

# LHC/ILC

Hitoshi Murayama (Berkeley)  
SLAC SSI, 7/27/2006

# Technicolor



Lykken: *“It doesn’t look good but is not going away”*



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# Outline

- $e^+e^-$  Linear Collider
- Reconstruction of the Lagrangian
  - mass, spin, quantum numbers, mixing, couplings
  - use supersymmetry as an example
- Physics Significance

$e^+e^-$  Linear Collider



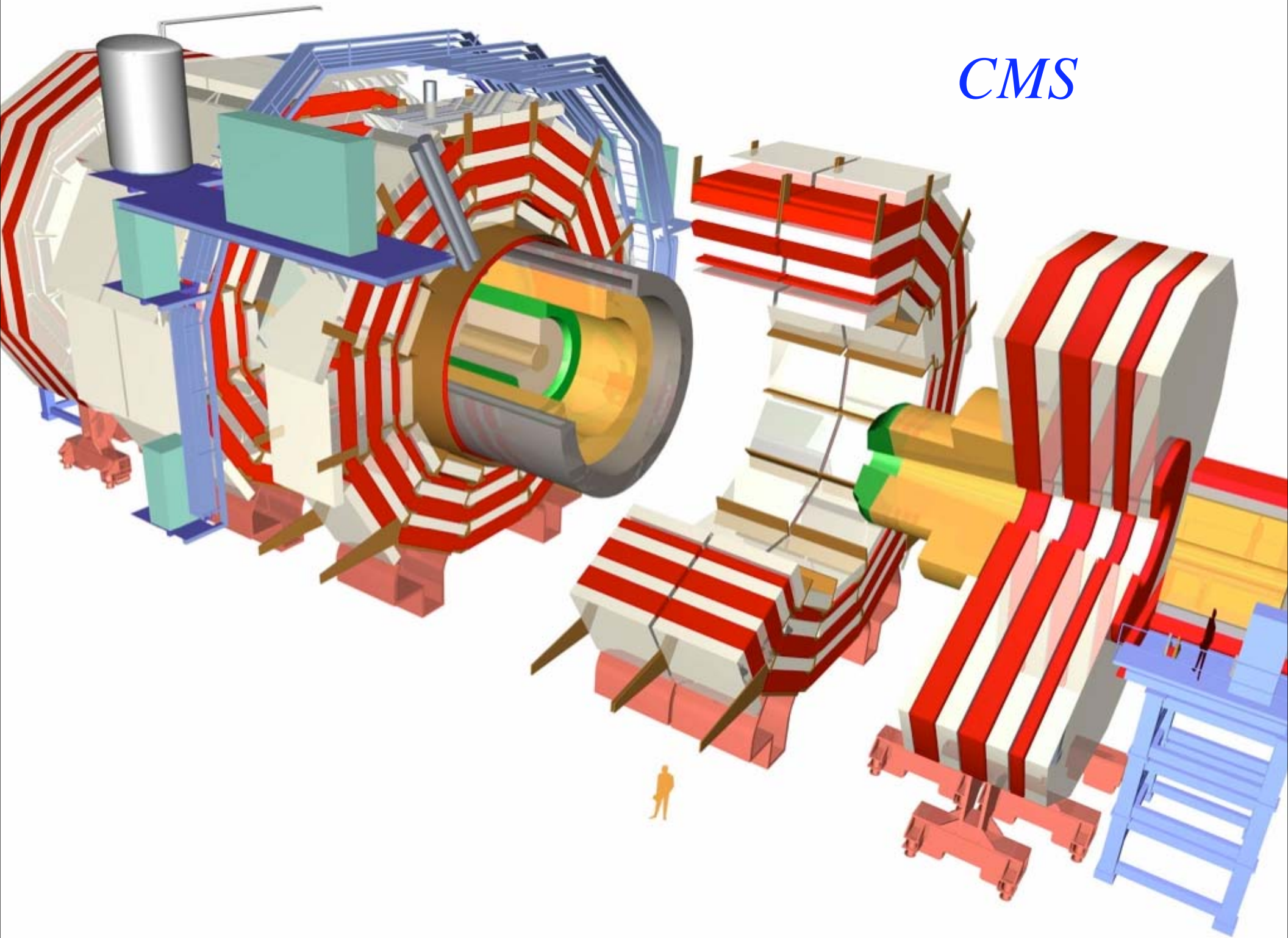
# Large Hadron Collider (LHC)

- proton-proton collider
- 14TeV energy  
(cf. 2TeV @ Fermilab)
- Under construction at  
CERN, Geneva
- Turn on in 2007
- Finally reaching the  
energy Fermi told us  
about in 1933!

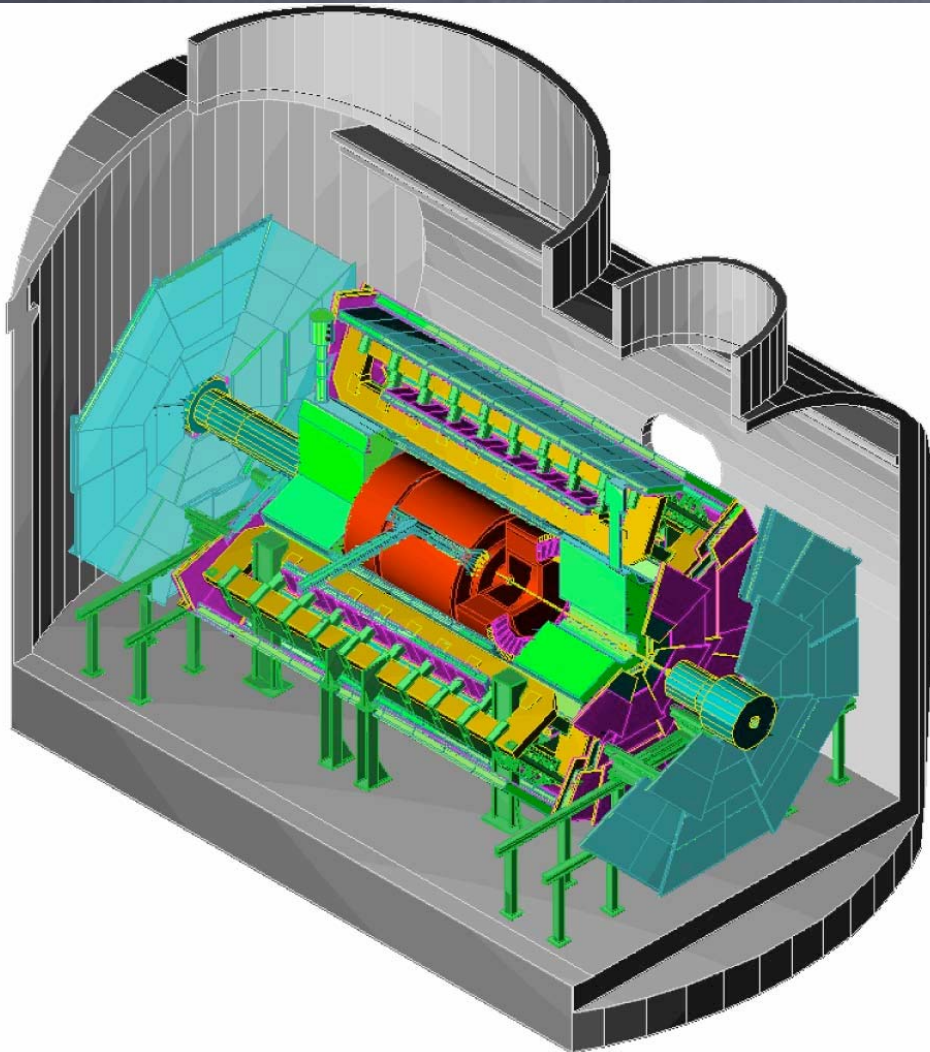




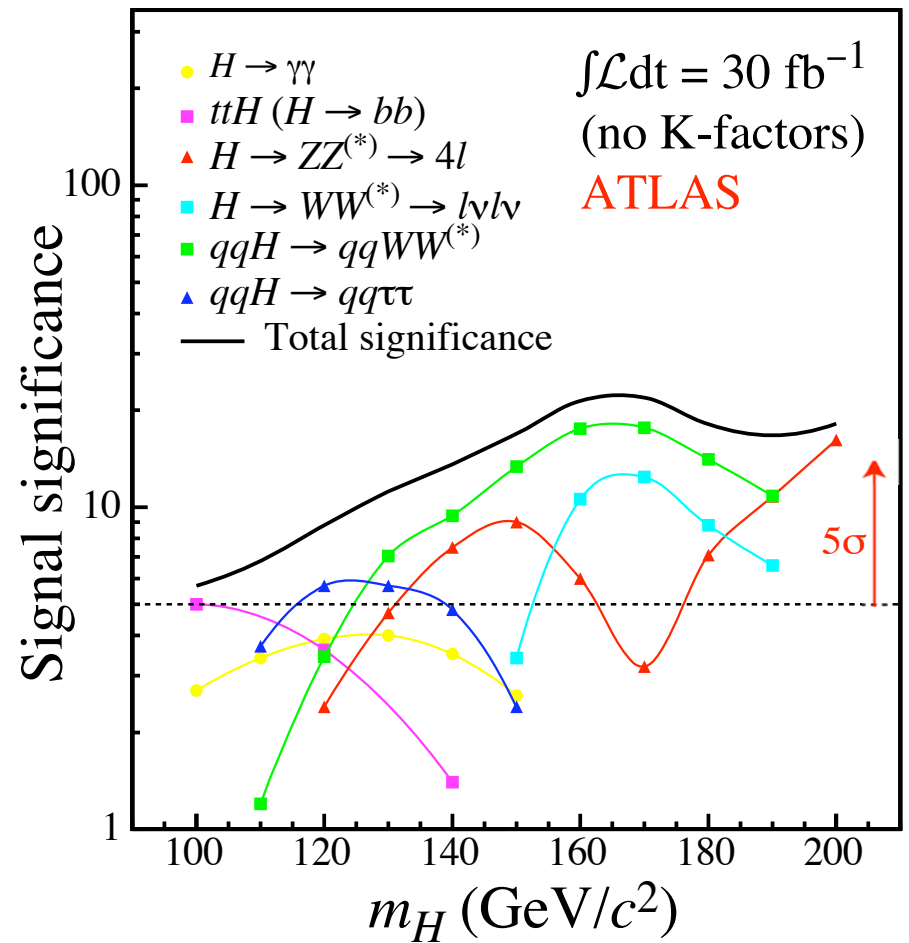
*CMS*



# Higgs at ATLAS



## Robust discovery

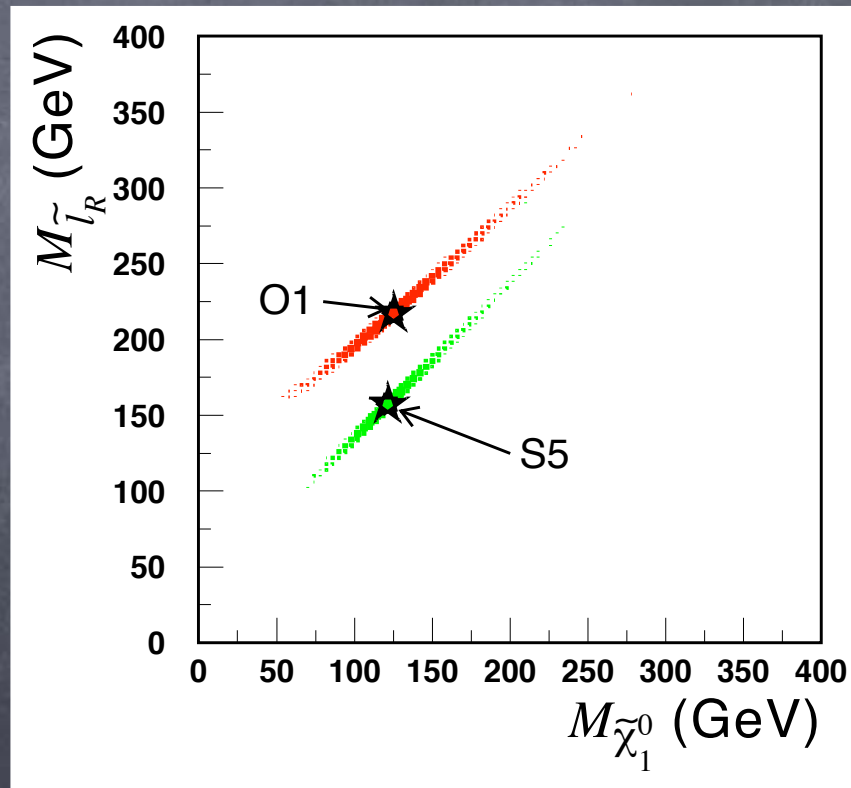
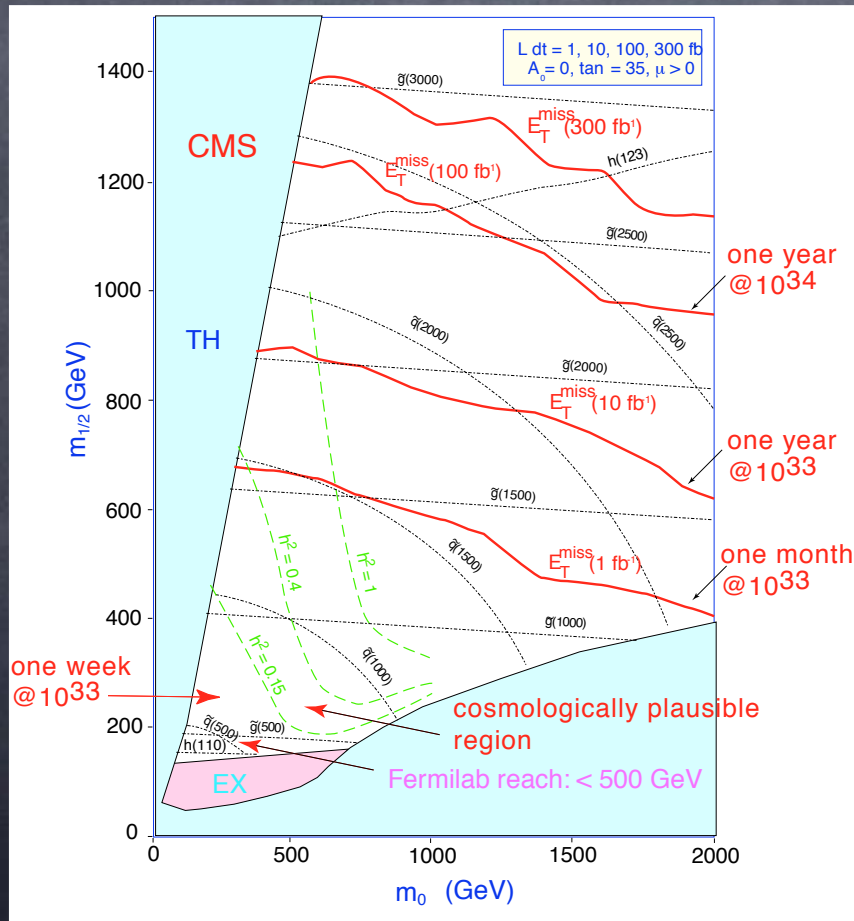




# Supersymmetry

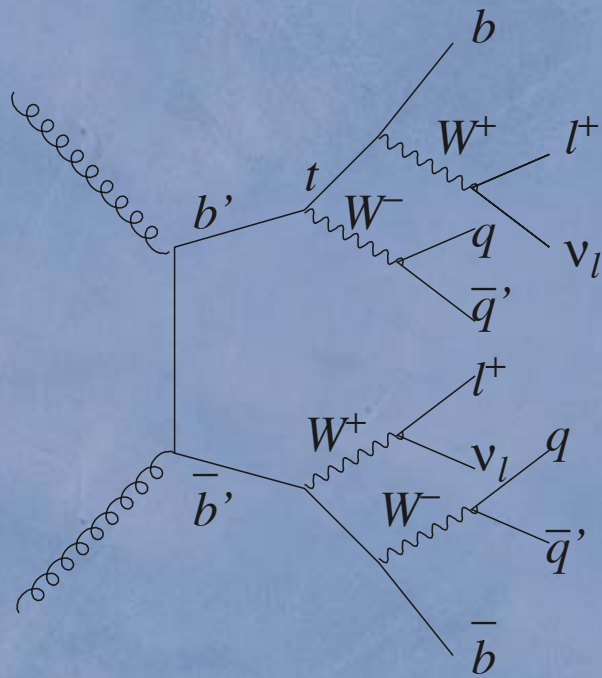
Tevatron/LHC will discover

Can do many precision measurements at LHC

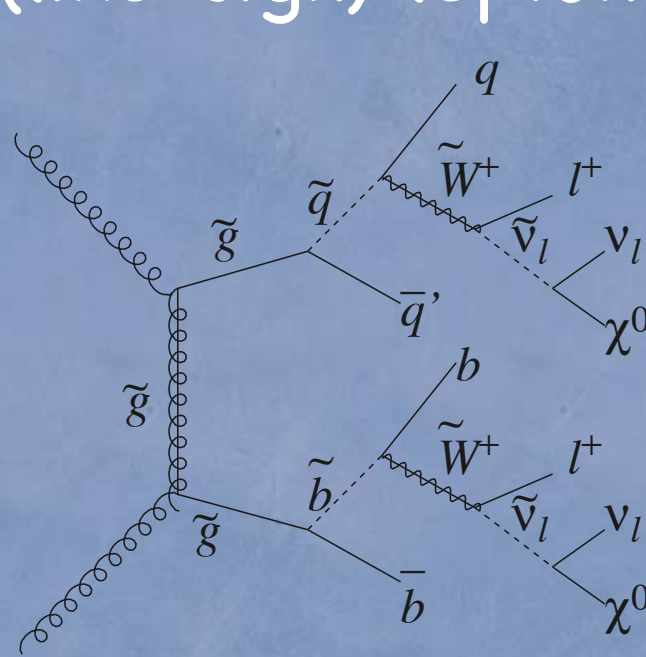


# New physics looks alike

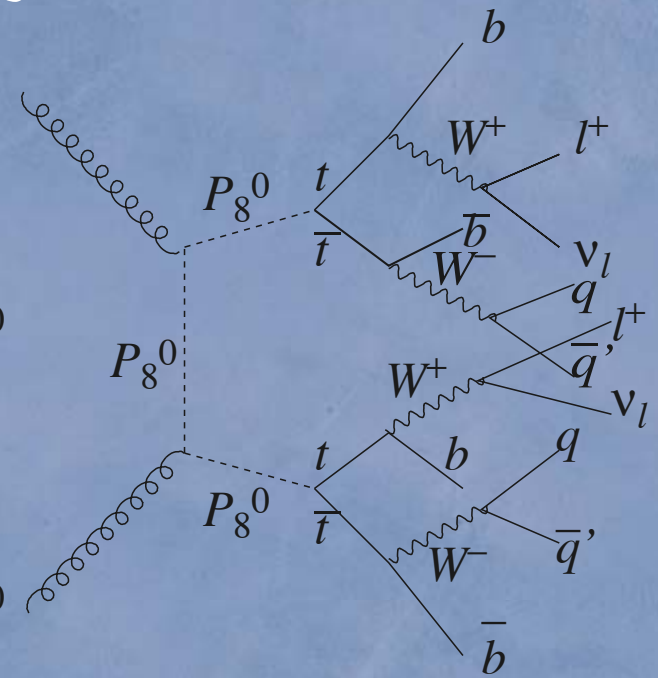
missing ET, multiple jets, b-jets,  
(like-sign) leptons



4th generation



SUSY



technicolor

+Universal extra dimension, little Higgs with T-parity



TC-TC composite Higgs  
hypercolor  
supercolor  
techni-GIM  
extended TC  
pseudo TC

NOT YET THOUGHT OF

effective SUSY  
SUSY  
MSSM + VR  
NMSSM  
unified SM  
axigloun

THOUGHT OF

NOT YET

6th gen  
5th gen  
4th gen

sterile  $\nu$   
heavy Majorana  
lepto quark  
fractionally charged

$Z_1, Z_2, Z_3, Z_4$   
 $Z_5, Z_6, Z_7, Z_8$   
 $X, Y$   
milli-charged  
mono-pole

vector-like family

S.T.U

shadow matter  
symmetry

NOT YET THOUGHT OF

highly charged doubly

triplet Higgs  
general 2HDSM Type 2  
spontaneous CP  
superweak

shadow matter  
symmetry

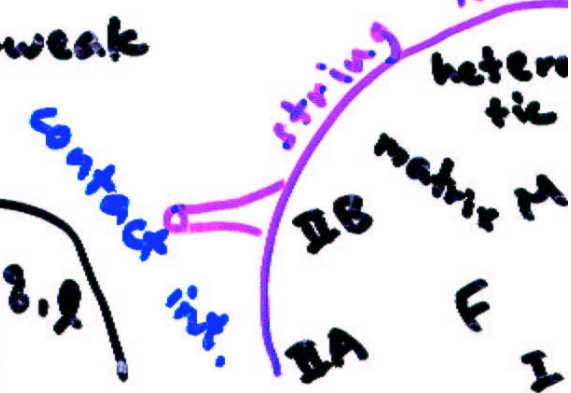
Majoron  
axion  
familon  
NGB

Weinberg's 3HD  
milli-weak

$N=2$   
 $N=4$   
 $N=8$



quintessence  
k-essence  
composite w, z



# Task

*Why do we live in a cosmic superconductor?*

- We can eliminate many possibilities at LHC
- But new interpretations necessarily emerge
- Race will be on:
  - **theorists** coming up with new interpretations
  - **experimentalists** excluding new interpretations
  - ⇒ A loooong process of elimination
- Crucial information is **in details**
  - ⇒ Reconstruct the theory from measurements



# The New York Times

July 23, 2008

## The Other Half of the World Discovered

Geneva, Switzerland

As an example, supersymmetry

“New-York Times level” confidence

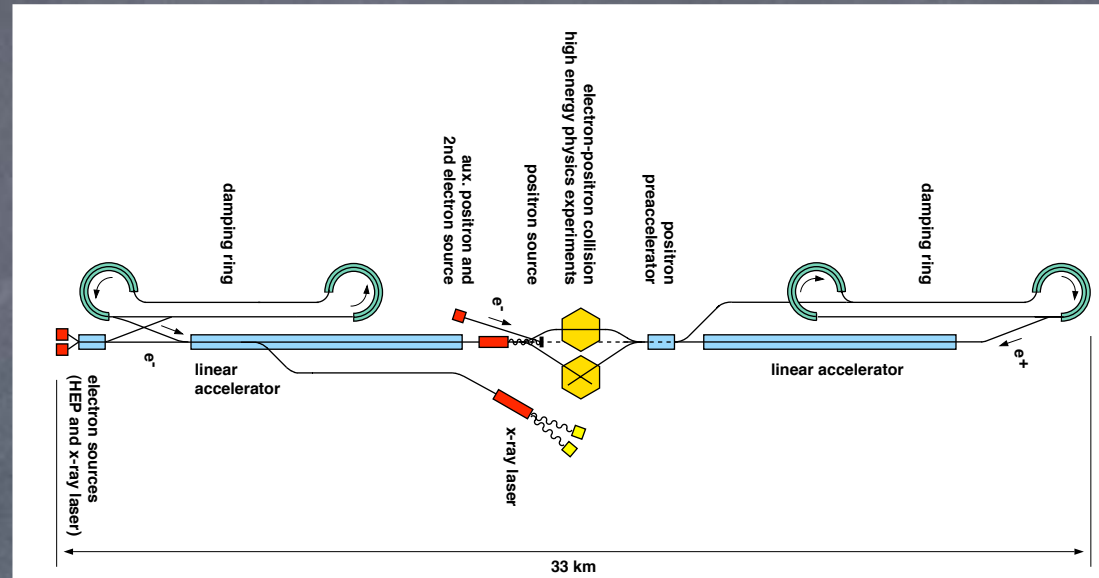
still a long way to

“Halliday-Resnick” level confidence

“We have learned that all particles we observe have unique partners of different spin and statistics, called superpartners, that make our theory of elementary particles valid to small distances.”

