massive to appear on-shell can contribute to loops
 - \diamond top quarks, $V_{t(s,d)}$ terms.
 - ♦ Higgs, SUSY, new structure.
- Interfering terms of comparable magnitudes expose phases, e.g., *CP* violating.
- Penguins complicate *CP* measurements; need to be understood for interpretation.







BABAR luminosity

- "Run 1", 1999-2000
- 20.7 fb⁻¹ on- $\Upsilon(4S)$
- 2.6 fb⁻¹ off- $\Upsilon(4S)$
- 22.7×10^6 produced *B* pairs

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Reconstruction of B decays at the $\Upsilon(4S)$

• $\Upsilon(4S) \rightarrow B^0 \bar{B}^0$, $B^+ B^- \sim \text{at rest} (p_B^* \simeq 325 \text{ MeV/c})$

•
$$m_B \simeq 5.3 \text{ GeV/c}^2 = m_{\text{recoil}}$$

• Energy, momentum conservation expressed as

$$\Delta E \equiv \sum E_i^* - E_{\text{beam}}^* = 0$$

$$M \equiv (m_{ES} \text{ or } m_{EC}) \simeq \sqrt{E_{\text{beam}}^{*2} - |\mathbf{p}_i^*|^2} = m_B$$

within resolution. (E_{beam}^* is much better measured than $\sum E_i^*$.) Typical resolution $M: 2.8 \text{ MeV}/c^2$, $\Delta E: 25 - 30 \text{ MeV}$

- For two-body
 - \diamond daughter $E^* \simeq 2.6 \text{ GeV}$
 - ◊ daughters nearly back-to-back

Background rejection

• $b \rightarrow c$ background: heavier daughter, lower energy recoil.

Not a problem for most modes.

- Continuum background (e.g. two x = .5 particles in opposite jets from $e^+e^- \rightarrow q\bar{q}$):
 - \diamond Control data with $E_{\rm Cm} < B\bar{B}$ threshold.
 - ♦ Distinguished by event shape.
 - * R_2 , ratio of 2nd to zeroth Fox-Wolfram moments (peaks higher for jetty continuum)
 - * B decay axis correlation with jet axis ...



Electroweak annihilation





$$B \to K^*\gamma, \ B \to (\rho, \omega)\gamma, \ b \to s\gamma, \ B \to K^{(*)}\ell^+\ell^-$$











(c)





$$B o K^* \gamma$$

- Include $K^{*+} \rightarrow K^0_{\scriptscriptstyle S} \pi^+, K^+ \pi^0$, $K^{*0} \rightarrow K^+ \pi^-, K^0_{\scriptscriptstyle S} \pi^0$
- Select events in $\diamond K^*$ mass peak $\diamond 2.30 < E_{\gamma}^* < 2.85 \text{ GeV}$ $\diamond \text{Track, photon, vertex quality}$ $\diamond \text{Positive kaon ID from DIRC}$
- Reduce background with
 - $\diamond \ |\!\cos\theta_{\rm T}| < 0.8$
 - $\diamond~{\rm For}~p_B^*~{\rm direction}$ wrt beam $\cos\theta_B^* < 0.75$
 - ♦ For K* decay helicity angle $\cos \theta_H^* < 0.75$ Both have (for different reasons) $\sin^2 \theta$ distribution, vs uniform for BG







Vigorous theoretical program to compute these.



$$\begin{aligned} \mathcal{B}(B^- \to K^{*-}\gamma) &= (3.92 \pm 0.62 \pm 0.21) \times 10^{-5} \\ \mathcal{B}(B^0 \to K^{*0}\gamma) &= (4.23 \pm 0.40 \pm 0.22) \times 10^{-5} \\ \text{theory estimate} & (7.1 \pm \sim 2.5) \times 10^{-5} \end{aligned}$$

In progress:
$$B
ightarrow (
ho, \omega) \gamma$$

From the ratios of $b \rightarrow d$ to $b \rightarrow s$ we determine the top quark CKM couplings:

$$\frac{\mathcal{B}(B^- \to \rho^- \gamma)}{\mathcal{B}(B^- \to K^{*-} \gamma)} = \frac{\mathcal{B}(B^- \to \rho^0 \gamma) + \mathcal{B}(B^- \to \omega^0 \gamma)}{\mathcal{B}(B^- \to K^{*-} \gamma)} = \left|\frac{V_{td}}{V_{ts}}\right| \xi \Omega$$

- $\xi = SU(3)$ -breaking factor
- $\Omega = \text{phase space factor}$ $\implies \text{bounds on } \left| \frac{V_{td}}{V_{ts}} \right|$



Signal is evident; results still to be quantified.

