

New Sub-Millimeter Dimensions,

The Hierarchy Problem,

and

Quantum Gravity at a TeV.

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The Hierarchy Problem

- There are two seemingly fundamental energy scales in nature:
 - Electroweak scale $m_{EW} \sim 10^3 \text{ GeV}$
 - Planck scale $M_{Pl} \sim G_N^{-1/2} \sim 10^{18} \text{ GeV}$.
- Hierarchy problem: why is $m_{EW} \ll M_{Pl}$, and why is it radiatively stable?
- Central motivation for extensions of SM, with two major frameworks: low-energy SUSY, technicolor.
- But the problem only exists if we believe that both m_{EW}^{-1} , M_{Pl}^{-1} correspond to physical short-distance scales

- That m_{EW}^{-1} is a physical scale is an experimental certainty, since electroweak interactions have been probed @ distances approaching m_{EW}^{-1} .

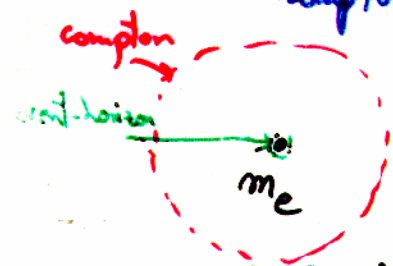
- But why do we believe M_{pl} is fundamental?
 - Gravity becomes strong (comparable to gauge interactions) @ energies $\sim M_{pl}$:

 $\sigma(E) \sim 1/E^2$

 $\sigma(E) \sim G_N^2 E^2 \sim E^2/M_{pl}^4$

- Quantum gravitational effects become important @ energies $\sim M_{pl}$: point particle of mass m has

$R_{\text{Compton}} \sim 1/m$, $R_{\text{event-horizon}} \sim G_N m$



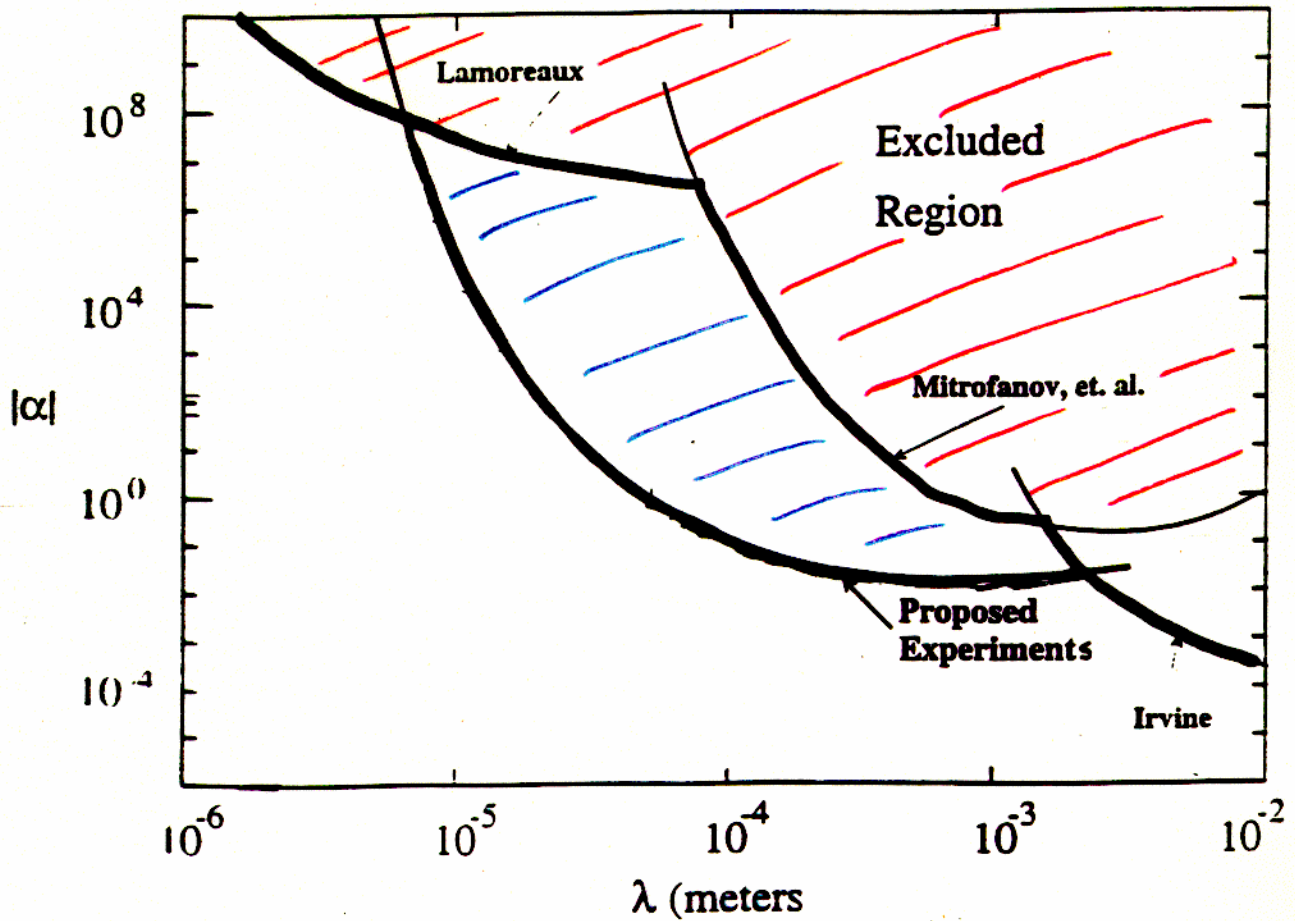
Quant. Gravity irrel.