# Linear Collider

- Electron-positron collider
- $e^-$ ,  $e^+$  point-like with no structure
- Well-understood environment
- Linear instead of ring to avoid synchrotron loss
- Super-high-tech machine
- Accelerate the beam over >15km
- Focus beam down to a few **nanometers** and make them collide



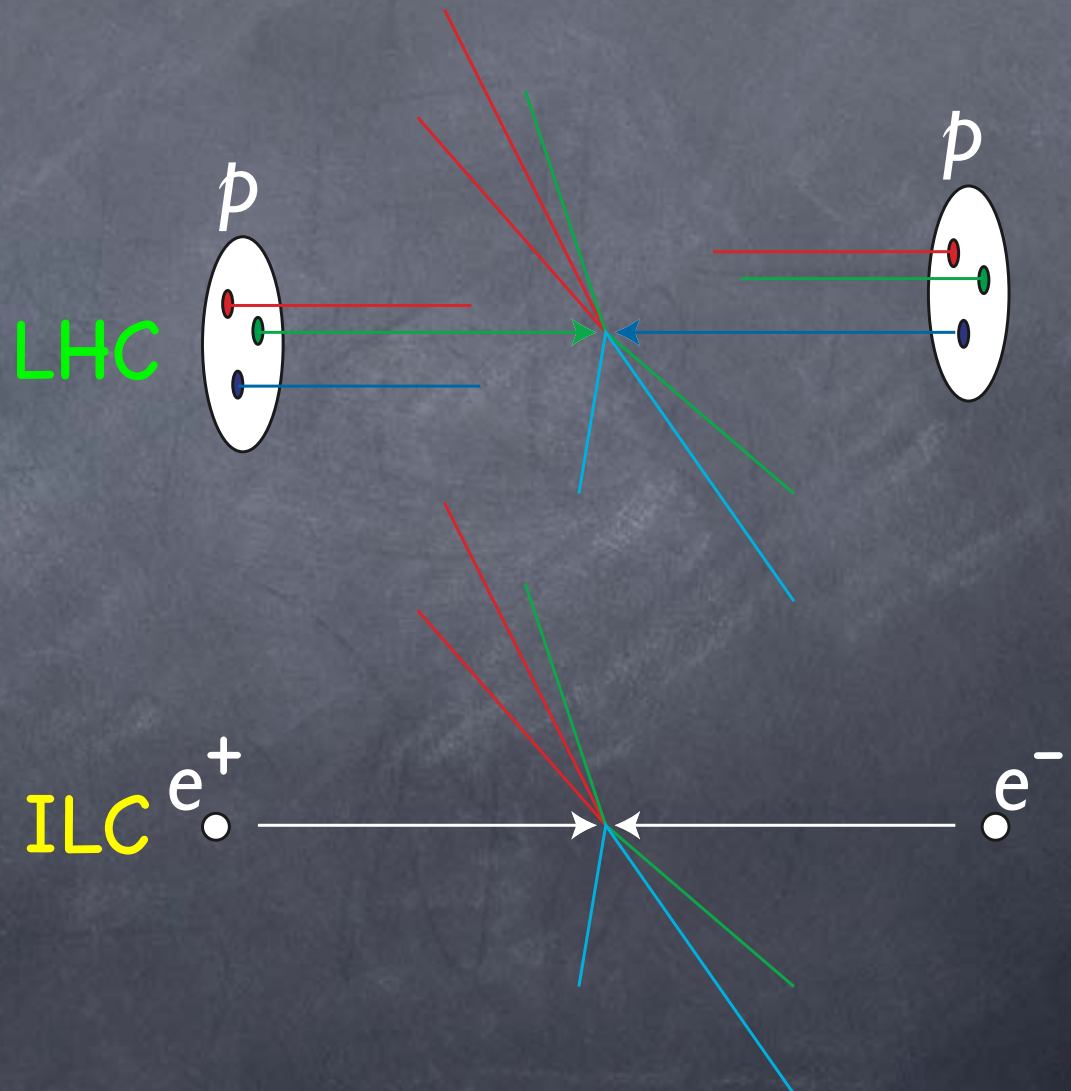
International Linear Collider (ILC)





# ILC

- elementary particles
- well-defined energy, angular momentum
- uses its full energy
- can produce particles democratically
- can capture nearly full information



# LHC vs ILC

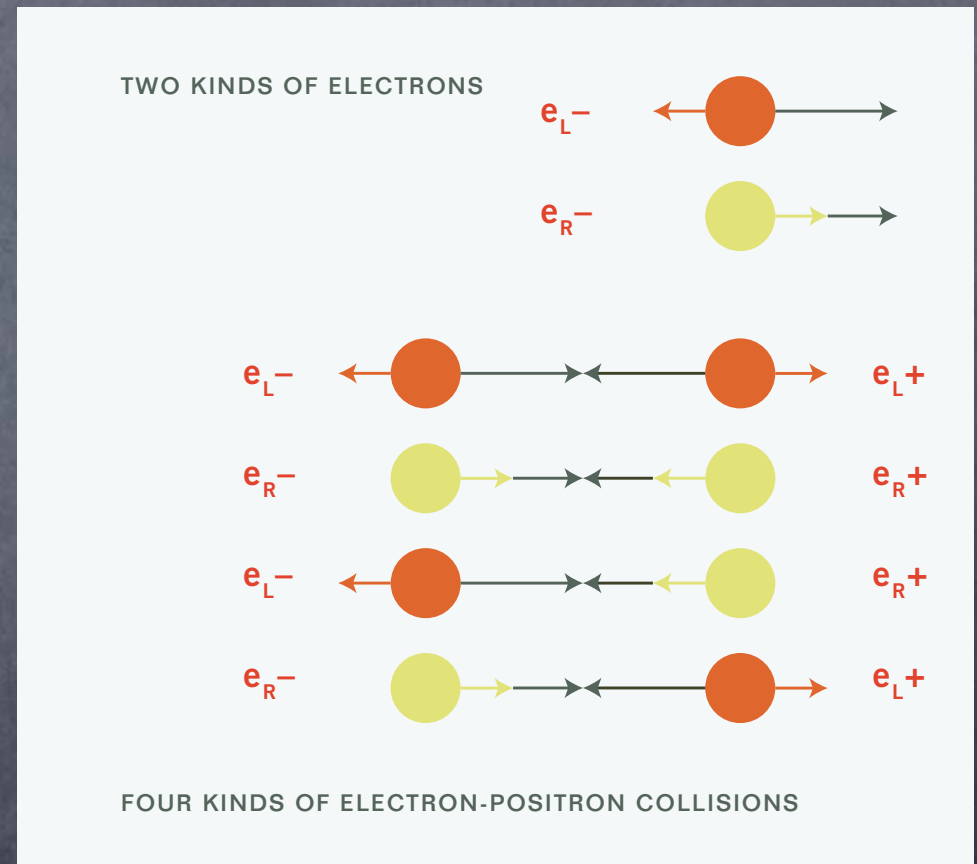
(oversimplified)

total energy	14TeV	0.5-1 TeV
usable energy	a fraction	full
beam	proton (composite)	electron (point-like)
signal rate	high	low
noise rate	very high	low
analysis	specific modes	nearly all modes
events	lose info along the beams	capture the whole
status	NEARLY READY!	needs to finish design



# Polarized beams

- serves as two different machines
- $e^+ e_R^-$
- $e^+ e_L^-$



# Polarized beams

• serves as two (nine?)  
different machines

•  $e_L^+ e_R^-$

•  $e_R^+ e_L^-$

•  $e_L^+ e_L^-$

•  $e_R^+ e_R^-$

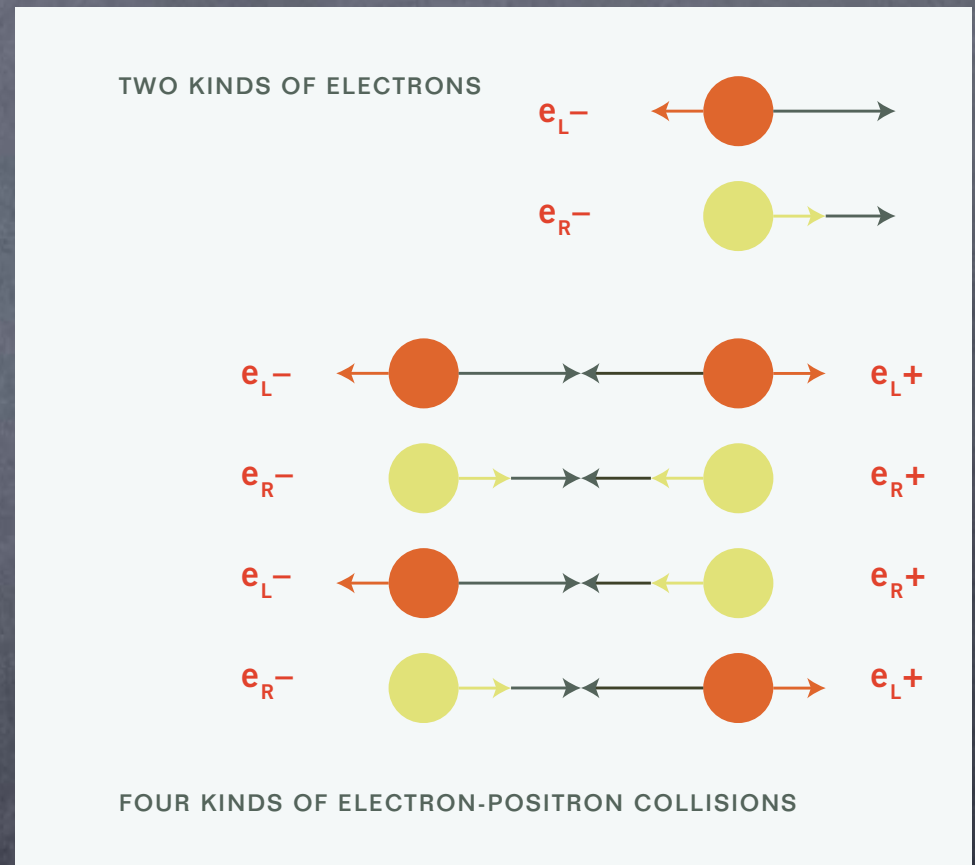
•  $e^- \gamma$

•  $\gamma \gamma$

•  $e_L^- e_R^-$

•  $e_L^- e_L^-$

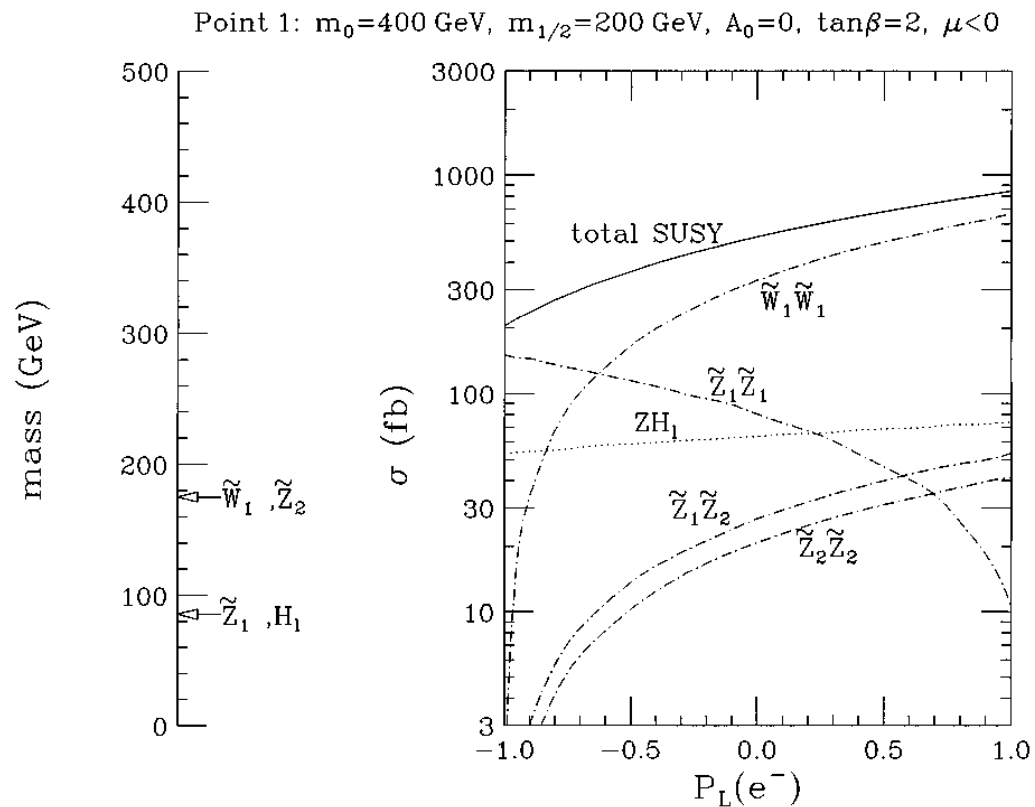
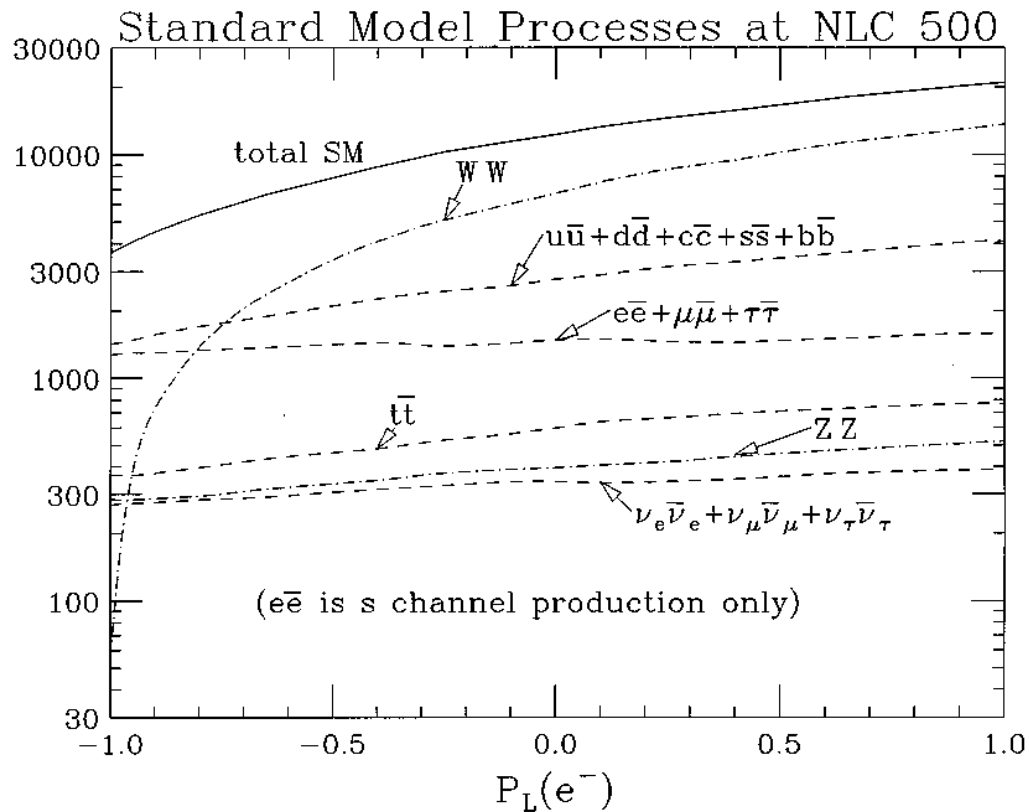
•  $e_R^- e_R^-$





# Polarized beams

- $e_L$  and  $e_R$  are really different particles at  $E \gg m_Z$



# Reconstruct Lagrangian from data

- Specify the fields
  - mass
  - spin  $\Rightarrow$  Klein-Gordon, Dirac, Majorana, gauge
  - $SU(3) \times SU(2) \times U(1)$  quantum numbers
  - mixing of states
- Specify their interactions
  - $SU(3) \times SU(2) \times U(1)$  quantum numbers determine gauge interactions
  - Yukawa couplings

