

Part II

Measuring neutrino parameters

Solar Neutrinos



electron neutrinos are produced

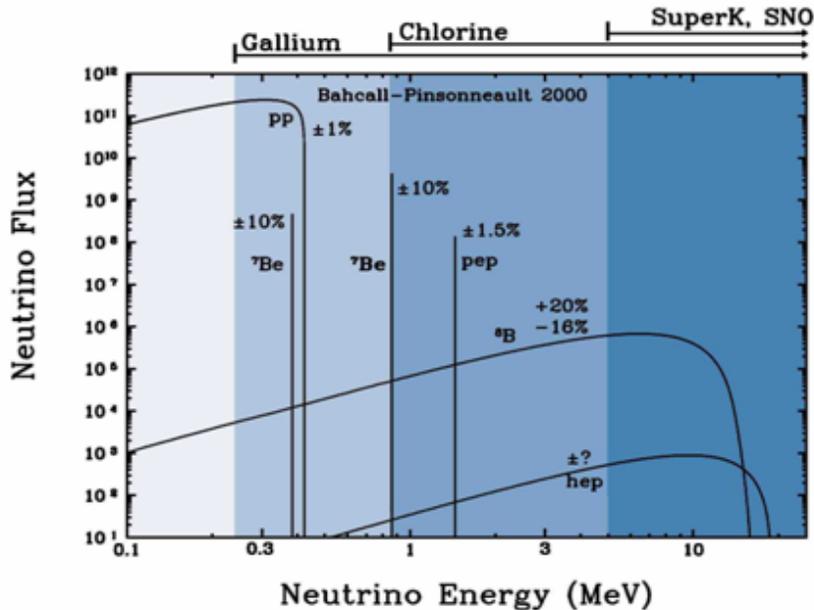
$$F = 6 \cdot 10^{10} \text{ cm}^{-2} \text{ c}^{-1}$$

total flux at the Earth



Adiabatic
conversion

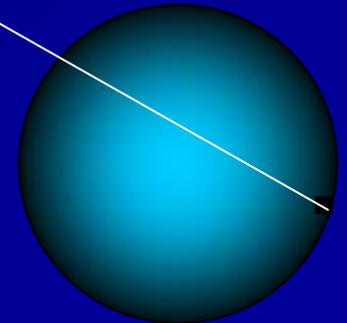
J.N. Bahcall



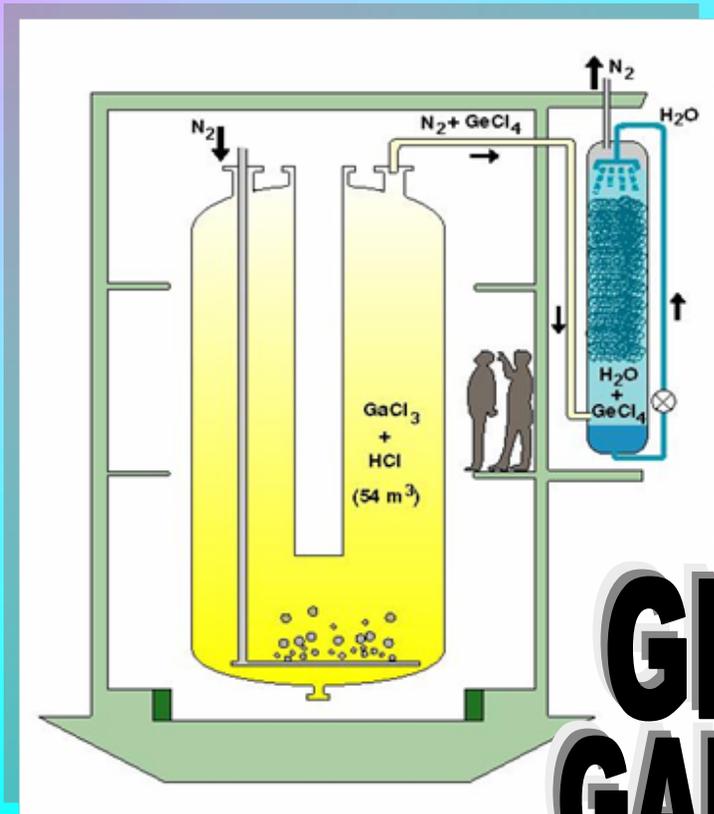
ν

$\rho : (150 \Rightarrow 0) \text{ g/cc}$

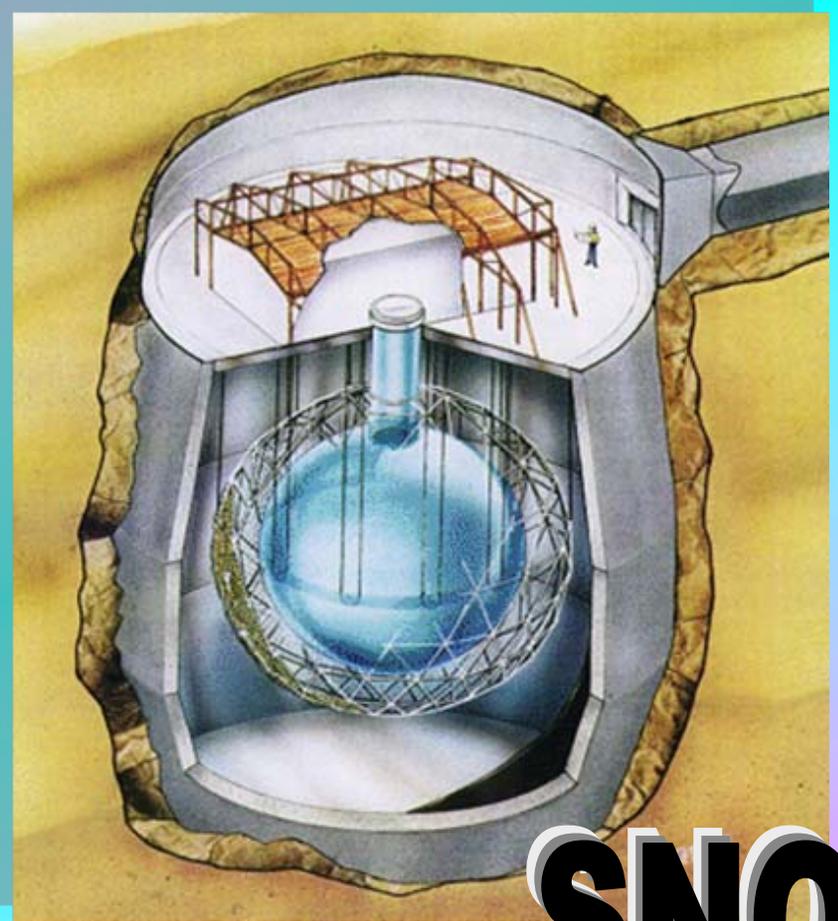
Oscillations
in matter
of the Earth



Homestake Kamiokande SAGE



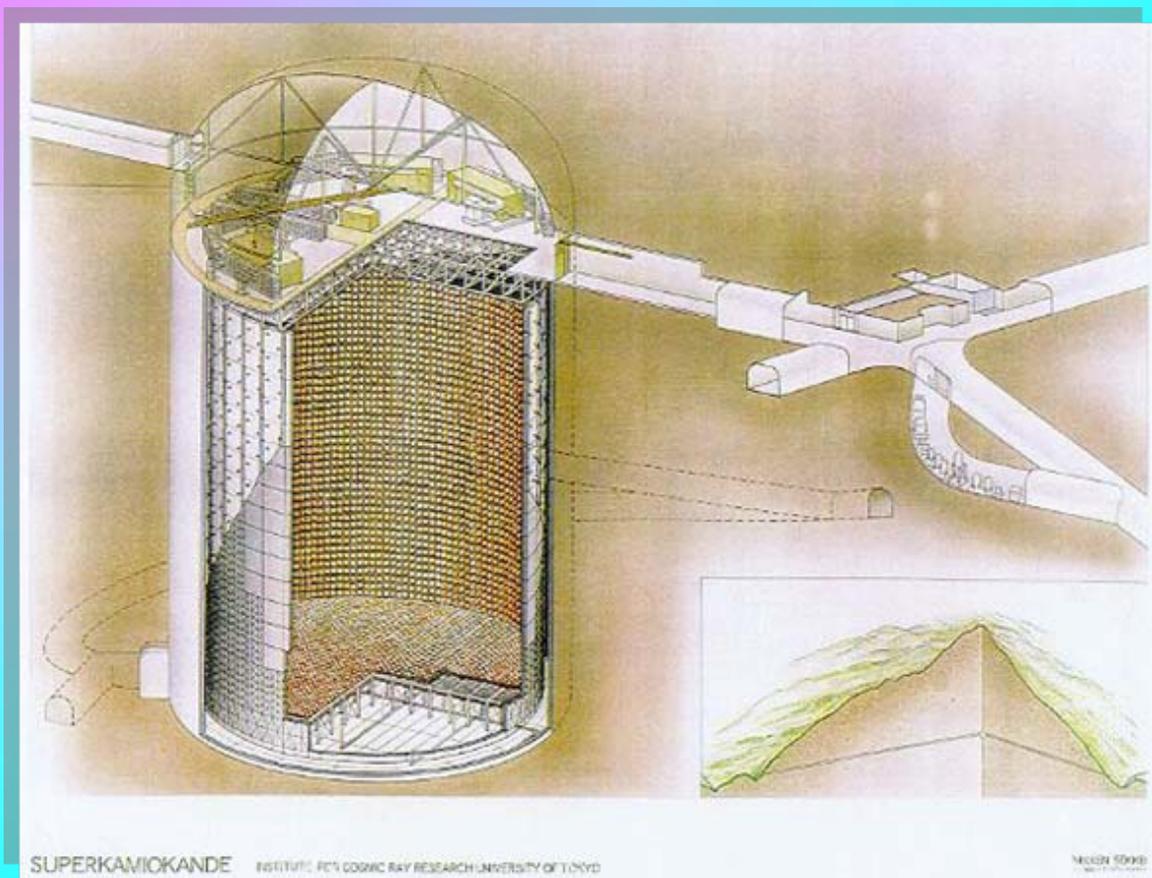
GNO
GALLEX



SNO

SuperKamiokande

SuperKamioKande



Water
Cherenkov detector

$$\nu_e \rightarrow \nu_e$$

Experiments and data

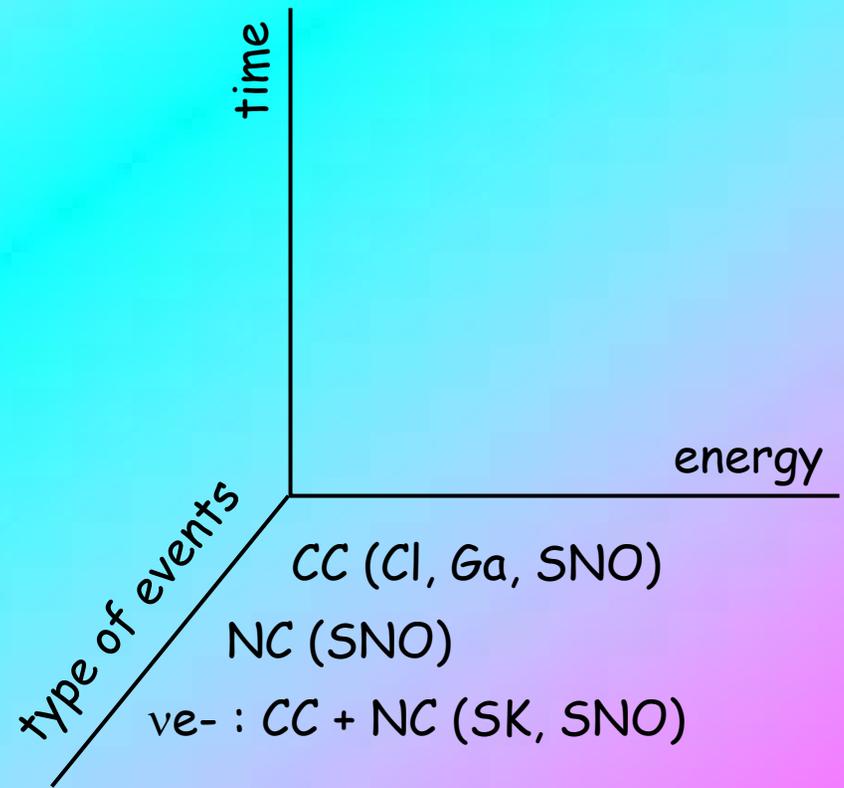
Homestake

SAGE

Galex, GNO

SuperKamiokande

SNO



Evidence of conversion

1). Deficit of signals in all experiments

2). Indirect evidence of the spectrum distortion - dependence of the suppression factor on energy:

Low energies (Ga): $R \sim 0.5 - 0.6$

High energies (Cl, SNO, SK): $R \sim 0.3$

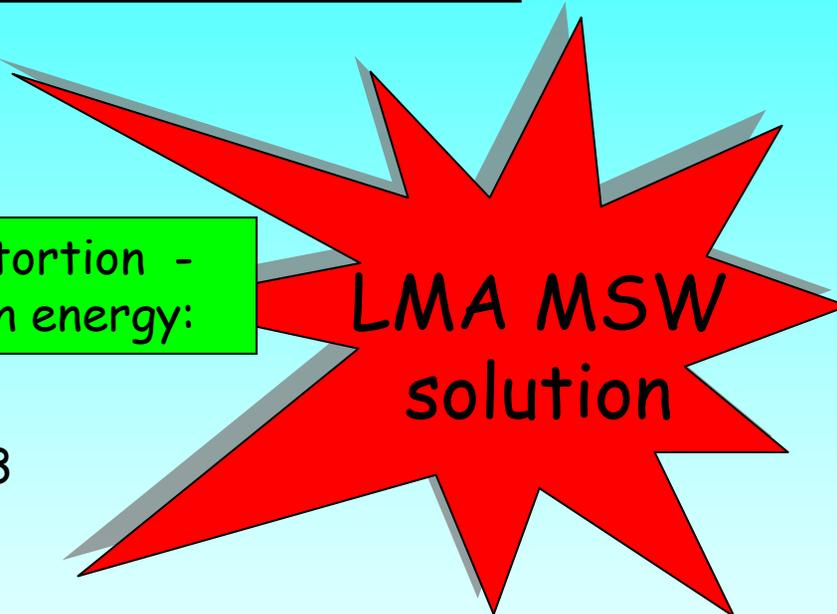
3). $CC/NC \sim 0.34$

No statistically significant

distortion of the boron neutrino spectrum (at $E > 5$ MeV) in SK, SNO

day-night effect

time variations apart from annual variations (due to eccentricity of the Earth orbit)



