

# Proton and Deuteron Spin Structure from SLAC Experiment E155

Lee Sorrell  
American University  
Representing the E155 Collaboration

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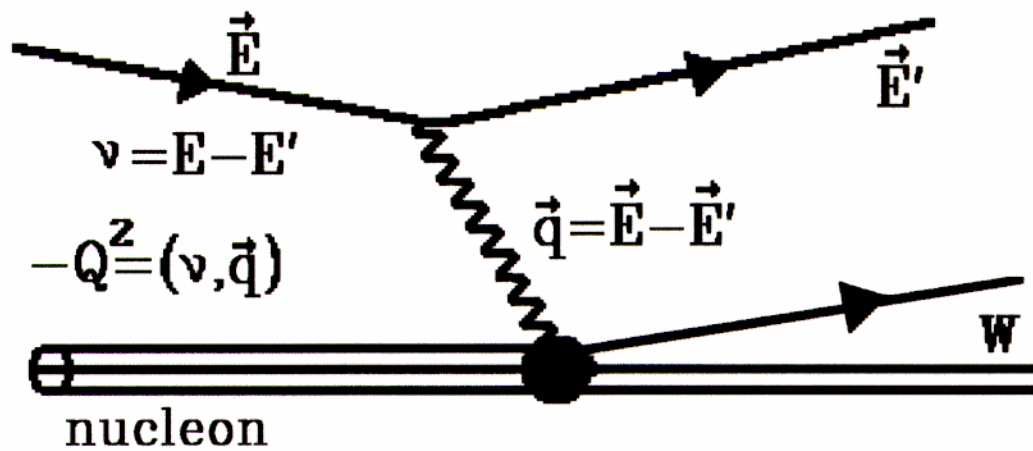
## E155 Collaboration

P. Anthony, R. Arnold, T. Averett, H. Band, C. Berisso, H. Borel, P. Bosted  
S. Bültmann, M. Buenerd, T. Chupp, S. Churchwell, G.R. Court, D. Crabb  
D. Day, P. Decowski, P. DePietro, R. Erbacher, R. Erickson, A. Feltham  
H. Fonvieille, E. Frlez, R. Gearhart, V. Ghazikhanian, J. Gomez, K. Griffioen  
C. Harris, M.A. Houlden, E. W. Hughes, C. Hyde-Wright, G. Igo, S. Incerti  
J. Jensen, J. Johnson, **P. King**, Yu. G. Kolomensky, S. Kuhn, R. Lindgren  
R. Lombard-Nelsen, J. Marroncle, W. Meyer, J. McCarthy, **P. McKee**  
**G. Mitchell**, J. Mitchell, M. Olson, S. Penttila, G. Peterson, G. G. Petratos  
R. Pitthan, D. Pocanic, R. Prepost, C. Prescott, L. M. Quin, B. Raue  
D. Reyna, L.S. Rochester, S. Rock, O. Rondon-Aramayo, F. Sabatie  
S. St. Lorant, I. Sick, L. Sorrell, T. Smith, F. Staley, L. M. Stuart, Z. Szalata  
Y. Terrien, L. Todor, **A. Tobias**, **T. Toole**, S. Trentalange, D. Walz  
R. Welsh, **F.R. Wesselmann**, T. Wright, C. Young, B. Youngman, M. Zeier  
H. Zhu, B. Zihlmann

American University  
Universität Basel  
University Blaise Pascal  
Ruhr-Universität Bochum  
California Inst. of Technology  
University of California  
PNC Saclay  
Florida International University  
Kent State University  
University of Liverpool  
LANL

University of Massachusetts  
University of Michigan  
Old Dominion University  
Smith College  
SLAC  
Stanford University  
TJNAF  
University of Virginia  
College of William and Mary  
University of Wisconsin

## Inclusive Deep Inelastic Scattering



- $\theta$  : scattering angle
- $Q^2$ : 4-momentum transfer squared
- $\nu$  : energy transfer
- $x = \frac{Q^2}{2M\nu}$  : Bjorken scaling variable

## Unpolarized Deep Inelastic Scattering

$$\frac{d\sigma}{d\Omega dE'} = \frac{4\alpha^2 E'^2 \cos^2(\theta/2)}{Q^4} \left[ \frac{F_2(x, Q^2)}{\nu} + \frac{2F_1(x, Q^2)}{M} \tan^2(\theta/2) \right]$$

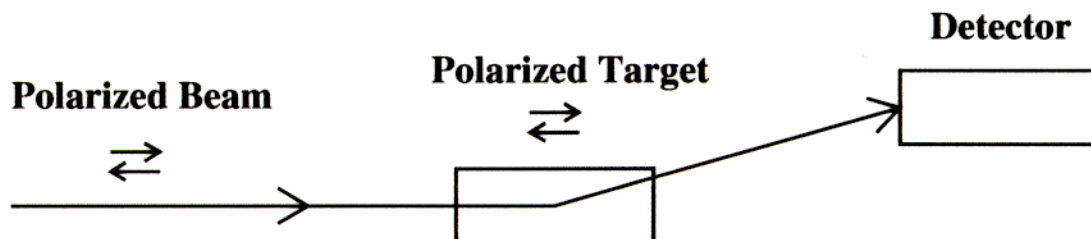
$F_1$  and  $F_2$  are the unpolarized structure functions.

In naive quark parton model  $F_1$  is related to the spin averaged quark distribution functions:

$$F_1(x, Q^2) = \frac{1}{2} \sum_i e_i^2 [q_i^\uparrow(x, Q^2) + q_i^\downarrow(x, Q^2)]$$

- Up (down) arrow denotes spin aligned (anti-aligned) with nucleon spin
- $i$  is over the quark and antiquark
- $e_i$  is the corresponding charge

## Polarized Deep Inelastic Scattering



$$A_{\parallel} = \frac{\sigma^{\uparrow\downarrow} - \sigma^{\uparrow\uparrow}}{\sigma^{\uparrow\downarrow} + \sigma^{\uparrow\uparrow}} = f_k \left[ g_1(x, Q^2) [E + E' \cos(\theta)] - \frac{Q^2}{\nu} g_2(x, Q^2) \right]$$

$$A_{\perp} = \frac{\sigma^{\downarrow\leftarrow} - \sigma^{\uparrow\leftarrow}}{\sigma^{\downarrow\leftarrow} + \sigma^{\uparrow\leftarrow}} = f_k E' \sin(\theta) \left[ g_1(x, Q^2) + \frac{2E}{\nu} g_2(x, Q^2) \right]$$

$g_1$  and  $g_2$  are the polarized structure functions.

- $A_{\parallel}$  is primarily sensitive to  $g_1$
- $A_{\perp}$  is more sensitive to  $g_2$
- $f_k$  includes contribution from unpolarized structure functions

## $g_1$ Structure Function

In QPM  $g_1$  is related to the spin distribution functions:

$$g_1(x, Q^2) = \frac{1}{2} \sum_i e_i^2 \Delta q_i$$

$$\Delta q_i = q_i^\uparrow(x, Q^2) - q_i^\downarrow(x, Q^2)$$

The index  $i$  is over the quark and antiquark flavors and  $e_i$  is the corresponding charge

However, in pQCD  $g_1$  is also sensitive to the polarized gluon distribution  $\Delta G$ :

$$g_1(x, Q^2) = \frac{1}{2} \sum_q e_q^2 [ C_q \otimes \Delta q + C_G \otimes \Delta G ]$$

where  $C_{q,G}$  are the perturbative coefficients



## Virtual Photon-Nucleon Asymmetries

$$A_1(x, Q^2) = \frac{\sigma_T^{1/2} - \sigma_T^{3/2}}{\sigma_T^{1/2} + \sigma_T^{3/2}} = \frac{g_1(x, Q^2) - \gamma^2 g_2(x, Q^2)}{F_1(x, Q^2)}$$

$\sigma_T^{1/2}$  and  $\sigma_T^{3/2}$  are transverse virtual photoabsorption cross sections for final helicity states of 1/2 and 3/2

$$\rightarrow \sigma_T^{1/2} \sim \frac{1}{2} \sum_i e_i^2 q_i^\uparrow \quad \text{and} \quad \sigma_T^{3/2} \sim \frac{1}{2} \sum_i e_i^2 q_i^\downarrow$$

$$\rightarrow A_1 \approx g_1/F_1 \text{ for } \gamma^2 = Q^2/\nu^2 = 0 \text{ (high energy)}$$

$$A_2(x, Q^2) = \frac{2\sigma_{TL}}{\sigma_T^{1/2} + \sigma_T^{3/2}} = \gamma \frac{g_1(x, Q^2) + g_2(x, Q^2)}{F_1(x, Q^2)}$$

$\rightarrow \sigma_{TL}$  is the transverse-longitudinal interference term.

