

Non-local Symmetry of the 3-Dimensional Burgers-Type Equation

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Abstract

Non-local transformation, which connects the 3-dimensional Burgers-type equation with a linear heat equation, is constructed. Via this transformation, nonlinear superposition formulae for solutions are obtained and the conditional non-local symmetry of this equation is studied.

The multidimensional generalization of the Burgers equation

$$L_1(u) = u_0 - u|\nabla u| - \Delta u = 0, \quad (1)$$

is called further the Burgers-type equation. This equation was suggested by W. Fushchych in [1]. We use here such notations:

$$\begin{aligned} \partial_\mu u &= \frac{\partial u}{\partial x_\mu}, & \{x_\mu\} &= (x_0, x_1, \dots, x_{n-1}), \\ \nabla u &= \|u_1, u_2, \dots, u_n\|^T, & (\mu &= 0, n-1), \\ |\nabla u| &= \sqrt{(\nabla u)^2}, & \Delta &= (\nabla)^2 = \partial_1^2 + \partial_2^2 + \dots + \partial_{n-1}^2. \end{aligned}$$

In the present paper, we construct the non-local transformation, which connects the 3-dimensional equation (1) with a linear heat equation. Via this transformation, we obtain nonlinear superposition formulae for solutions of equation (1). Also we investigate the conditional non-local symmetry of equation (1) and obtain formulae generating solutions of this equation.

1. Conditional non-local superposition

Let us consider the 3-dimensional scalar heat equation

$$L_2(w) = w_0 - \Delta w = 0. \quad (2)$$

For the vector-function H , such equations are fulfilled:

$$H = 2\nabla \ln w, \quad H = \|h^1, h^2, h^3\|^T, \quad (3)$$

$$\frac{1}{2}(H)^2 + (\nabla \cdot H) = \partial_0 \ln w, \quad \nabla \times H = 0. \quad (4)$$

From the integrability condition for equations (3), (4), it follows that equation (2) is connected with the vector equations

$$H_0 - \frac{1}{2}(H)^2 - \nabla(\nabla \cdot H) = 0, \quad \nabla \times H = 0. \tag{5}$$

Let $|H|^2 = (h^1)^2 + (h^2)^2 + (h^3)^2 = u^2$ for $H = \theta \cdot u$, where $|\theta| = 1$, $\theta = |\theta^1, \theta^2, \theta^3|^T$. Then we obtain from (5) that $\nabla \times \theta = 0$ and the equality

$$\theta [u_0 - u|\nabla u| - \Delta u] = u [\theta_0 - 2\nabla \ln u(\nabla \cdot \theta) - \nabla(\nabla \cdot \theta)].$$

Let relations

$$L_1(u) = u_0 - u|\nabla u| - \Delta u = 0, \tag{6}$$

$$\theta_0 - 2\nabla \ln u(\nabla \cdot \theta) - \nabla(\nabla \cdot \theta) = 0, \tag{7}$$

$$\nabla \times \theta = 0.$$

be fulfilled on the some subset of solutions of equation (5). System (3), (4) in new variables, which connect equations (2) and (6), has the form

$$\nabla \ln w = \frac{1}{2}\theta u, \tag{8}$$

$$\partial_0 \ln w = \frac{1}{4}u^2 + \frac{1}{2}(\nabla \cdot \theta u). \tag{9}$$

So via transformations (8), (9), PDE (6) is reduced to the linear equation (2). The corresponding generalization of the Cole-Hopf transformation (substitution) is obtained in the form

$$u = 2\sqrt{(\nabla \ln w)^2}. \tag{10}$$

Theorem 1. *The non-local substitution (10) is a linearization of the 3-dimensional Burgers-type equation (1), and the operator equality*

$$\partial_0(\nabla \ln w) - 2(\nabla \ln w) \cdot \Delta \ln w - \nabla(\Delta \ln w) = [w^{-1}\nabla - w^{-2} \cdot \nabla w] \cdot (w_0 - \Delta w) = 0$$

is fulfilled.

Example 1. The function

$$u = 2\sqrt{[\nabla \ln \varphi(x_0, \omega)]^2}, \quad \omega = \alpha \cdot x = \alpha_a x_a, \quad (a = 1, 2, 3), \tag{11}$$

is a non-Lie solution of the nonlinear equation (1). Let us substitute this non-local ansatz into equation (1). We find the condition on φ :

$$[\varphi^{-1} \cdot \partial_\omega - \varphi^{-2} \cdot \varphi_\omega] \cdot [\varphi_0 - \varphi_{\omega\omega}] = 0. \tag{12}$$

So, we obtain the non-local reduction of equation (1) to the form (12). Analogously, with the ansatz

$$u = 2\sqrt{[\nabla \ln \{2(n-1)x_0 + x\}]^2}, \quad x \equiv (x_1, x_2, x_3), \quad (n = 4), \tag{13}$$

one obtains

$$[w^{-1}\nabla - w^{-2}\nabla w] \cdot [\partial_0 - \Delta] \cdot \{2(n-1)x_0 + x\}. \quad (14)$$

Here, $w \equiv 2(n-1)x_0 + x$.