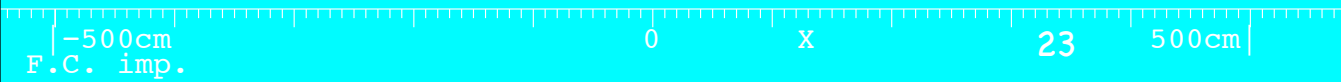
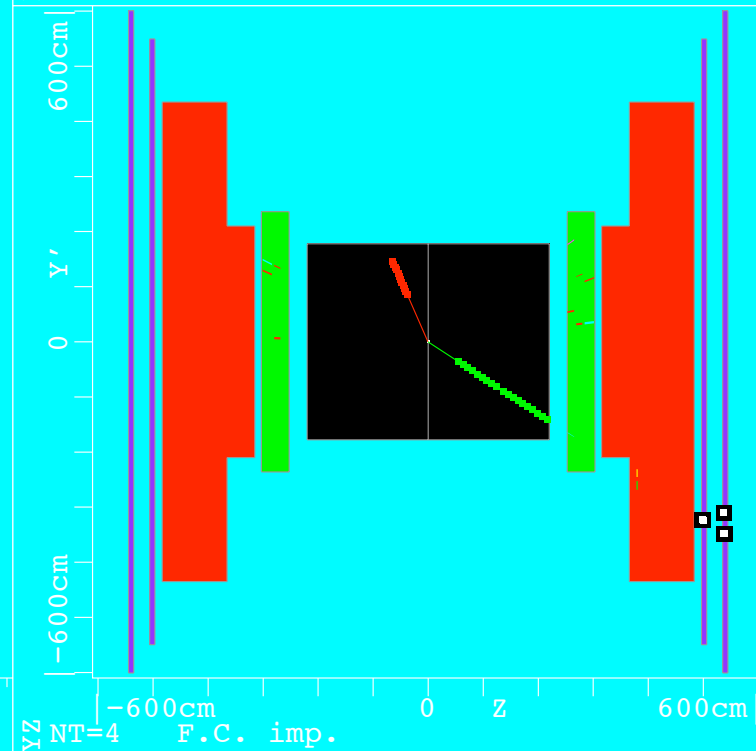
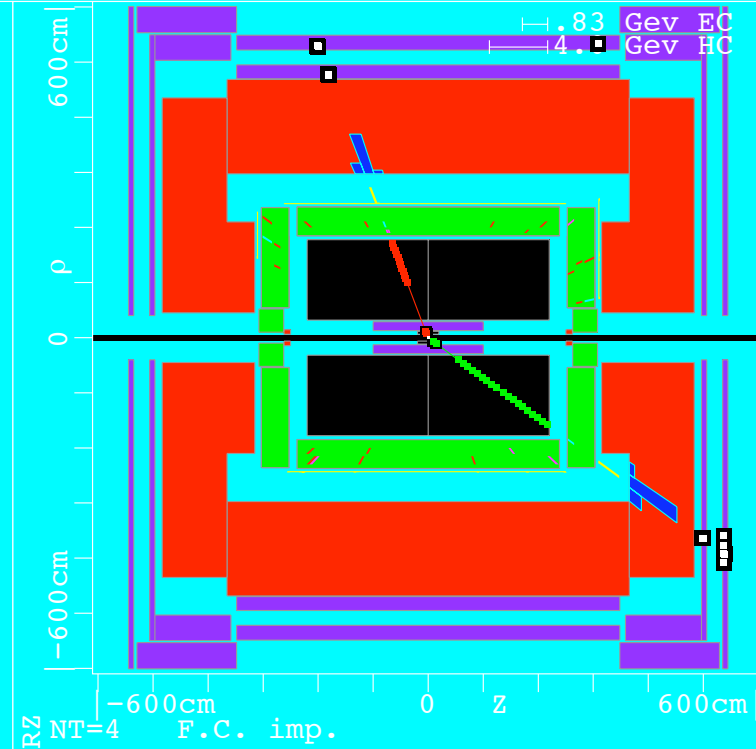
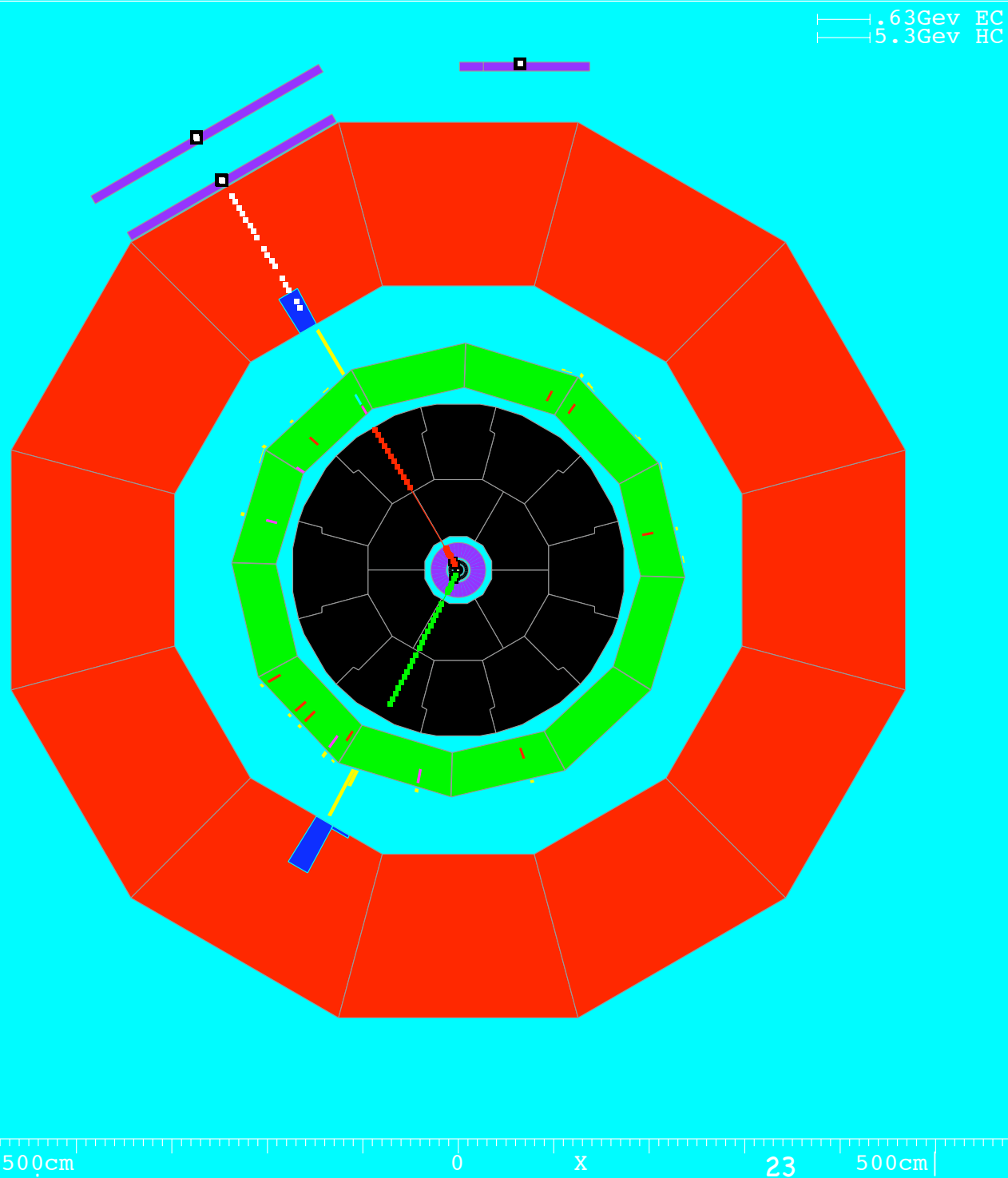


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- Specify the fields
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mass





# Two-body kinematics

- In the CM frame of two particles of mass  $m_1$  and  $m_2$

$$E_1 = \frac{\sqrt{s}}{2} \left( 1 + \frac{m_1^2}{s} - \frac{m_2^2}{s} \right)$$

$$E_2 = \frac{\sqrt{s}}{2} \left( 1 + \frac{m_2^2}{s} - \frac{m_1^2}{s} \right)$$

$$p_1 = p_2 = \frac{\sqrt{s}}{2} \sqrt{1 - \frac{2(m_1^2 + m_2^2)}{s} + \frac{(m_1^2 - m_2^2)^2}{s^2}}$$



$$\tilde{\mu} \rightarrow \mu \chi^0$$

• In the smuon rest frame  $\hat{p}_\mu = \frac{m_{\tilde{\mu}}}{2} \left( 1 - \frac{m_{\tilde{\chi}^0}^2}{m_{\tilde{\mu}}^2} \right) (1, \sin \hat{\theta}, 0, \cos \hat{\theta})$

• In the lab frame  $\gamma_{\tilde{\mu}} = \frac{E_{\tilde{\mu}}}{m_{\tilde{\mu}}} = \frac{\sqrt{s}}{2m_{\tilde{\mu}}}$   $\beta_{\tilde{\mu}} = \sqrt{1 - \frac{4m_{\tilde{\mu}}^2}{s}}$

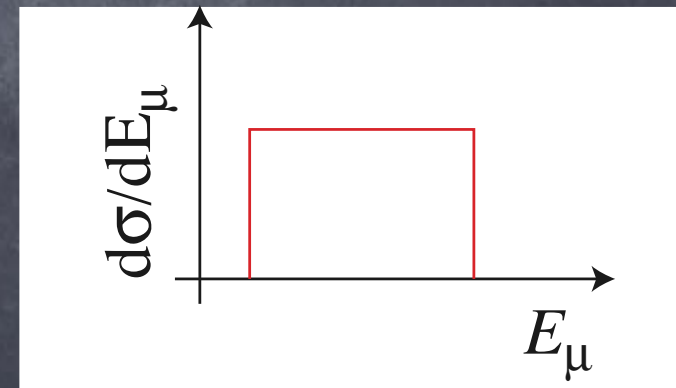
• muon momentum in the lab frame

$$p_\mu = \frac{m_{\tilde{\mu}}}{2} \left( 1 - \frac{m_{\tilde{\chi}^0}^2}{m_{\tilde{\mu}}^2} \right) (\gamma_{\tilde{\mu}} + \gamma_{\tilde{\mu}} \beta_{\tilde{\mu}} \cos \hat{\theta}, \sin \hat{\theta}, 0, \gamma_{\tilde{\mu}} \cos \hat{\theta} + \gamma_{\tilde{\mu}} \beta_{\tilde{\mu}})$$

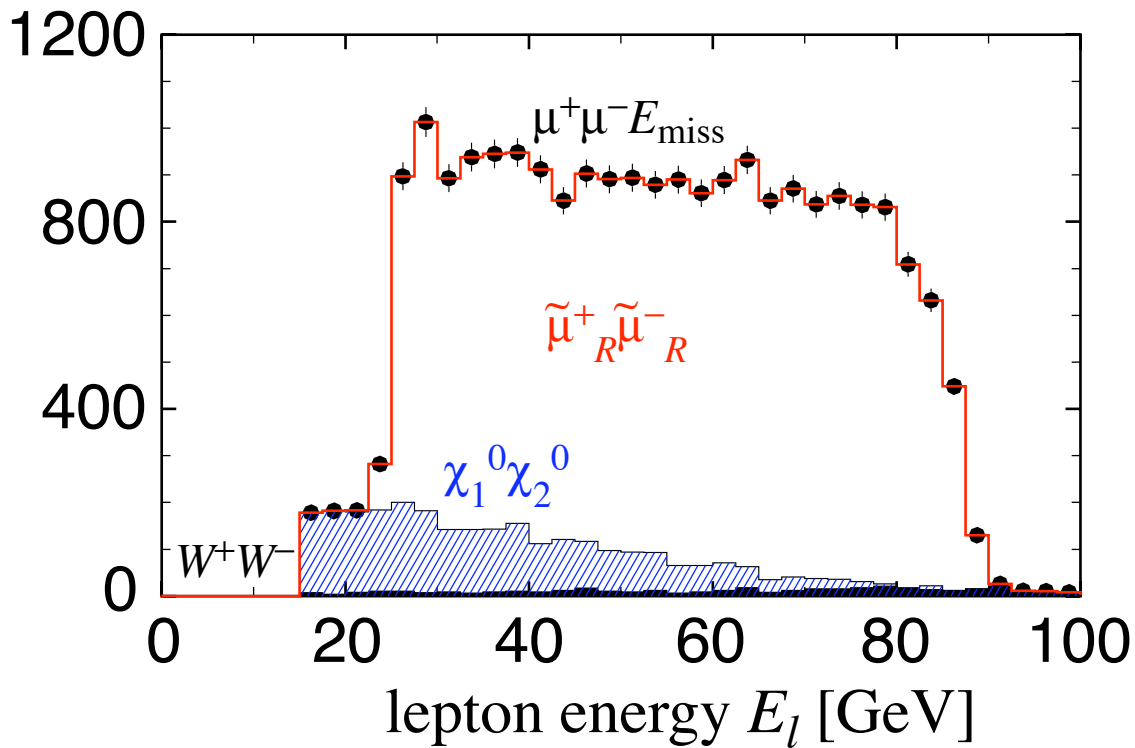
• Therefore, the muon energy is

$$\frac{\sqrt{s}}{4} \left( 1 - \frac{m_{\tilde{\chi}^0}^2}{m_{\tilde{\mu}}^2} \right) (1 - \beta_{\tilde{\mu}}) < E_\mu < \frac{\sqrt{s}}{4} \left( 1 - \frac{m_{\tilde{\chi}^0}^2}{m_{\tilde{\mu}}^2} \right) (1 + \beta_{\tilde{\mu}})$$

$$\frac{d\sigma}{dE_\mu} \propto \frac{d\sigma}{d \cos \hat{\theta}} = \text{constant}$$



$$\tilde{\mu} \rightarrow \mu \chi^0$$



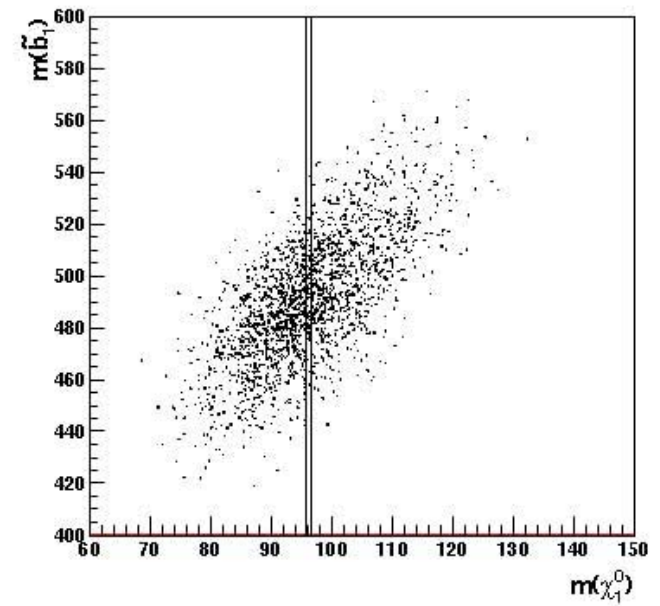
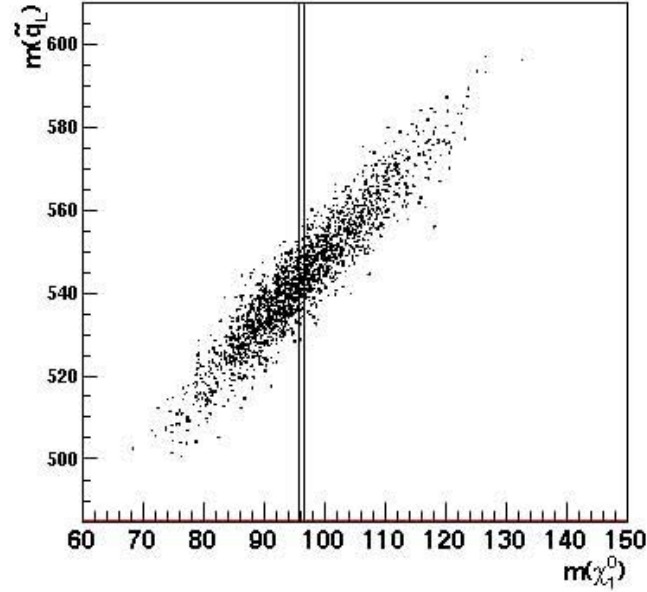
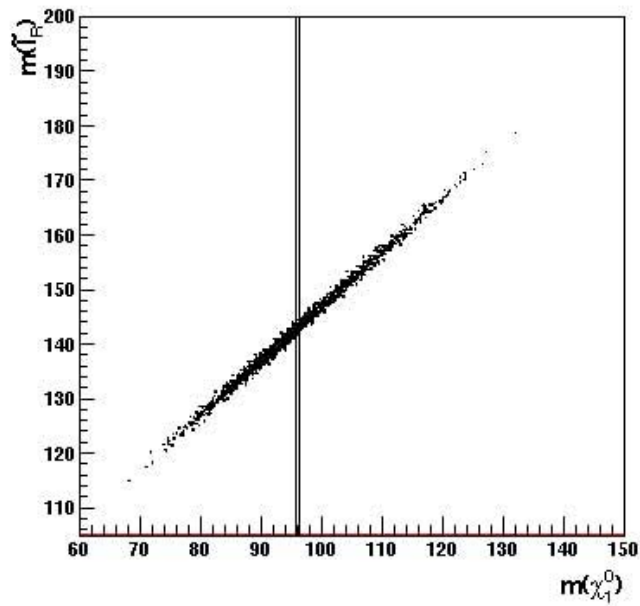
fit to the kinetic distribution

$$m_{\tilde{\mu}} = 132.0 \pm 0.3 \text{ GeV}$$

$$m_{\tilde{\chi}^0} = 71.9 \pm 0.1 \text{ GeV}$$



# LHC/LC synergy



# Reconstruct Lagrangian from data

- Specify the fields
  - **mass**
  - **spin**  $\Rightarrow$  Klein-Gordon, Dirac, Majorana, gauge
  - $SU(3) \times SU(2) \times U(1)$  quantum numbers
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Spin

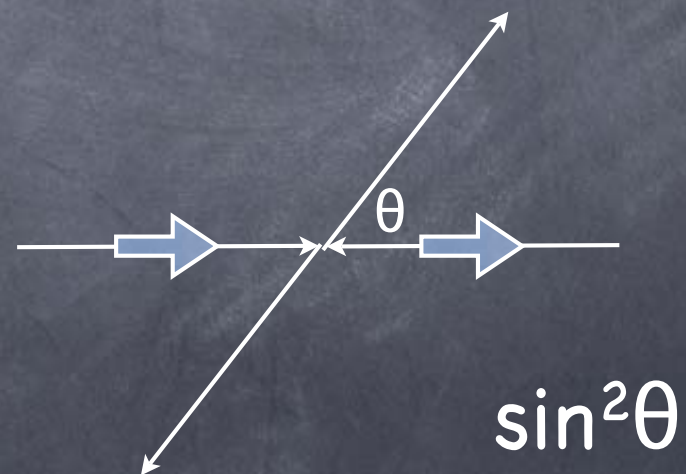
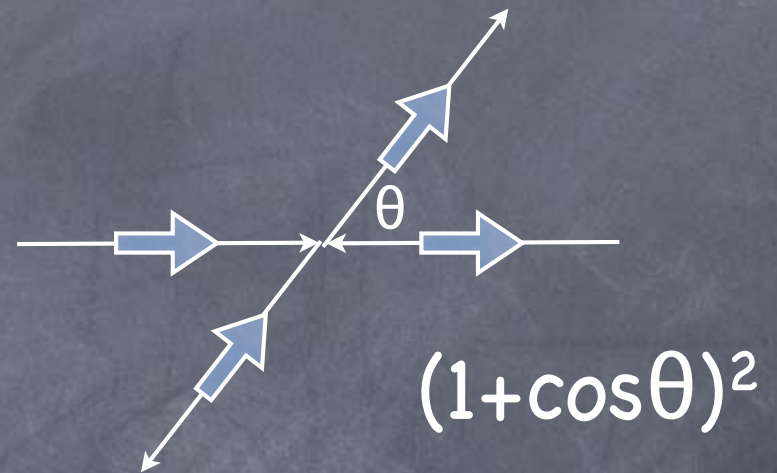


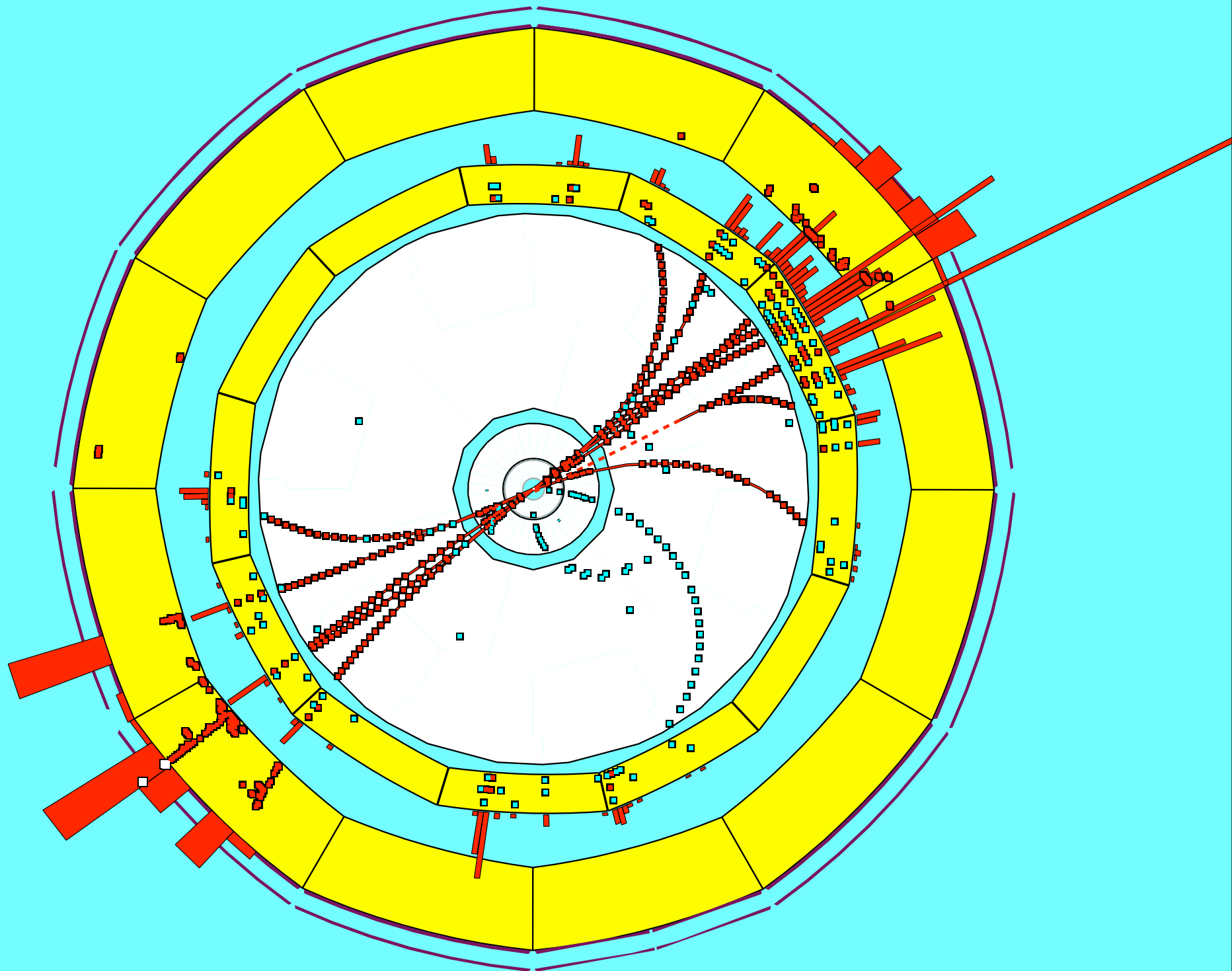
# Spin

- production angle distribution well above the threshold:

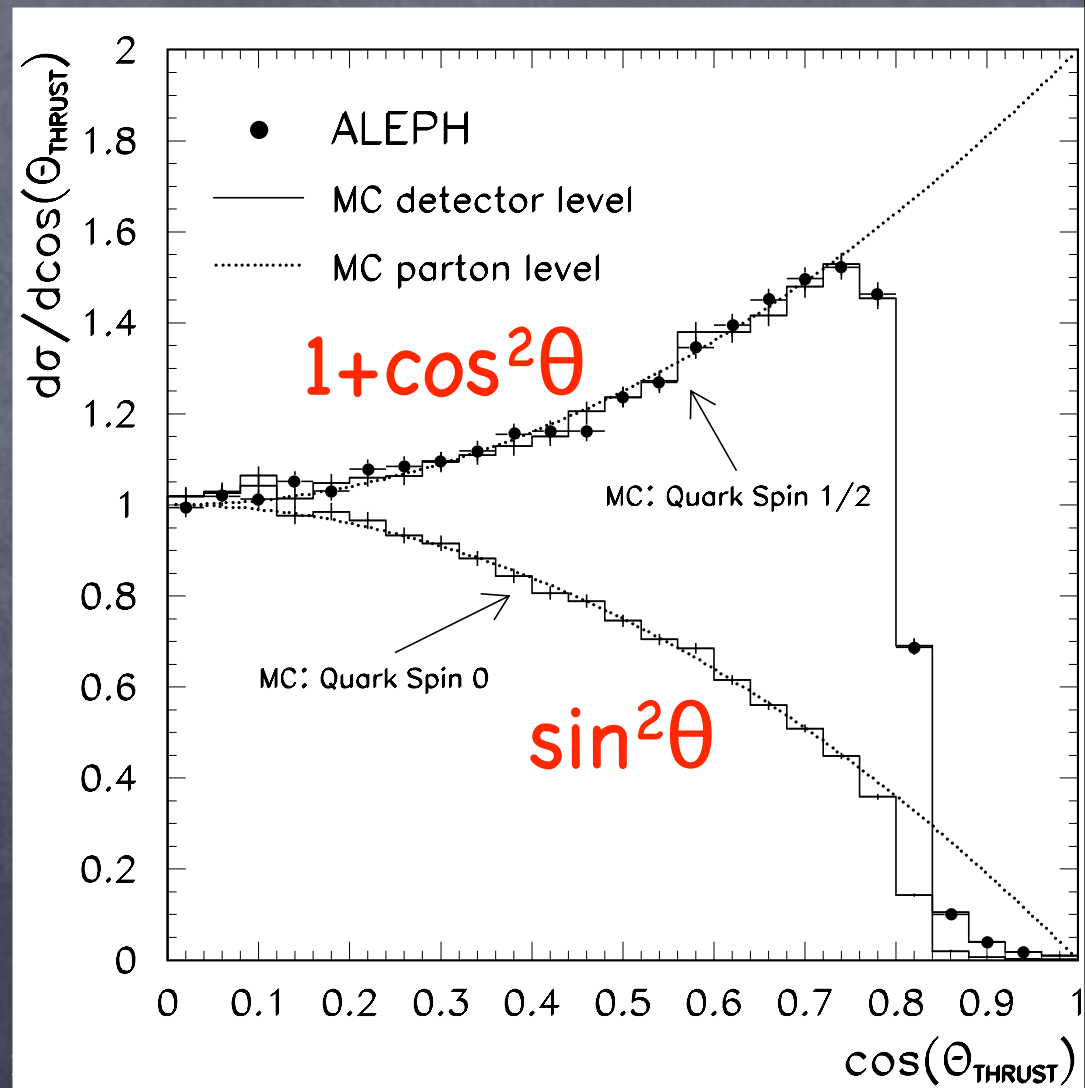
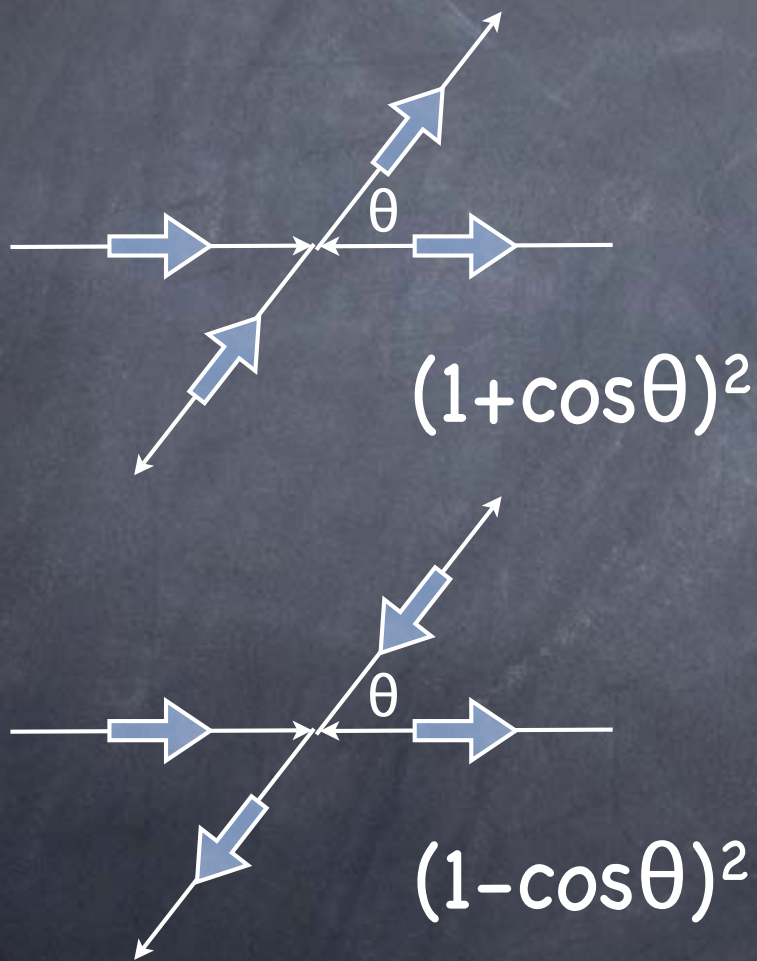
- spin 1/2

- spin 0

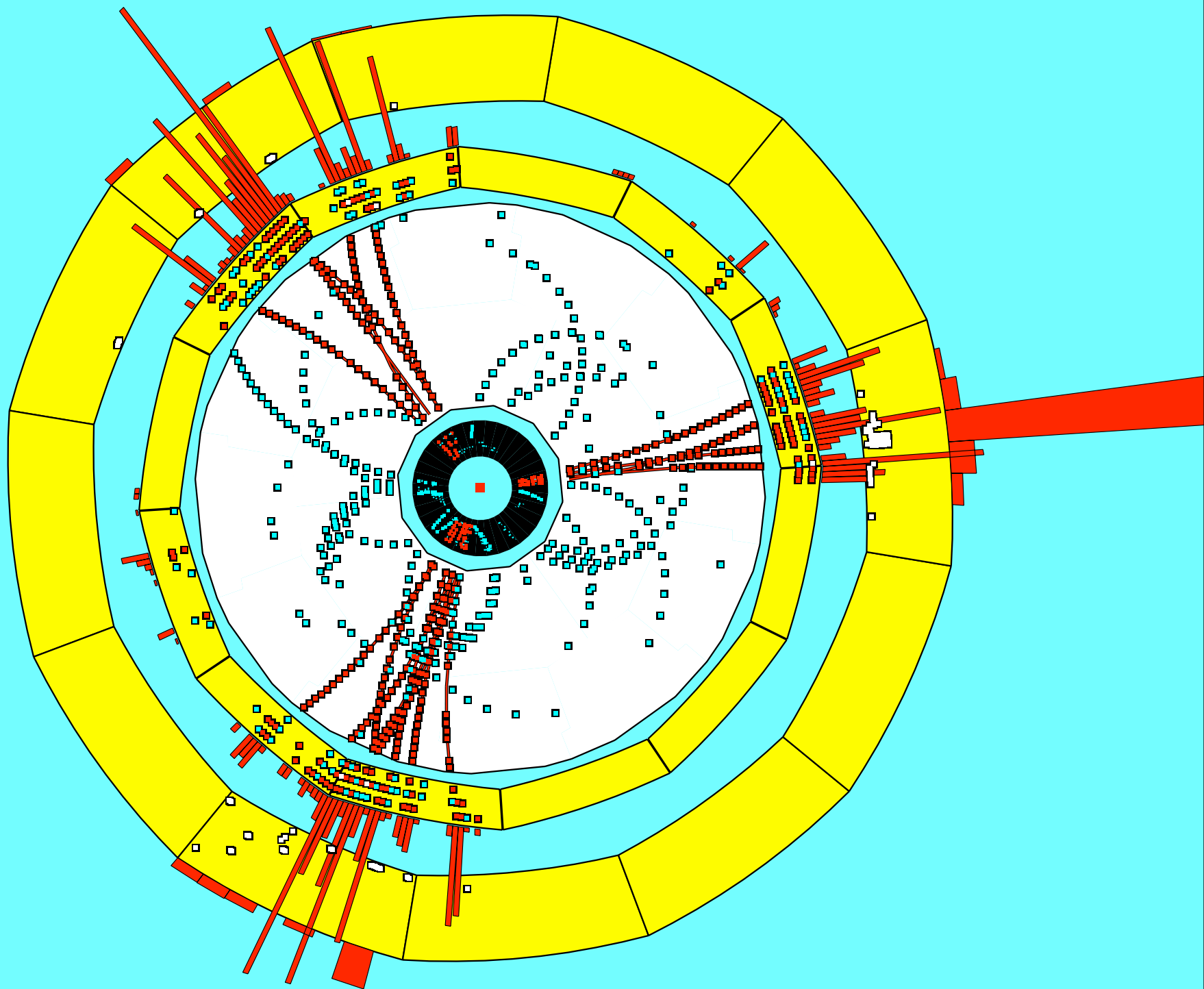




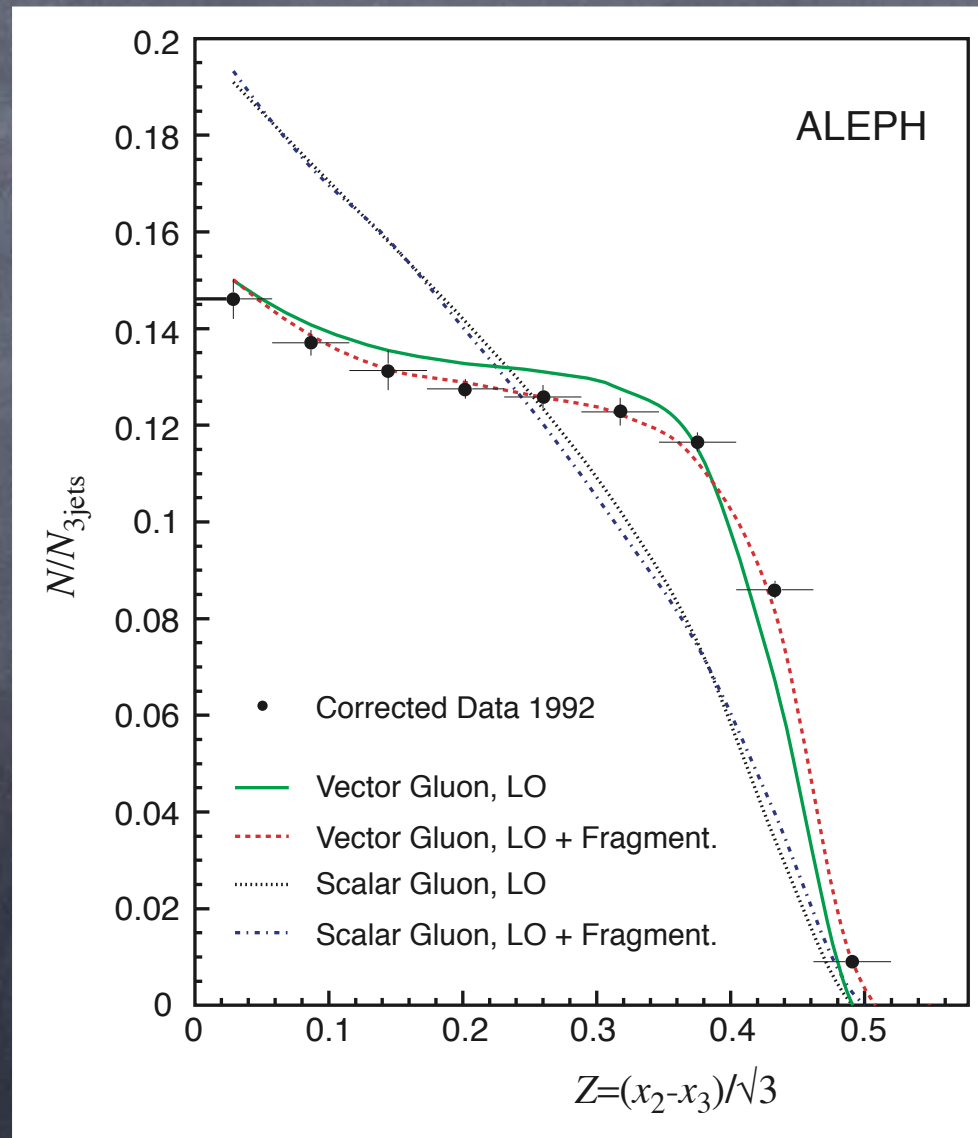
# "New particle" has spin 1/2

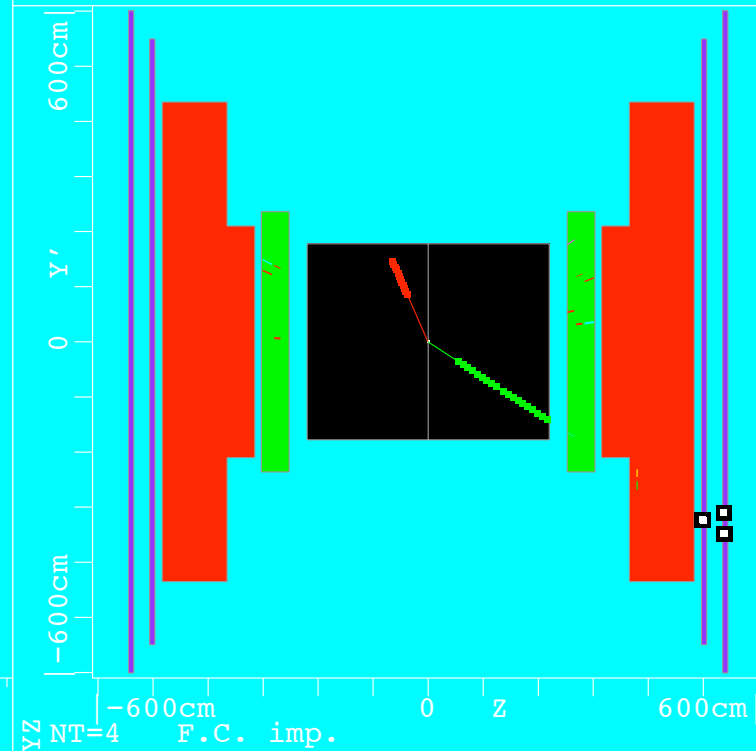
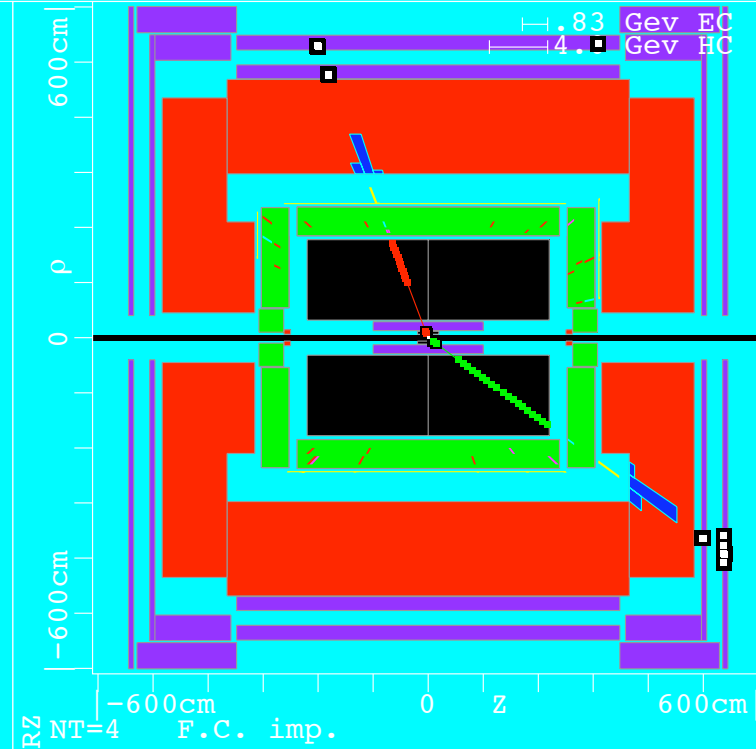
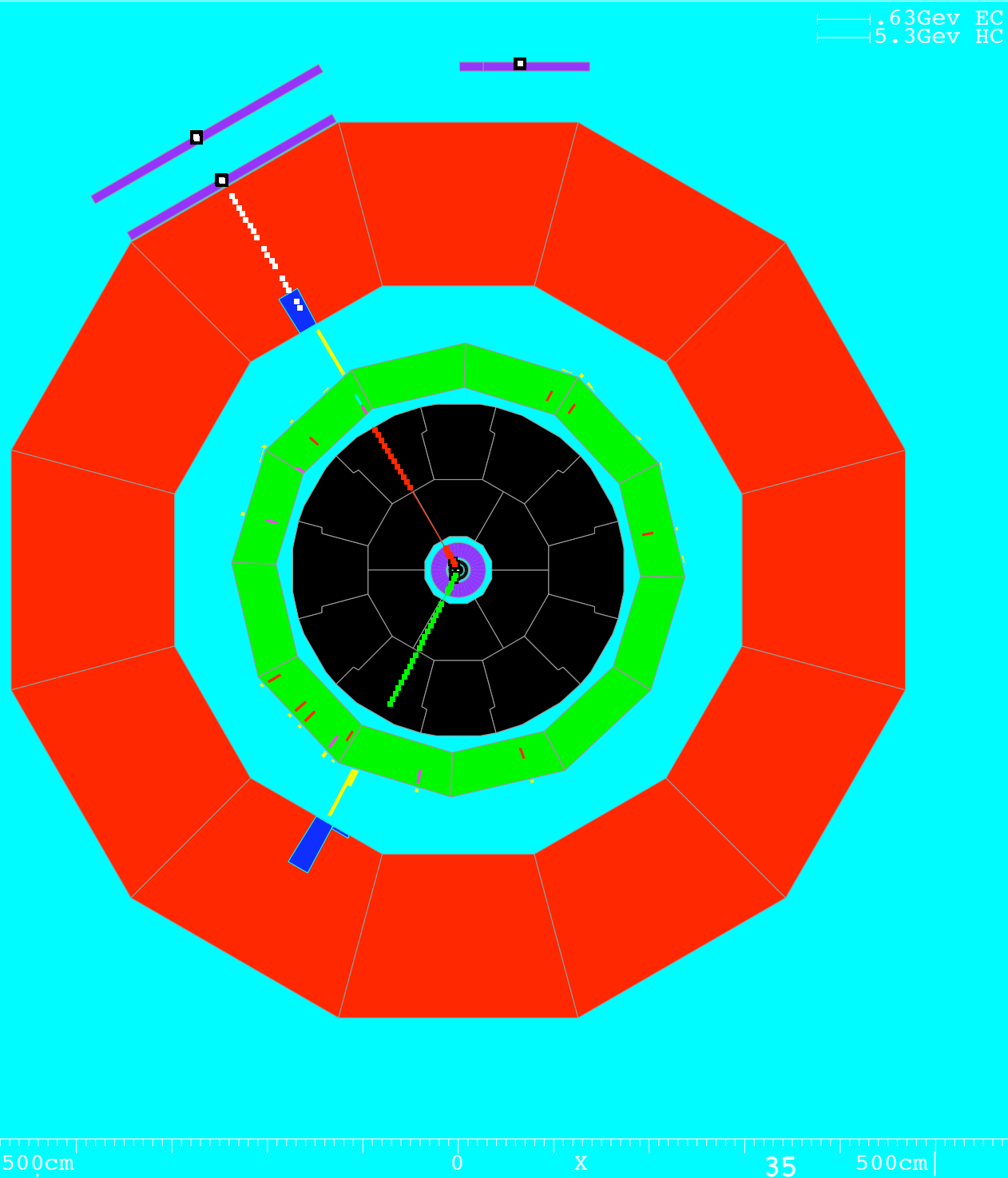