LMA MSW solution

Physics of conversion

1. Adiabaticity

Adiabaticity
parameter

$$\gamma(x) = \frac{l_m(x)}{4\pi h(x)} \sim 10^{-4}$$

l_m - oscillation length
in matter;
 h - height of the density
profile

- Corrections are γ^2
- Also relevant for oscillation in the Earth where affect $\sim \gamma$

2. Loss of coherence

3. Oscillations in the matter of the Earth

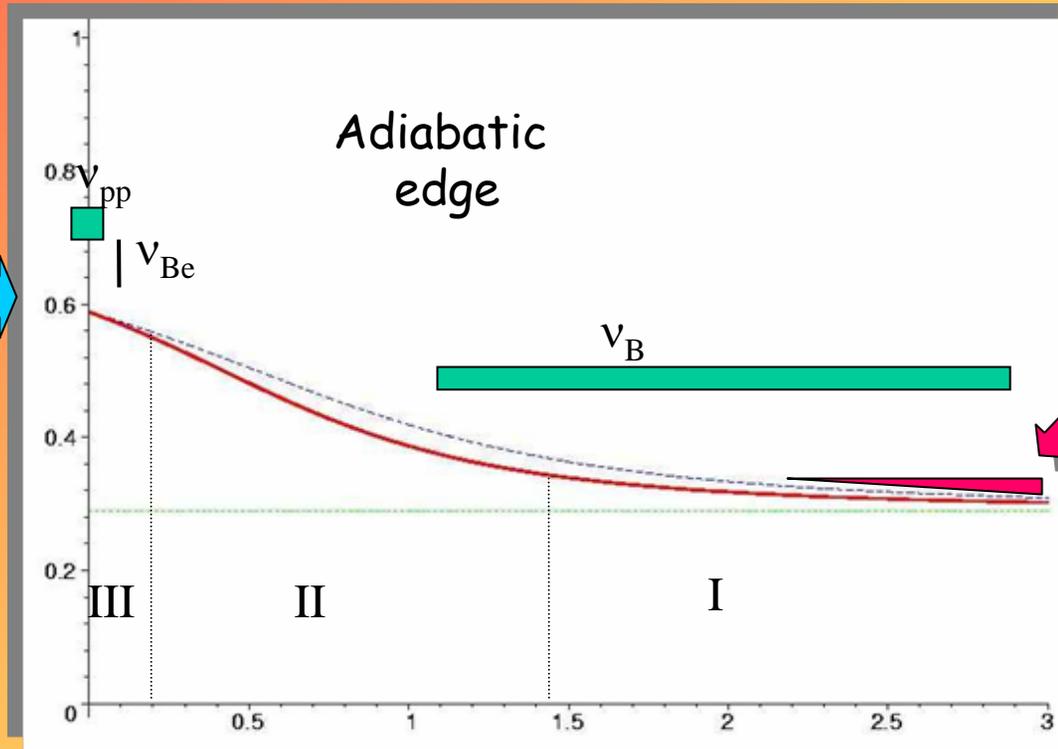
Physical picture

LMA MSW: profile of the effect

Synthetic solution

Survival probability

$1 - 0.5\sin^2 2\theta$



Earth matter effect

$\sin^2\theta$

$l_v/l_0 \sim E$

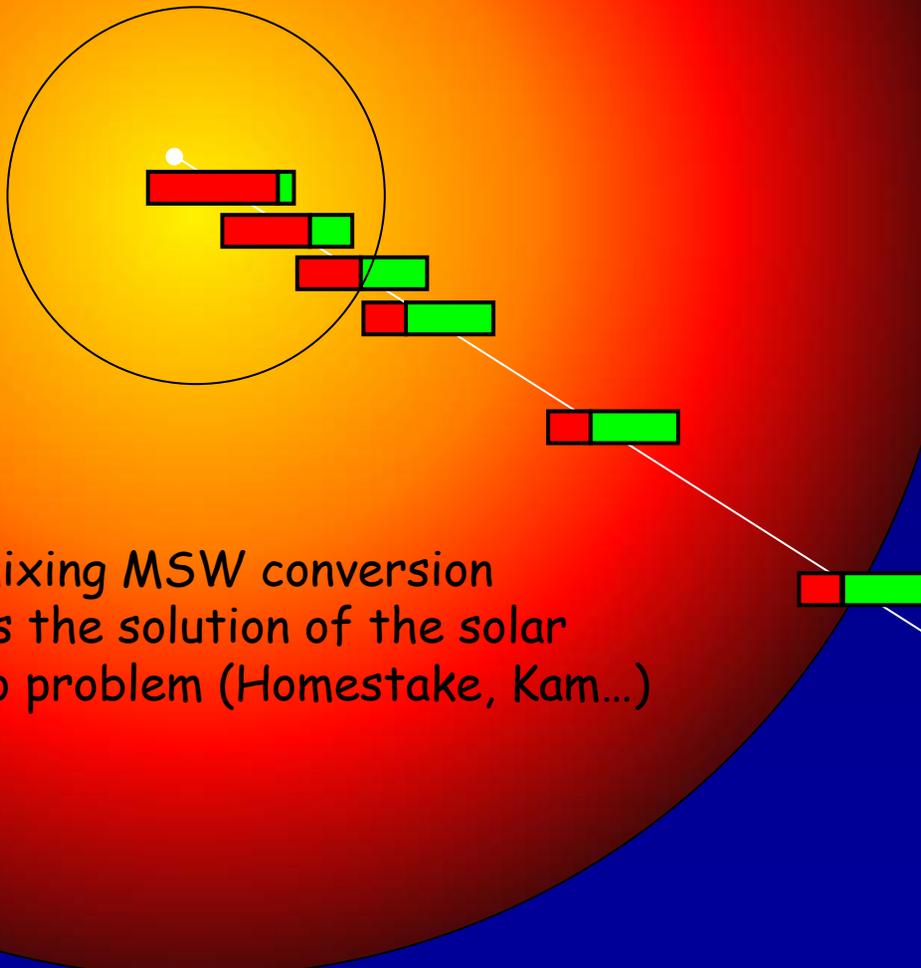
Averaged oscillations with small matter effect

Conversion + oscillations

Conversion with small oscillation effect

Non-oscillatory transition

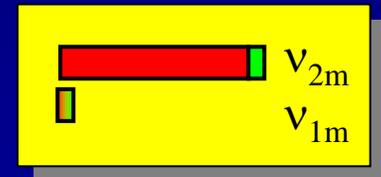
Solar neutrinos



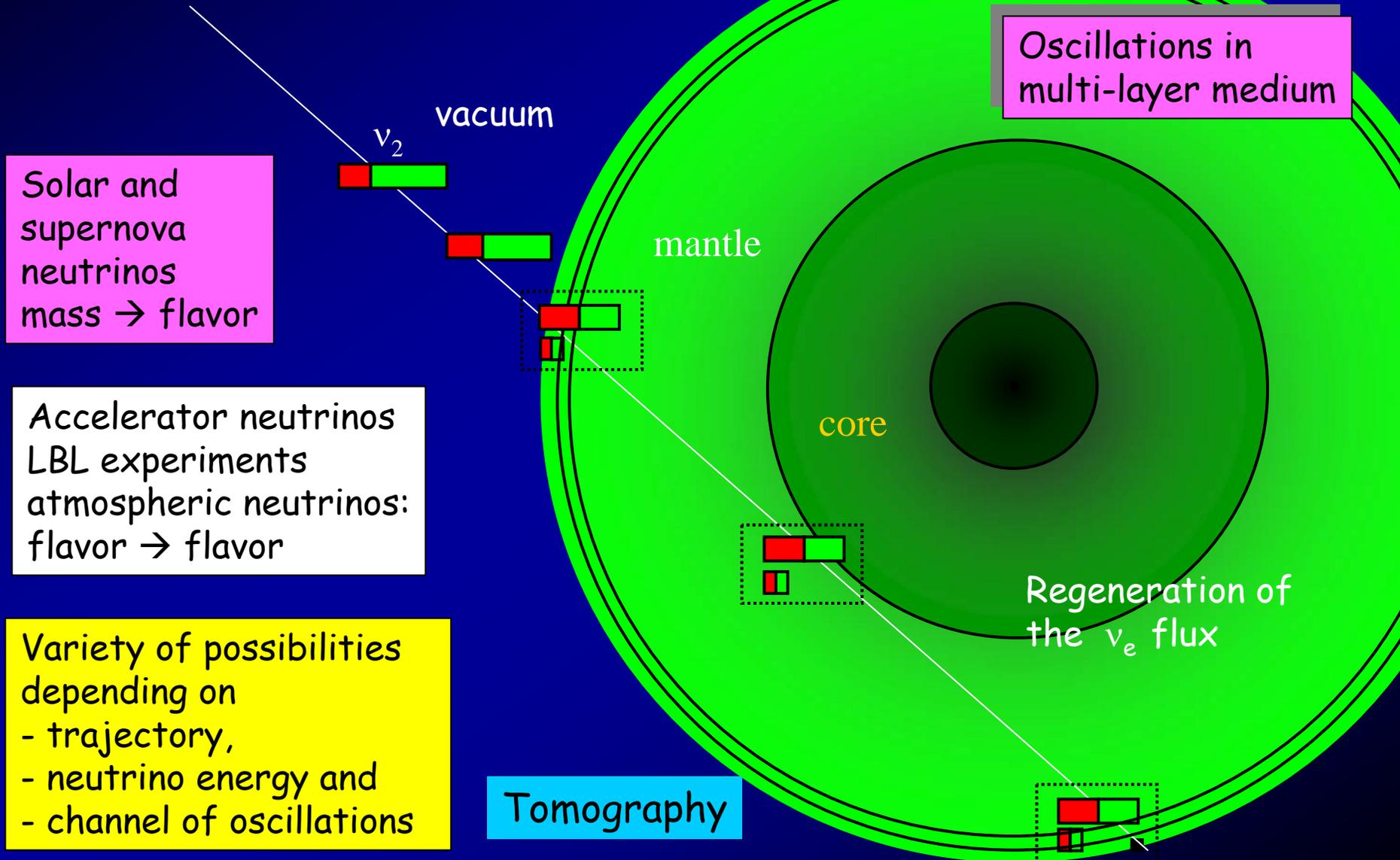
Large mixing MSW conversion provides the solution of the solar neutrino problem (Homestake, Kam...)

$E = 10 \text{ MeV}$
Resonance layer:
 $n_R Y_e = 20 \text{ g/cc}$
 $R_R = 0.24 R_{\text{sun}}$

In the production point:
 $\sin^2 \theta_m^0 = 0.94$
 $\cos^2 \theta_m^0 = 0.06$



Oscillations inside the Earth

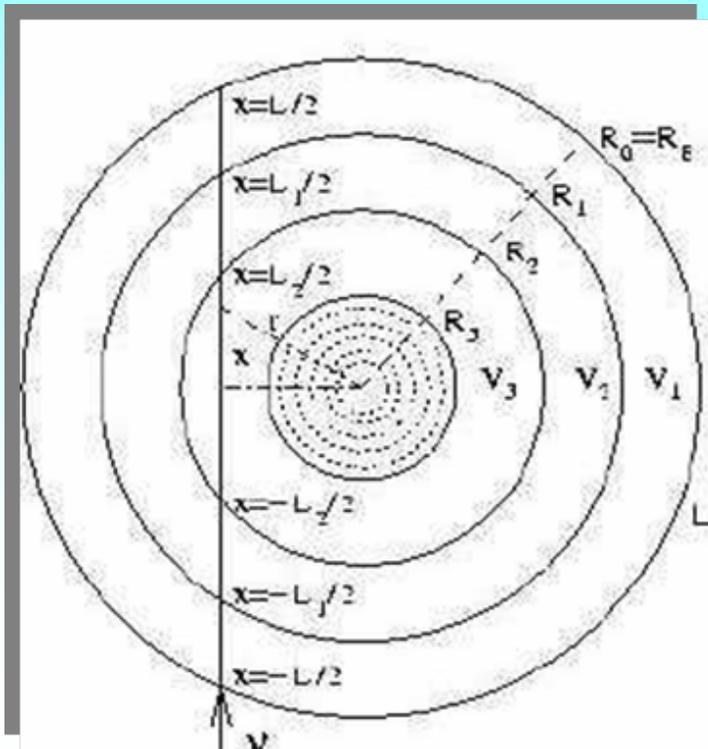


Oscillations inside the Earth

1). Incoherent fluxes of ν_1 and ν_2 arrive at the surface of the Earth

2). In matter the mass states oscillate

3). the mass-to-flavor transitions, e.g. $\nu_2 \rightarrow \nu_e$ are relevant



Regeneration factor:

$$P_{2e} = \sin^2\theta + f_{\text{reg}}$$

$$P_{ee} = 0.5[1 + \cos 2\theta_m^0 \cos 2\theta] - \cos 2\theta_m^0 f_{\text{reg}}$$

4). The oscillations proceed in the weak matter regime:

$$\varepsilon(x) = \frac{2EV(x)}{\Delta m^2} \ll 1$$

Regeneration factor

$$f_{\text{reg}} = P_{2e} - \sin^2\theta$$

$$P(\nu_2 \rightarrow \nu_e) = |\langle \nu_e | U(\theta_{\text{mR}}) S(x_0 \rightarrow x_f) U^\dagger(\theta_{\text{mR}}) U(\theta) | \nu_2 \rangle|^2$$

θ_{mR} - mixing angle at the surface of the Earth

Adiabatic perturbation theory

Liao Wei, P de Holanda, A.S.

$$f_{\text{reg}} = \varepsilon(R) \sin^2 2\theta \sin^2 [\Phi^{\text{m}}(x_0 \rightarrow x_f)/2] + \sin 2\theta \text{Re}\{c(x_0 \rightarrow x_f)\}$$