Quant. Gravity Imp.

Is M_{pl} Fundamental?

Experimental Status of Gravitational-Strength Forces in the Sub-Centimeter Regime Figure 1
(Price et al. '98)



- Price, Long, Chan et. al.
- Kapitulnik, Keay et. al.

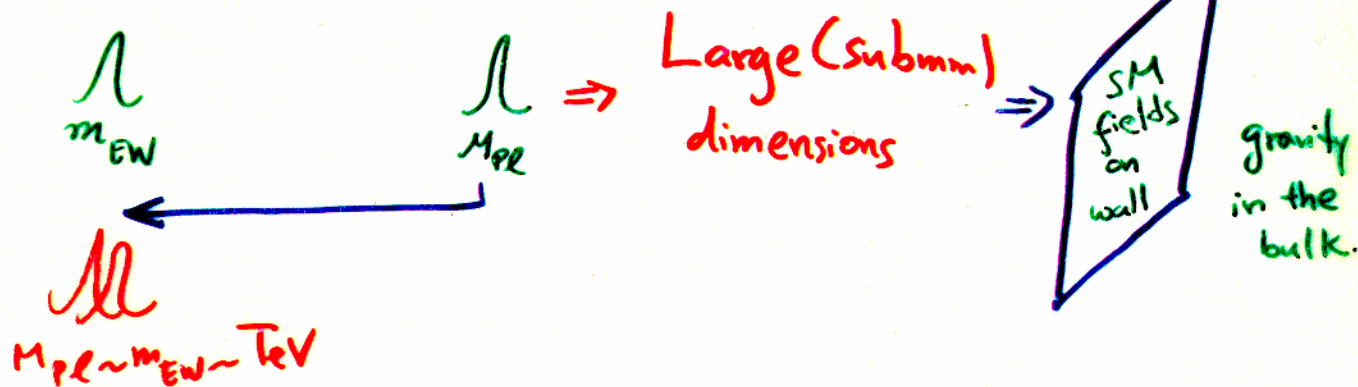
- In sharp contrast to electroweak interactions, gravity has not remotely been measured at distances approaching $M_{\text{Pl}}^{-1} \sim 10^{-33}$ cm.

Gravity has only been directly measured @ ~ 1 cm.

- Interpretation of M_{Pl} as a fundamental energy scale is based on assumption that gravity is unmodified over 33 orders of magnitude, from where it is measured down to M_{Pl}^{-1} .
- It is worthwhile to question this extrapolation

Outline

(1) Our Framework



(2) Is it alive?

- Lab bounds
- Astrophysics { Sun, Red Giants, SN 1987A
- Cosmology.

(3) Experimental signatures

- sub-mm forces
- Quantum gravity at LHC, NLC.

(4) Model-building issues (proton decay, gauge coupling unification, ...)

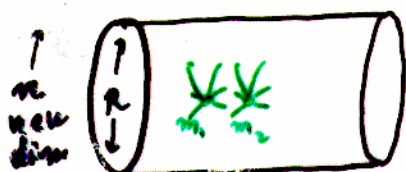
Our Scenario

• Simplest "solution" to hierarchy problem:

gravity becomes strong near the electroweak scale, i.e. $M_{Pl \text{ fund.}} \sim m_{EW}$.