## Spin Structure Measurements

Year	Exp.	Data	Target	Publications
1978	E80	$A_{\parallel}^p$	Butanol	PRL 37, 1258 (76) PRL 37, 1261 (76) PRL 41, 70 (78) PRL 45, 200 (80)
1983	E130	$A_{\parallel}^p$	Butanol	PRL 51, 1135 (83)
1988	EMC	$A_{\parallel}^p$	NH <sub>3</sub>	PLB 206, 364 (88) NPB 328, 1 (89)
1991	SMC	$A_{\parallel}^d$	Butanol	PLB 302, 533 (93)
1992	E142	$A_{\parallel}^n, A_{\perp}^n$	<sup>3</sup> He	PRL 71, 959 (93) PRD 54, 6620 (96)
1993	SMC	$A_{\parallel}^p, A_{\perp}^p$	Butanol	PLB 329, 399 (94) PLB 336, 125 (94)
1994	E143	$A_{\parallel}^p, A_{\perp}^p, A_{\parallel}^d, A_{\perp}^d$	<sup>15</sup> NH <sub>3</sub> , <sup>15</sup> ND <sub>3</sub>	PRL 74, 346 (95) PRL 75, 25 (95) PLB 364, 61 (95) HEP-PH 9802357 PRL 76, 587 (96) PRL 78, 815 (97)
1994	SMC	$A_{\parallel}^d$	Butanol	PRB 357, 248 (95)
1995	SMC	$A_{\parallel}^d$	Butanol	PRB 396, 338 (97)
1995	HERMES	$A_{\parallel}^n$	<sup>3</sup> He	PLB 404, 383 (97)
1995	E154	$A_{\parallel}^n, A_{\perp}^n$	<sup>3</sup> He	PRL 79, 26 (97) PLB 404, 377 (97) PLB 405, 180 (97)
1996	SMC	$A_{\parallel}^p, A_{\parallel}^d$	NH <sub>3</sub>	CERN-EP 98-085
1996	HERMES	$A_{\parallel}^p$	H <sub>2</sub>	HEP-EX 980715
1997	E155	$A_{\parallel}^p, A_{\perp}^p, A_{\parallel}^d, A_{\perp}^d$	<sup>15</sup> NH <sub>3</sub> , LiD	



## Physics Asymmetries

$$A_{\parallel}(\text{or } A_{\perp}) = C_1 \left( \left( \frac{N_L - N_R}{N_L + N_R} \right) \frac{1}{f P_b P_t} + C_2 \right) + A_{RC}$$

- $N_L$  and  $N_R$  are rates for L and R beam helicity, corrected for pair-symmetric contributions and electronics dead time.
- $C_1$  and  $C_2$  are for nuclear corrections and/or target contamination corrections.
- $f$  is the dilution factor ( $\approx 0.15$  for  $\text{NH}_3$ ,  $\approx 0.35-0.40$  for LiD)
- $P_b$ ,  $P_t$  are the beam and target polarizations
- $A_{RC}$  radiative correction

## Neutron Extraction

$$g^n(x, Q^2) = \frac{2g^d(x, Q^2)}{(1 - 1.5\omega_d)} - g^p(x, Q^2)$$

- $\omega_d = 0.05 \pm 0.01 =$  probability deuteron is in D-state

## E155 Keys

- Decrease statistical uncertainties
  - Collected  $\approx 200$  million deep inelastic events
  
- Reduce systematic uncertainties
  
- Cover wide  $x$  and  $Q^2$  range in single experiment
  - Use 48.3 GeV electron beam
  - Add third spectrometer
  
- Improve precision on  $Q^2$  dependence of  $g_1$  to constrain  $\Delta G$

## Polarized Electron Beam

- Intense 48.3 GeV Beam
- Highly polarized :  $P_b = 0.813 \pm 0.020$
- Beam polarization measured with two independent Møller detector systems

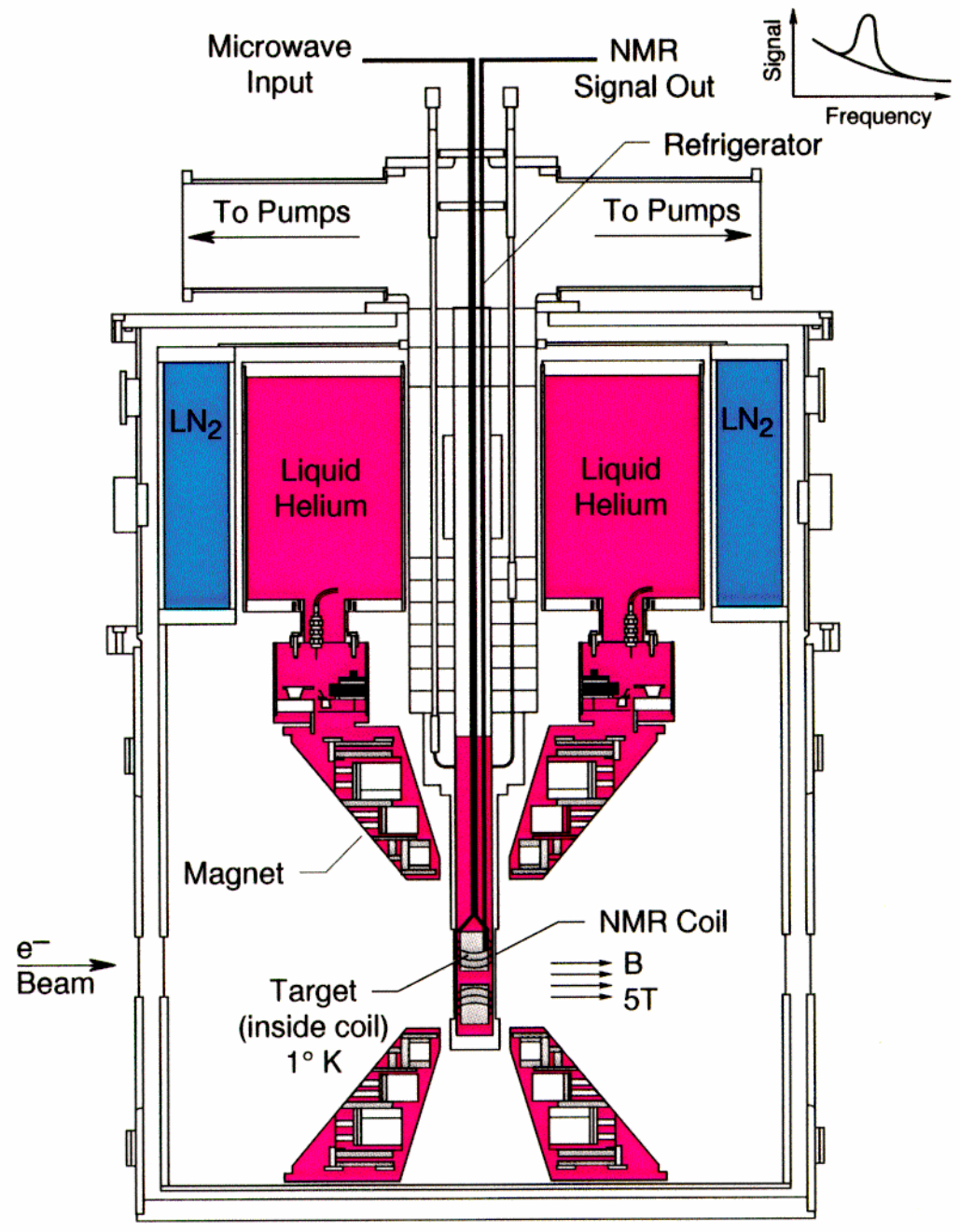
$$\begin{aligned}
 A &= \frac{N_{e^-}(\uparrow\uparrow) - N_{e^-}(\uparrow\downarrow)}{N_{e^-}(\uparrow\uparrow) + N_{e^-}(\uparrow\downarrow)} \\
 &= -P_{beam}P_{foil} \frac{(7 + \cos^2\theta_{cm})\sin^2\theta_{cm}}{(3 + \cos^2\theta_{cm})^2}
 \end{aligned}$$

- Single arm system and double arm coincidence system agree within 1%
- 20-50  $\mu\text{m}$  thick target foils
- Helicity varied randomly from pulse to pulse (120 Hz)