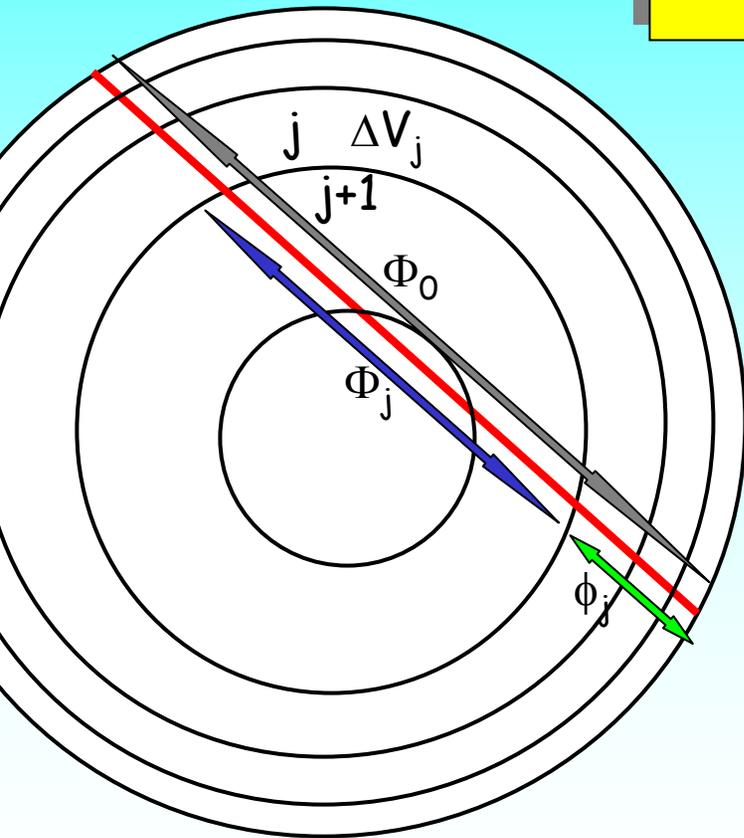
amplitude of
jump probability

$$\varepsilon(R) = \frac{2EV(R)}{\Delta m^2}$$

- If adiabaticity is realized, the regeneration depends on the potential $V(R)$ at the surface and on the total adiabatic phase
- Non-adiabatic conversion appears as the interference term and therefore - linearly (in contrast to conversion in the Sun)

Analytic result

$$f_{\text{reg}} = \frac{2E \sin^2 2\theta}{\Delta m^2} \sin \Phi_0 / 2 \sum_{j=0 \dots n-1} \Delta V_j \sin \Phi_j / 2$$



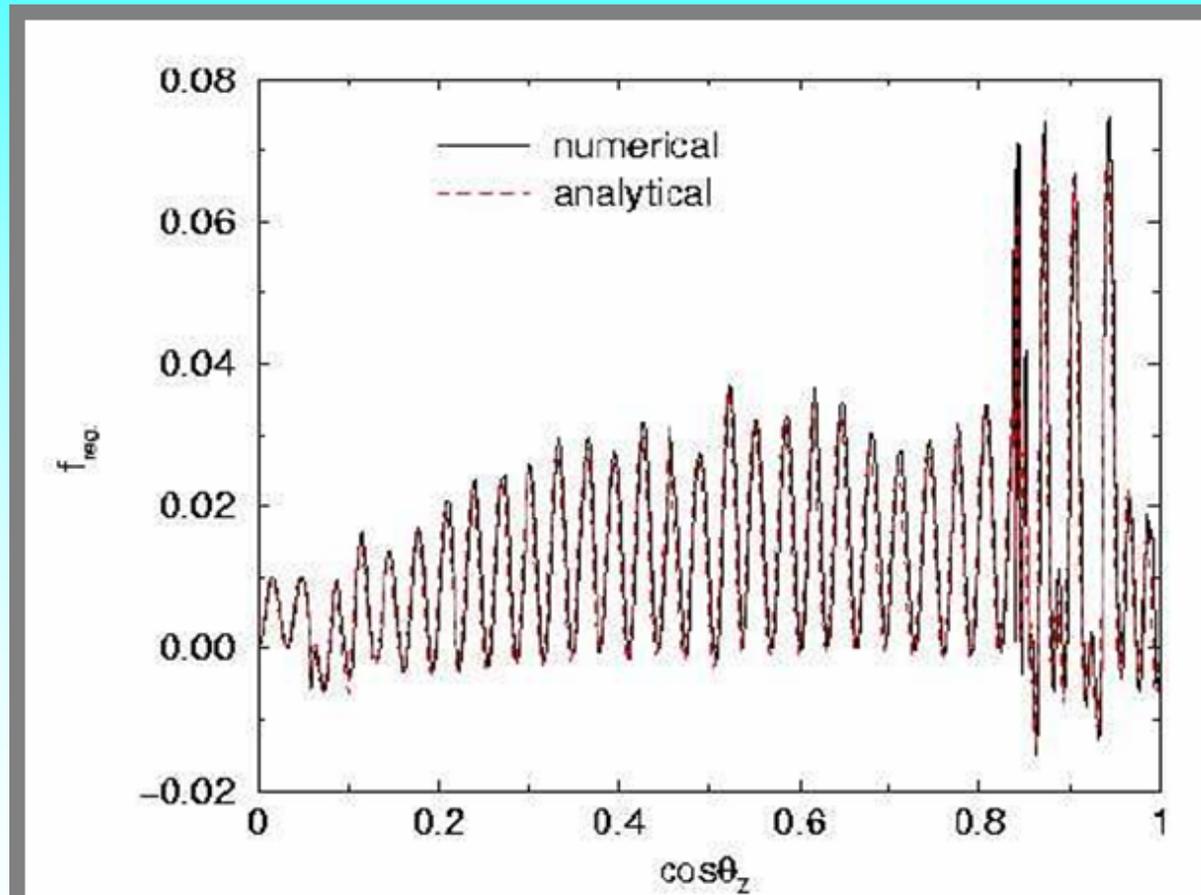
Defining $\phi_j = 0.5(\Phi_0 - \Phi_j)$

$$f_{\text{reg}} = \frac{2E \sin^2 2\theta}{\Delta m^2} \times$$

$$\sum_{j=0 \dots n-1} \Delta V_j [\sin^2 \Phi_0 / 2 \cos \phi_j - 0.5 \sin \Phi_0 \sin \phi_j]$$

If ϕ_j is large - averaging effect.
This happens for remote structures, e.g. core

Analytic vs. numerical results



*A. Ioanissian,
A.S.*

Regeneration factor as function of the zenith angle
 $E = 10 \text{ MeV}$, $\Delta m^2 = 6 \cdot 10^{-5} \text{ eV}^2$, $\tan^2\theta = 0.4$

Attenuation effect

Important for the oscillation tomography of the Earth with solar and atmospheric neutrinos

Related to finite energy resolution of detectors
→ integration (averaging) over the energy resolution intervals

Neutrino microscopes

Integral formula

$$P_{2e} = \sin^2\theta + f_{\text{reg}}$$

ε - perturbation theory

Regeneration factor

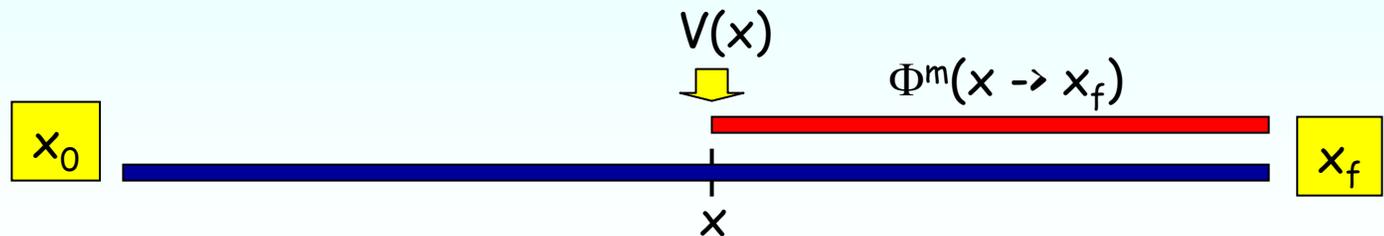
$$f_{\text{reg}} = 0.5 \sin^2 2\theta \int_{x_0}^{x_f} dx V(x) \sin \Phi^m(x \rightarrow x_f)$$

A. Ioannisian, A.S.

Explicitly:

$$f_{\text{reg}} = 0.5 \sin^2 2\theta \int_{x_0}^{x_f} dx V(x) \sin \left[\frac{\Delta m^2}{2E} \int_x^{x_f} dy \sqrt{\left[\cos 2\theta - \frac{2EV(y)}{\Delta m^2} \right]^2 - \sin^2 2\theta} \right]$$

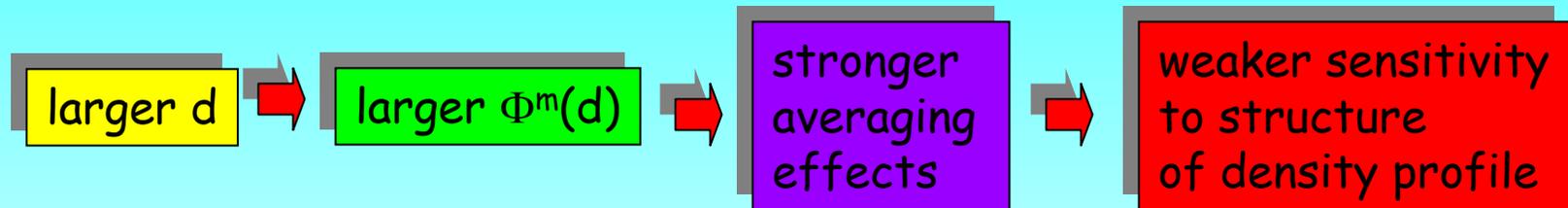
Integration limits:



The phase is integrated from a given point to the final point

Sensitivity to density profile

For mass-to-flavor transition $V(x)$ is integrated with $\sin \Phi^m(d)$
 $d = x_f - x$ the distance from structure to the detector



Integration with the energy resolution function $R(E, E')$:

$$\overline{f_{\text{reg}}} = \int dE' R(E, E') f_{\text{reg}}(E')$$

The effect of averaging:

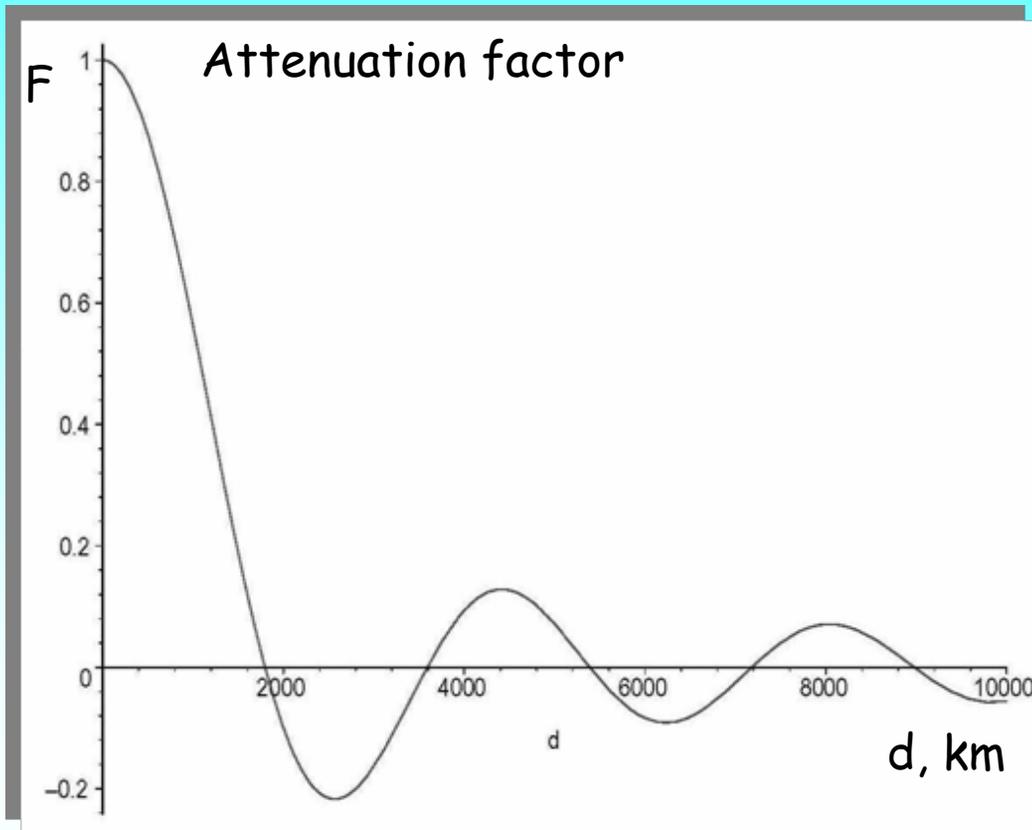
$$\overline{f_{\text{reg}}} = 0.5 \sin^2 2\theta \int_{x_0}^{x_f} dx V(x) F(x_f - x) \sin \Phi^m(x \rightarrow x_f)$$

← averaging factor

For box-like $R(E, E')$ with width ΔE :

$$F(d) = \frac{I_\nu E}{\pi d \Delta E} \sin \left(\frac{\pi d \Delta E}{I_\nu E} \right)$$

Attenuation effect



The width of the first peak

$$d < l_v E / \Delta E$$

l_v is the oscillation length

The sensitivity to remote structures is suppressed:

- ➔ Effect of the core of the Earth is suppressed
- ➔ Small structures at the surface can produce stronger effect
- ➔ The better the energy resolution, the deeper penetration