# "New particle" has spin 1







# Smuon production

- $e^+e^- \rightarrow \tilde{\mu}^+\tilde{\mu}^- \rightarrow (\mu^+\tilde{\chi}_1^0)(\mu^-\tilde{\chi}_1^0)$
- once masses known, you can solve kinematics up to a two-fold ambiguity
- muon momenta measured:  $p_{1,2}^\mu = (E_{1,2}, \vec{p}_{1,2})$
- neutralino momenta:  $q_{1,2}^\mu = \left( \frac{\sqrt{s}}{2} - E_{1,2}, \vec{q}_{1,2} \right)$
- neutralino mass constraint:  $\vec{q}_{1,2}^2 = \left( \frac{\sqrt{s}}{2} - E_{1,2} \right)^2 - m_{\tilde{\chi}}^2$
- smuon mass constraint:  $\left( \frac{\sqrt{s}}{2} \right)^2 - (\vec{p}_1 + \vec{q}_1)^2 = m_{\tilde{\mu}}^2 \longrightarrow \vec{p}_1 \cdot \vec{q}_1$

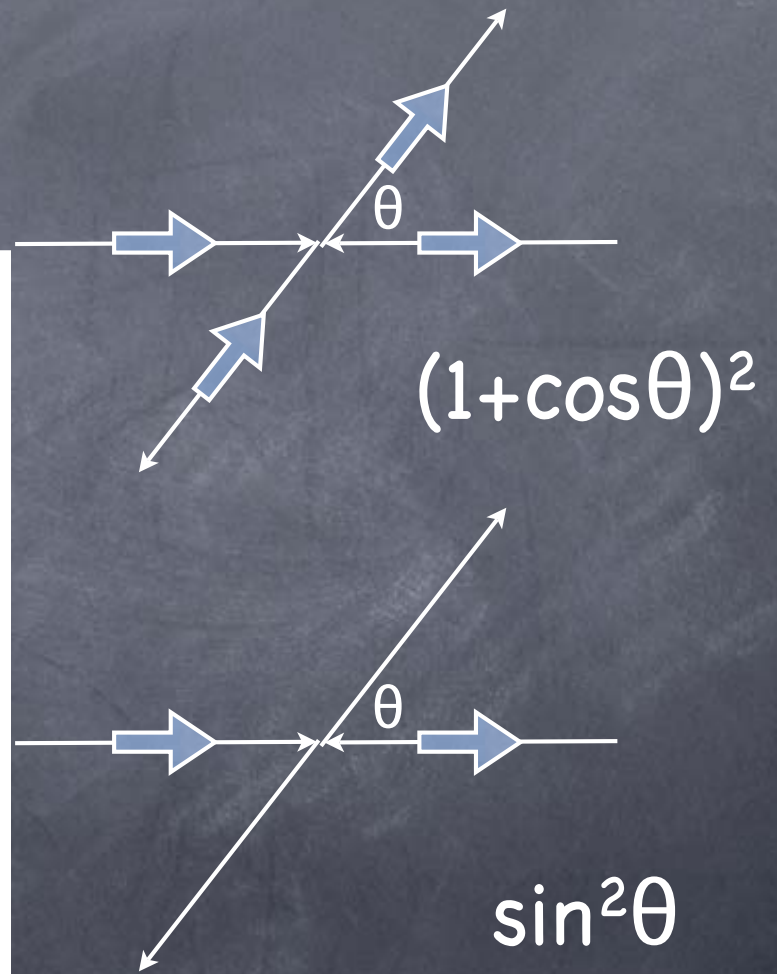
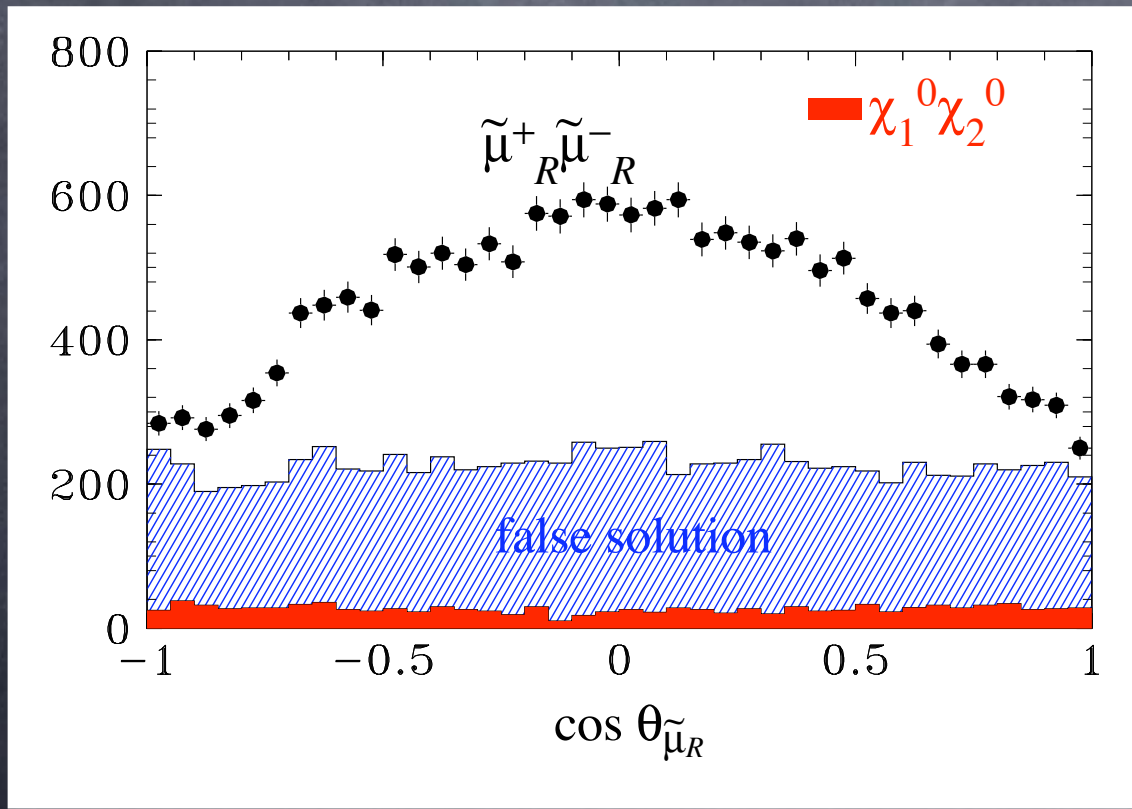
- momentum conservation:

$$\vec{q}_2^2 = (\vec{p}_1 + \vec{p}_2 + \vec{q}_1)^2 = (\vec{p}_1 + \vec{p}_2)^2 + \vec{q}_1^2 + 2\vec{p}_1 \cdot \vec{q}_1 + 2\vec{p}_2 \cdot \vec{q}_1 \longrightarrow \vec{p}_2 \cdot \vec{q}_1$$

- Now know  $|\vec{q}_1|$ ,  $\vec{p}_1 \cdot \vec{q}_1$ ,  $\vec{p}_2 \cdot \vec{q}_1$

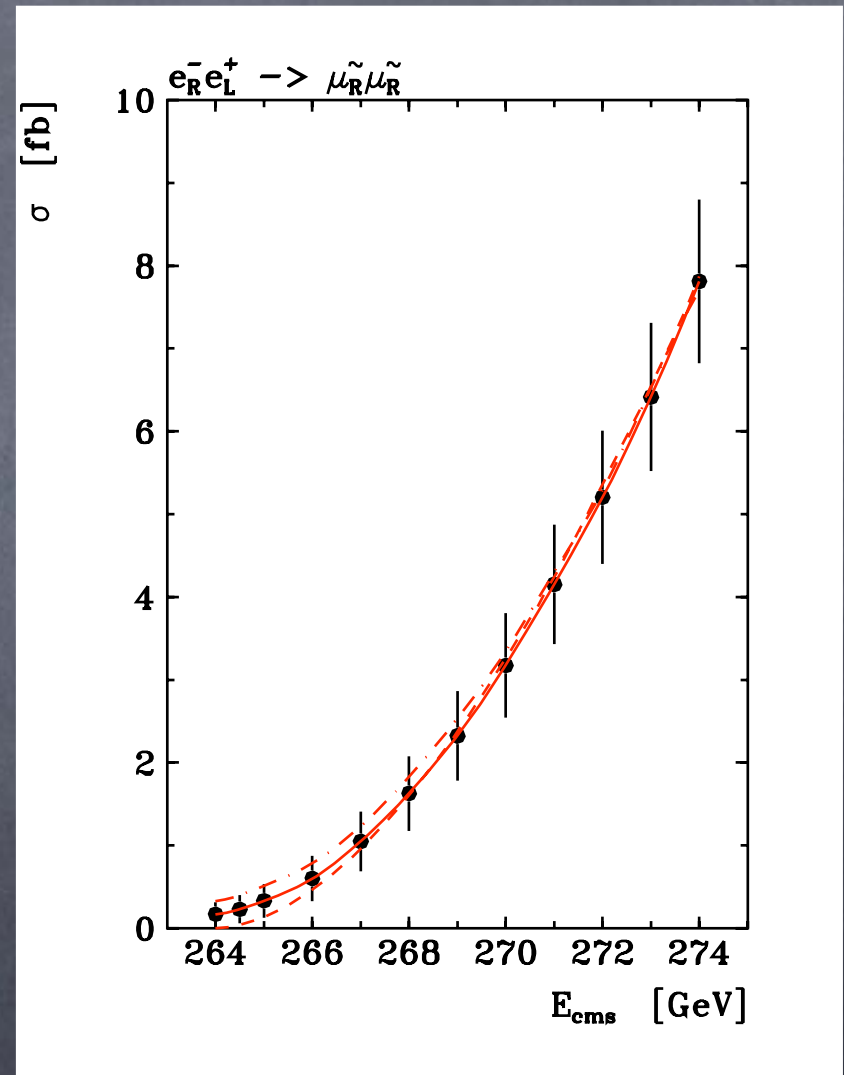
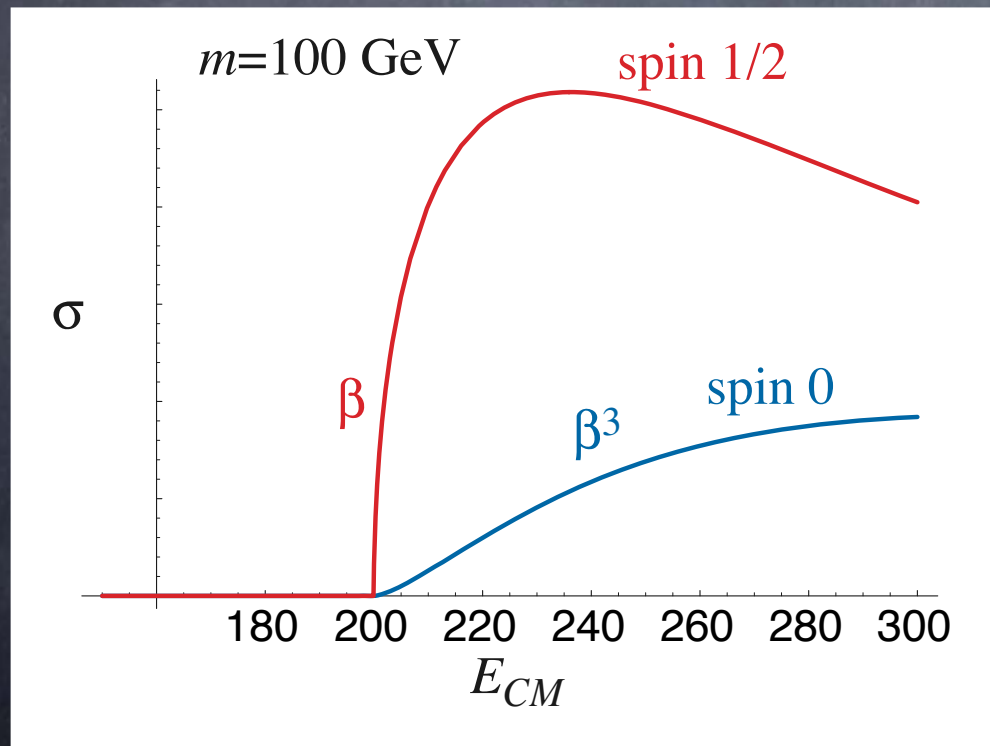
- Know  $\vec{q}_1$  up to a two-fold ambiguity

# Smuon has spin 0



# Spin

- threshold behavior  
non-relativistic limit:  
L, S separately  
conserved



$$m_{\tilde{\mu}} = 132.0 \pm 0.09 \text{ GeV}$$
$$m_{\tilde{\chi}^0} = 71.9 \pm 0.05 \text{ GeV}$$



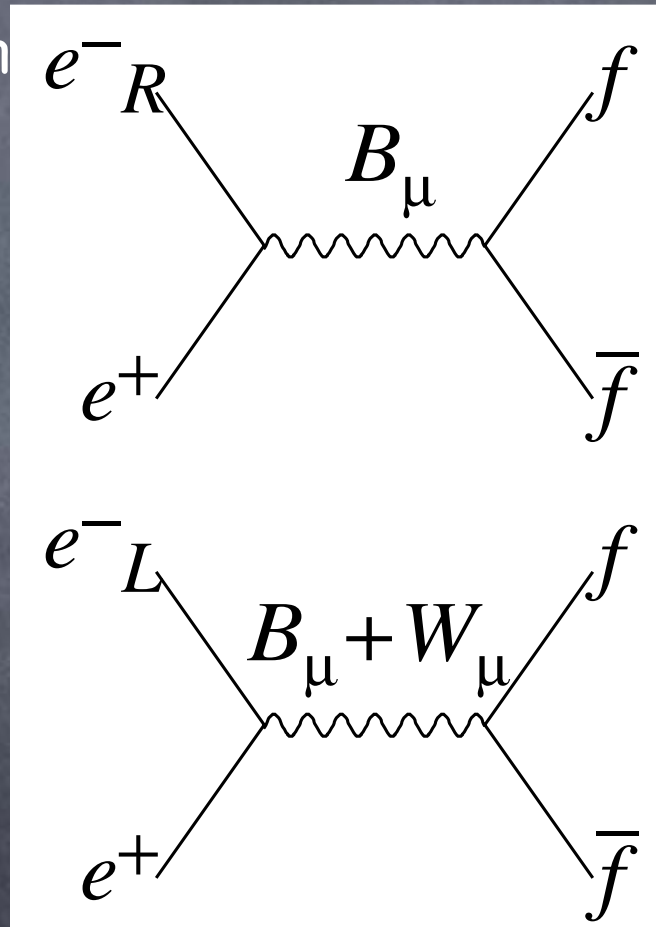
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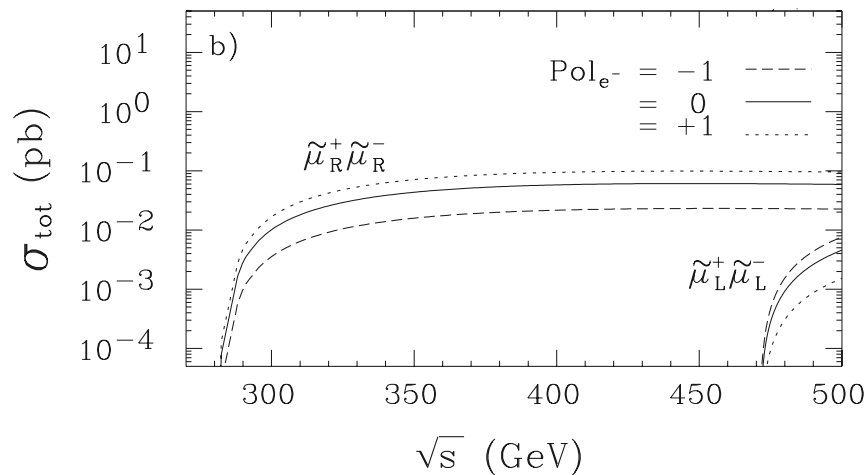
gauge quantum numbers

# polarization

- Use polarized electron beam
- can ignore  $m_Z^2 \ll s$
- $e_R$  couples only to  $B_\mu$
- $e_L$  couples to  $B_\mu + W_\mu^0$



$$\propto (g'^2 Y_f)^2$$



$$\propto (g'^2 Y_f + g^2 I_{3f})^2 / 4$$



# Reconstruct Lagrangian from data

- Specify the fields
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  - spin  $\Rightarrow$  Klein-Gordon, Dirac, Majorana, gauge
  - $SU(3) \times SU(2) \times U(1)$  quantum numbers
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Disentangle mixings

# gauginos, higgsinos

- charged ones "charginos"

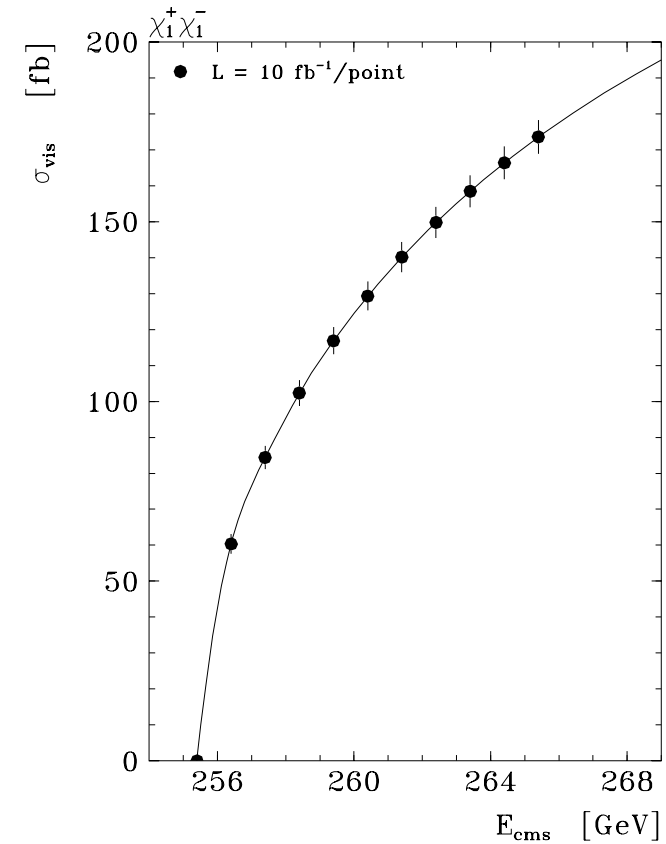
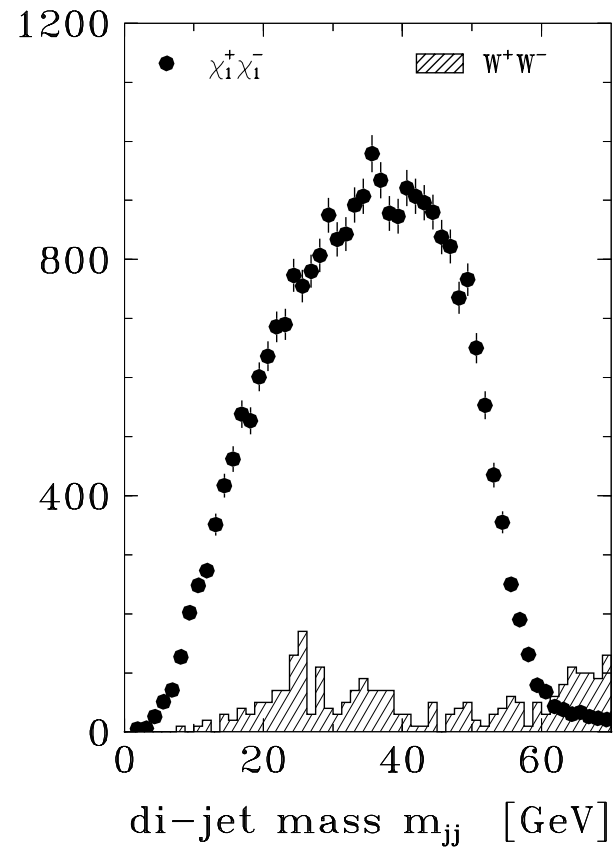
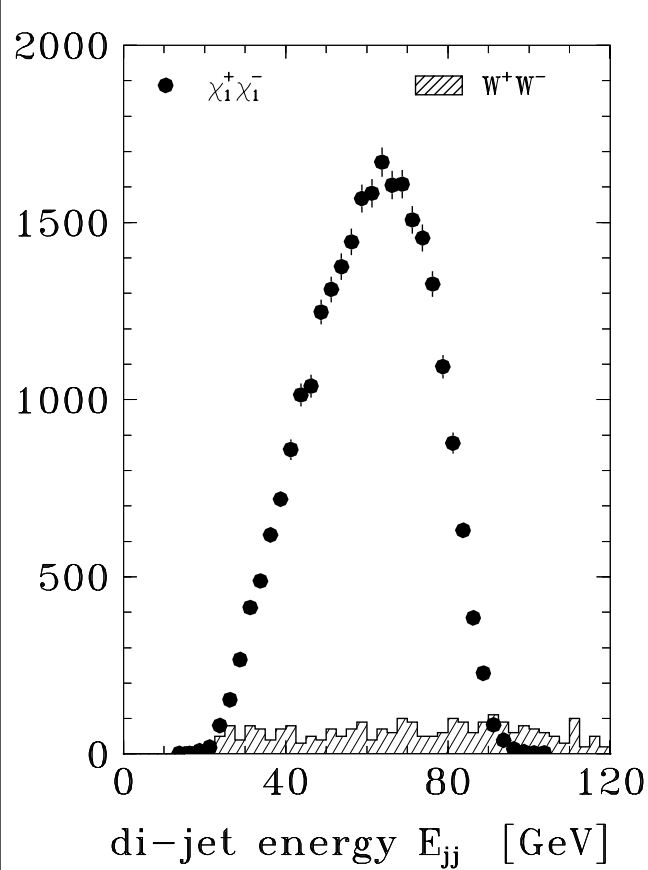
$$(\tilde{W}^- \tilde{H}_d^-) \begin{pmatrix} M_2 & \sqrt{2}m_W \sin\beta \\ \sqrt{2}m_W \cos\beta & \mu \end{pmatrix} \begin{pmatrix} \tilde{W}^+ \\ \tilde{H}_u^+ \end{pmatrix}$$