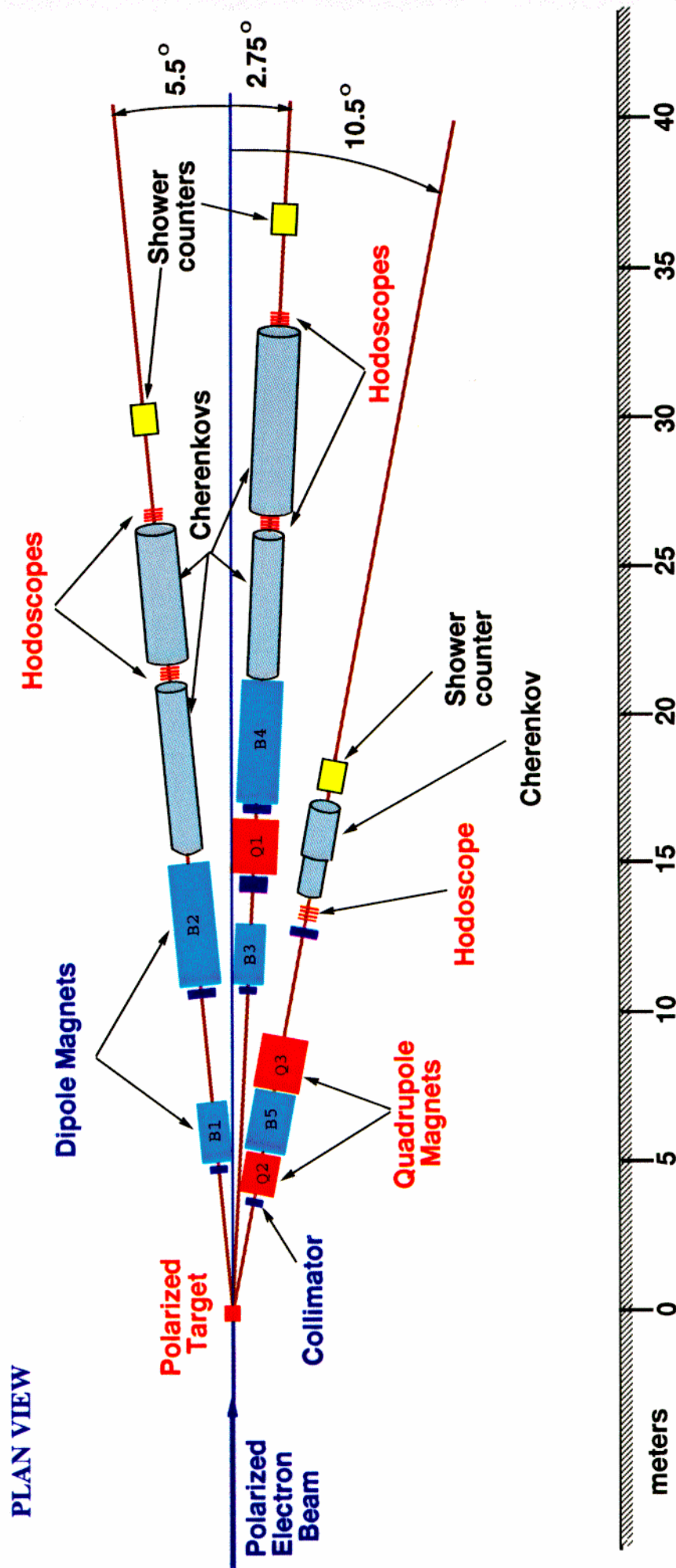
## Polarized Targets

- Solid cryogenic targets
- Dynamic nuclear polarization using microwaves
- $^{15}\text{NH}_3$  used for proton target
- Improved  $\mu$ wave power delivery →  $\langle P_t \rangle \approx 80\%$
- $^6\text{LiD}$  used for deuteron target
- $\text{LiD} \approx 50\%$  improvement in dilution over  $^{15}\text{ND}_3$
- $\text{LiD}$  6 times as radiation resistant
- $\langle P_t \rangle \approx 20\%$