More on continuum rejection

- θ_T (introduced above)
- θ_B^* (introduced above)
- B decay angle wrt beams θ^*_{TB}
- Energy flow about the decay axis
 - \diamond Conical bins
 - ◊ Legendre moments
- Global event shape (R_2 , sphericity)

As these are correlated, combine in a Fisher discriminant, \mathcal{F} (or neural network). Typically include conical bins, with or without θ_B^* , θ_{TB}^* .



Search for
$$B \to K^{(*)} \ell^+ \ell^-$$

- Tight PID cuts on all tracks
- Continuum suppression with a Fisher $\mathcal{F}(R_2, \theta_B, \theta_{BT}, m_{K\ell})$



• Control samples $B \to K^{(*)}J/\psi, \ B \to K^{(*)}\psi(2S)$



Search for $B \to K^{(*)}\ell^+\ell^-$, cont.



Perform a maximum likelihood fit in these two variables

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Search for
$$B o K^{(*)} \ell^+ \ell^-$$
, cont.

The Standard Model predictions are based on light quark sum rules, quark models, and QCD sum rules. Find (90% CL upper limits)

$$\mathcal{B}(B \to K\ell^+\ell^-) < 0.8 \times 10^{-6} \quad (\text{theory } 0.3 \text{ to } 0.75 \times 10^{-6})$$

$$\mathcal{B}(B \to K^*\ell^+\ell^-) < 2.5 \times 10^{-6} \quad (1.0 \text{ to } 3.0 \times 10^{-6})$$

Limits are close to prediction; improved experimental sensitivity should soon make this decay obervable.

Two body hadronic B decays

Can make ~ 70 combinations among members of lowest-lying pseudoscalar, vector nonets



$K, \ \pi$ separation

For modes with a charged track recoiling against a light particle, lab momentum is uniformly distributed between about 1.6 and 4.4 GeV/c. DIRC gives K, π separation (measured with $D^0 \rightarrow K^+\pi^-$):



Maximum likelihood fit

Results (most analyses) come from maximum of unbinned extended likelihood:

$$\mathcal{L} = \frac{e^{-(\sum n_j)}}{N!} \prod_{i=1}^N \mathcal{L}_i$$

Allowing for final states KX and πX ,

$$\mathcal{L}_{i} = n_{\pi} \mathcal{P}_{\pi S}(\mathbf{x}_{i}) + n_{K} \mathcal{P}_{KS}(\mathbf{x}_{i})$$
 Signals

$$+ n_{C} \left[f_{KC} \mathcal{P}_{KC}(\mathbf{x}_{i}) + (1 - f_{KC}) \mathcal{P}_{\pi C}(\mathbf{x}_{i}) \right]$$
 Background

- Input observables for each event include $\mathbf{x}_i = \{M, \Delta E, \mathcal{F}, \text{ resonance masses, helicity angles } \mathcal{H}, \text{ PID pulls} \}.$
- PDF's $\mathcal{P}_{\pi S}$, etc., are products of functions of uncorrelated observables.
- PDF's determined with sideband data and MC.





$$B
ightarrow (K,\pi)\pi$$



 $B^+ \to \pi^+ \pi^0 \simeq \text{pure } T$ $B^+ \to \pi^+ K^0 \simeq \text{pure } P'$

Results for K, π modes

Mode	N_S	$S(\sigma)$	$B(10^{-6})$
$\pi^+\pi^-$	$41 \pm 10 \pm 7$	4.7	$4.1\pm1.0\pm0.7$
$K^+\pi^-$	$169\pm17\pm13$	15.8	$16.7 \pm 1.6 \pm 1.3$
K^+K^-	$8.2^{+7.8}_{-6.4} \pm 3.5$	1.3	< 2.5 (90% C.L.)
$\pi^+\pi^0$	$37 \pm 14 \pm 6$	3.4	< 9.6 (90% C.L.)
$K^+\pi^0$	$75\pm14\pm7$	8.0	$10.8^{+2.1}_{-1.9} \pm 1.0$
$K^0\pi^+$	$59^{+11}_{-10} \pm 6$	9.8	$18.2^{+3.3}_{-3.0} \pm 2.0$
$\overline{K}{}^{0}K^{+}$	$-4.1^{+4.5}_{-3.8} \pm 2.3$	—	< 2.4 (90% C.L.)
$K^0\pi^0$	$17.9^{+6.8}_{-5.8} \pm 1.9$	4.5	$8.2^{+3.1}_{-2.7} \pm 1.2$

Predictions for pure tree, penguin modes (BBNS, current inputs): $\mathcal{B}(B^+ \to \pi^+ \pi^0) = 5.3^{+0.8}_{-0.4} \pm 0.3$ $\mathcal{B}(B^+ \to \pi^+ K^0) = 14.1^{+6.4}_{-4.0}$



$$H_{\text{eff}} = -4\frac{G}{\sqrt{2}}V_{tb}V_{ts}^* \sum_{i=1}^{10} C_i(\mu)Q_i(\mu) + (s \to d)$$