• How can this be reconciled with observed weakness of gravity at distances $\gtrsim \text{cm}$?

New dimensions large compared to m_{EW} :



$$V \sim \frac{G_{N(4+n)} m_1 m_2}{r^{n+1}} \quad (r \ll R)$$

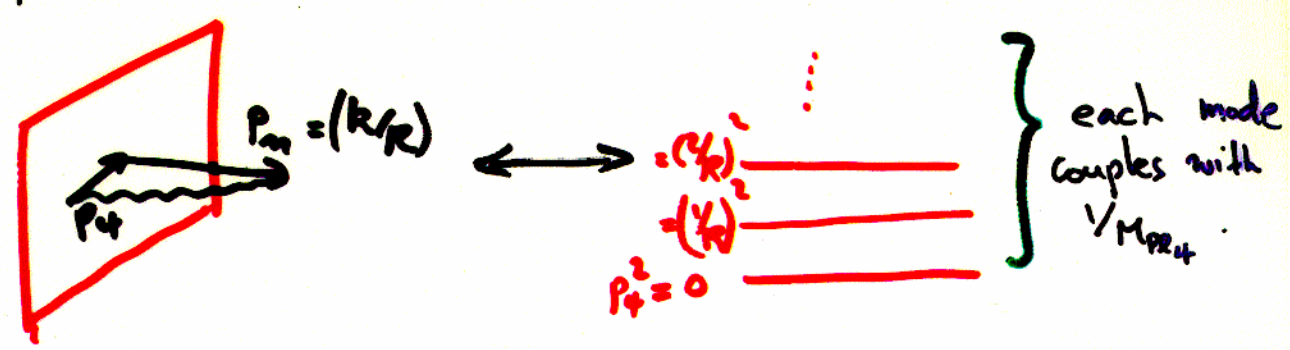
$$V \sim \frac{G_{N(4+n)} m_1 m_2}{R^n r} \quad (r \gg R)$$

$$\Rightarrow M_{Pl,4}^2 = M_{Pl,(4+n)}^{n+2} R^n \quad \text{Gauss' Law}$$

• Alternatively:

$$\int d^{4+n} x F_g M_{Pl(4+n)}^{n+2} R \rightarrow \int d^4 x F_g (M_{Pl(4+n)}^{n+2} V_n) R = \int d^4 x F_g M_{Pl,4}^2 R \quad \text{Kaluza-Klein}$$

- Still another way: From 4-d viewpoint, we have the graviton + its Kaluza-Klein excitations for each extra dimension with masses evenly spaced by $\frac{1}{R}$:



- $$V(r) = \sum_{KK} G_{N_4} m_1 m_2 \frac{e^{-m_{KK} r}}{r}$$

For $r \gg R$, only the usual massless graviton is relevant and $V(r) = G_{N_4} \frac{m_1 m_2}{r}$.

For $r \ll R$, $\sim (R/r)^n$ KK excitations are unsuppressed, so $V(r) \sim G_{N_4} \frac{m_1 m_2}{r} \times (R/r)^n \sim \frac{G_{N_4} m_1 m_2}{r^{n+1}}$.

\Rightarrow $M_{Pl,4}^2 = M_{Pl,(n+1)}^{n+2} R^n$

Since $M_{Pl_4}^2 = M_{Pl_{(4+n)}}^{n+2} R^n$, Gauss

Can be enormous \nearrow

$\sim \text{TeV}$ \nearrow

large compared to TeV^{-1} \nearrow

Enormity of M_{Pl_4} (\Leftrightarrow weakness of gravity @ long distances) does not have to reflect a fundamental short-distance scale, it could merely reflect large new spatial dimensions.