Attenuation length

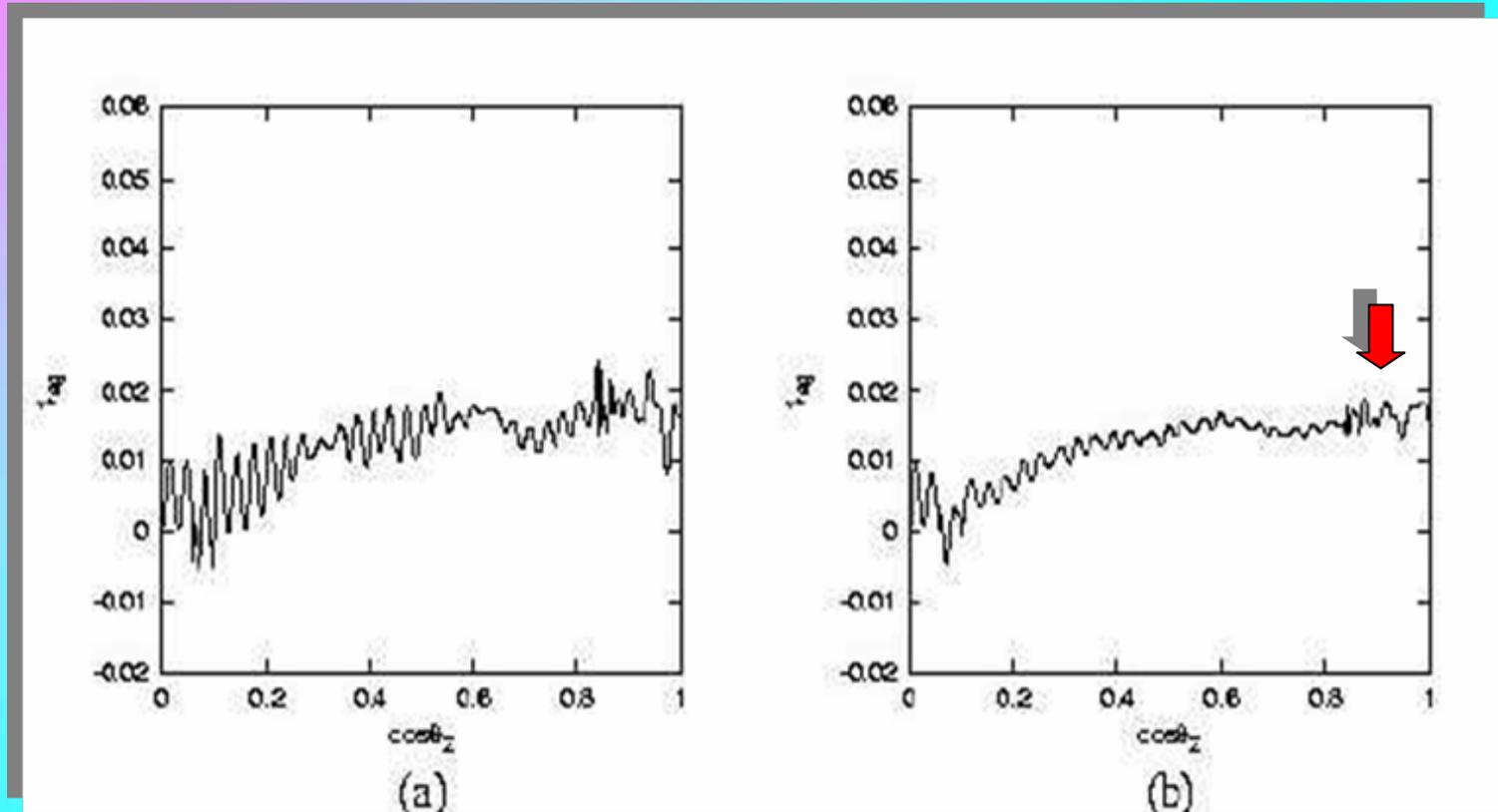
$$d = l_{\nu} E / \Delta E$$

$$d \sim 4\pi E E / \Delta E \Delta m^2$$

	l_{ν}	$\Delta E/E$	$d, \text{ km}$	Core
Solar neutrinos $E = 10 \text{ MeV}$	300 km	10 - 20%	1500 - 3000	can not be seen
Supernova neutrinos $E = 30 \text{ MeV}$	900 km	10%	9000	can be seen

l_{ν} is the oscillation length

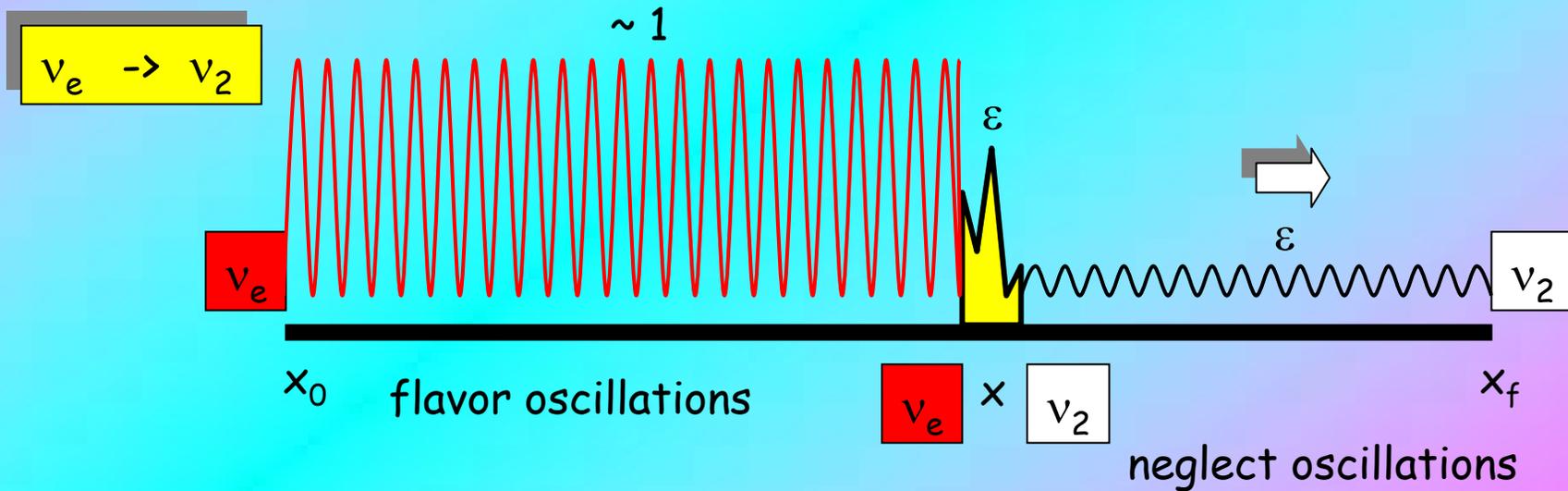
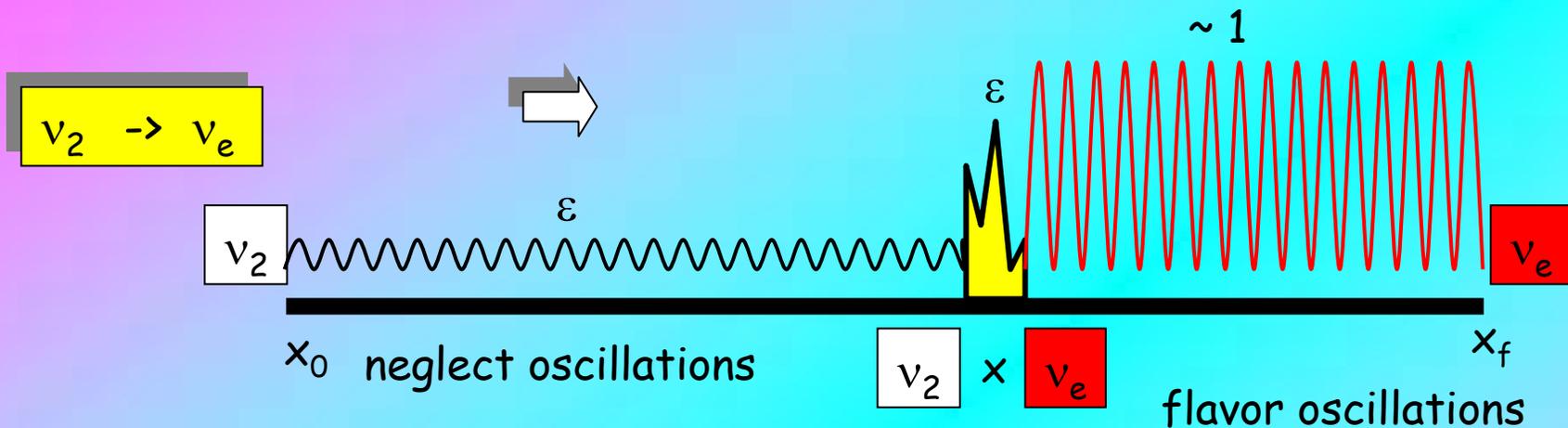
Averaging regeneration factor



Regeneration factor averaged over the energy intervals $E = (9.5 - 10.5) \text{ MeV}$ (a), and $E = (8 - 10) \text{ MeV}$ (b).

No enhancement for core crossing trajectories in spite of larger densities

In ϵ order



LMA MSW

SNO Collaboration

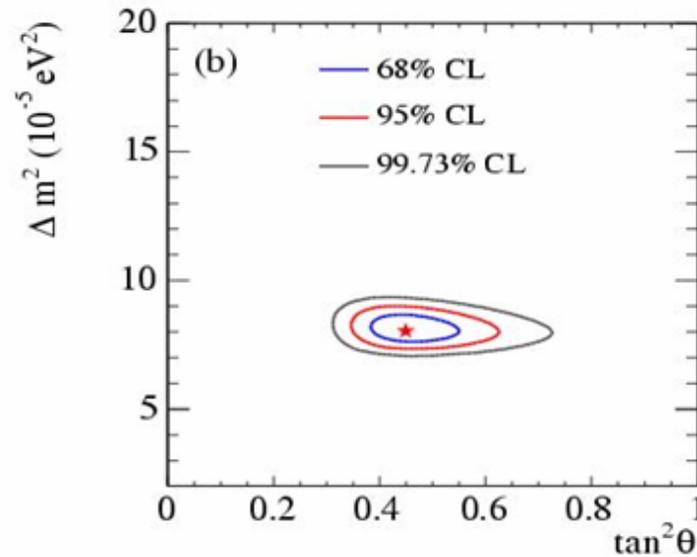
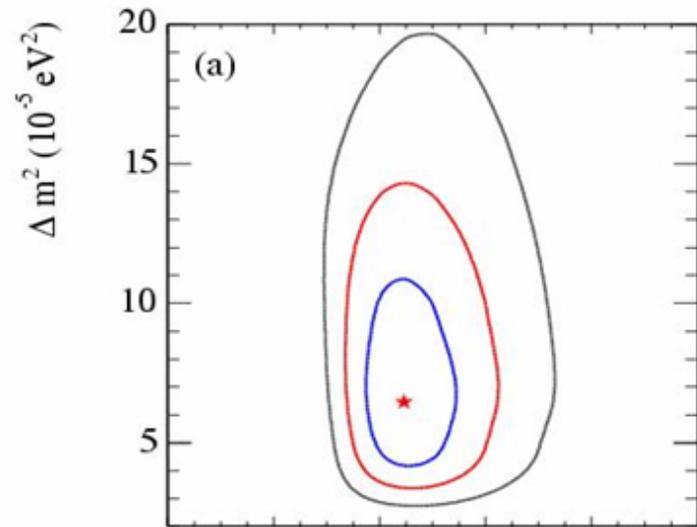
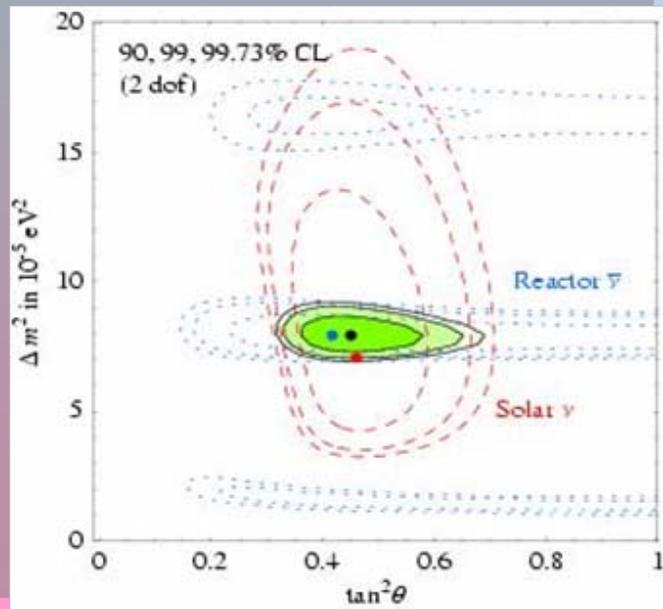
Global
Solar:

$$\Delta m^2 = 6.5 \cdot 10^{-5} \text{ eV}^2$$
$$\tan^2\theta = 0.45_{-0.08}^{+0.09}$$

Solar
+ KL:

$$\Delta m^2 = 7.9 \cdot 10^{-5} \text{ eV}^2$$
$$\tan^2\theta = 0.45$$

A. Strumia, F. Vissani



Stability of bfp

Status and perspectives

Expect new SNO data: phase-3 ^3He counters of neutrons

better identification of the NC events
-> better measurements of CC/NC ratio
-> better determination of θ_{13}

BOREXINO, KamLAND: measurements of fluxes in the intermediate energy range (Be neutrino flux)

SAGE calibration result: $\text{obs/exp} = 0.79 \pm 0.09 / - 0.10$

Lower cross-section - higher pp-flux
-> larger survival probability? Tension?

Monitoring of the solar neutrino flux
-> searches for time variations, periodicity

Analysis of data:
periodicity

Solar models, abundance of heavy elements at the surface??
Boron neutrino cross-section