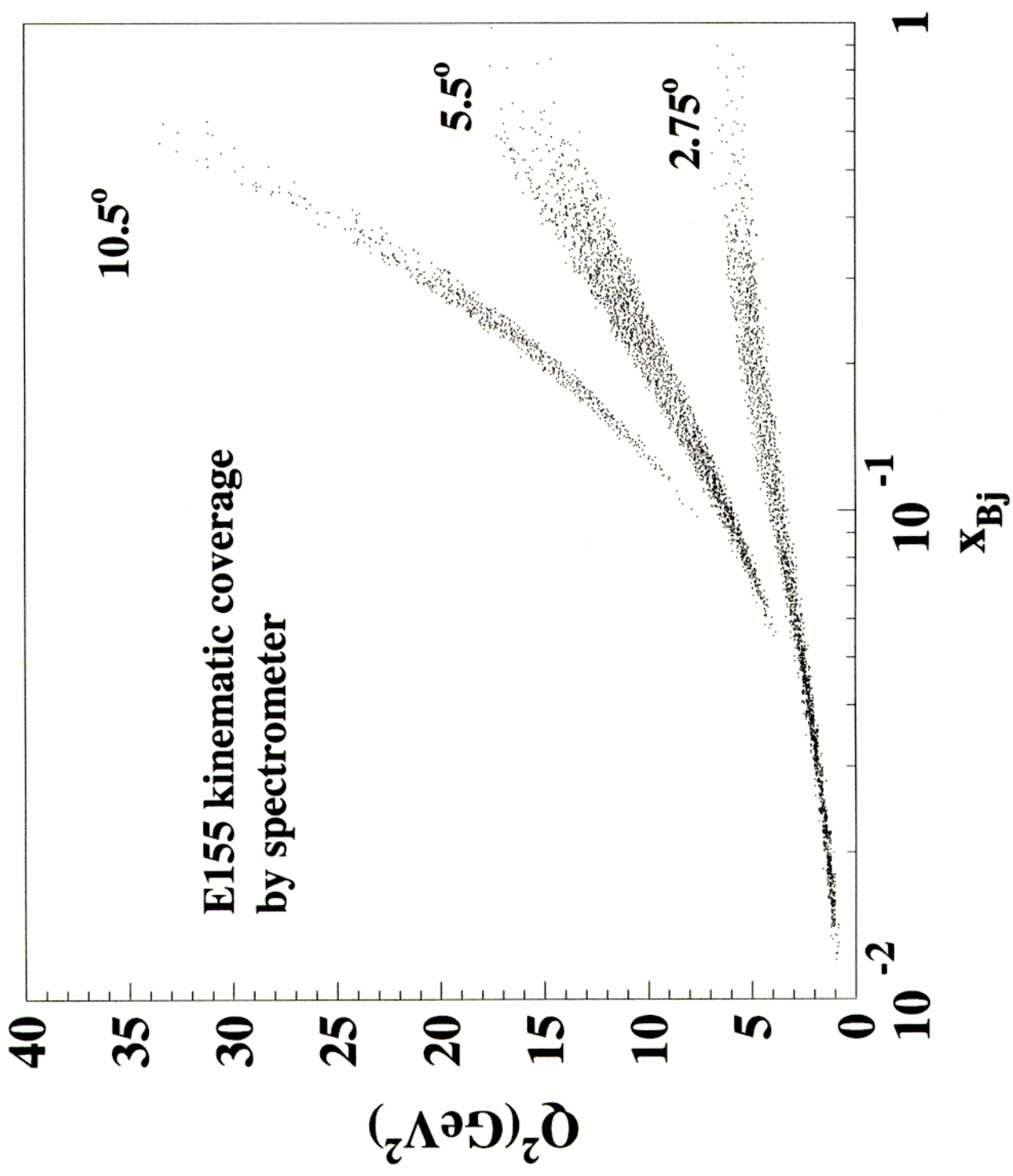
# E143/E155 Target



# E155 Spectrometers



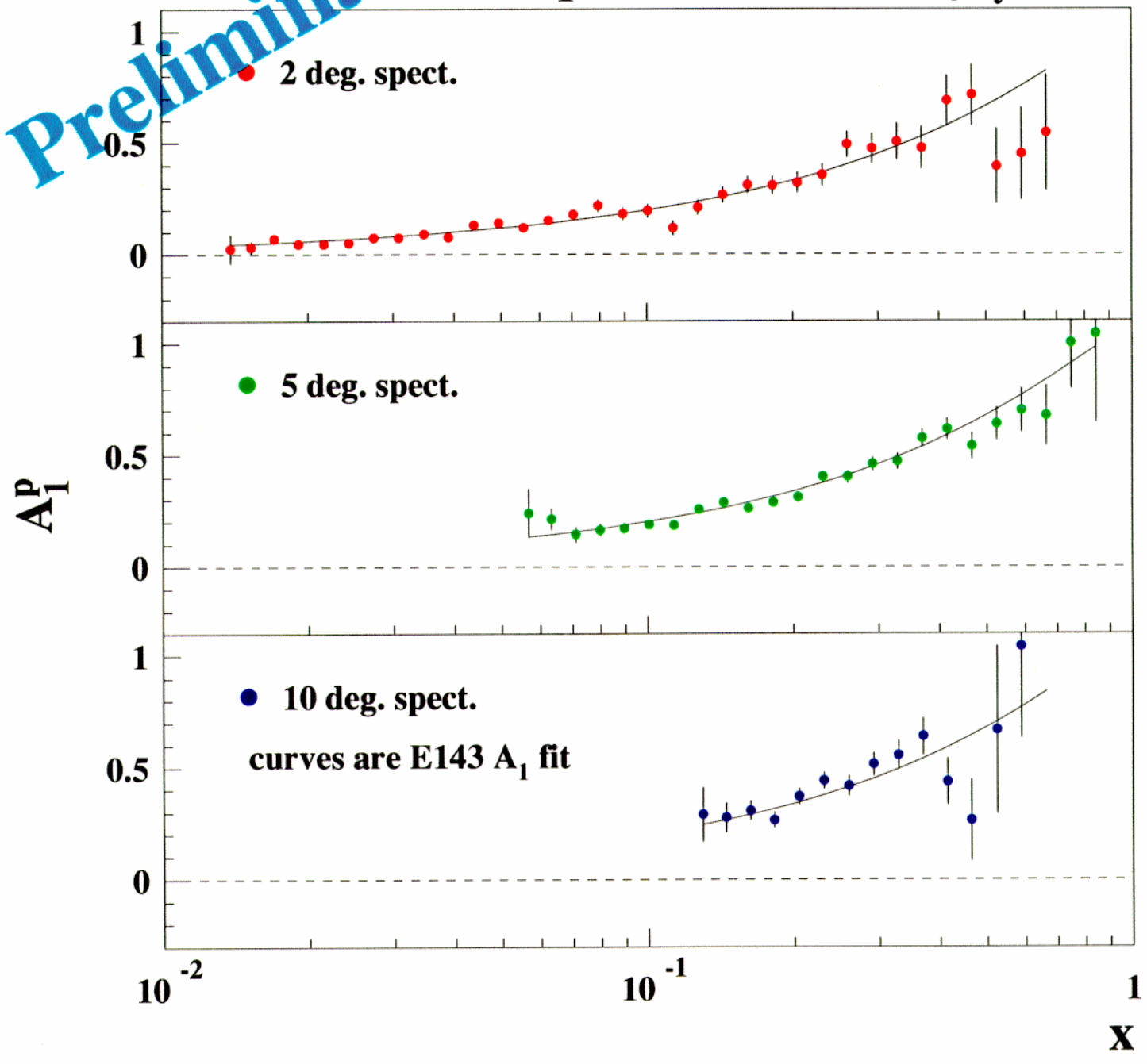




Preliminary

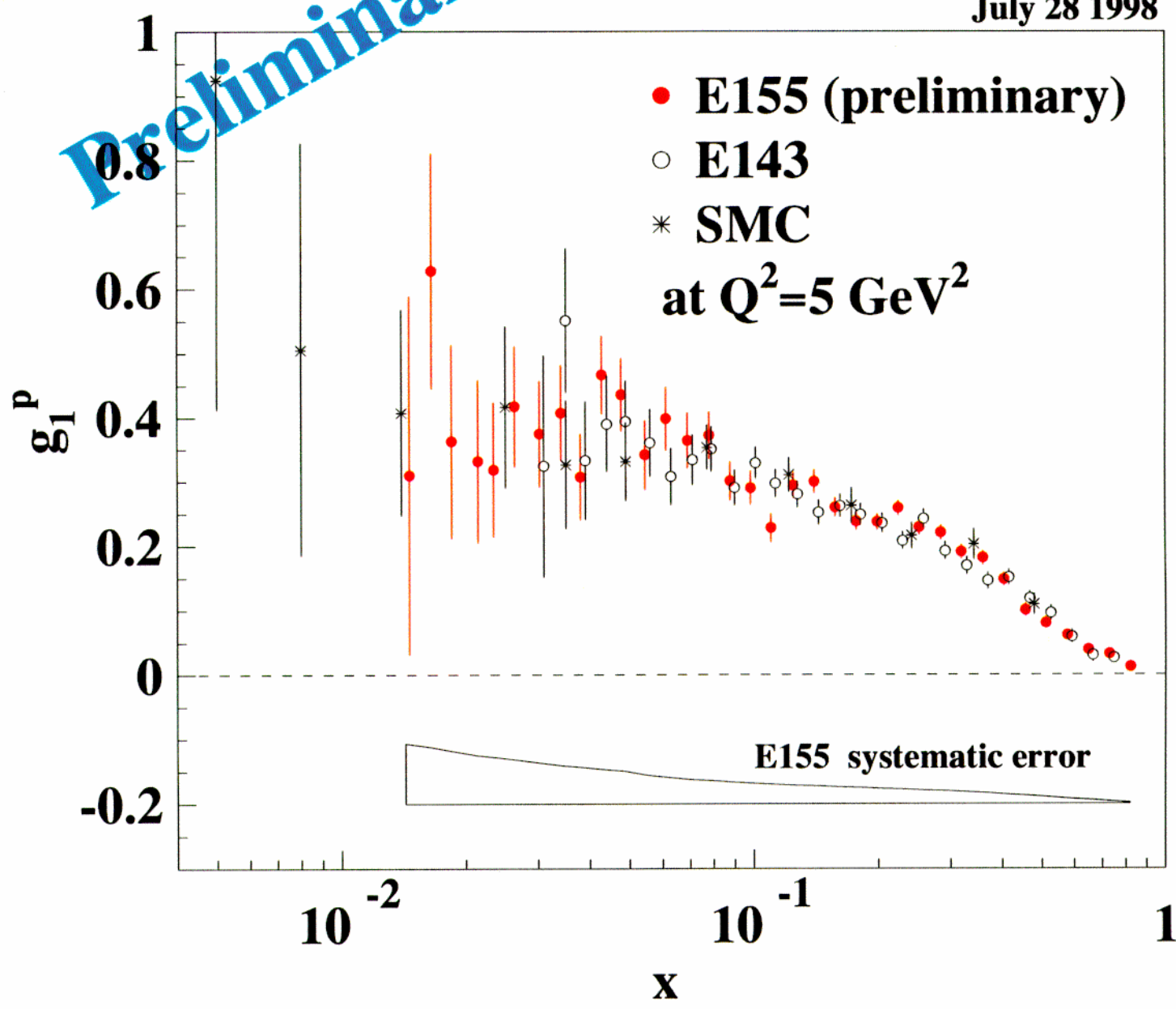
# E155 $A_1$ proton

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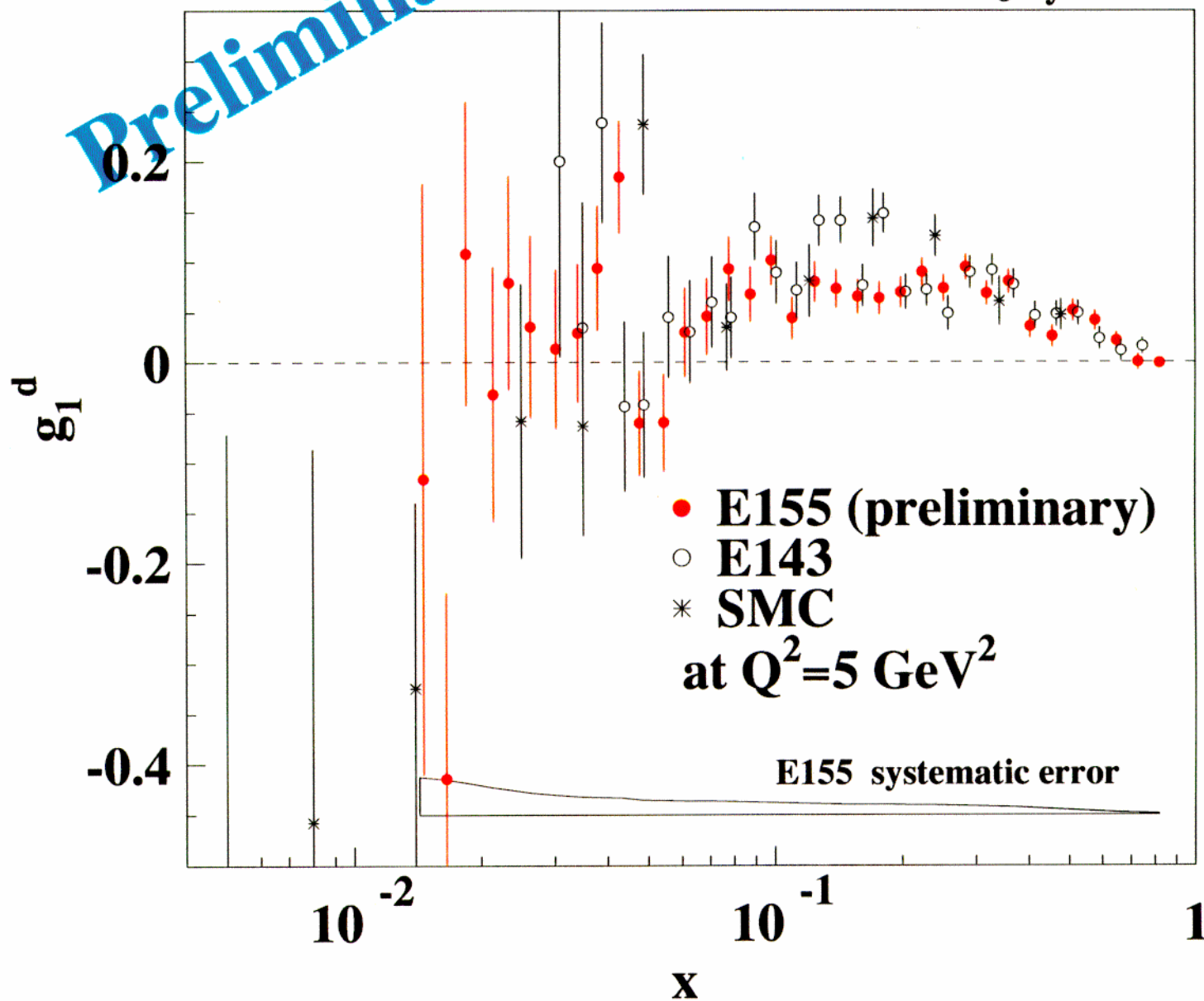
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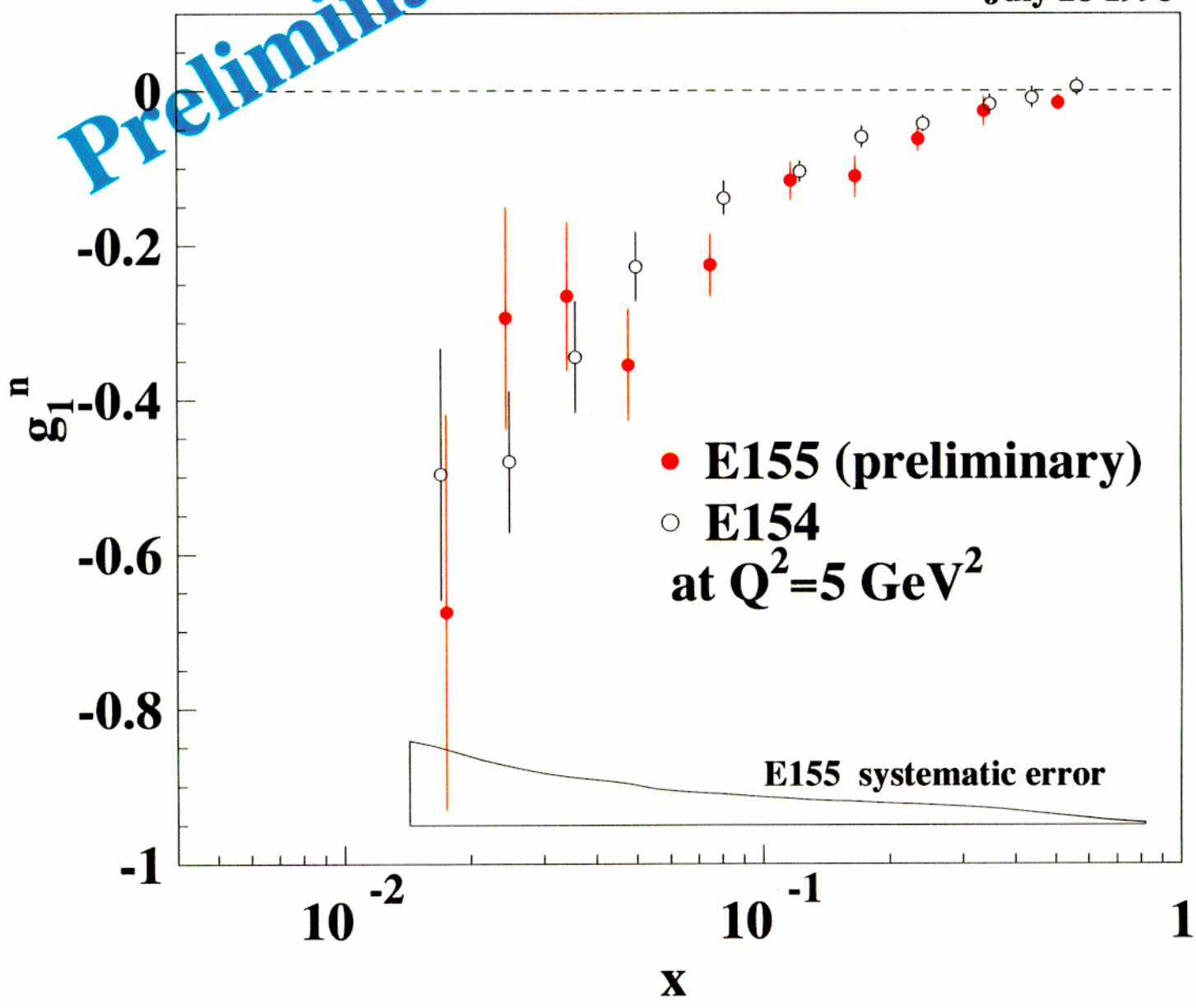
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Preliminary