$R \sim 10^{\frac{30}{n} - 17} \text{ cm} \times \left(\frac{1 \text{ TeV}}{M_{Pl_{(4+n)}}} \right)^{\frac{n+2}{n}}$

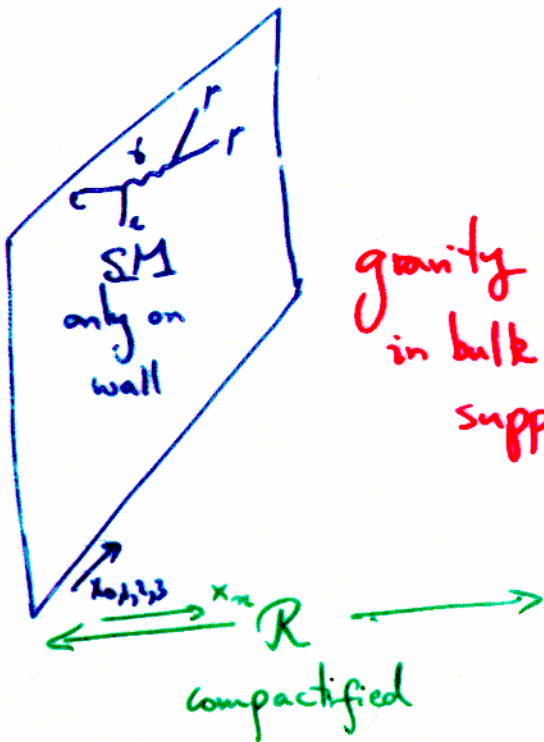
$n=1, R \sim 10^{13} \text{ cm} \times$
 $n=2, R \sim \text{mm}!$
 \vdots
 $n=6, R \sim (10 \text{ MeV})^{-1}$

Note as $n \rightarrow \infty, R \rightarrow M_{Pl_{(4+n)}}^{-1}$

Horava-Witten $10^3 \text{ GeV} \rightarrow 10^{16} \text{ GeV} \quad R \sim (10^{16} \text{ GeV})^{-1} \sim 10^{-27} \text{ cm}$

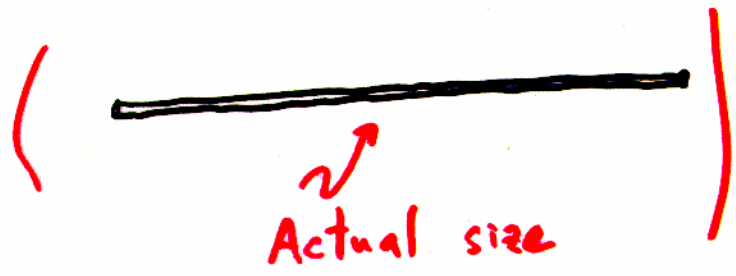
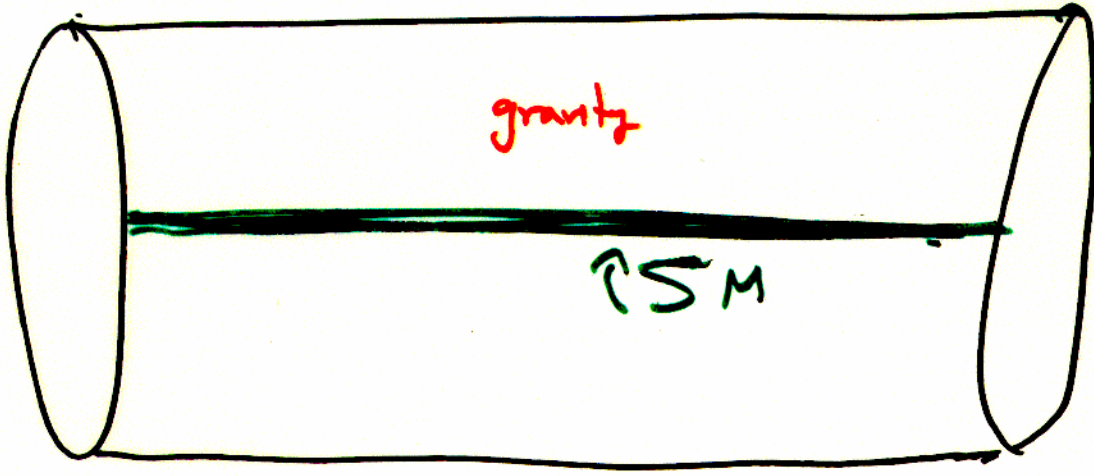
• While gravity has not been probed @ distances \lesssim mm, QED, QCD, electroweak interactions certainly have been.