Wilson operators Q_i , coefficients C_i , renormalization scale μ

Effective theory, cont.

$$H_{\text{eff}} = -4\frac{G}{\sqrt{2}}V_{tb}V_{ts}^* \sum_{i=1}^{10} C_i(\mu)Q_i(\mu) + (s \to d)$$

 $i = \text{odd}(\text{even}) \leftrightarrow \text{external(internal)}$ $i = 1, 2 : W^{\pm}(L)$ i = 3 - 6 : gluonic penguins(L, R)i = 7 - 10 : electroweak penguins(L, R)

- Renormalization group evolution $\mu = m_W \rightarrow m_B$ mixes the operators
- $\sum C_i(\mu)Q_i(\mu)$ scale independent, except for neglected terms in perturbation series
- Inclusive $\langle X \mid Q_i \mid B \rangle \simeq \langle s \mid Q_i \mid b \rangle$

Factorization

$$\langle m_1 m_2 \mid H_2 \mid B \rangle = \langle m_1 \mid J^{\mu} \mid B \rangle \langle m_2 \mid J_{\mu} \mid 0 \rangle$$

Uncertainties:

- Unmeasurable non-color singlet terms $\propto 1/N_c \rightarrow \xi$, a free parameter.
- $\langle m_1 \mid J^{\mu} \mid B \rangle$ depends on form factor(s) $F_{B \to m_1}$
- Average 4-momentum of virtual gluons and photons in loops
- quark masses

 $\mathcal{B}(B^0 \to K^+\pi^-)$ Ali, Kramer, Lü prediction (14 to 21) × 10⁻⁶ BABAR:

 $(16.7 \pm 1.6 \pm 1.3) \times 10^{-6}$



Bounds on $\gamma_{ m CKM}$

(Fleisher, Mannel; others): $B^{\pm} \rightarrow \overset{(-)}{K^0} \pi^{\pm}$ is pure penguin.

$$\mathcal{A}_{CP}^{\mathrm{dir}} \equiv \frac{\mathcal{B}\left(B^{0} \to K^{+}\pi^{-}\right) - \mathcal{B}\left(\bar{B}^{0} \to K^{-}\pi^{+}\right)}{\mathcal{B}\left(B^{0} \to K^{+}\pi^{-}\right) + \mathcal{B}\left(\bar{B}^{0} \to K^{-}\pi^{+}\right)}$$
$$R \equiv \frac{\mathcal{B}\left(\overset{(-)}{B^{0}} \to K^{\pm}\pi^{\mp}\right)}{\mathcal{B}\left(\bar{B}^{\pm} \to \overset{(-)}{K^{0}}\pi^{\pm}\right)}$$
$$r \equiv \frac{(\mathrm{tree})}{(\mathrm{penguin})}; \quad \delta \equiv \phi_{(\mathrm{tree})} - \phi_{(\mathrm{penguin})}$$

$$\mathcal{A}_{CP}^{\text{dir}} = 2\frac{r}{R}\sin\delta\sin\gamma$$
$$R = 1 - 2r\cos\delta\cos\gamma + r^2$$

 $\sin^2 \gamma \ < \ R = 0.65 \pm 0.38$

 K^{+}, K^{*+}

u

U

u

 η, η'

 W^+

 B^+

u

$$B
ightarrow \eta^{(\prime)} K^{(*)}$$

The tree is Cabibbo suppressed:



enhance $B \to \eta K^*$, $B \to \eta' K$

suppress $B \to \eta K$, $B \to \eta' K^*$



$$B o \eta K^*$$











Results for
$$B o \eta^{(\prime)} K^{(*)}$$

Mode	Signal yield	$S(\sigma)$	$B/10^{-6}$	(90% CL)
ηK^{*0}	20.5 ± 6.0	5.4	$19.8^{+6.5}_{-5.6} \pm 1.7$	
ηK^{*+}	14.3 ± 6.6	3.2	$22.1^{+11.1}_{-9.2} \pm 3.3$	(< 33.9)
$\eta'_{\eta\pi\pi}K^+$	$49.5_{-7.3}^{+8.1}$	15	63^{+10}_{-9}	
$\eta'_{\rho\gamma}K^+$	$87.6^{+13.4}_{-12.5}$	11	80^{+12}_{-11}	
$\eta' K^+$		17	$70\pm8\pm5$	
$\eta_{\eta\pi\pi}' K^0$	$6.3^{+3.3}_{-2.5}$	4.7	28^{+15}_{-11}	
$\eta'_{ ho\gamma}K^0$	$20.8_{-6.5}^{+7.4}$	4.2	61^{+22}_{-19}	
$\eta' K^0$		5.9	$42^{+13}_{-11}\pm 4$	
$\eta'_{\eta\pi\pi}\pi^+$	$5.7^{+3.8}_{-2.8}$	3.2	$7.1^{+4.8}_{-3.5}$	
$\eta'_{ ho\gamma}\pi^+$	$-0.9^{+7.8}_{-6.2}$	0.1	$-0.7\substack{+6.7 \\ -5.3}$	
$\eta'\pi^+$		2.8	$5.4^{+3.5}_{-2.6}\pm0.8$	(< 12)