KamLAND

Kamioka Large Anti-Neutrino Detector

- Reactor long baseline experiment
150 - 210 km

Liquid scintillation detector

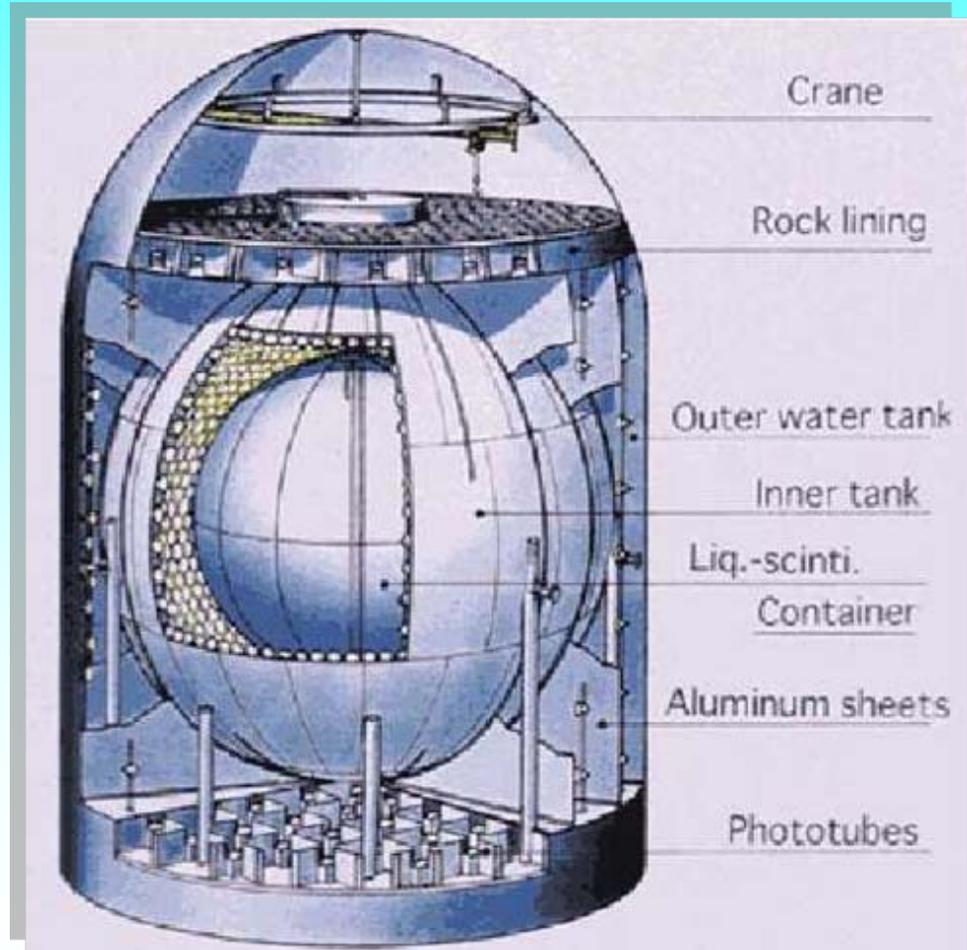


$$E_{pr} > 2.6 \text{ MeV}$$

- Data: total rate
energy spectrum of events

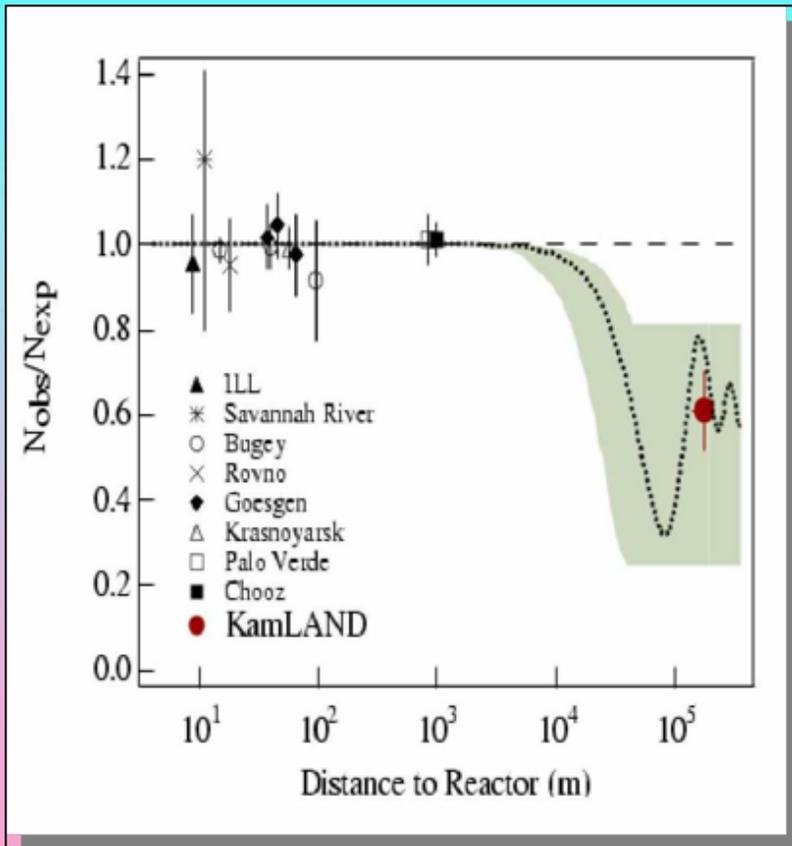
- LMA 
precise determination of
the oscillation parameters
10% accuracy

- Detection of the Geo-neutrinos
 $E_{pr} > 1.3 \text{ MeV}$



1 kton of LS

KamLAND results



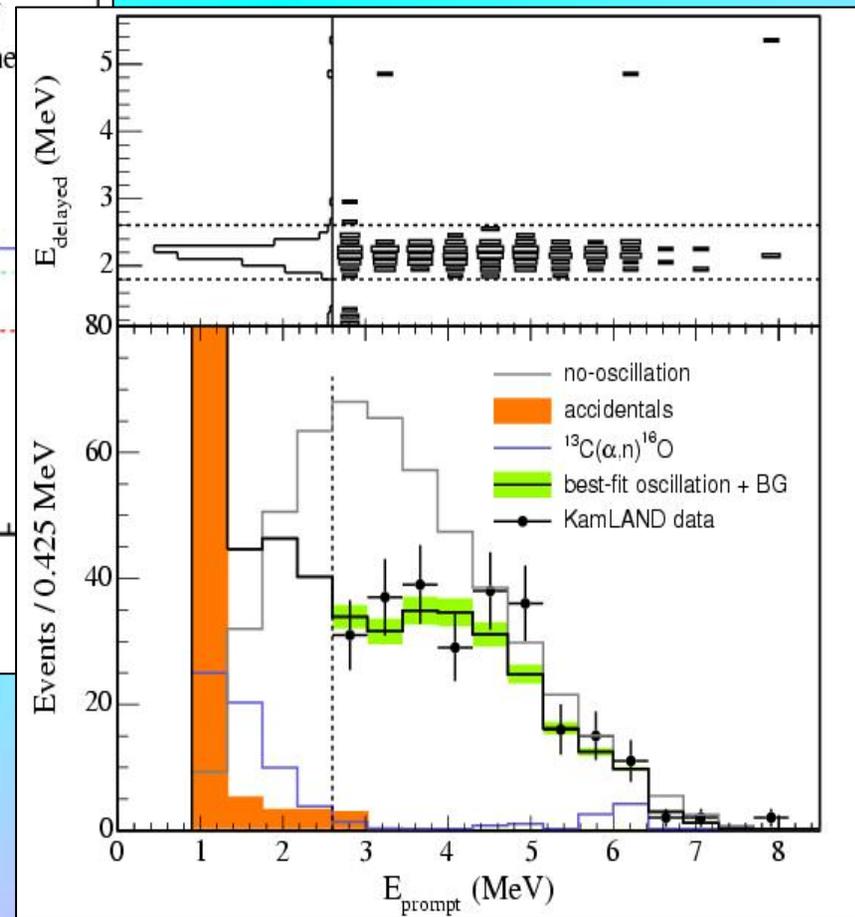
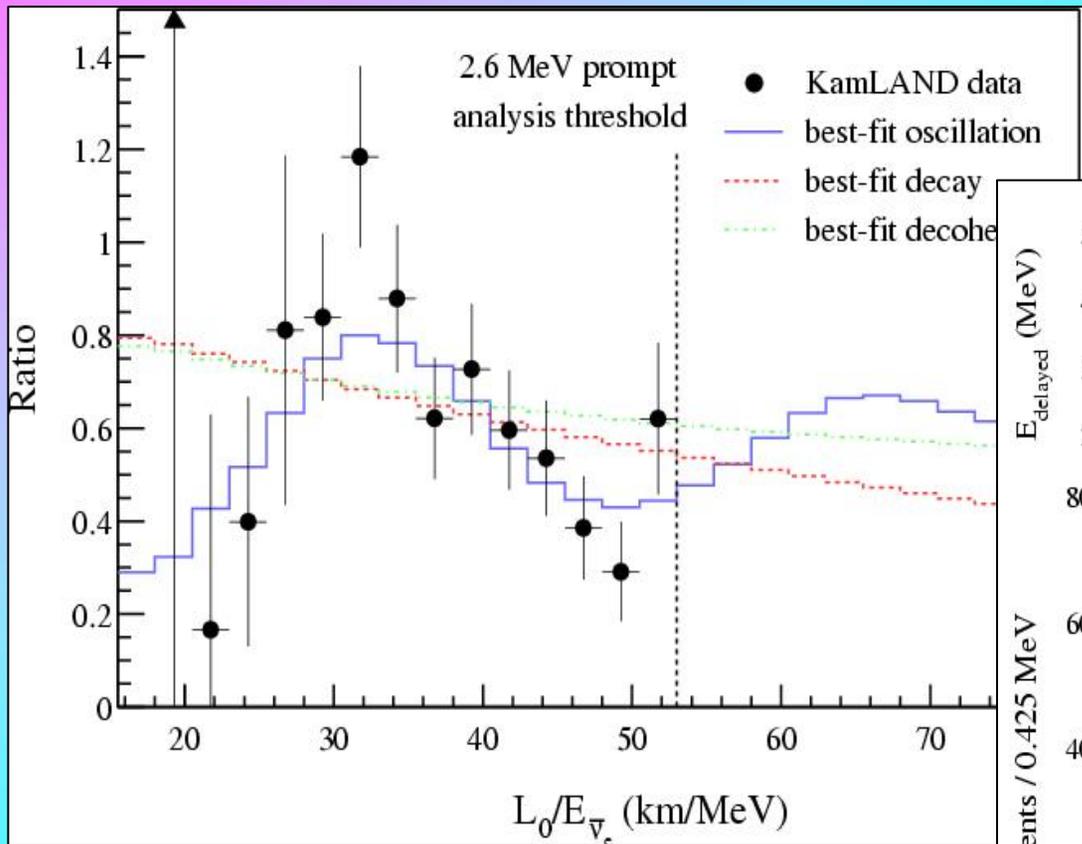
Δm^2 is determined from
phase of oscillations
(energy spectrum distortion)

According to LMA the range of Δm^2
is restricted

- from above - by spectrum distortion
(up turn)
- from below by increasing day-night
asymmetry

No dependence on phase

KamLAND results



PRL 94, 081801 (2005)

Testing Theory of oscillations

Solar neutrinos

- Adiabatic conversion (MSW)
 - Matter effect dominates (at least in the HE part)
 - Non-oscillatory transition, or averaging of oscillations
- the oscillation phase is irrelevant

Adiabatic conversion formula

KamLAND

- Vacuum oscillations
- Matter effect is very small
- Oscillation phase is crucial for observed effect

Vacuum oscillations formula

$\Delta m^2, \theta$

Coincidence of these parameters determined from the solar neutrino data and from KamLAND results testifies for the correctness of the theory (phase of oscillations, matter potential, etc..)

Testing the MSW effect in the Sun

SK+SNO:

$$\langle P_{ee} \rangle = 0.34 \pm 0.12 (3\sigma)$$

$$\langle P_{ee} \rangle < 0.46 (3\sigma)$$

2ν - vacuum oscillations:

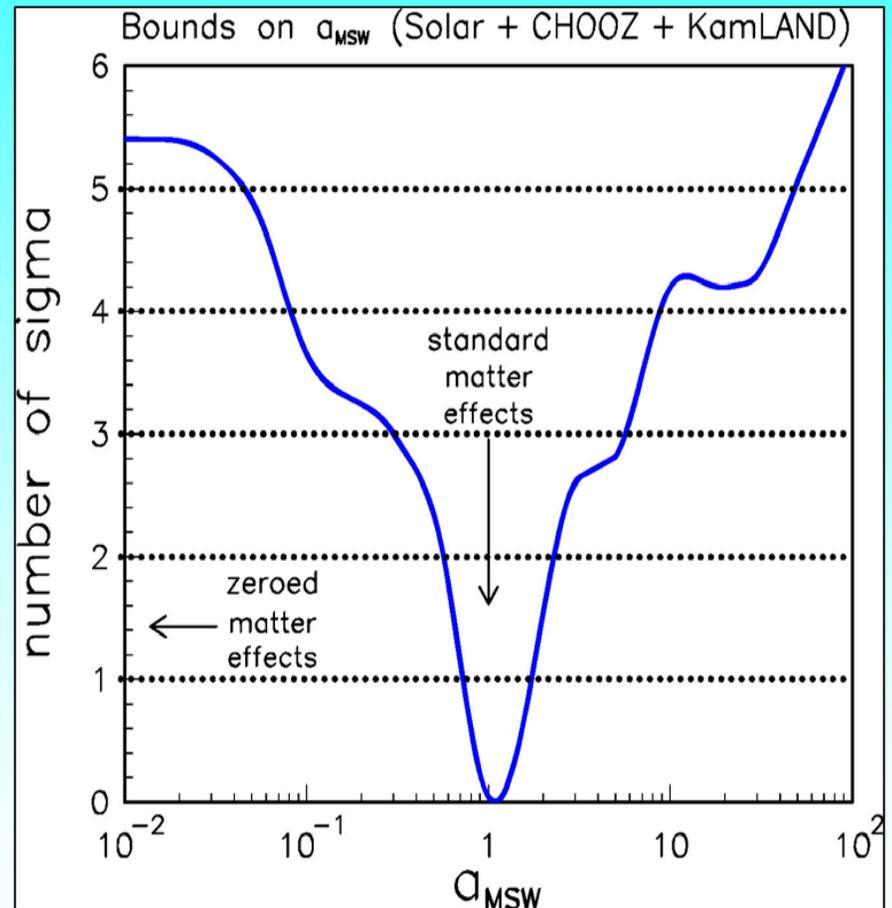
$$\langle P_{ee} \rangle > 0.5$$

$$V \rightarrow a_{MSW} V$$

free parameter

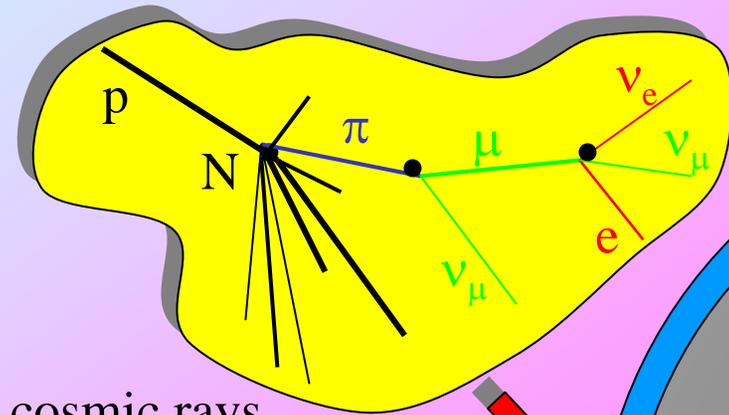
Global analysis of solar and reactor (KamLAND+CHOOZ) data with Δm_{12}^2 , $\sin^2\theta_{12}$, a_{MSW} unconstrained:

G. L. Fogli, E. Lisi, A. Marrone, A. Palazzo, A. M. Rotunno, hep-ph/0506307



Notice: in the best fit point $a_{MSW} > 1$ reflects mismatch of the b.f. values from solar and KamLAND analysis