• Therefore, @ energies \lesssim TeV, SM fields must be localised to a wall in the extra dimensions:



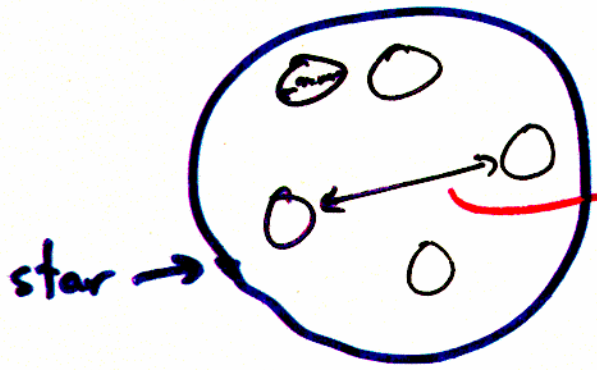
gravity in bulk, coupling suppressed by $1/M_{pl(4+m)} \sim 1/TeV$

Perfectly reasonable effective theory at energies \lesssim TeV, @ TeV, Quant. gravity is revealed



Is this scenario alive?

- Most gravitational effects are macroscopic + do not sense modification of gravity at sub-mm distance, even if typical interparticle distances are smaller than mm.



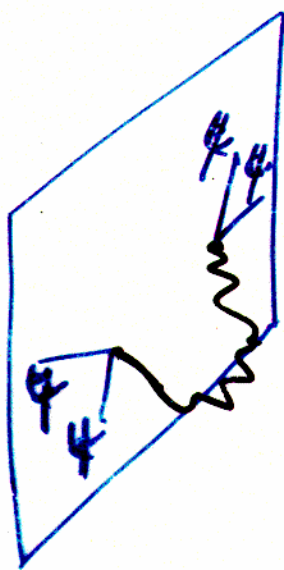
Normal Newtonian gravity between pairs, normal macroscopic energetics

$$\frac{E_{\text{grav}}}{M} = \underbrace{\int_{r_n}^R d^3r \frac{G_{N4}}{r} \rho}_{\text{Normal piece, dominated by large distances}} + \underbrace{\int_0^{r_n} d^3r \frac{G_{N(4+n)}}{r^{n+1}} \rho}_{\text{new piece.}}$$

$$\frac{\delta E_{\text{grav}}}{E_{\text{grav}}} \sim \frac{M_{\text{pl}}^2}{M_{\text{pl}}^{n+2} r_{\text{typ}}^{n-2} R^2} \lesssim \left(\frac{1 \text{ mm}}{R}\right)^2 \text{ irrelevant}$$

• Lab bounds

- The theory above $\sim \text{TeV}$ will produce higher-dimension operators suppressed by $\sim \text{TeV}$, constrained by compositeness searches which are at most $\sim 4 \text{ TeV}$ for electron compositeness. Since we don't know the theory above TeV we don't know the coefficients of these operators...

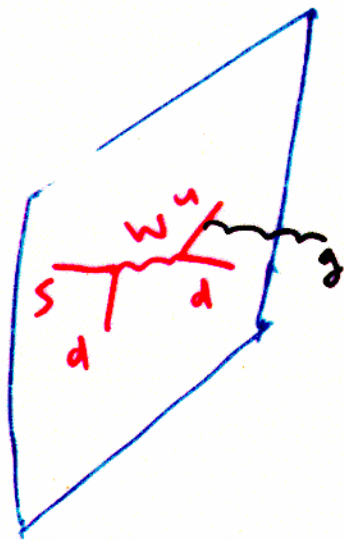


$$\rightarrow \mathcal{O} \sim \frac{E^2}{M_{\text{electron}}^4} (\bar{\psi}\psi)^2$$

Same even for

$$M_{\text{electron}} \lesssim 1 \text{ TeV}$$

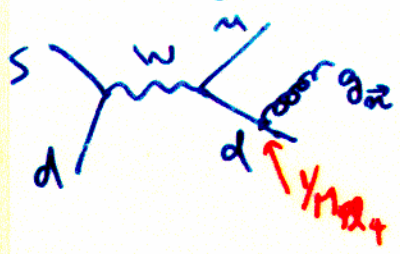
- Much more important: we have a massless $1/\text{TeV}$ coupled particle, the graviton.
- We know that other light particles (axions, familons), coupled with $1/\text{TeV}$, are a disaster from rare decays and stellar cooling in astrophysics — must check. $\rightarrow \frac{1}{10^{10} \text{ GeV}}$
- Start with a very rare decay:



$$K \rightarrow \pi \text{ graviton.}$$

[Note that since translation invariance is broken in extra dimensions, there is no \vec{p}_n conservation, but energy must still be conserved].