Large rate for
$$B
ightarrow \eta' K$$

Has prompted conjectures

- η^\prime approximates a flavor singlet state
- QCD anomaly, glue coupling to η^\prime
- "Charming penguins" c enhanced in loop:





$$B o \phi K^{(*)}$$



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$$B
ightarrow \omega\pi$$
 , $B
ightarrow \omega K$









Ali, Kramer, Lü





Dotted line is CP symmetry axis

- ρ^\pm along left, bottom
- ρ^0 along diagonal
- $\rho^{\prime\pm}$, $\rho^{\prime0}$ along inner bands
- Spin alignment $\implies \cos^2 \theta_H$ peaking at low, high m
- Charged, neutral $(f^0(400 1200))$ or " σ " scalar in band centers
- Non-resonant $\pi^+\pi^-\pi^0$ in central region



$$B^0 o \pi^+ \pi^- \pi^0$$

Cuts:
$$\pi^0 \rightarrow \gamma\gamma$$
 mass,
 $\cos \theta_{\mathrm{T}}, \mathcal{F}, \mathcal{H}$
Signal in $B^0 \rightarrow \rho^+ \pi^-$
 $m_{ES}, \Delta E$:
 $p^{\circ} \pi^- 42.8 \pm 11.5 \pm 4.0$
 $\rho^- \pi^+ 46.2 \pm 11.4 \pm 3.8$
 $\rho^\pm \pi^\mp 5.01 28.9 \pm 5.4 \pm 4.3$
 $\rho^0 \pi^+ 6.1 \pm 5.8 \pm 2.8 0.96 3.6 \pm 3.5 \pm 1.7$
 $(< 10.6, 90\% \text{ CL})$

$$B^0
ightarrow a_0^+ \pi^-$$

Going beyond the lowest-lying meson nonets, $a_0(980)(\Gamma = 50 \text{ to } 100 \text{ MeV}/c^2) \rightarrow \eta\pi$ ("dominant"); $\eta \rightarrow \gamma\gamma$ Continuum rejection Neural Net

- R_2 , $\cos \theta_{\mathrm{T}}$, $\cos \theta_B^*$
- three more topological angles
- Legendre moments $L_j^{(c,n)} \equiv \sum_{i_{(c,n)}} p_i |\cos \theta_i|^j$, j = 0, 2, 6c: tracks; n: calorimeter clusters,
 - θ_i wrt B candidate thrust axis



$$B^0
ightarrow a_0^+ \pi^-$$
, cont.

Max. likelihood fit with NN, $m(\gamma\gamma)$, ΔE , $m_{\rm EC}$ (no $m(a_0)$)



$$\mathcal{B}(B^0 \to a_0^+ (a_0^+ \to \eta \pi^+) \pi^-) = (6.2^{+3.0}_{-2.5} \pm 1.1) \times 10^{-6}$$

with a statistical significance, including systematics, of 3.7 σ . At 90% CL, $\mathcal{B}(B^0 \to a_0^+(a_0^+ \to \eta \pi^+)\pi^-) < 11.5 \times 10^{-6}$



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Summary for η, η', ω modes







Near future focus of theory, expt effort

- Reliability of predictions for radiative decays
 - ♦ Constrain new physics in loops
- Reliability of "triangle" relations for CP angles
 - ♦ Control of penguin polution
 - ◊ "Pure" tree, penguin modes
- Charge asymmetries direct *CP*
- Time dependent asymmetries CKM angles
 - ♦ Control of penguin polution
- Role of QCD anomaly, charm enhancement, ...
 - ♦ Large rates for $B \to \eta' K$, inclusive $B \to \eta'$

Conclusions

- Experiments are approaching the $\mathcal{O}(10^{-6})$ decade of branching fractions in B decays.
- ~ 20 exclusive modes observed so far:
 - $\diamond \ K^* \gamma$
 - $\diamond \ K\pi$
 - $\diamond \pi\pi$
 - $\diamond \ K\eta'$
 - $\diamond \ K^*\eta$
 - $\diamond \ \pi^+ \rho^-$
 - $\diamond \ \pi^+ \omega$
 - $\diamond \ K^{(*)}\phi$
- Measurements have stimulated much theoretical interest and effort
- More data in hand; much more in the next few years.