Atmospheric neutrinos

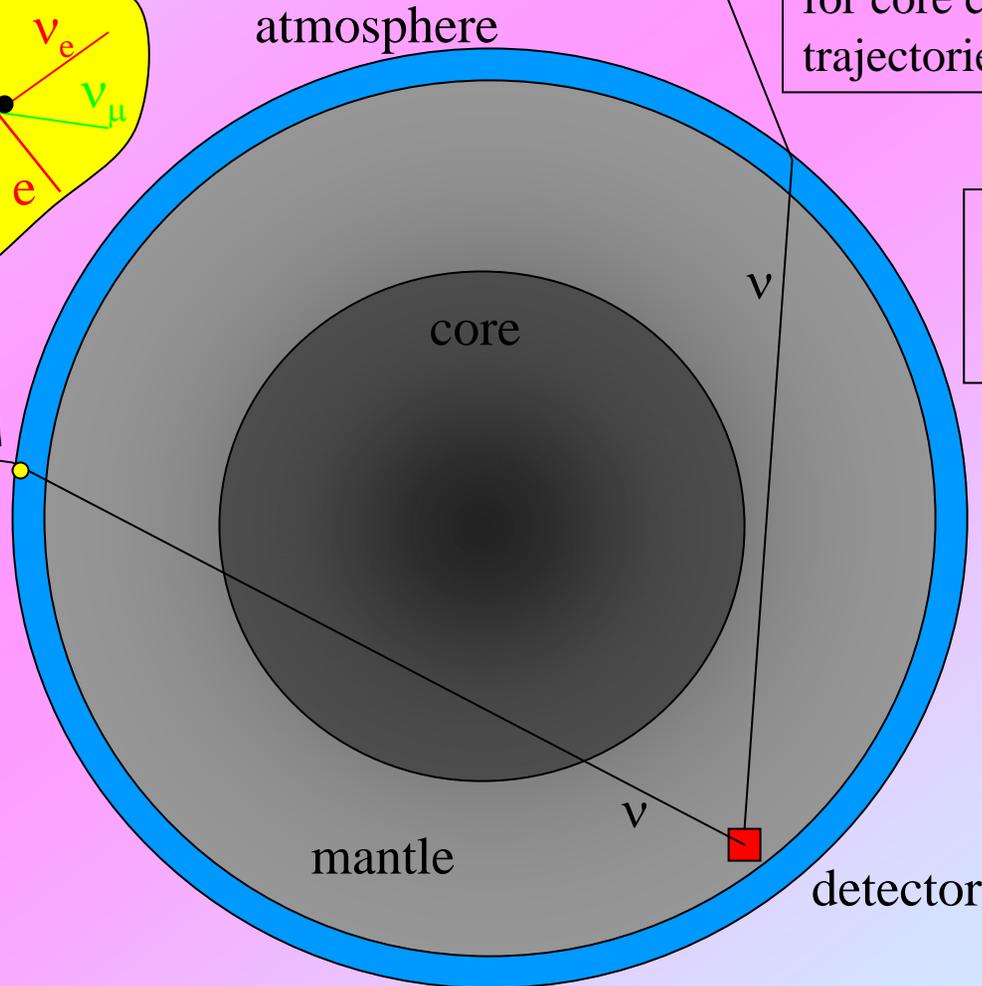


cosmic rays

At low energies:

$$r = F_{\mu} / F_e = 2$$

$\nu_{\mu} - \nu_{\tau}$
vacuum oscillations



Parametric effects
in $\nu_{\mu} - \nu_e$ oscillations
for core crossing
trajectories

$\nu_{\mu} - \nu_e$
oscillations
in matter

detector

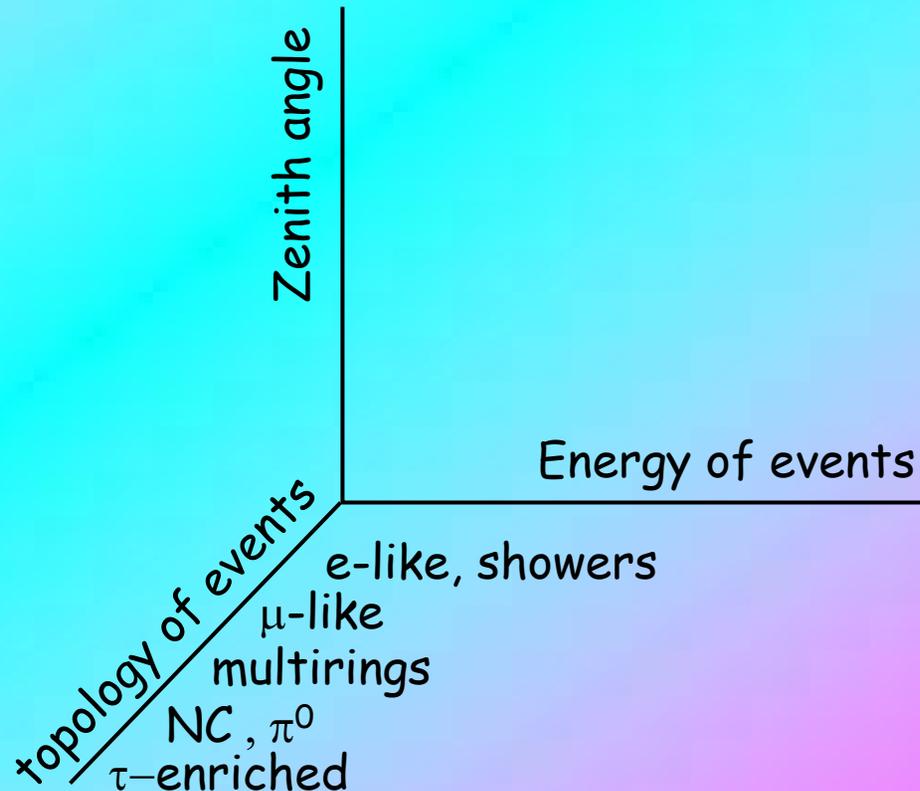
Experiments and data

SuperKamiokande

MACRO

SOUDAN

MINOS



Evidence of oscillations

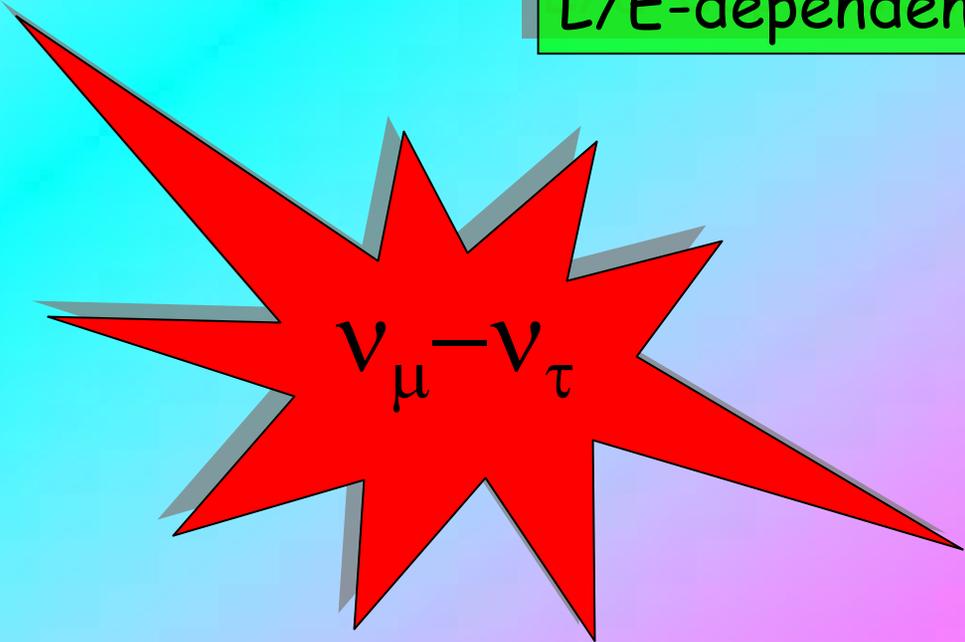
Zenith angle
distribution
of μ -like events

Appearance
of τ -like events

L/E-dependence

Up-down
asymmetry

Ratio of numbers
of μ -like to e-like
events

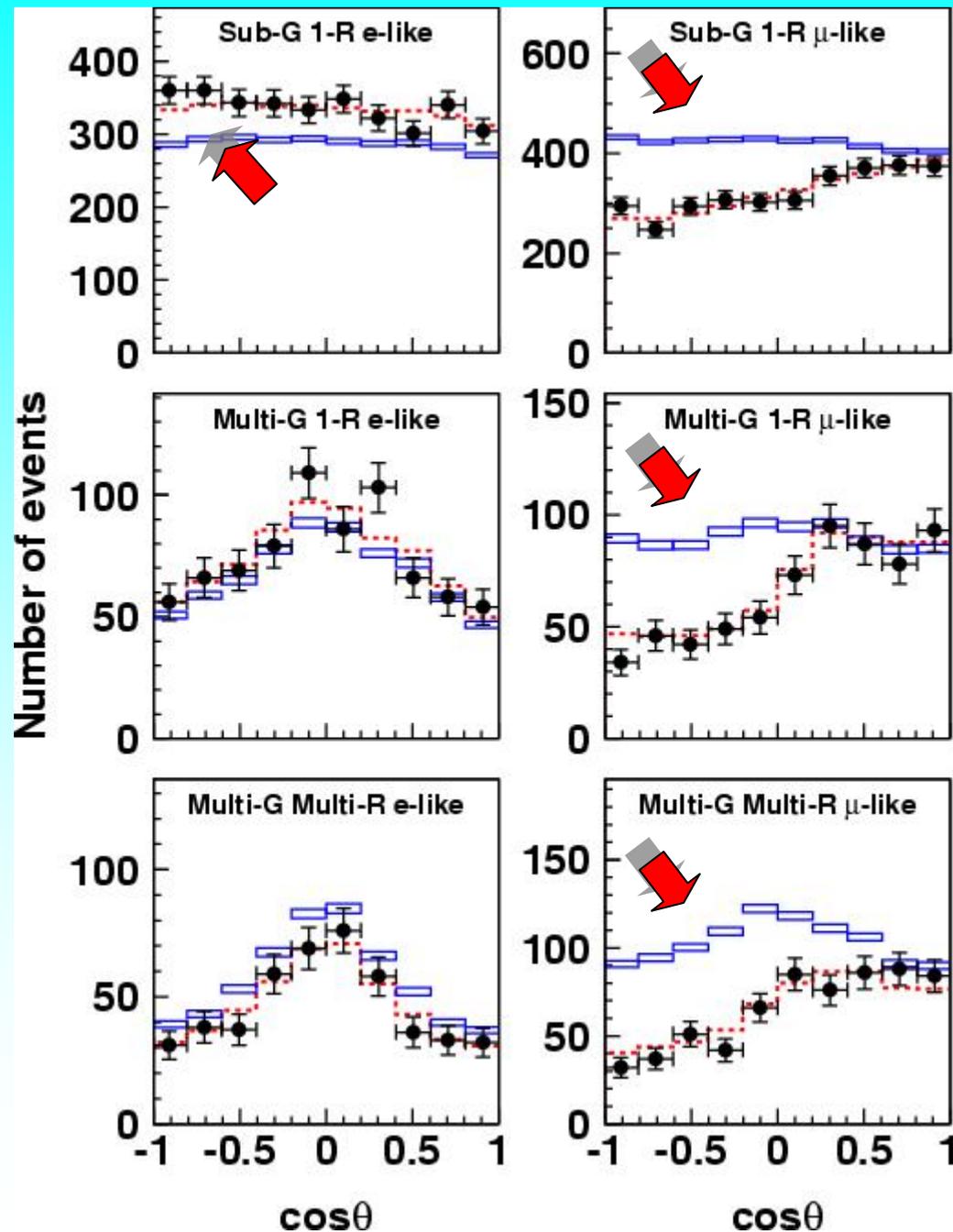

$$\nu_{\mu} - \nu_{\tau}$$

Zenith angle distribution

SuperKamiokande
hep-ex/0604011

Deficit

and excess



mu/e-ratio

$$R = \frac{(\mu/e)_{\text{data}}}{(\mu/e)_{\text{MC}}}$$

$$R \sim 1 - \sin^2 2\theta \langle \sin^2 \phi/2 \rangle_z$$

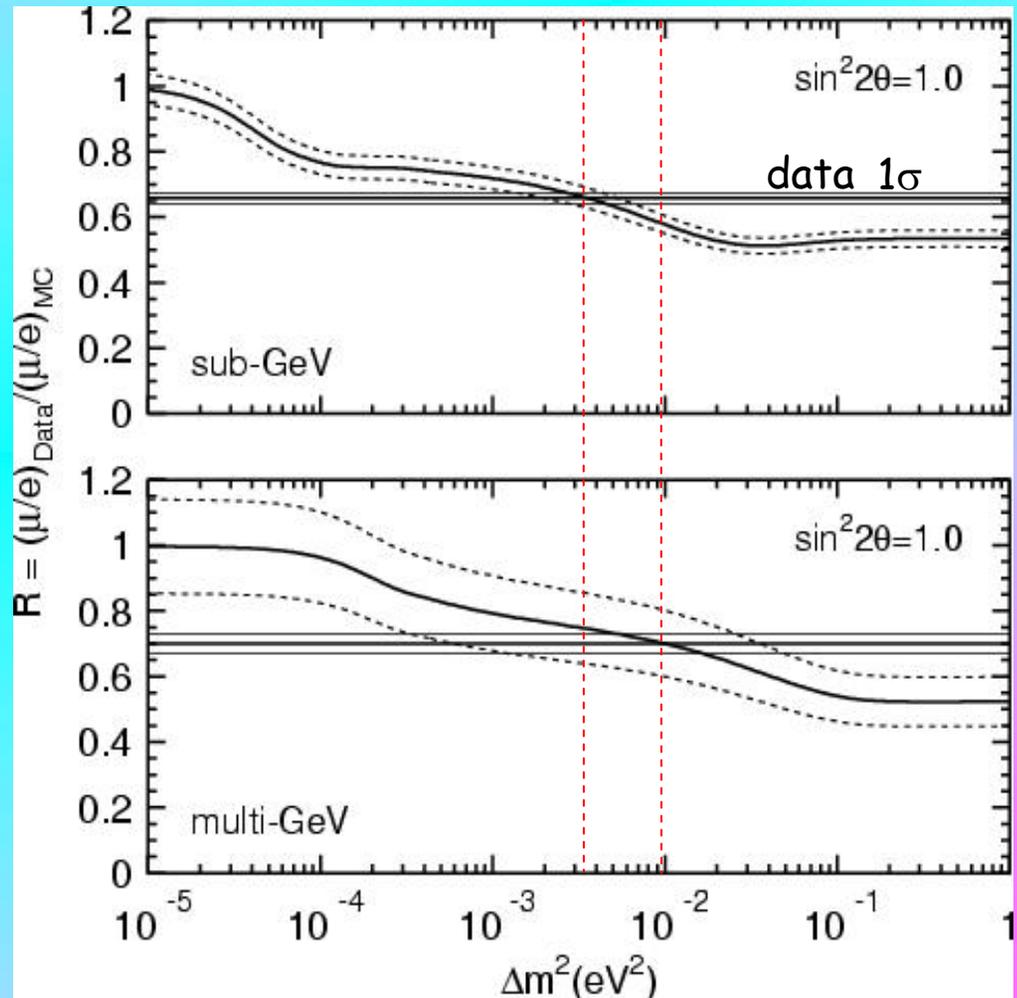

 oscillation phase
 factor integrated
 over zenith angle
 (decreases with energy)

for multi-GeV:

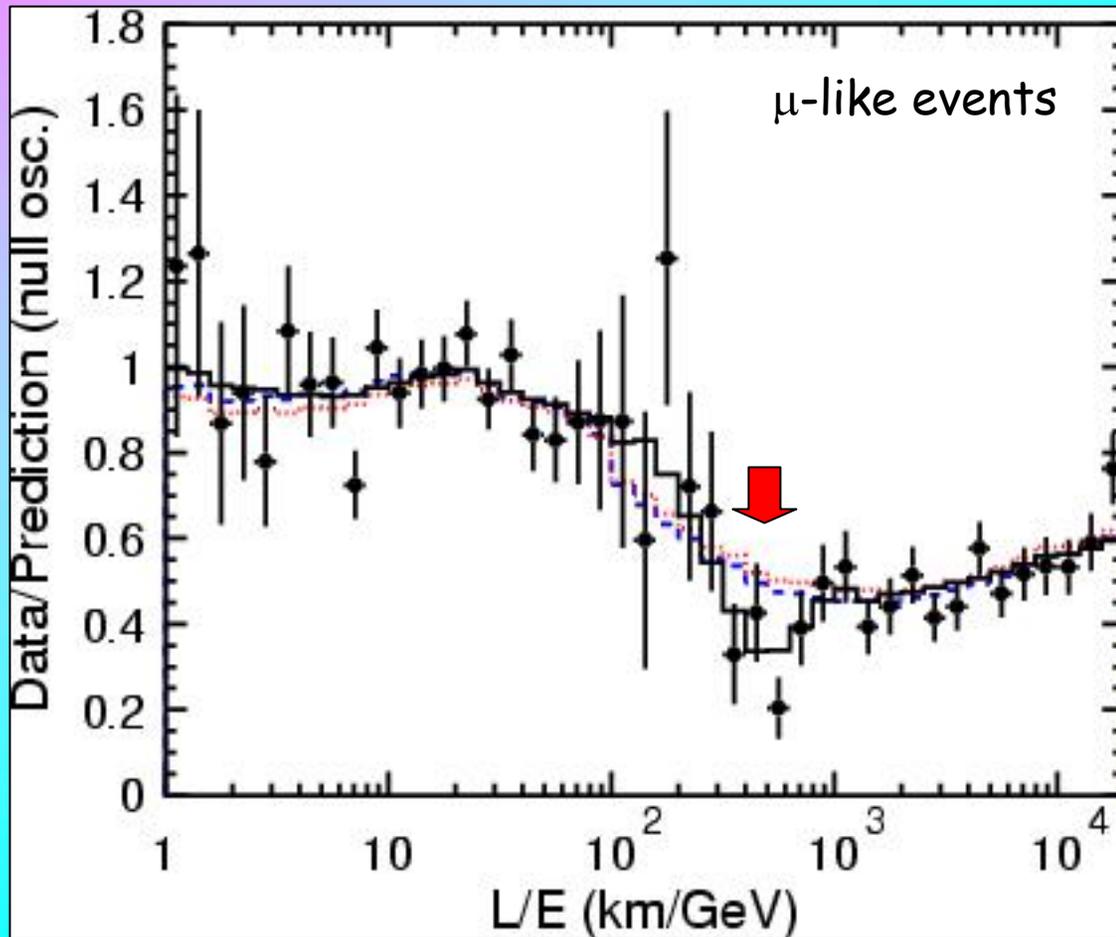
$$\langle \sin^2 \phi/2 \rangle_z \sim 0.25 - 0.30$$

from data

$$\sin^2 2\theta \sim 1$$



L/E dependence



- oscillations
- - - decay
- ⋯ decoherence

In the first oscillation maximum (dip): $\phi = \pi$



$$\Delta m^2 = 2\pi E/L$$

data: $L/E \sim 500 \text{ km/GeV}$



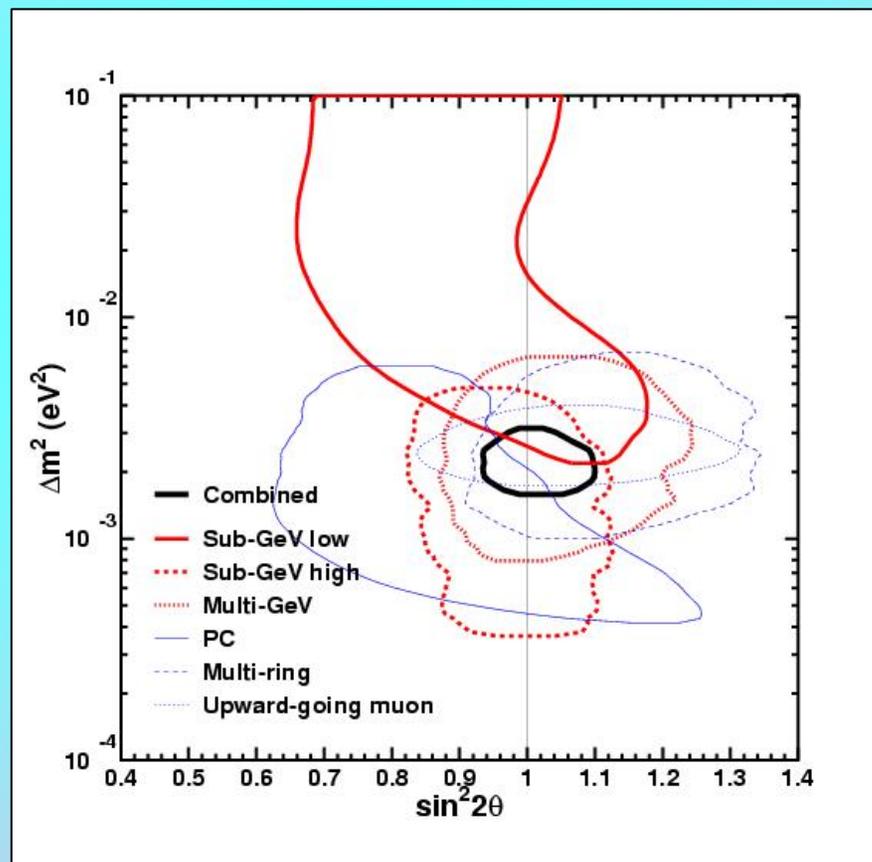
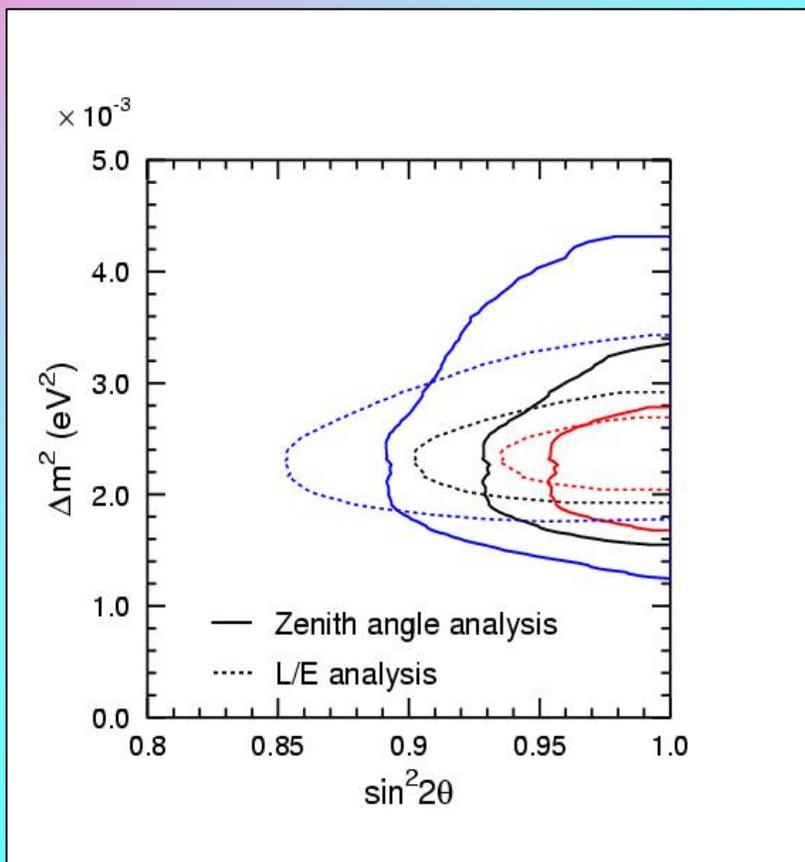
$$\Delta m^2 \sim 2.5 \cdot 10^{-3} \text{ eV}^2$$

Super-Kamiokande
PRL, 93 101801 (2004)

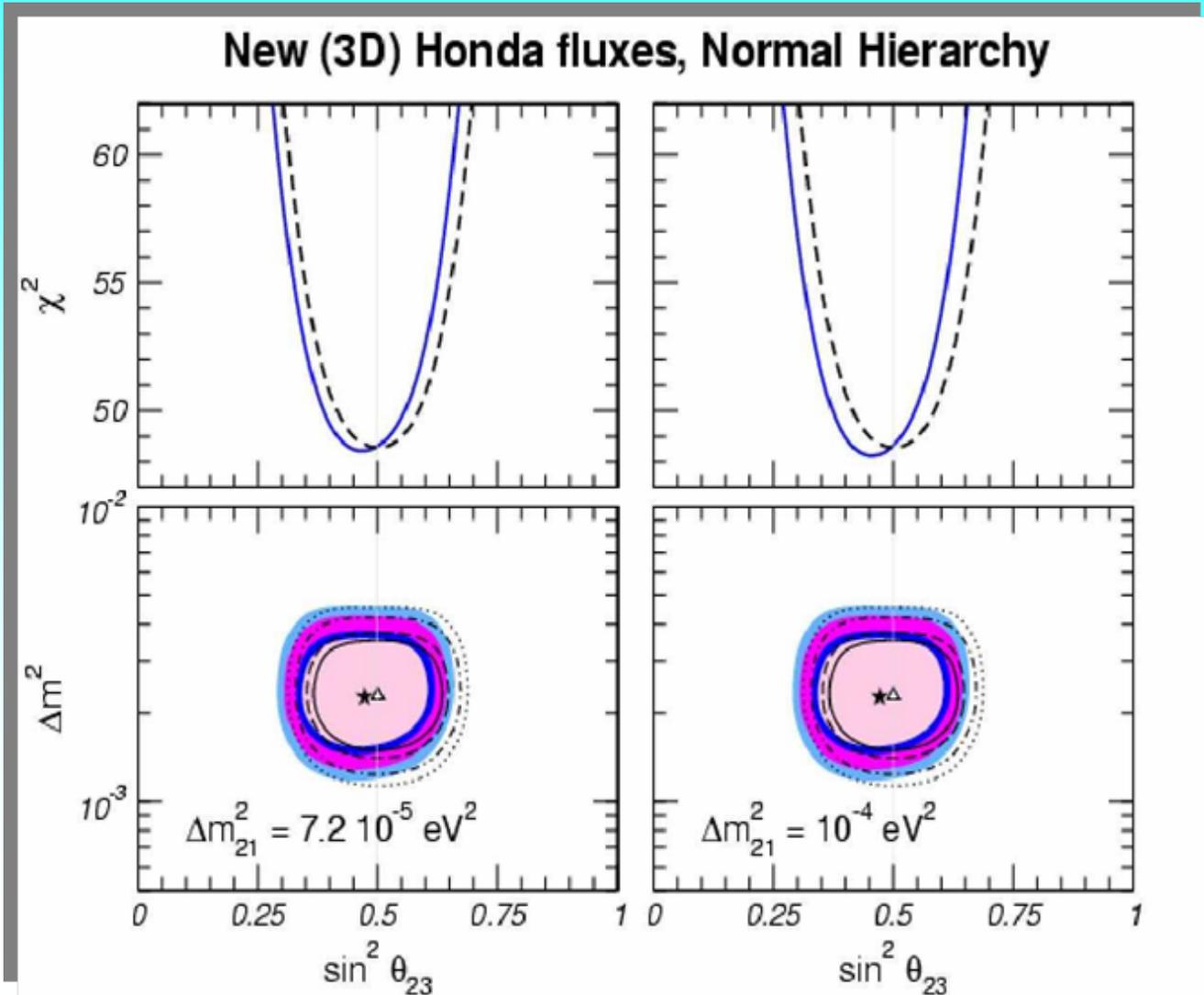
Oscillation parameters

Super-Kamiokande
hep-ex/0501064

$$1.5 \times 10^{-3} < \Delta m^2 < 3.4 \times 10^{-3} \text{ eV}^2$$
$$0.92 < \sin^2 2\theta \quad \text{at 90\% C.L.}$$



Deviation of 2-3 mixing from maximal



----- $\Delta m^2 = 0$
——— $\Delta m^2 \neq 0$

$\sin^2 \theta_{23} = 0.45 - 0.47$

M. C. Gonzalez-Garcia
M. Maltoni, A.S.

Status and perspectives

Still complete 3ν analysis of the atmospheric neutrino data should be done by the experimental collaboration (SK)

Physics potential of future detectors (next generation experiments)

Rich source:

- various flavors ν_e and ν_μ
- neutrinos and antineutrinos
- huge interval of energies: $E \sim 0.1 - 10^4$ GeV

“Variable baseline”:

$L \sim 10 - 10^4$ km (zenith angles)
sensitivity to both Δm^2
opens a possibility (in principle)
to establish the CP violation

Drawback:

- small statistics
- uncertainties in the predicted fluxes

Physics and Detectors

- determination of mass hierarchy
- 1-3 mixing
- deviation of 2-3 mixing from maximal (and quadrant)
- CP violation
- Earth tomography

■ INO - Indian Neutrino observatory

50 kton iron calorimeter

■ HyperKamiokande

0.5 Megaton water Cherenkov detectors

■ Underwater detectors NEMO, ANTARES

■ Icecube (1000 Mton)

■ TITAND (Totally Immersible Tank Assaying Nuclear Decay)