• $K \rightarrow \pi$ grav. from 4-d point of view



$$Br(K \rightarrow \pi g_{\pi}) \sim \left(\frac{m_K^2}{M_{Pl,4}^2} \right)$$

But, there are $\sim (m_K / 1/R)^n = (m_K R)^n$ energetically available KK modes, so

$$Br(K \rightarrow \pi g) \sim \left(\frac{m_K^2}{M_{Pl,4}^2} \right) \times (m_K R)^n \sim \left(\frac{m_K}{M_{Pl(4+n)}} \right)^{n+2}$$

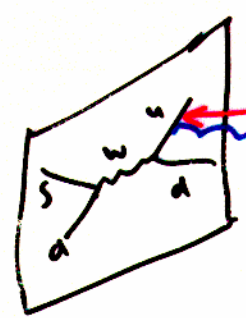
$\sim 10^{-12}$ even for $n=2$,
 $M_{Pl(4+n)} \sim 500 \text{ GeV}$
SAFE.

(similar FZ 10^{19} GeV)

• $K \rightarrow \pi$ grav. from $(4+n)$ d point of view:

$$g_{AB} = \gamma_{AB} + \frac{h_{AB}}{\sqrt{M_{Pl(4+n)}^{n+2}}}$$

← canonically normalized in $(4+n)$ d.



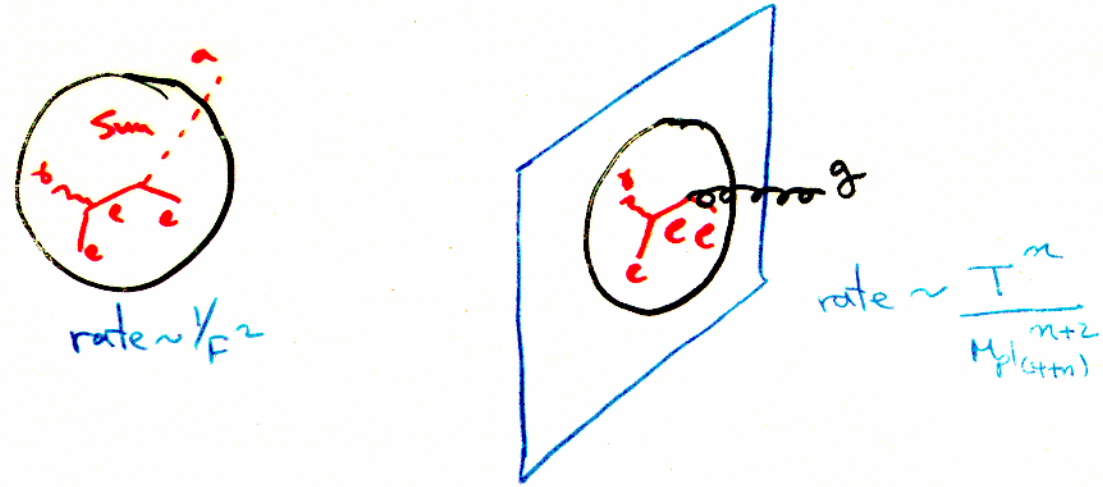
$\frac{1}{g_{Pl(4+n)}}$

$$\rightarrow \frac{1}{M_{Pl(4+n)}^{n+2}}$$

in the rate ✓

Astrophysical constraints

Analogous to axions, worry is that stars can lose energy too rapidly, by emitting gravitons into the extra dimensions:



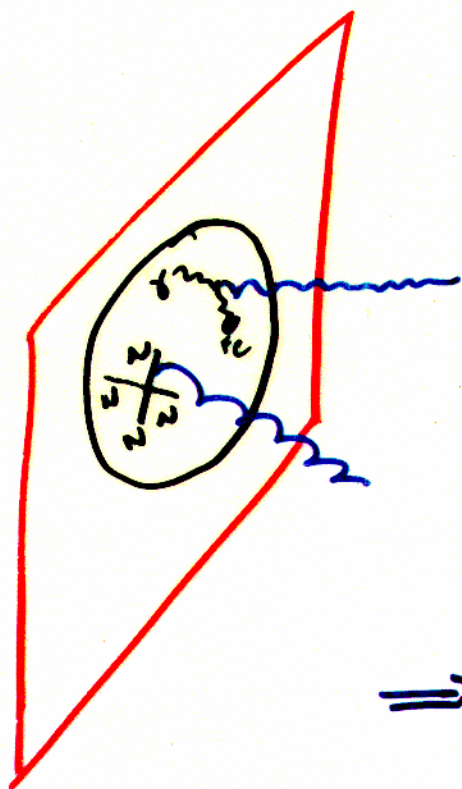
[Note: if $n=7$, the first KK excitation is @ $1/r \sim 100$ MeV, too heavy to be produced even in SN 1987A where $T_{SN} \sim 20-60$ MeV [$T_{\text{edgiant}} \sim 10$ keV, $T_{\text{sun}} \sim$ keV]]