## Global pQCD Fitting

- Use pQCD expansion for  $g_1$ :

$$g_1(x, Q^2) = \frac{1}{2} \sum_q e_q^2 [C_q \otimes \Delta q + C_G \otimes \Delta G]$$

→  $C_q$  contribution starts at order 1,  $C_G$  at order  $\alpha_S^1$

- Evolve  $\Delta q$  and  $\Delta G$  from initial  $Q_0^2$  to  $Q^2$  using NLO DGLAP equations.
- Parameterize non-perturbative inputs.

$$EG : \quad \Delta f(x, Q_0^2) = A_f x^{\alpha_f} f(x, Q_0^2)$$

where  $f =$  valence  $\Delta u_v$  and  $\Delta d_v$ , sea, gluons

- Factorization Scheme Dependence

→ Common schemes are  $\overline{MS}$  and Adler-Bardeen

→ Ambiguity in definition of contributions

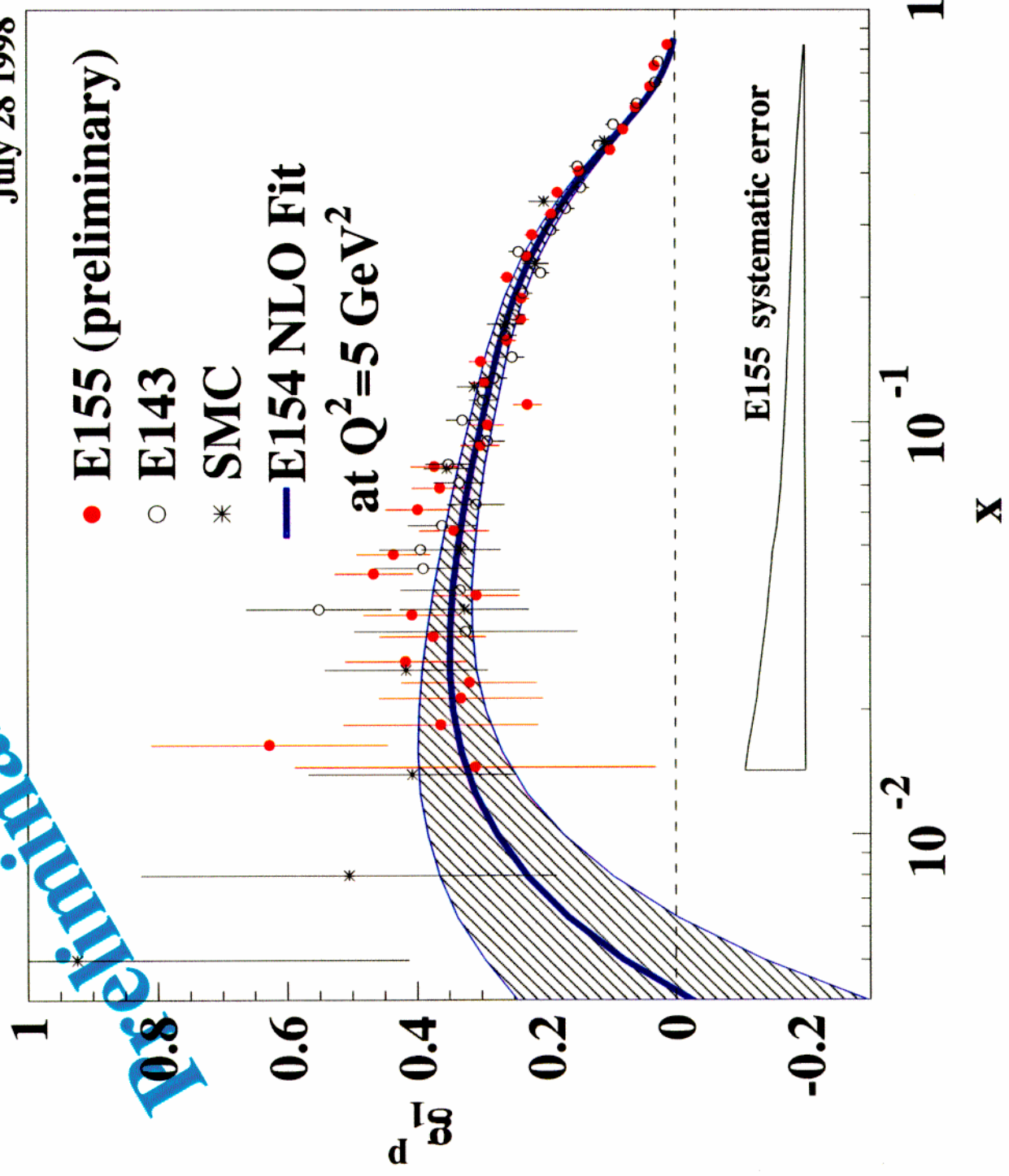
$$\Delta\Sigma(AB) = \Delta\Sigma(\overline{MS}) + \frac{N_f}{2\pi} \alpha_S \Delta G$$

$\alpha_S(Q^2)\Delta G(Q^2)$  is independent of  $Q^2$  to  $O(\alpha_S^2)$



Preliminary

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## $g_2(x, Q^2)$ Structure Function

- Wandzura-Wilczek  $g_2$  expression (twist-2 only)

$$g_2^{WW}(x, Q^2) = -g_1(x, Q^2) + \int_x^1 \frac{dy}{y} g_1(y, Q^2)$$

- In general there are additional contributions.

$$\begin{aligned}
 & \begin{array}{ccc} \text{twist-2} & & \text{twist-2} \quad \text{twist-3} \\ \Downarrow & & \Downarrow \quad \Downarrow \end{array} \\
 g_2(x, Q^2) &= g_2^{WW}(x, Q^2) - \int_x^1 \frac{\partial}{\partial y} \left( \frac{m}{M} h_T(y, Q^2) + \xi(y, Q^2) \right) \frac{dy}{y} \\
 &= g_2^{WW}(x, Q^2) + \bar{g}_2(x, Q^2)
 \end{aligned}$$

- $\bar{g}_2$  includes quark-gluon correlations ( $\xi$ ) inside the nucleon
- Twist-2 term (quark transverse spin distribution  $h_T$ ) in  $\bar{g}_2$  suppressed by  $\frac{m}{M} \ll 1$
- Burkhardt-Cottingham (BC) sum rule from virtual Compton scattering dispersion relations in  $Q^2 \rightarrow \infty$  limit:  $\int_0^1 g_2(x, Q^2) dx = 0$

## Operator Product Expansion (OPE) For $g_1$ and $g_2$

OPE used in QCD to separate the physics into a perturbative piece (easily treated) and a non-perturbative piece (unknown matrix elements).

$$\Gamma_1^{(n)} = \int_0^1 x^n g_1(x, Q^2) dx = \frac{a_n}{2}, \quad n = 0, 2, 4, \dots$$