- neutral ones "neutralinos"

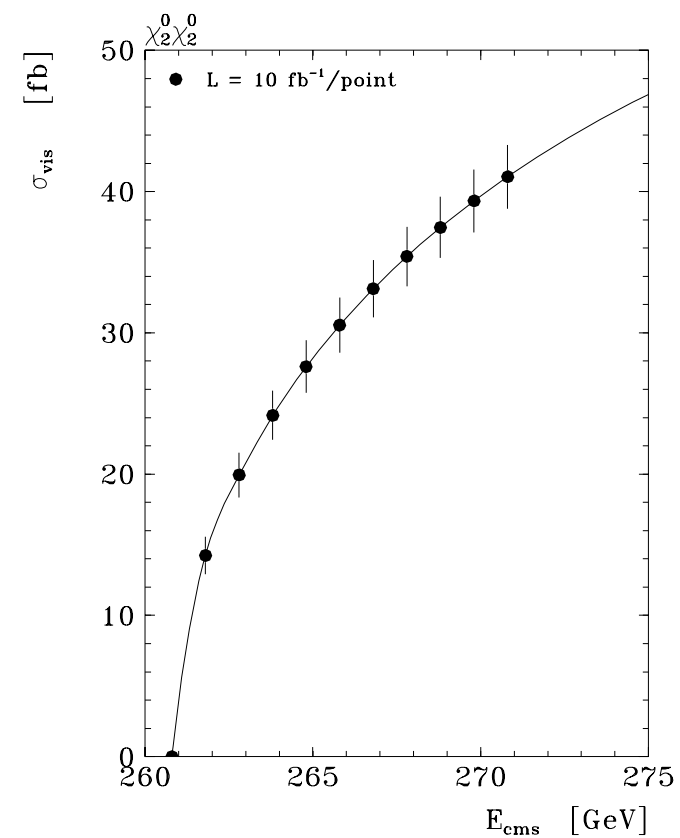
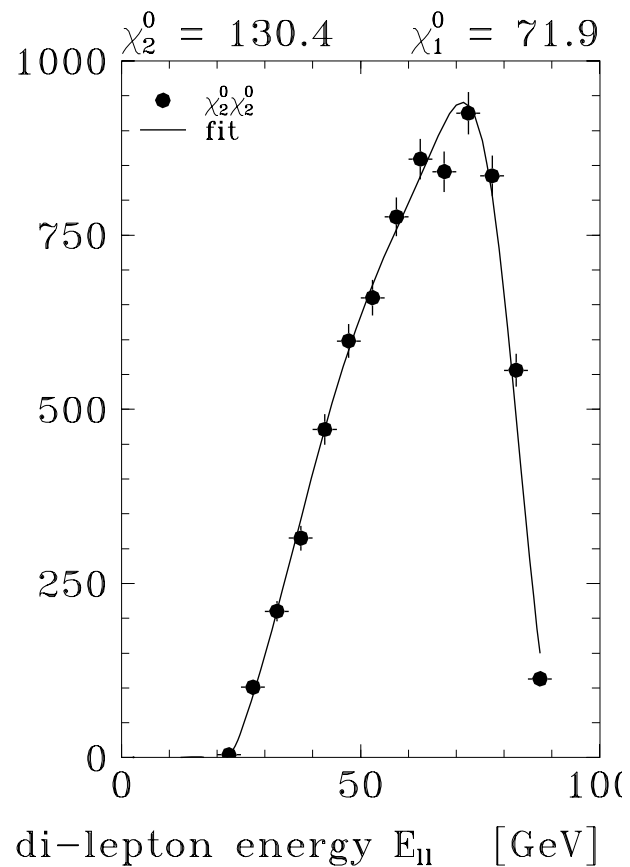
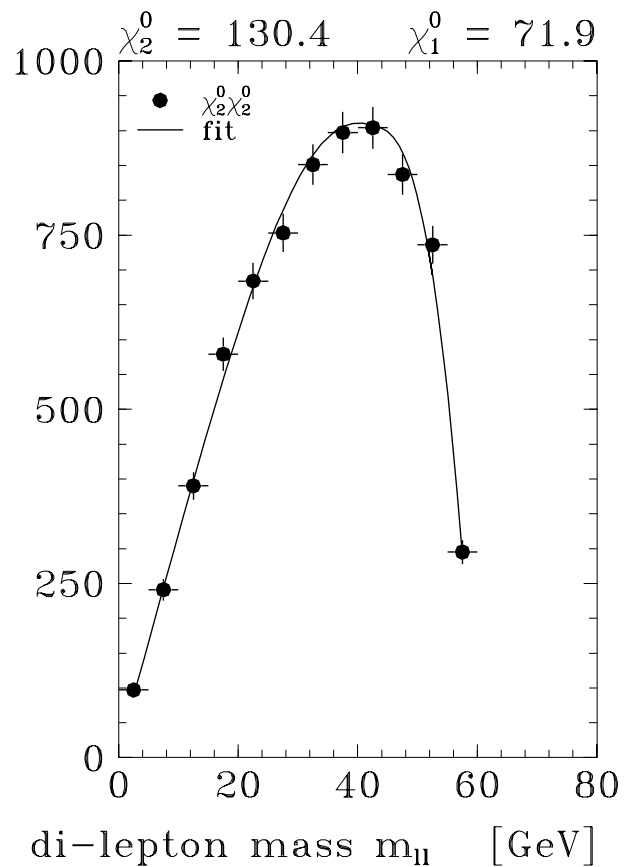
$$(\tilde{B}, \tilde{W}^0, \tilde{H}_d^0, \tilde{H}_u^0) \begin{pmatrix} M_1 & 0 & -m_Z s_W c_\beta & m_Z s_W s_\beta \\ 0 & M_2 & m_Z c_W c_\beta & -m_Z c_W s_\beta \\ -m_Z s_W c_\beta & m_Z c_W c_\beta & 0 & -\mu \\ m_Z s_W s_\beta & -m_Z c_W s_\beta & -\mu & 0 \end{pmatrix} \begin{pmatrix} \tilde{B} \\ \tilde{W}^0 \\ \tilde{H}_d^0 \\ \tilde{H}_u^0 \end{pmatrix}$$



$$e^+e^- \rightarrow \tilde{\chi}_1^+ \tilde{\chi}_1^- \rightarrow (\tilde{\chi}_1^0 l^\pm \nu_l) (\tilde{\chi}_1^0 q \bar{q}')$$

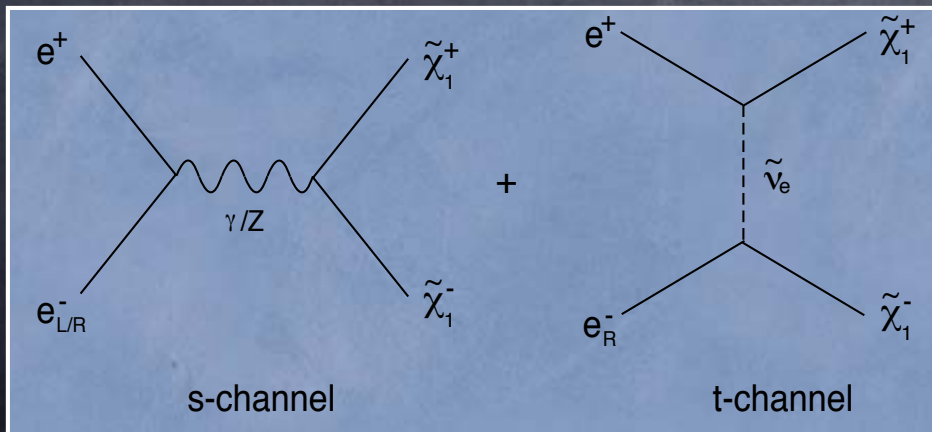


$$e^+e^- \rightarrow \tilde{\chi}_2^0\tilde{\chi}_2^0 \rightarrow (\tilde{\chi}_1^0 l^+ l^-) (\tilde{\chi}_1^0 l'^+ l'^-)$$



# Model-independent parameter determination

- Chargino/neutralino mass matrices have four parameters  $M_1, M_2, \mu, \tan\beta$
- Can measure 2+4 masses
- can measure 10x2 neutralino cross sections
 
$$\sigma_{L,R}(e^+e^- \rightarrow \tilde{\chi}_i^0 \tilde{\chi}_j^0) \quad \sigma_{L,R}(e^+e^- \rightarrow \tilde{\chi}_i^+ \tilde{\chi}_j^-)$$
- can measure 3x2 chargino cross sections
- depend on masses of  $\tilde{\nu}_e, \tilde{e}_L, \tilde{e}_R$

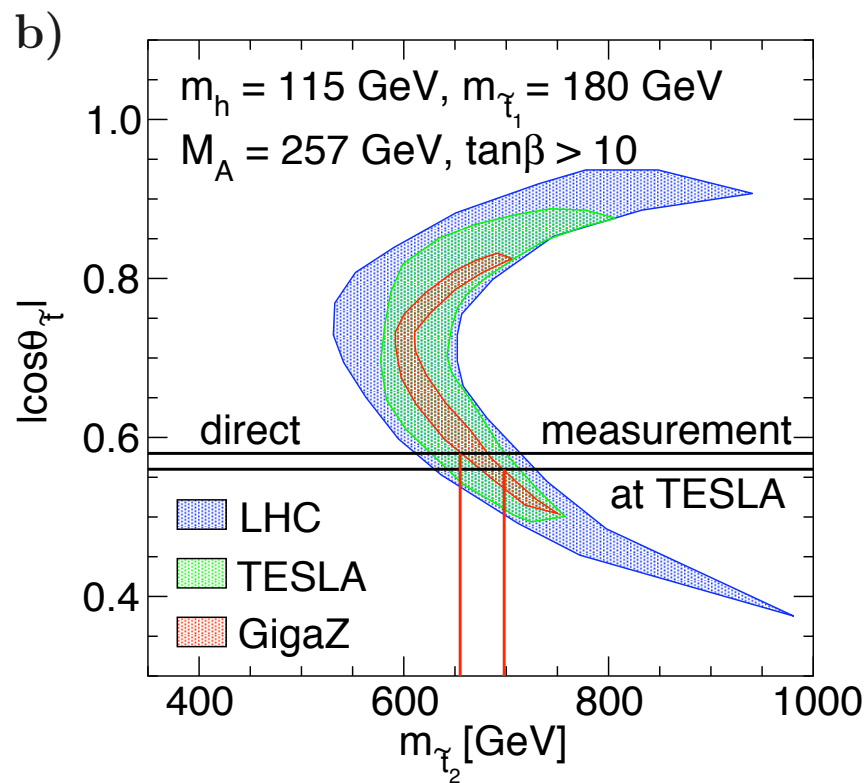
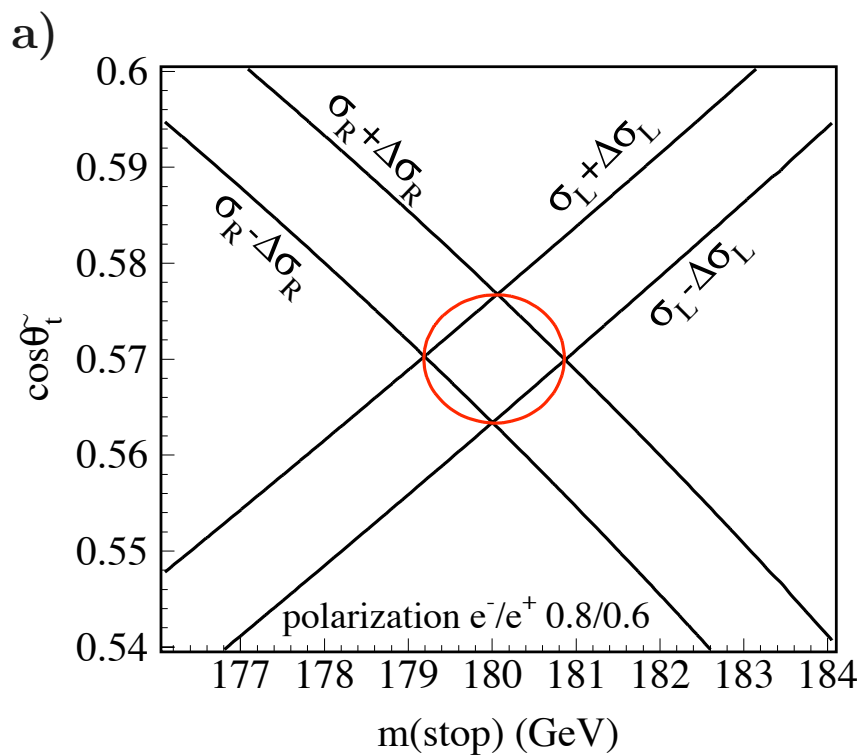


	input	fit
$M_2$	152 GeV	$152 \pm 1.8$ GeV
$\mu$	316 GeV	$316 \pm 0.9$ GeV
$\tan\beta$	3	$3 \pm 0.7$
$M_1$	78.7 GeV	$78.7 \pm 0.7$ GeV



# Stop

$$\begin{pmatrix} \tilde{t}_L^* & \tilde{t}_R^* \end{pmatrix} \begin{pmatrix} m_{\tilde{Q}_3}^2 + m_t^2 & (A_t - \mu^* \cot \beta) m_t \\ (A_t^* - \mu \cot \beta) m_t & m_{\tilde{t}}^2 + m_t^2 \end{pmatrix} \begin{pmatrix} \tilde{t}_L \\ \tilde{t}_R \end{pmatrix}$$



# Reconstruct Lagrangian from data

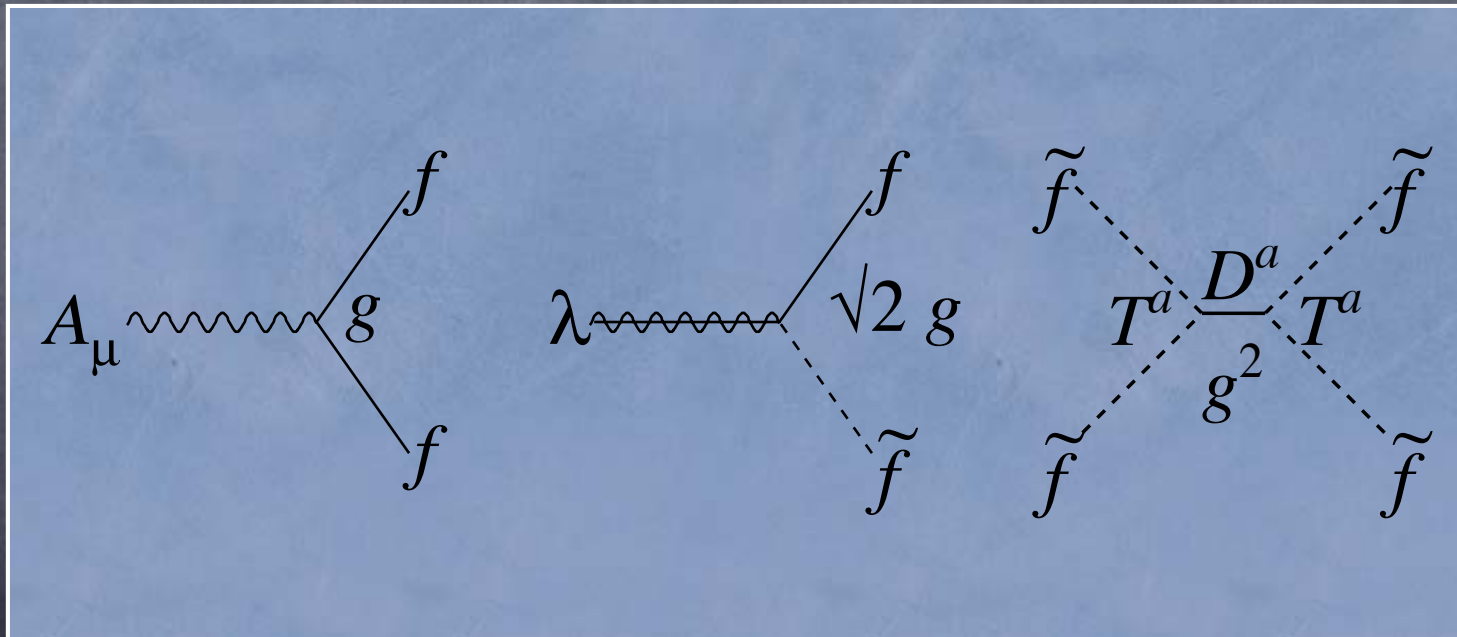
- Specify the fields
  - mass
  - spin  $\Rightarrow$  Klein-Gordon, Dirac, Majorana, gauge
  - $SU(3) \times SU(2) \times U(1)$  quantum numbers
  - mixing of states
- Specify their interactions
  - $SU(3) \times SU(2) \times U(1)$  quantum numbers determine gauge interactions
  - Yukawa couplings

# Interaction

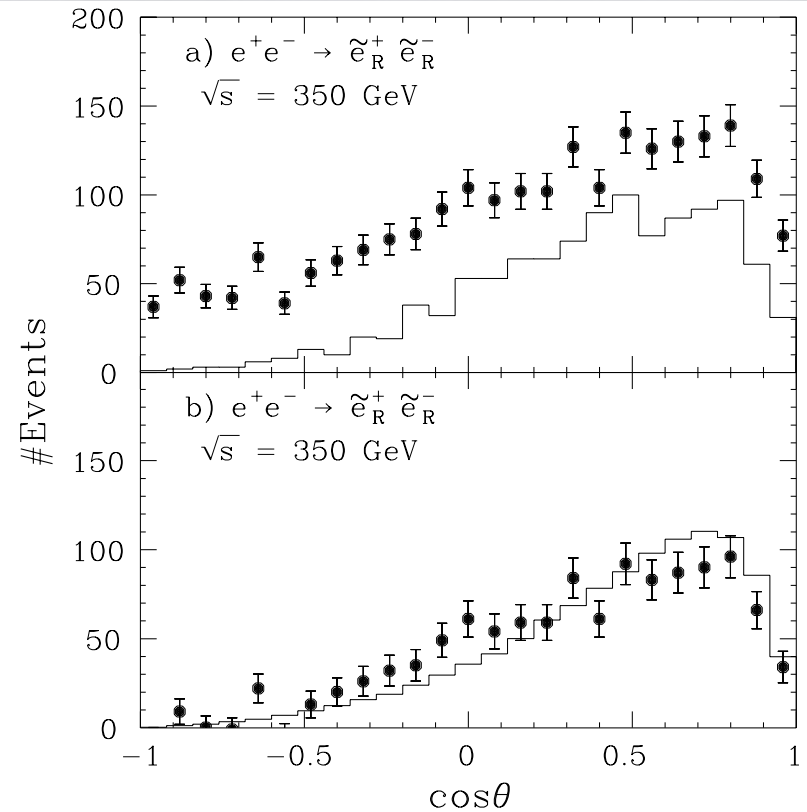
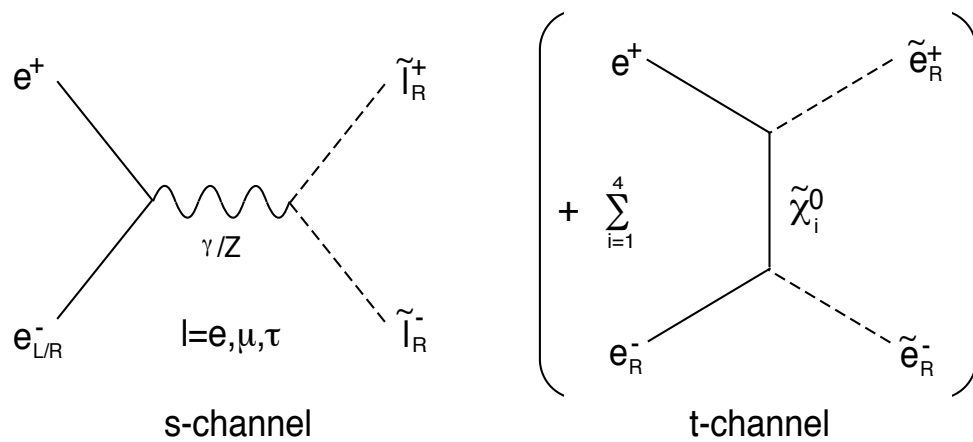