So we have a very rough dictionary:

$$\frac{1}{r^2} \longleftrightarrow \frac{T^n}{M_{pl(4+n)}^{n+2}} \Rightarrow \begin{aligned} &\text{For } n \geq 2, \text{ everything OK for } M_{pl(4+n)} \gtrsim 1 \text{ TeV} \\ &\text{For } n \geq 2, \text{ need } M_{pl(4+n)} \gtrsim 10 \text{ TeV} \end{aligned}$$

($M_s \sim 10$ TeV)

Strongest bounds come from SN1987A:



$$\mathcal{L}_{\text{SN1987A}} \sim 10^{53} \text{ egr s}^{-1} \sim 10^{26} \text{ TeV}^2$$

$$\mathcal{L}_{\text{grav. (Primakoff)}} \sim 10^{57} \frac{T^{n+4}}{M_{\text{pl}}^{n+2}} < \mathcal{L}_{\text{SN}}$$

$$\Rightarrow M_6 \gtrsim 10 \text{ TeV} \quad (n=2)$$

$$< 1 \text{ TeV} \quad n > 2$$

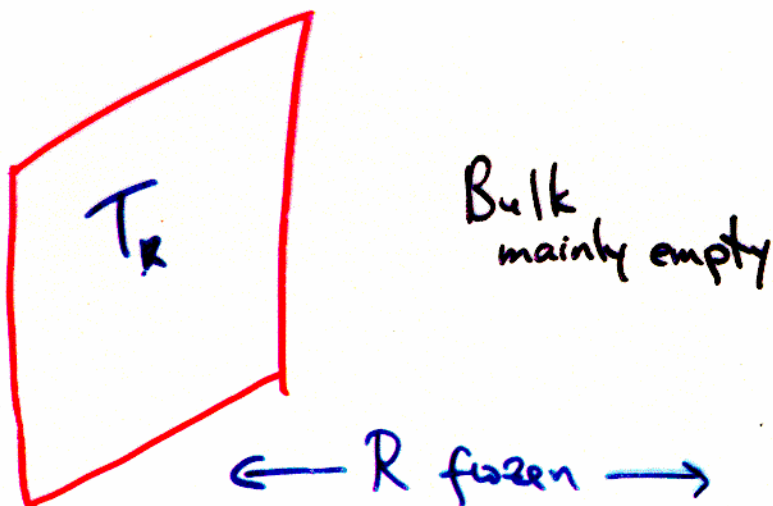
OK.

Cosmology.

Of course all of early-universe cosmology must be rethought in this framework.

As long as extra dimensions have stabilised @ temperatures $\sim 1 \text{ MeV}$ (by nucleosynthesis) the universe may be normal enough for everything to work.

Parametrise ignorance in following question: what is the highest temperature T_* at which the universe can be "normal"?