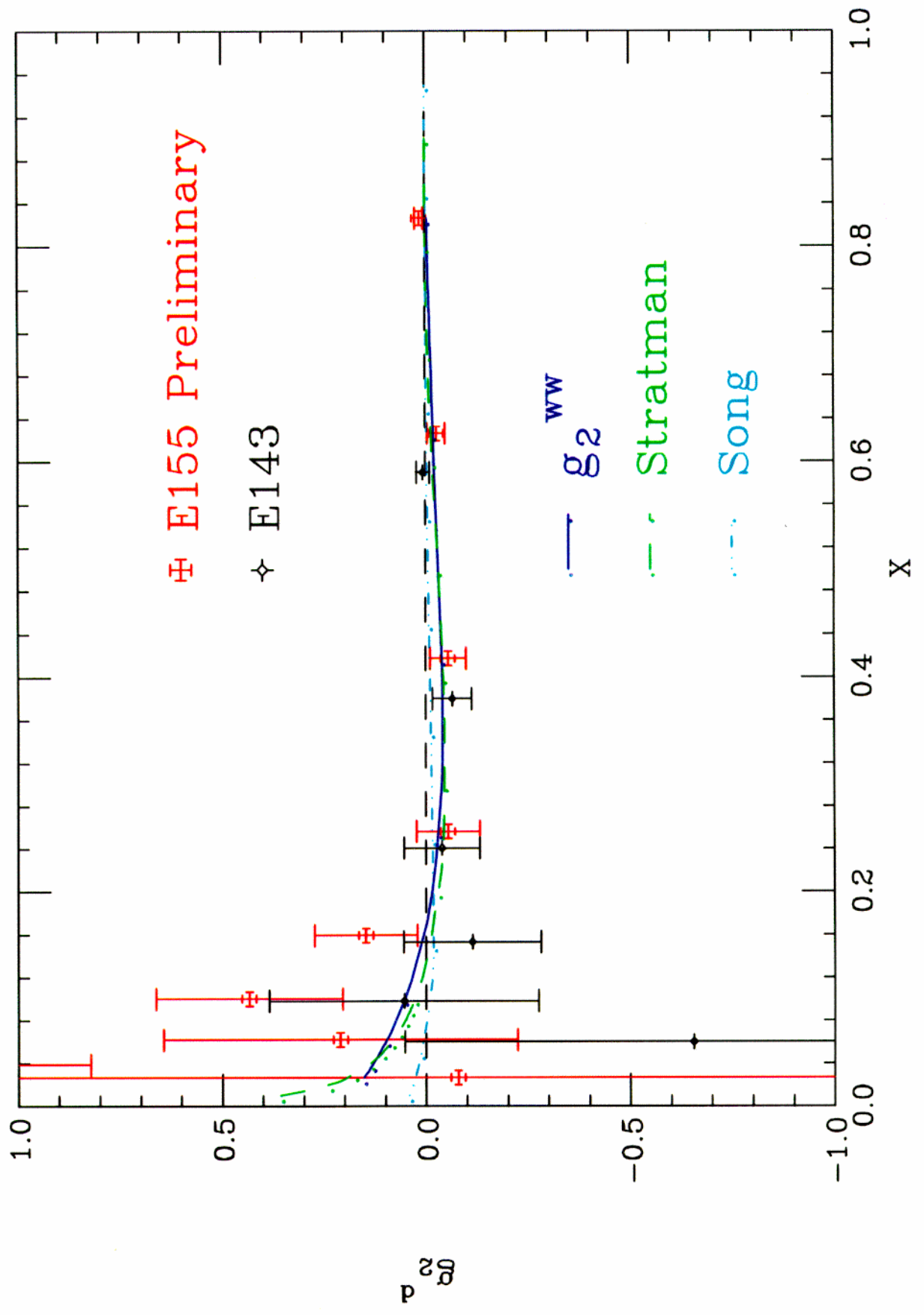
$$\Gamma_2^{(n)} = \int_0^1 x^n g_2(x, Q^2) dx = \frac{1}{2} \frac{n}{n+1} (d_n - a_n)$$

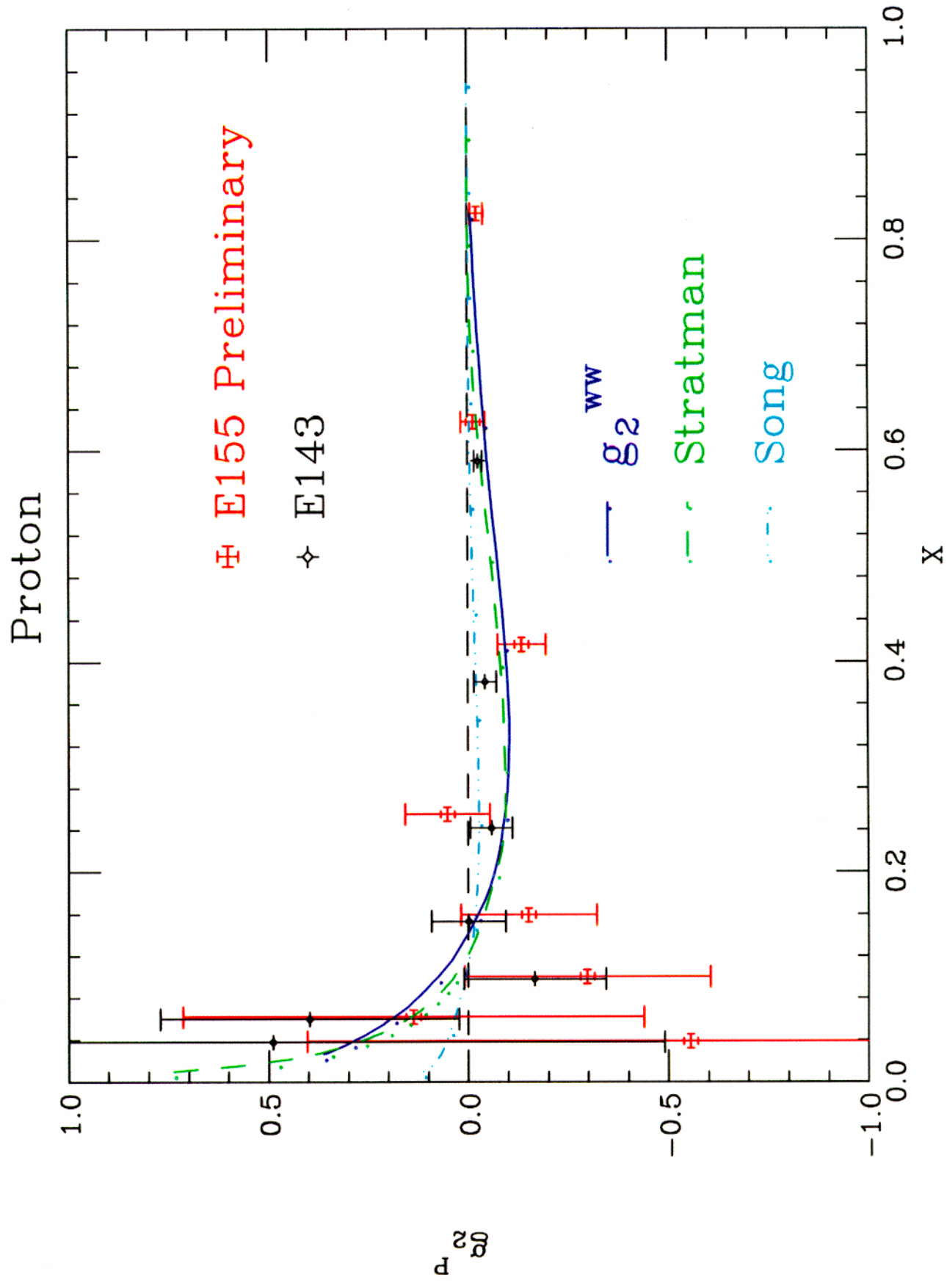
$$n=2,4,\dots$$

$a_n$  are the twist-2 and  $d_n$  are the twist-3 matrix elements.