# Feynman rules

- Single gauge coupling constant gives all of these Feynman vertices

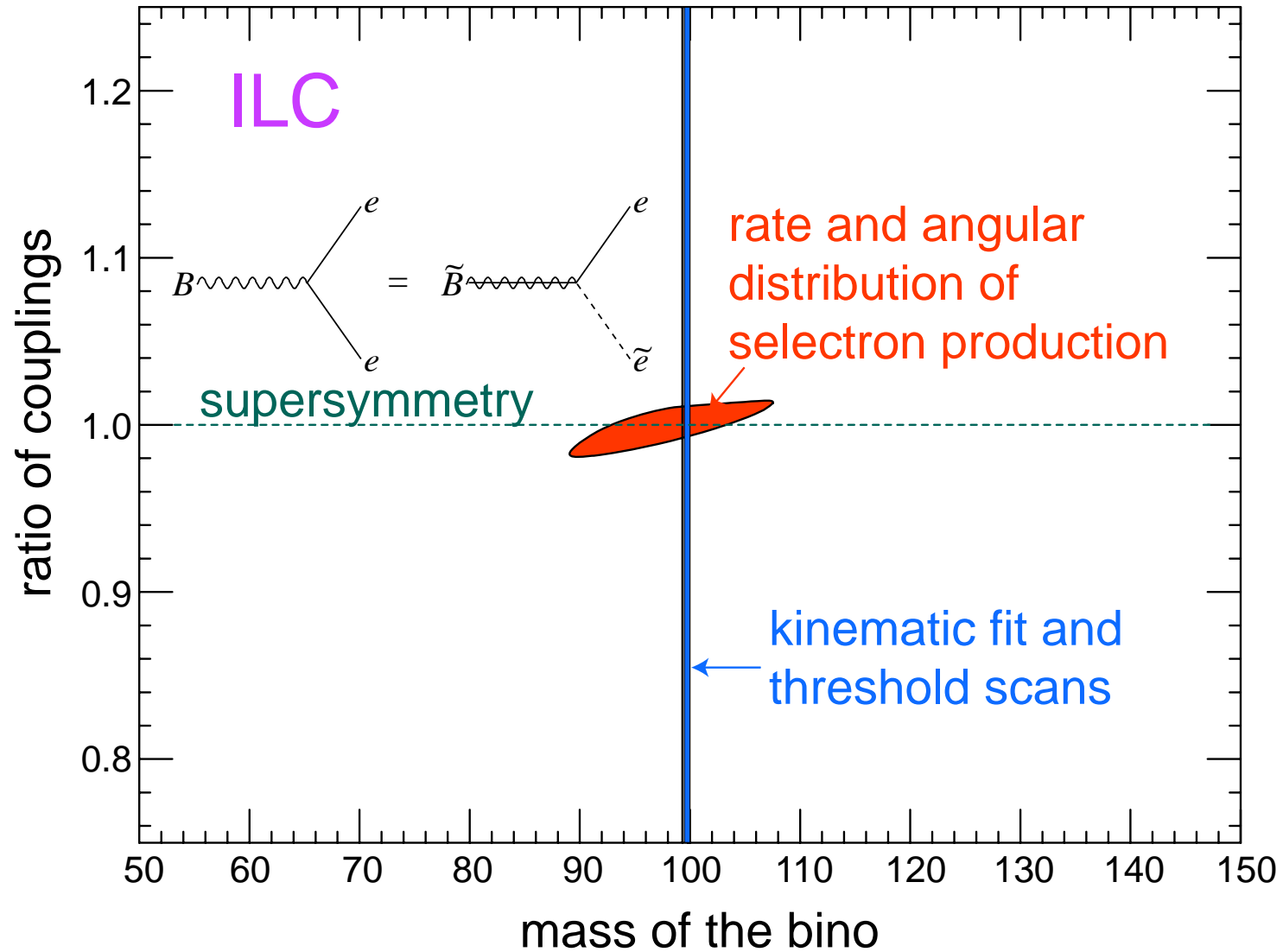


# selectron production



$$\mathcal{M} \propto \sin\theta \left[ 1 - \frac{4Y_{\tilde{B}}^2}{1 - 2\cos\theta\beta_f + \beta_f^2 + 4M_1^2/s} \right] \quad Y_{\tilde{B}} = \frac{g_{e_R\tilde{e}_R\tilde{B}}}{\sqrt{2}g'}$$

# gaugino coupling





# Reconstruct Lagrangian from data

- Specify the fields
  - mass
  - spin  $\Rightarrow$  Klein-Gordon, Dirac, Majorana, gauge
  - $SU(3) \times SU(2) \times U(1)$  quantum numbers
  - mixing of states
- Specify their interactions
  - $SU(3) \times SU(2) \times U(1)$  quantum numbers determine gauge interactions
  - Yukawa couplings

# Proof of supersymmetry

- This way, you can show:
  - new particle has the same gauge quantum numbers as one of the SM particle
  - their spins differ by  $1/2$
  - it has a Yukawa coupling whose size is  $\sqrt{2}$  times the known gauge coupling
- You have reconstructed the supersymmetric Lagrangian from data!

# The New York Times

July 23, 2008

## The Other Half of the World Discovered

Geneva, Switzerland

As an example, supersymmetry

“New-York Times level” confidence

still a long way to

“Halliday-Resnick” level confidence

“We have learned that all particles we observe have unique partners of different spin and statistics, called superpartners, that make our theory of elementary particles valid to small distances.”



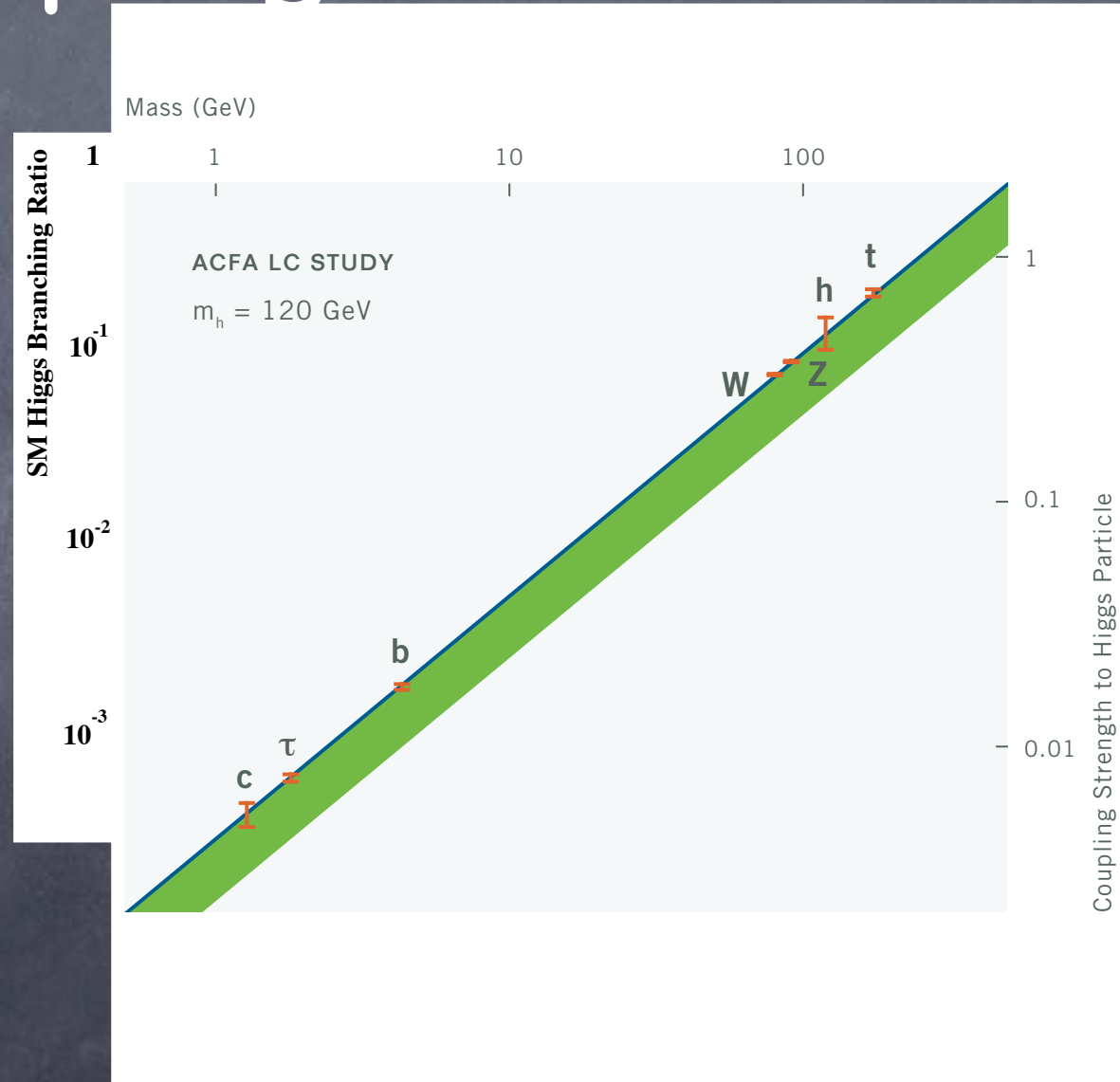
# Physics Significance

# Prove Higgs coupling $\propto$ mass

- Branching Fractions test the relation

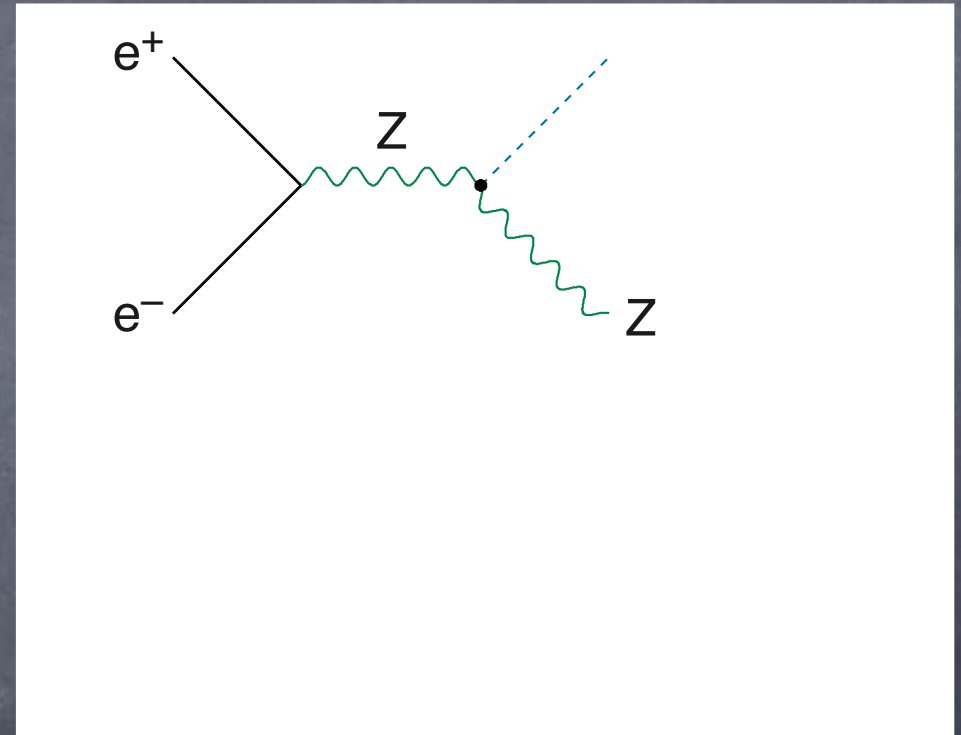
coupling  $\propto$  mass

$\Rightarrow$  proves that Higgs Boson is the Mother of Mass



# Prove it is condensed

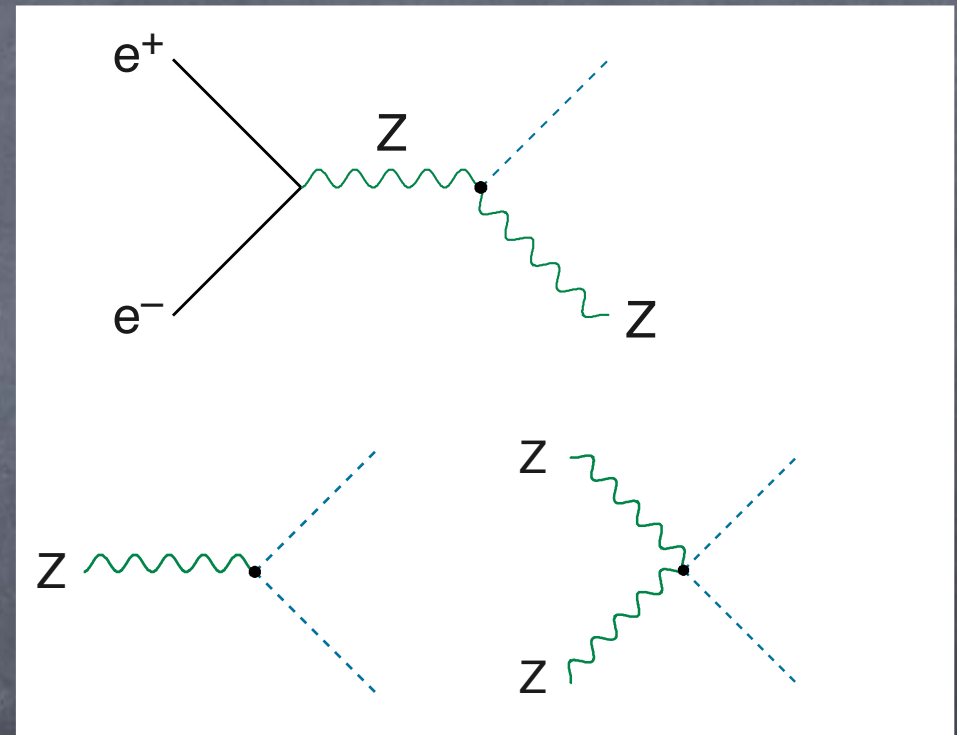
- ZH final state
- Prove the ZZH vertex





# Prove it is condensed

- ZH final state
- Prove the ZZH vertex
- We know Z: gauge boson, H: scalar boson  
 $\Rightarrow$  only two types of vertices

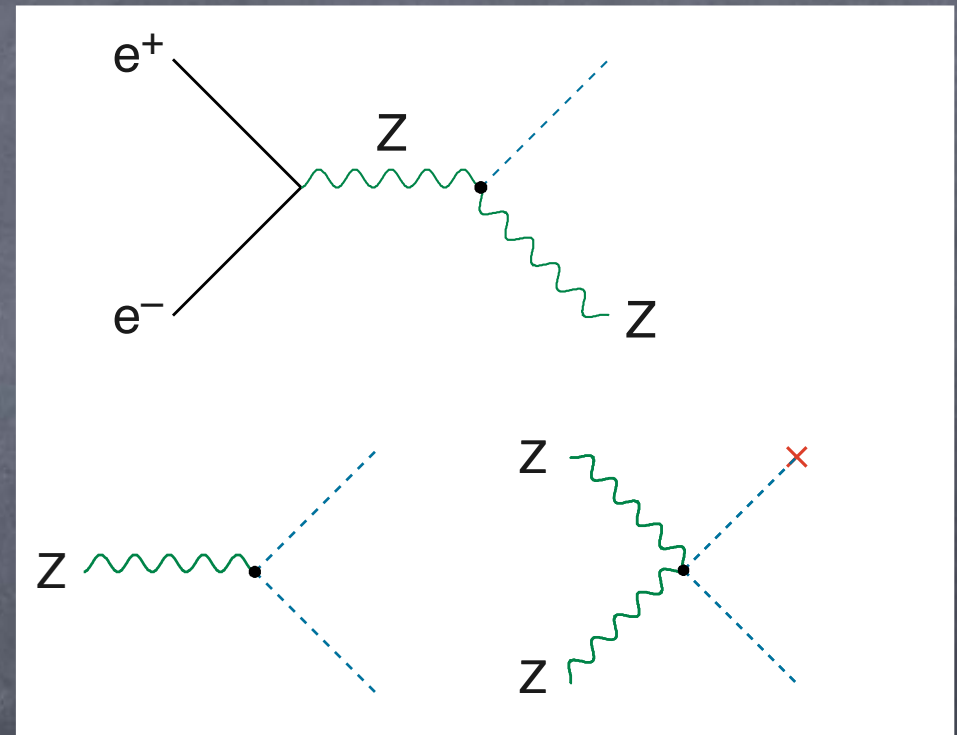


# Prove it is condensed

- ZH final state
- Prove the ZZH vertex
- We know Z: gauge boson, H: scalar boson  
 $\Rightarrow$  only two types of vertices
- Need a condensate to get ZZH vertex

$\Rightarrow$  proves it is condensed in Universe

HM, hep-ex/9606001

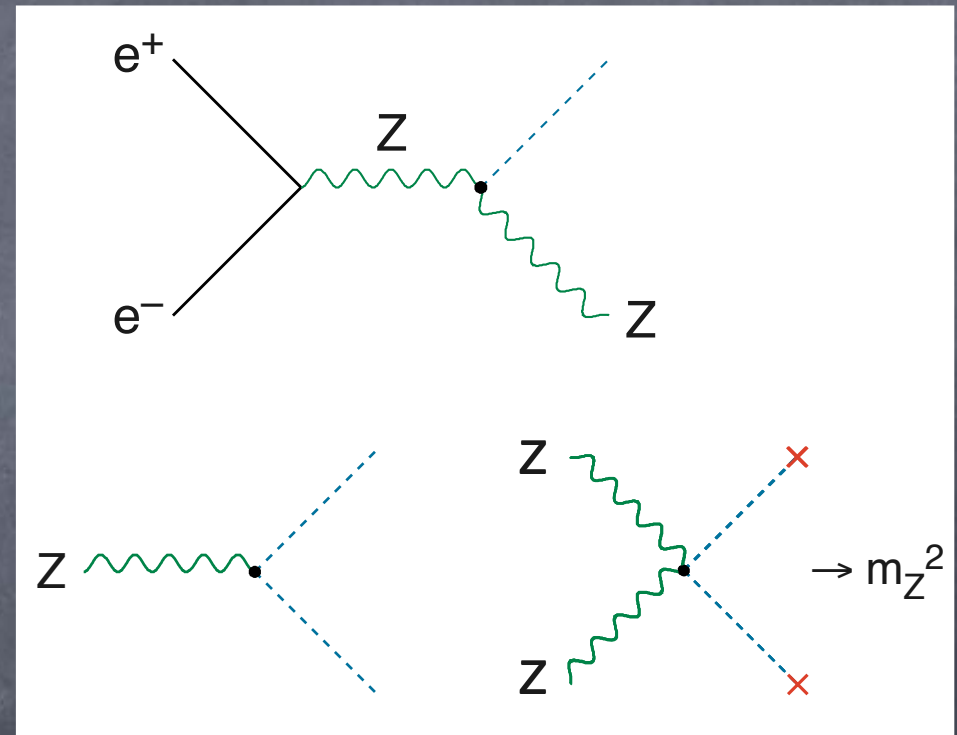


# Prove it is condensed

- ZH final state
- Prove the ZZH vertex
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*$\Rightarrow$  proves it is condensed in Universe*

HM, hep-ex/9606001





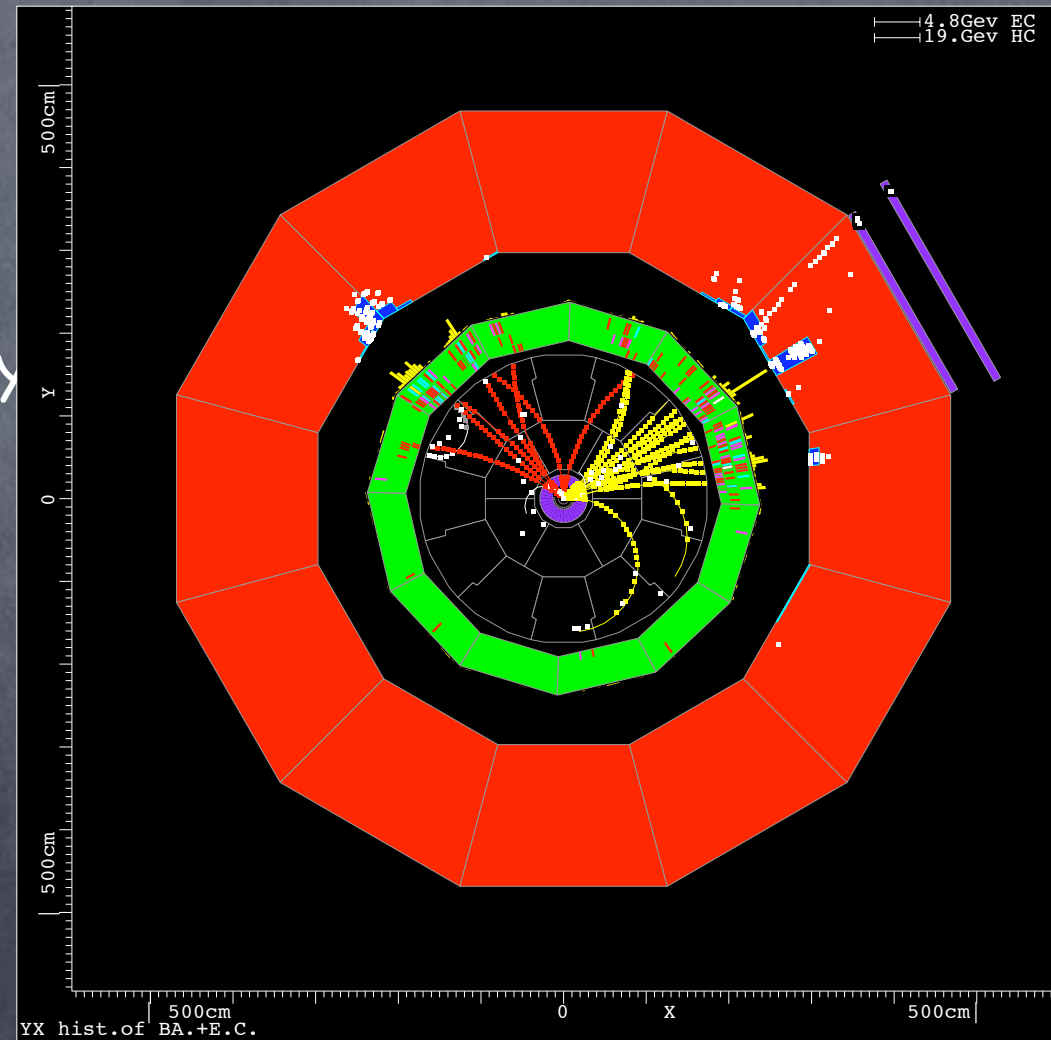
# Producing Dark Matter in the laboratory