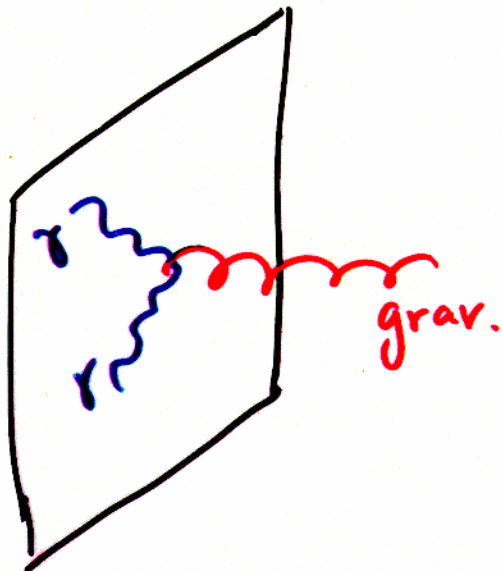
Cosmological Time Scales

Normal: Cooling by adiabatic expansion

$$\dot{T}/T \sim \dot{R}/R = H \sim T^2/M_{pl} \equiv 1/t_H.$$

dominates at late times

New: Cooling by evaporation



$$\dot{T}/T \sim \frac{T^{n+3}}{M_{(4+n)}^{n+2}} \equiv 1/t_{ev.}$$

dominates at early times

For normalcy, we want $t_{ev} \gg t_H$

$$\Rightarrow T_* \lesssim 10^{\frac{6n-9}{n+1}} \text{ MeV}$$

$$\left(\underset{n=2}{10 \text{ MeV}} - \underset{n=6}{10 \text{ GeV}} \right)$$

$$(M_{4+n} \sim \text{TeV})$$

- Are we "pushing the limits of parameter space"?
- Experimentally: Yes! Theory is testable, with all sorts of new physics at dangerously low energies.
- Theoretically: No. All of our scales are motivated by solving the hierarchy problem.

old string scenario

↓
susy
 $\sim 10^3 \text{ GeV}$

↓
string
scale
 $\sim 10^{18} \text{ GeV}$

New scenario

↓
susy
@ string
scale
 $\sim 10^3 \text{ GeV}$ large
radii

Experimental Signals

- Sub-mm measurements of gravitational strength forces in next 1-2 years, perhaps down to ~ 1 micron.
 \Rightarrow Could observe transition from $\gamma_{r=2} \rightarrow \gamma_{r=4}$ Newtonian gravity for $n=2$ new dimension.
- For any n , in a string scenario, there could be light particles of mass $\sim \frac{M_{EW}^2}{M_{Pl,4}}$ $\sim (1\text{mm})^{-1}$ mediating observable grav. strength Yukawa forces at sub-mm distances.
- For any n , repulsive forces $10^6 - 10^8 \times$ gravity, from gauge fields in the bulk coupled to lin. comb. of $B, L \#$.

Quantum Gravity @ LHC, NLC

- Since $M_{Pl(4+n)} \sim TeV$, LHC + NLC are Planck-scale machines, and could probe quantum-gravity.

- Example: Strings @ a TeV.



New phenomena include: production of higher string vibrational modes (Regge excitation) for every SM particle.

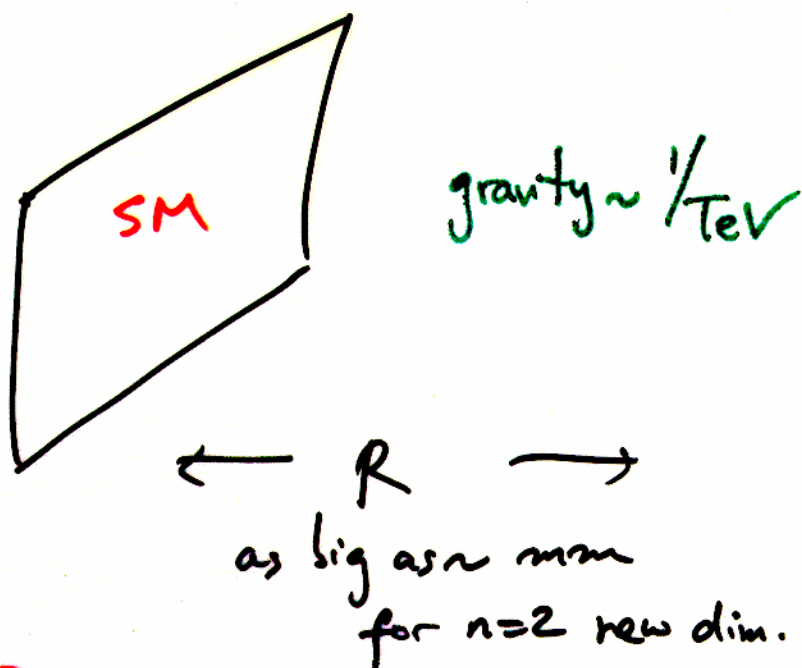
e.g.



Spin 2 "heavy gluon".

Conclusions

- There is a new framework for solving the hierarchy problem which does not rely on SUSY or technicolor.



- Raises exciting prospect that LHC, NLC are quant. gravity machines, and that new submm forces may soon be discovered.
- Many open questions (origin of R ? Unification? Early universe?)