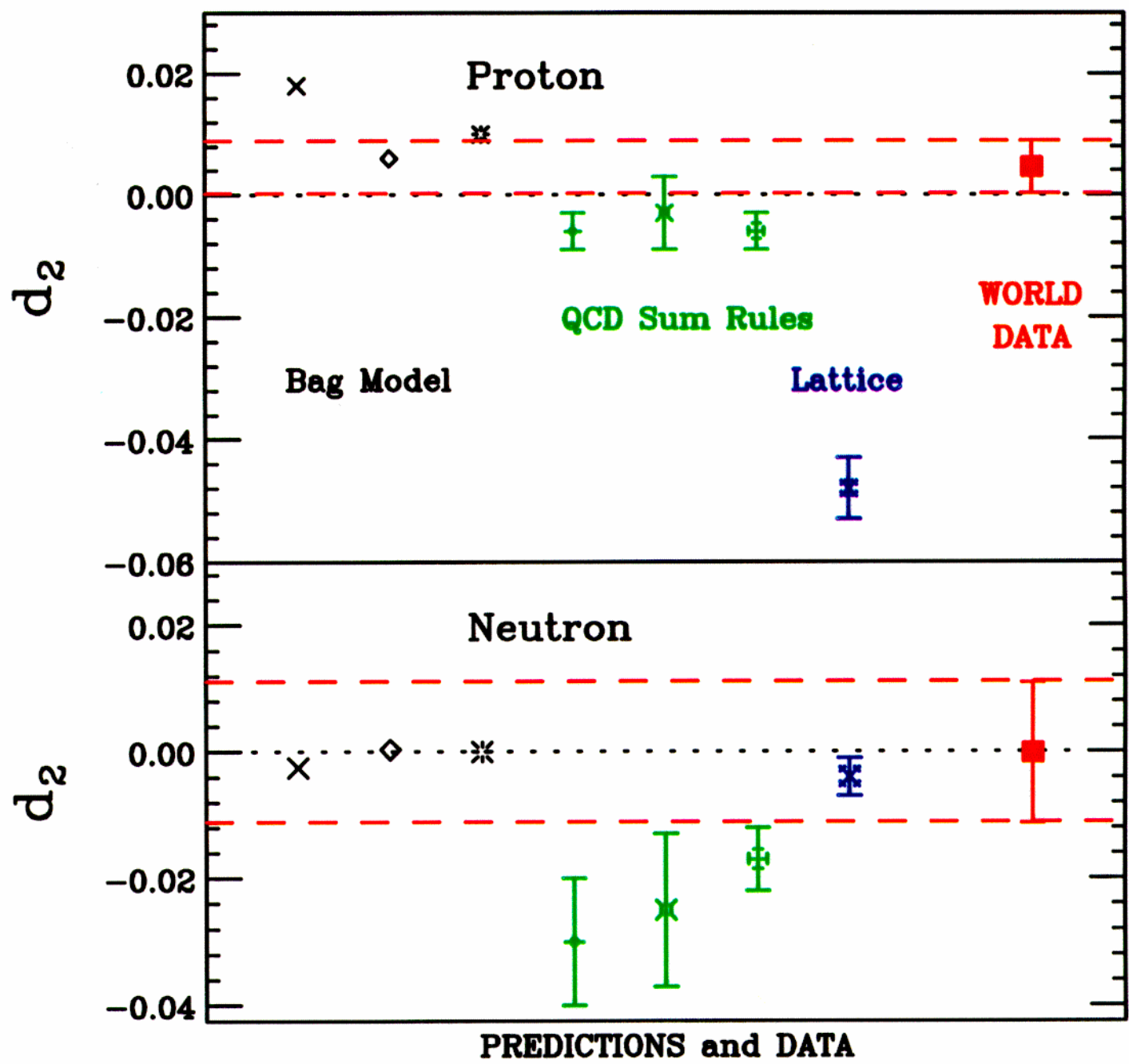
$$\begin{aligned} d_2 &= \int_0^1 x^2 [2g_1(x, Q^2) + 3g_2(x, Q^2)] dx \\ &= 3 \int_0^1 x^2 \overline{g}_2(x, Q^2) dx \end{aligned}$$