The linearization of the differential equation (1) makes it possible via a linear superposition of solutions of equation (2) and the non-local transformation (8), (9) to construct the principle of nonlinear superposition for solutions of equation (1).

Theorem 2. *The superposition formula for a subset of solutions ${}^{(k)}u(x)$, ($k = 1, 2$) of equation (1) has the form*

$$\begin{aligned} {}^{(3)}u(x) &= 2 \left[\frac{{}^{(1)(1)(1)}\tau u \theta + {}^{(2)(2)(2)}\tau u \theta}{{}^{(1)}\tau + {}^{(2)}\tau} \right], \\ {}^{(3)(3)}\theta u(x) &= 2 \left({}^{(1)}\tau + {}^{(2)}\tau \right)^{-1} \left({}^{(1)(1)(1)}\tau u \theta + {}^{(2)(2)(2)}\tau u \theta \right), \\ {}^{(k)}\tau_0 &= \frac{1}{2} \frac{{}^{(k)}\tau}{\tau} \left[u^2 + (\nabla \cdot \theta) u \right], \\ \nabla \frac{{}^{(k)}\tau}{\tau} &= \frac{1}{2} \frac{{}^{(k)}\tau u \theta}{{}^{(k)}\tau} \cdot \nabla \times \theta, \\ \partial_0 \ln \left(\frac{{}^{(1)}\tau + {}^{(2)}\tau}{\tau} \right) &= \left[\frac{{}^{(1)(1)(1)}\tau u \theta + {}^{(2)(2)(2)}\tau u \theta}{{}^{(1)}\tau + {}^{(2)}\tau} \right] + \left[\nabla \frac{{}^{(1)(1)(1)}\tau u \theta + {}^{(2)(2)(2)}\tau u \theta}{{}^{(1)}\tau + {}^{(2)}\tau} \right]. \end{aligned} \quad (15)$$

Here, ${}^{(k)}u$, ($k = 1, 2$) are known solutions of equation (1), ${}^{(s)}u$, ($s = 1, 2, 3$) and ${}^{(k)}\tau$, ($k = 1, 2$) are functional parameters, and $\frac{{}^{(k)}\tau}{\tau} - \Delta \frac{{}^{(k)}\tau}{\tau} = 0$.

2. Conditional non-local invariance

Let us consider the potential hydrodynamic-type system in $R(1, 3)$ of independent variables

$$H_0 + (H, \nabla)H - \Delta H = 0, \quad \nabla \times H = 0, \quad (16)$$

or, in the form (5),

$$H_0 + \frac{1}{2}\nabla(H)^2 - \nabla(\nabla \cdot H) = 0, \quad \nabla \times H = 0. \quad (17)$$

Let $|H| = u$ or, in other words

$$|H|^2 = (H^1)^2 + (H^2)^2 + (H^3)^2 = u^2.$$

Then we have such equalities:

$$H = \theta \cdot |H| = \theta \cdot u, \quad (\theta = H \cdot |H|^{-1}). \quad (18)$$

So, θ is a unit vector collinear with H . Let $\nabla \times \theta = 0$. It makes one possible to obtain

$$\begin{aligned} [\nabla \times \theta u] &= u \cdot [\nabla \times \theta] + [\nabla u \times \theta] \longrightarrow \nabla u \times \theta = 0, \\ \nabla u &= \theta \cdot |\nabla u|, \quad (\nabla \cdot \theta u) = u(\nabla \cdot \theta) + \theta \cdot \nabla u, \\ \nabla(\nabla \cdot \theta u) &= u \cdot \nabla(\nabla \cdot \theta) + 2\nabla u \cdot (\nabla \cdot \theta) + \theta \Delta u. \end{aligned} \quad (19)$$

Substituting H of the form (18) into equation (17), we find the operator equation

$$\theta[u_{0+}u|\nabla u| - \Delta u] = -u[\theta_0 - 2\nabla \ln u(\nabla \cdot \theta) - \nabla(\nabla \cdot \theta)].$$

Let now pick out the subset of solutions of this equation, which consists of solutions of the system

$$u_{0+}u|\nabla u| - \Delta u = 0, \quad (20)$$

$$\theta_0 - 2\nabla \ln u(\nabla \cdot \theta) - \nabla(\nabla \cdot \theta) = 0. \quad (21)$$

Variables $\theta = \|\theta^1, \theta^2, \theta^3\|^T$ are playing the role of supplementary parameters. Notice that, for a known solution u of equation (20), equation (21) is linear.