- Collision of high-energy particles mimic Big Bang
- We hope to create Dark Matter particles in the laboratory
- Look for events where energy and momenta are unbalanced

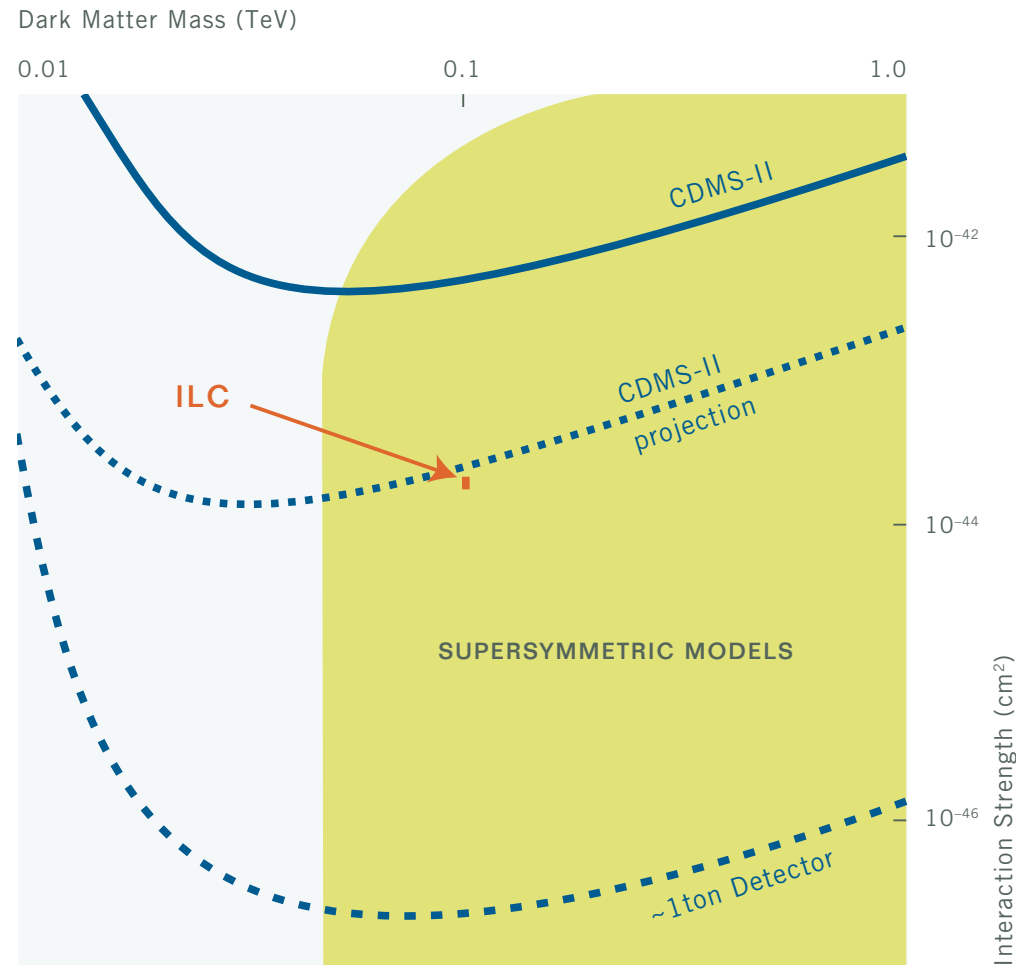
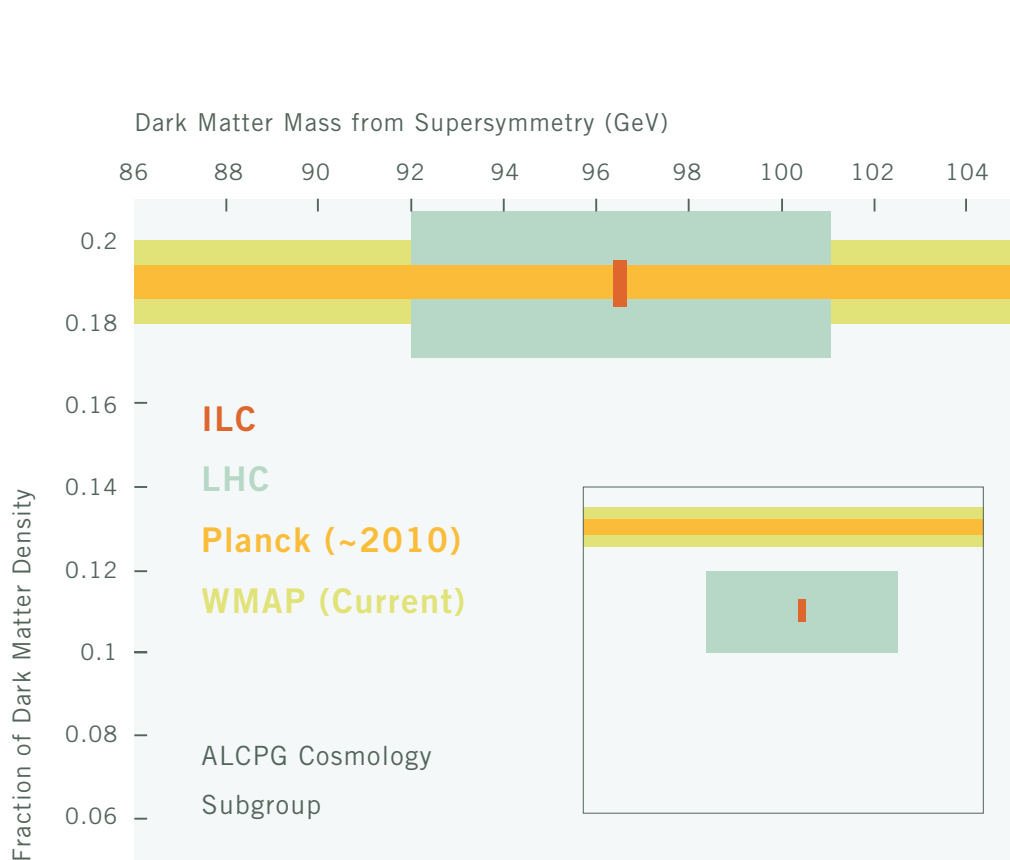
“missing energy”  $E_{\text{miss}}$

- **Something** is escaping the detector
- electrically neutral, weakly interacting

⇒ **Dark Matter!?**

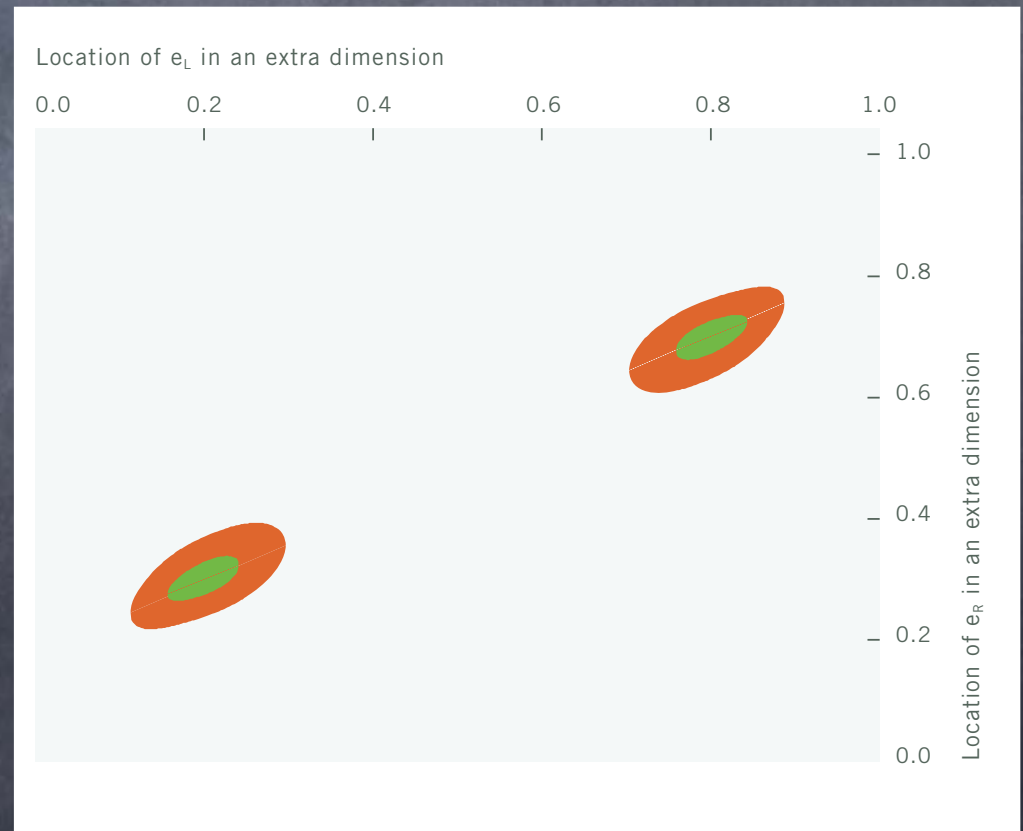
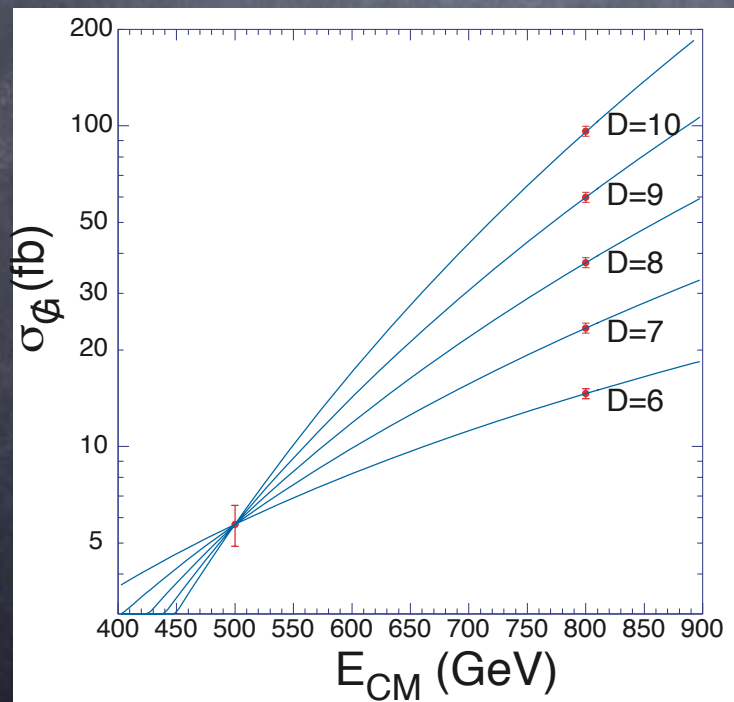
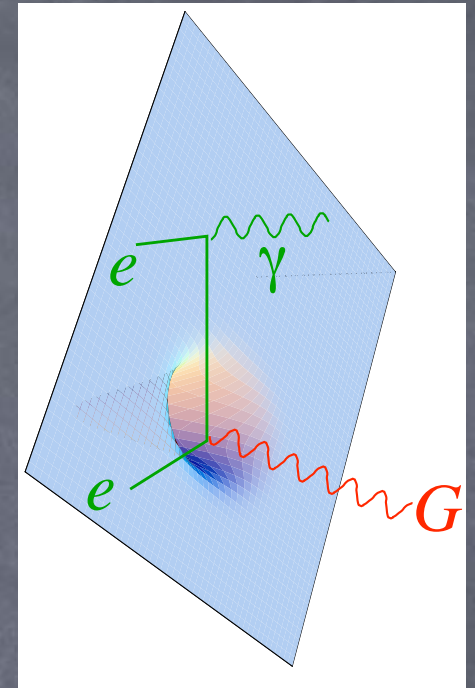


# Dark Matter



# Extra D

- measure the number of dimensions
- location of the wave functions

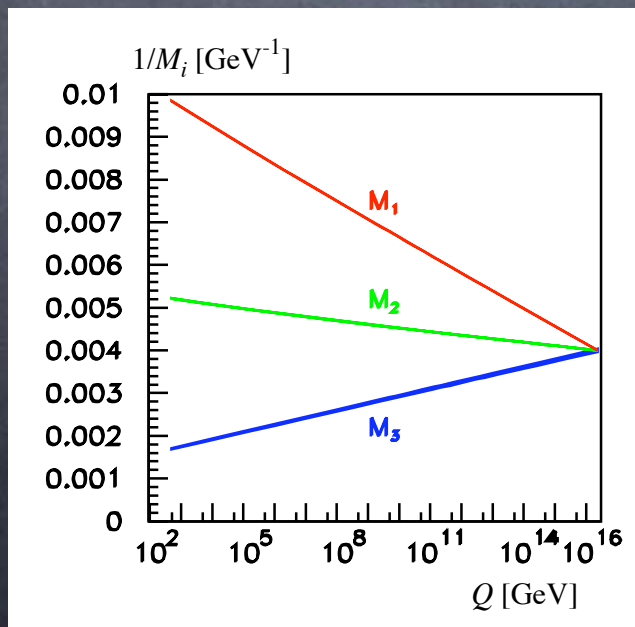




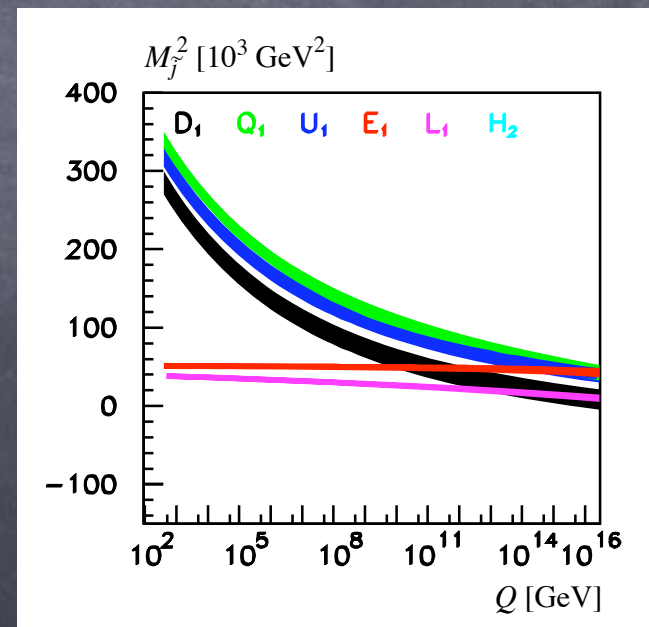
# Unification

Do the forces and matter unify?  
We know coupling constants appear  
to unify with supersymmetry

gaugino masses



scalar masses

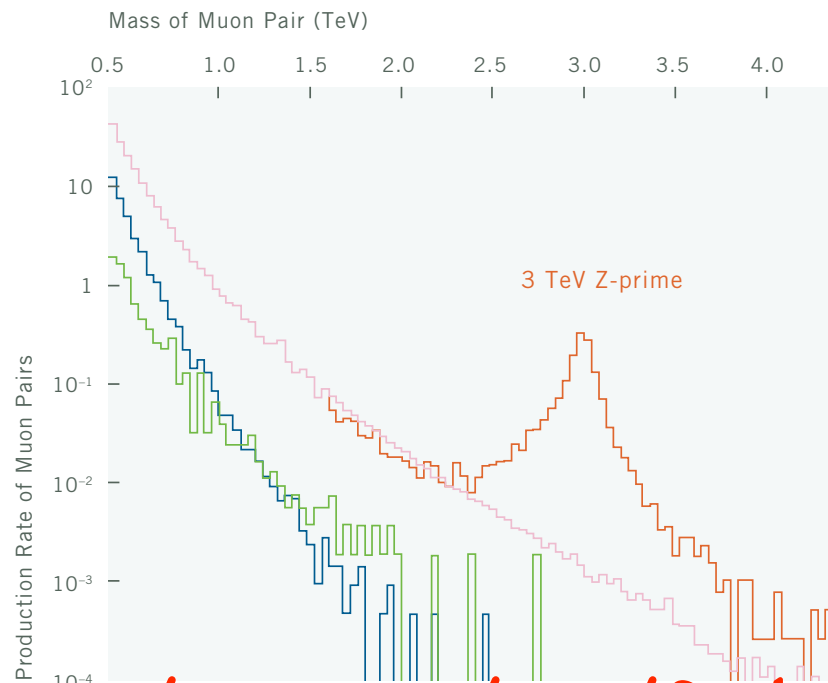


# Implications

- There is indeed unification!
- No gauge non-singlet particles below  $10^{14}\text{GeV}$
- Neutrino mass must come from gauge singlet exchange (i.e. seesaw!)
- Constraints on baryogenesis models (strong preference to leptogenesis by  $\nu_R$ ), axion models

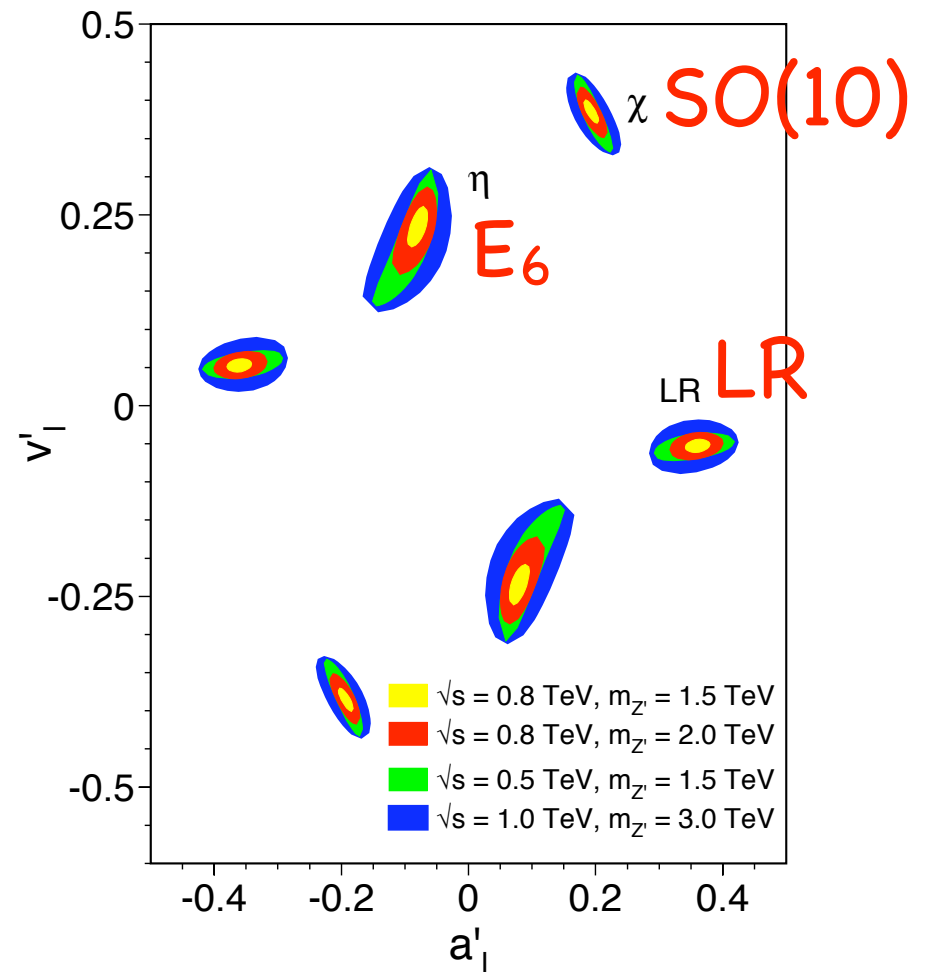
Buckley, HM

# New force: $Z'$



$\sim 1/2$  event/bin/fb $^{-1}$

What kind of force?





# Einstein's Telescope

- With both LHC and ILC, we hope to see way beyond the energy scale we can probe directly, i.e. GUT and string scales