# Deuteron





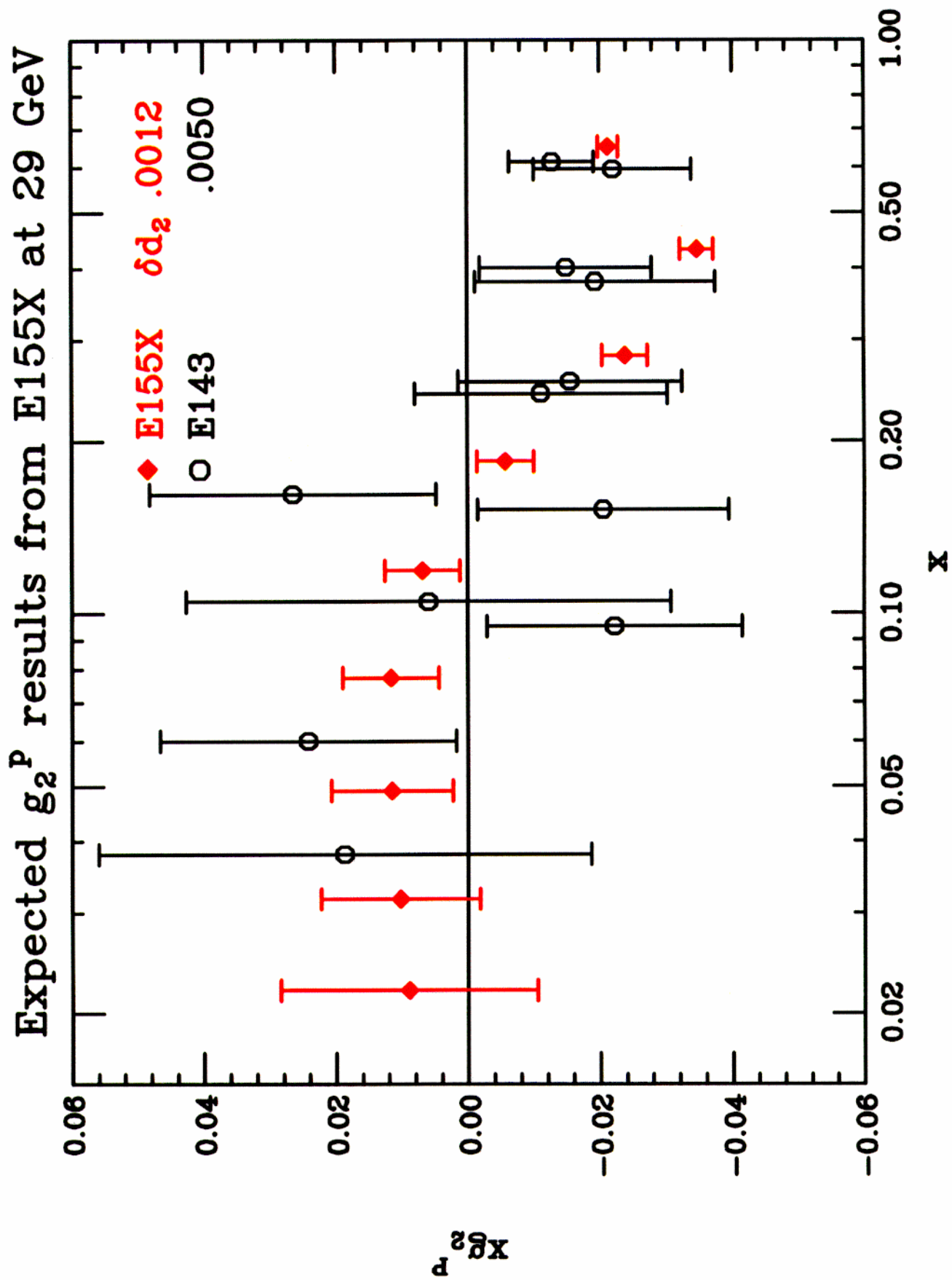
### Twist 3 Matrix Element $d_2$



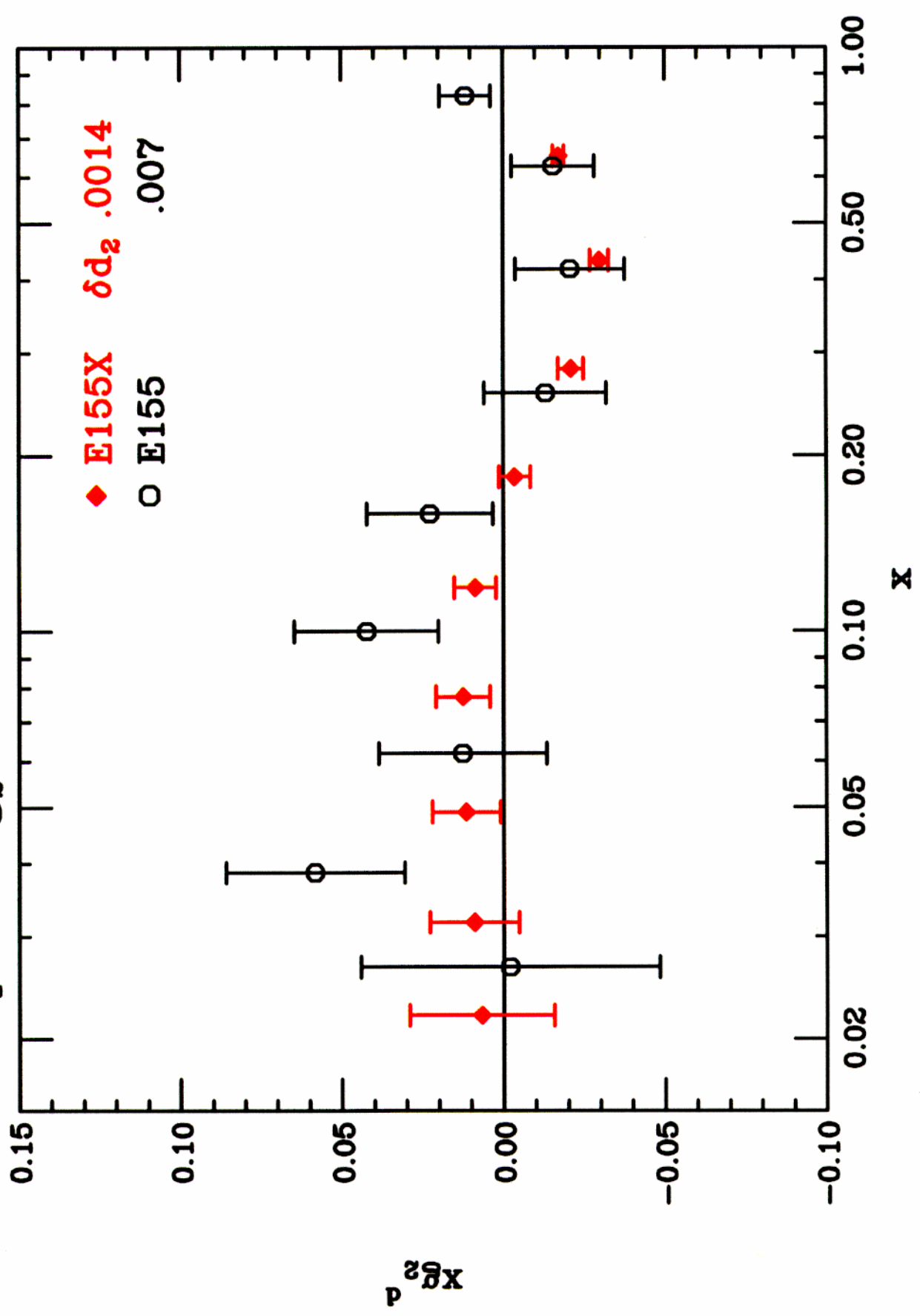


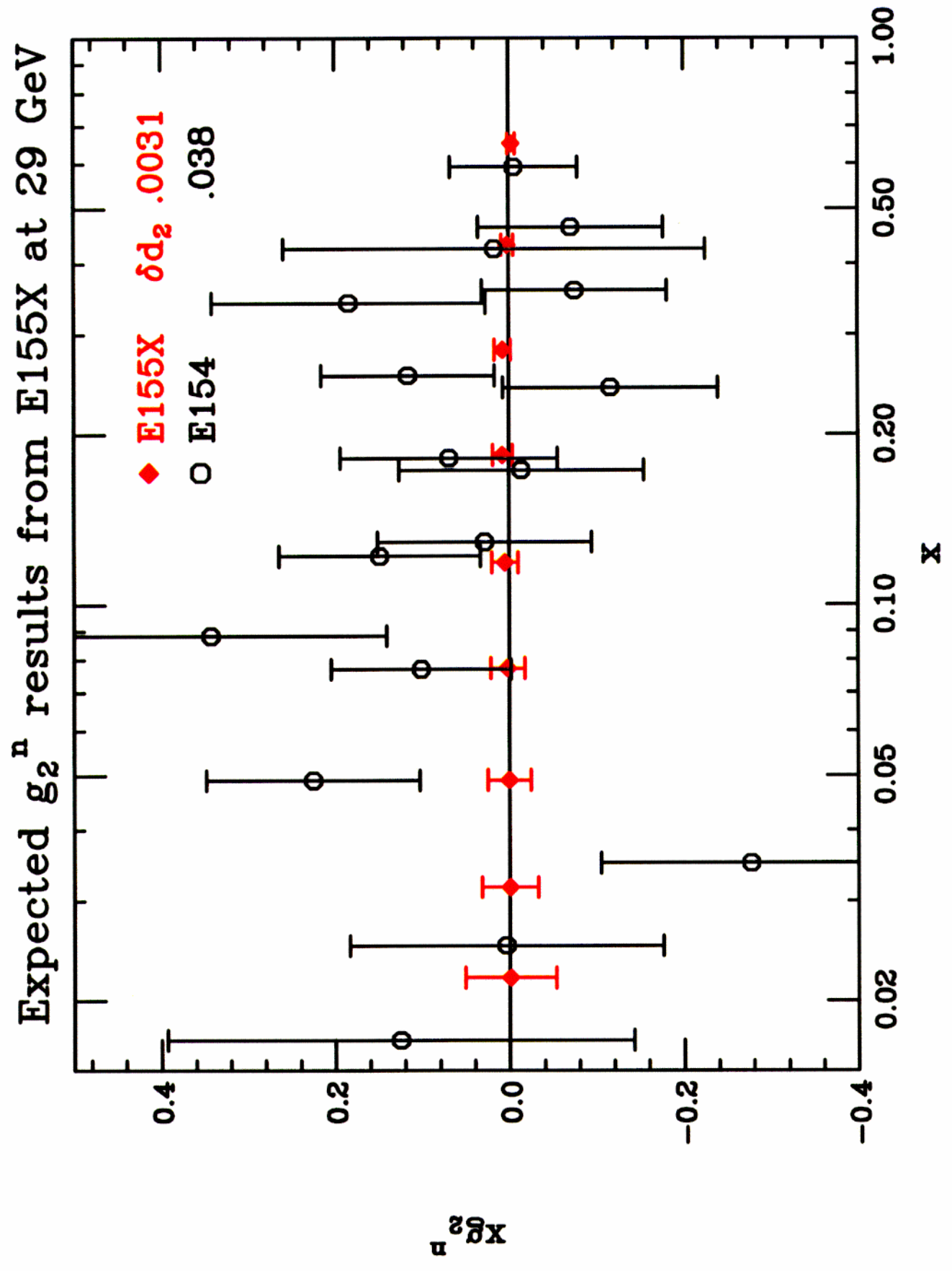
## E155x

- An extension to E155 has been approved to precisely measure  $g_2$ .
- The experiment is planned for February and March 1999.
- The goal is to measure  $g_2$  (and  $A_2$ ) for the proton and deuteron over a wide range in  $x$  and  $Q^2$ .
- Significant improvement in sensitivity to contributions beyond  $g_2^{WW}$

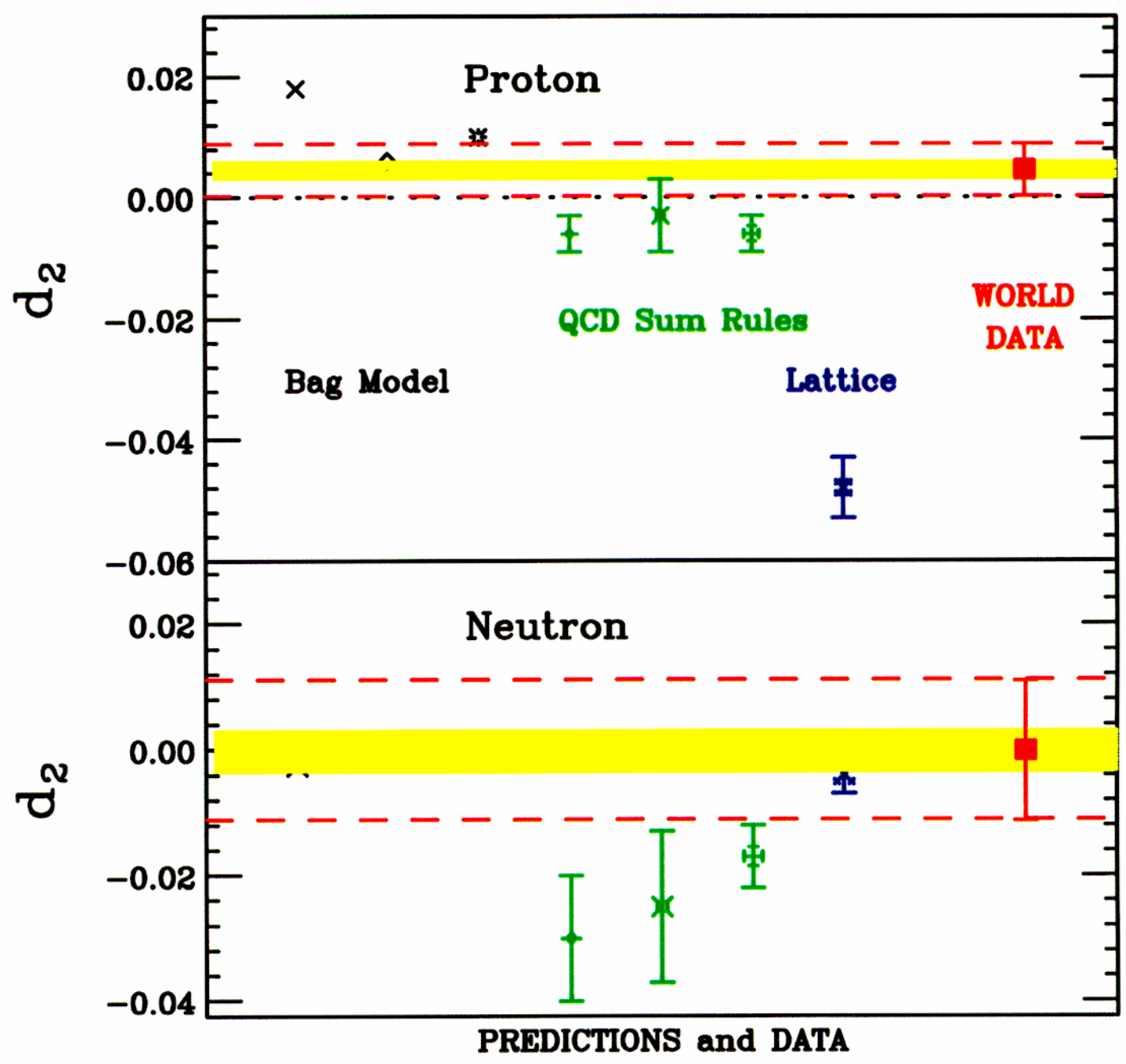


# Expected $g_2^d$ results from E155X at 29GeV





# E155X sensitivity for $d_2$



## Summary

- E155 has precision data on  $g_1$  over a wide  $x$  and  $Q^2$  range for the proton and deuteron
- NLO analysis should yield improved constraints on gluon polarization
- Extension run will provide precision  $g_2$  data on proton and deuteron
- More direct information on gluon polarization will have to come from RHIC, COMPASS, SLAC E156?, polarized HERA?