Let us take another copy of equation (17):

$$Q_0 + \frac{1}{2}\nabla(Q)^2 - \nabla(\nabla \cdot Q) = 0, \quad \nabla \times Q = 0. \quad (22)$$

Let $Q = \tau \cdot w$, where $w = |Q|$ and $\nabla \times \tau = 0$. The differential equation (22) can be substituted now by the system

$$w_0 + w|\nabla w| - \Delta w = 0, \quad (23)$$

$$\tau_0 - 2\nabla \ln w \cdot (\nabla \cdot \tau) - \nabla(\nabla \cdot \tau) = 0. \quad (24)$$

Assume now such equalities:

$$\begin{aligned} -2\nabla \ln w &= H - Q, \\ -\partial_0 \ln w &= \frac{1}{2}(H \cdot Q) - \left[\frac{1}{2}(H)^2 - (\nabla \cdot H) \right]. \end{aligned} \quad (25)$$

Excluding w from this system by cross-differentiation, we obtain

$$H_0 + \frac{1}{2}\nabla(H)^2 - \nabla(\nabla \cdot H) = Q_0 + \frac{1}{2}\nabla(H \cdot Q). \quad (26)$$

So, with the condition

$$Q_0 + \frac{1}{2}\nabla(H \cdot Q) = 0, \quad (27)$$

we get equation (17). Let

$$\frac{1}{2}(H \cdot Q) = \frac{1}{2}(Q)^2 - (\nabla \cdot Q). \quad (28)$$

Then (27) becomes equation (22). As one can see, condition (28) is the necessary condition for reducing equations (20), (21) via the non-local transformation (25), (28) to system (23), (24). Substituting

$$H = \tau \cdot w - 2\nabla \ln w$$

into (28), one obtains

$$w_0 + w \cdot |\nabla w| - \Delta w = \frac{1}{2}w\{|\nabla w| - (\nabla \cdot Q)\}.$$

Let us put the condition

$$|\nabla w| = (\nabla \cdot Q). \quad (29)$$

Since

$$(\nabla \cdot \tau w) = w(\nabla \cdot \tau) + |\nabla w|,$$

it follows from condition (29) that

$$(\nabla \tau) = -4|\nabla \ln w|.$$

Introducing notations $u \equiv \overset{(2)}{u}$, $w \equiv \overset{(1)}{u}$, we can formulate such a theorem:

Theorem 3. Let $\overset{(1)}{u}$ be a known solution of equation (1)

$$u_0 + u \cdot |\nabla u| - \Delta u = 0.$$

Then its new solution $\overset{(2)}{u}$ is defined by the formulae

$$\overset{(2)}{u} = (\theta \cdot \tau) \overset{(1)}{u} - 2 \cdot \theta \cdot \nabla \ln \overset{(1)}{u},$$

$$\frac{1}{2}(\theta \cdot \tau) \overset{(1)}{u} \overset{(2)}{u} + 2\theta_0 \ln \overset{(1)}{u} = \frac{1}{2} \overset{(2)}{u}^2 - (\nabla \cdot \theta \overset{(2)}{u}),$$

$$\theta \overset{(2)}{u} = \tau \overset{(1)}{u} - 2 \cdot \nabla \ln \overset{(1)}{u},$$

$$\theta_0 - 2 \cdot \nabla \ln \overset{(2)}{u} (\nabla \cdot \theta) - \nabla(\nabla \cdot \theta) = 0,$$

$$\tau_0 - 2 \cdot \nabla \ln \overset{(1)}{u} (\nabla \cdot \tau) - \nabla(\nabla \cdot \tau) = 0,$$

$$\nabla \times \theta = 0, \quad \nabla \times \tau = 0, \quad (\nabla \cdot \tau) = -4|\nabla \ln w|.$$

The last equalities are additional conditions for the non-local invariance of equation (1) with respect to the non-local transformation (25), (26).

References

- [1] Fushchych W.I., Symmetry in mathematical physics problems, in: Algebraic-Theoretical Studies in Mathematical Physics, Inst. of Math. Acad. of Sci. Ukraine, Kyiv, 1981, 6–28.