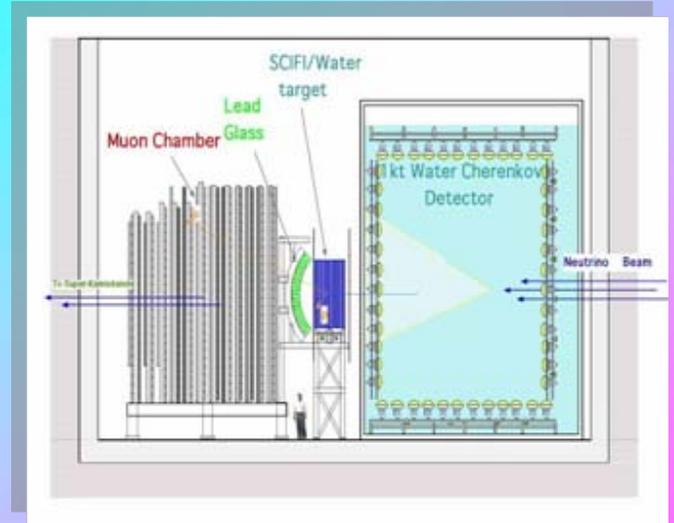
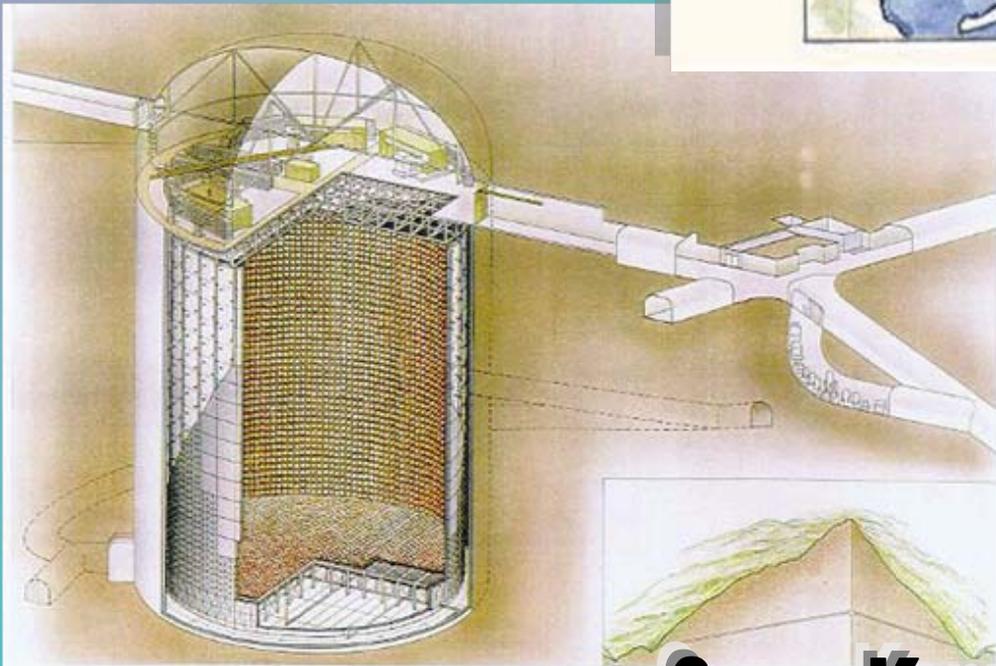
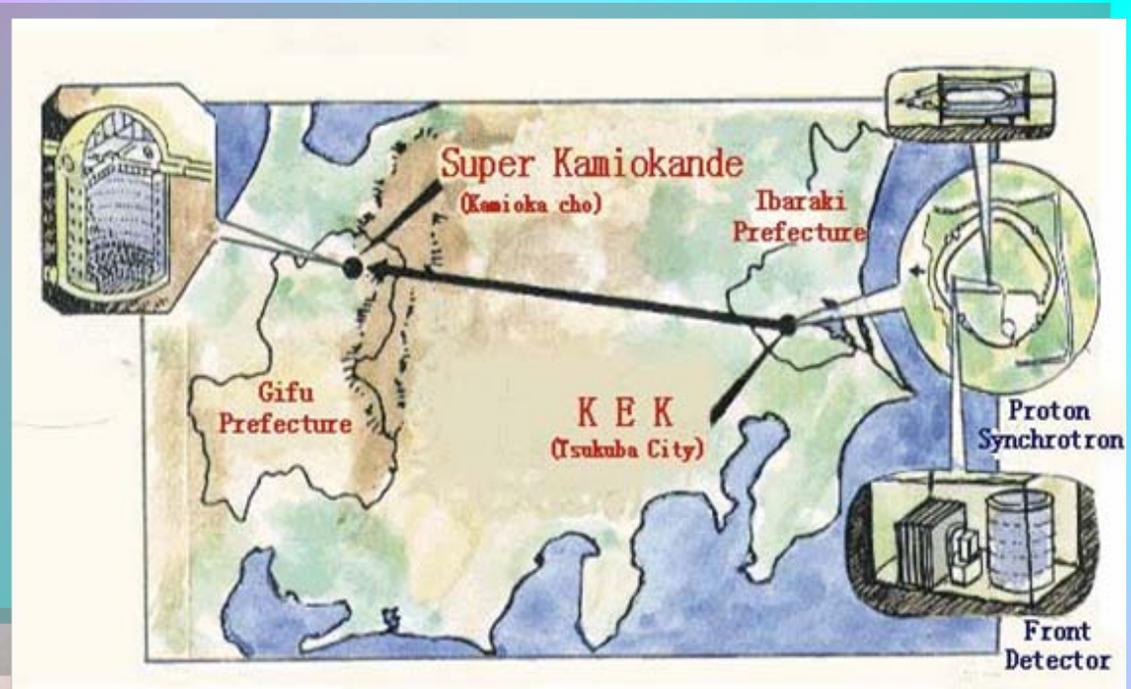
2 Mt and more

Y. Suzuki.

K2K

KEK to Kamioka

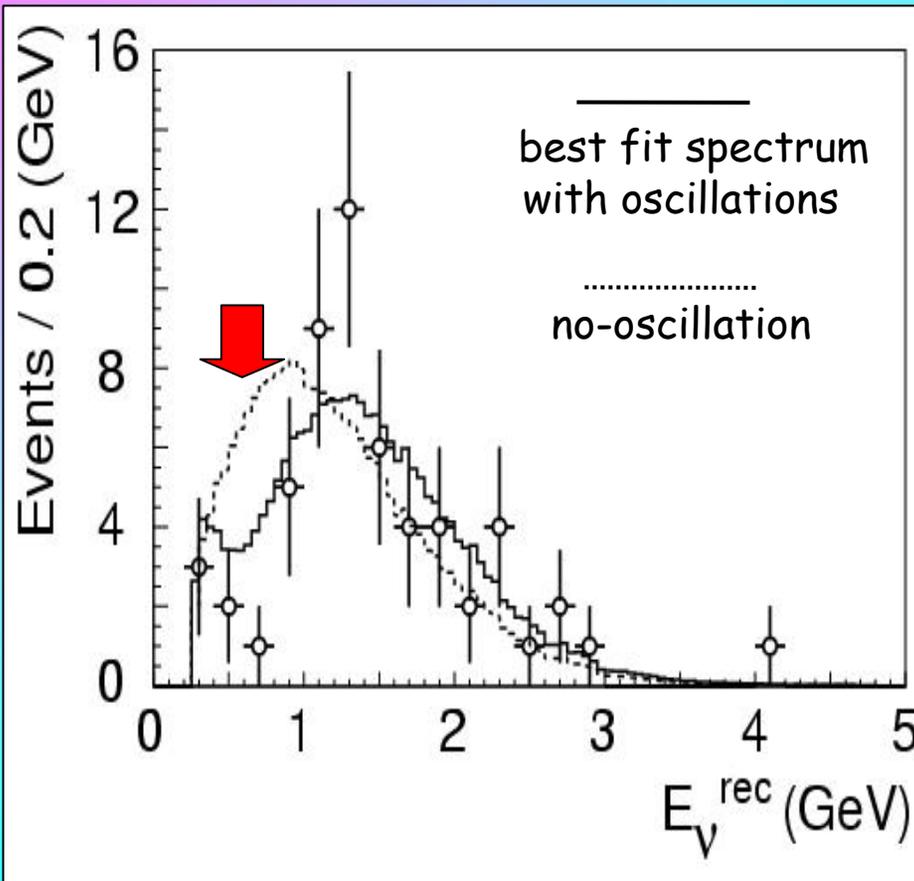
$$\nu_{\mu} \rightarrow \nu_{\mu}$$



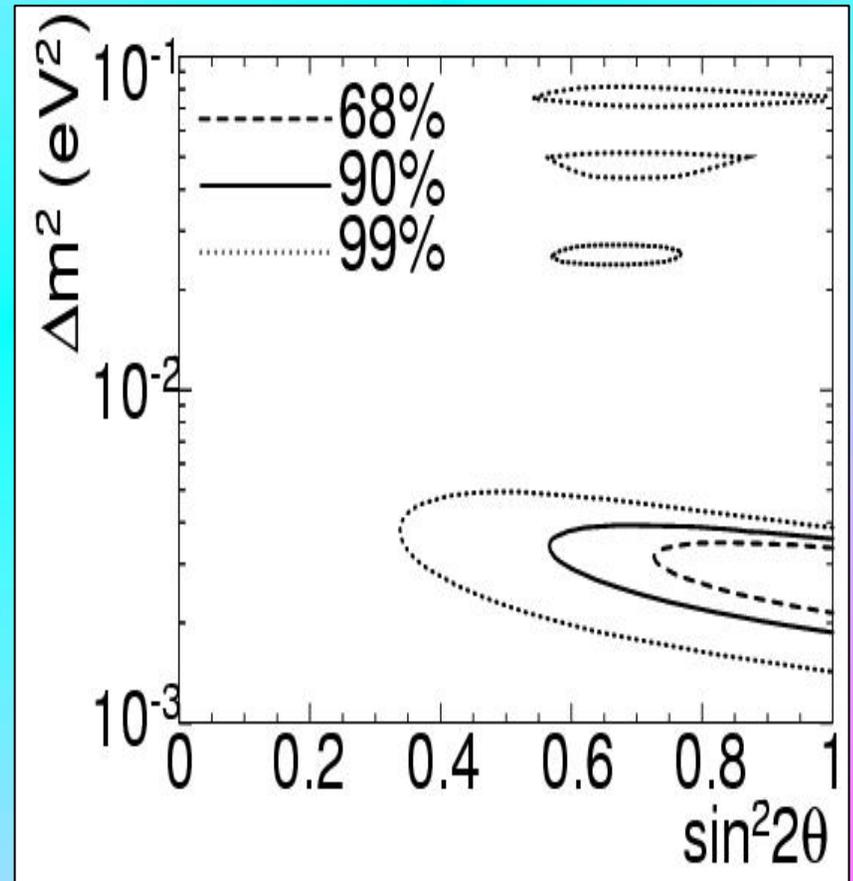
SuperKamiokande

K2K results

107 events observed
151 +12/-10 -expected



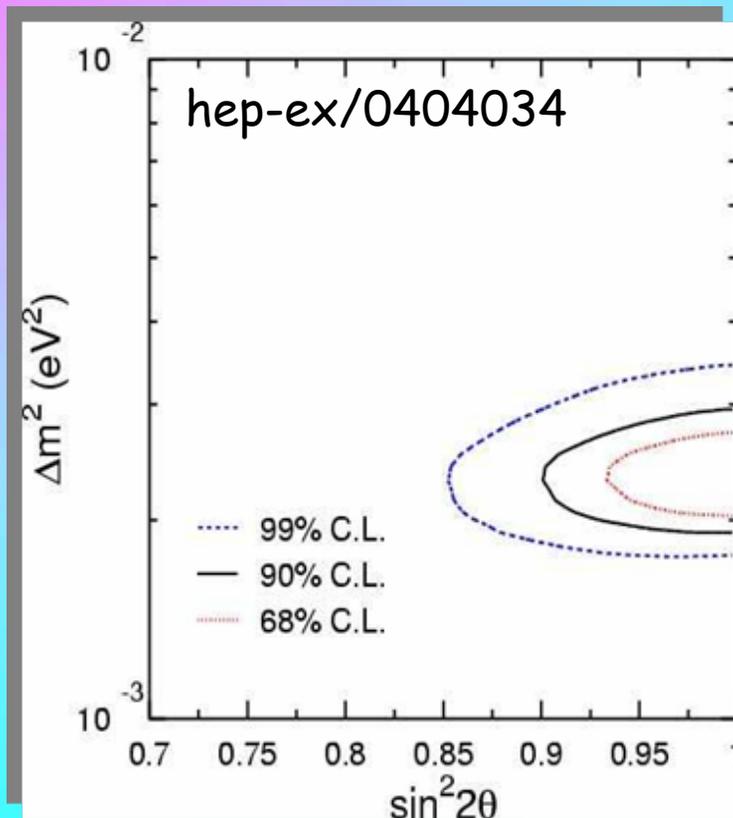
Normalized (57 events) reconstructed neutrino energy spectra



Allowed regions of parameters

2-3 masses and mixing

SuperKamiokande (atmospheric)



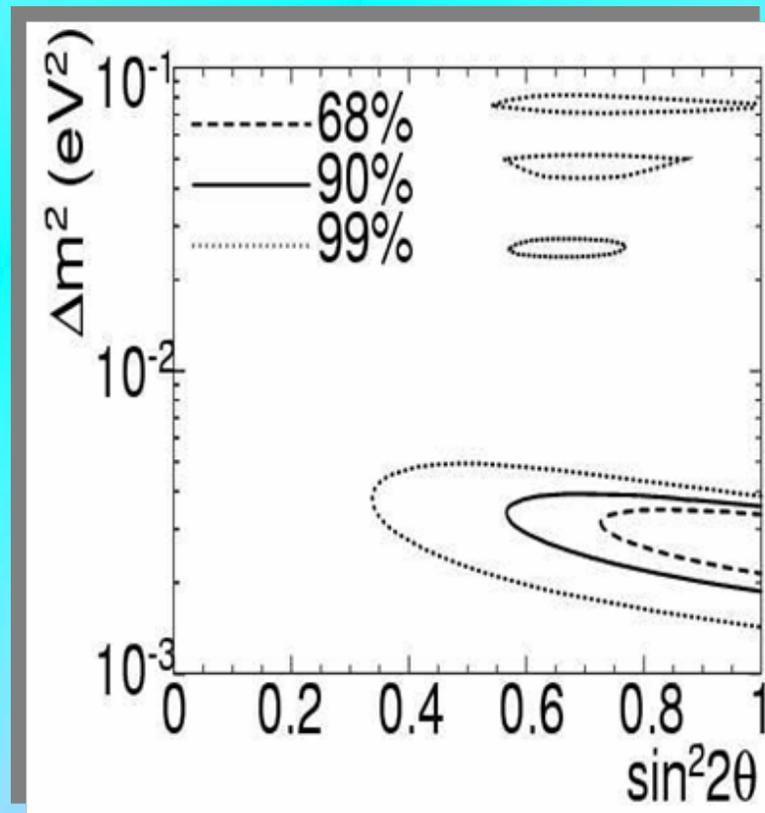
$$\Delta m_{23}^2 = (2.4 \pm 0.6 / -0.5) \cdot 10^{-3}, \text{ eV}^2$$

90 % C.L.

$$\sin^2 2\theta_{23} = 1.0$$

K2K

hep-ex/0411038



$$\Delta m_{23}^2 = (1.9 - 3.6) \cdot 10^{-3}, \text{ eV}^2$$

90 % C.L.

MINOS

Main Injector Neutrino
Oscillation Search

LBL: Fermilab - SOUDAN mine



Near detector (1km):
1 kton

Far detector (735 km)
5400 t, steel,
sampling calorimeter

Beam: 120 GeV protons
 $2.5 \cdot 10^{20}$ p/year
→ 1- 10 GeV neutrinos

MINOS results

Expected: 177 +/- 11 events

Observed: 92 events

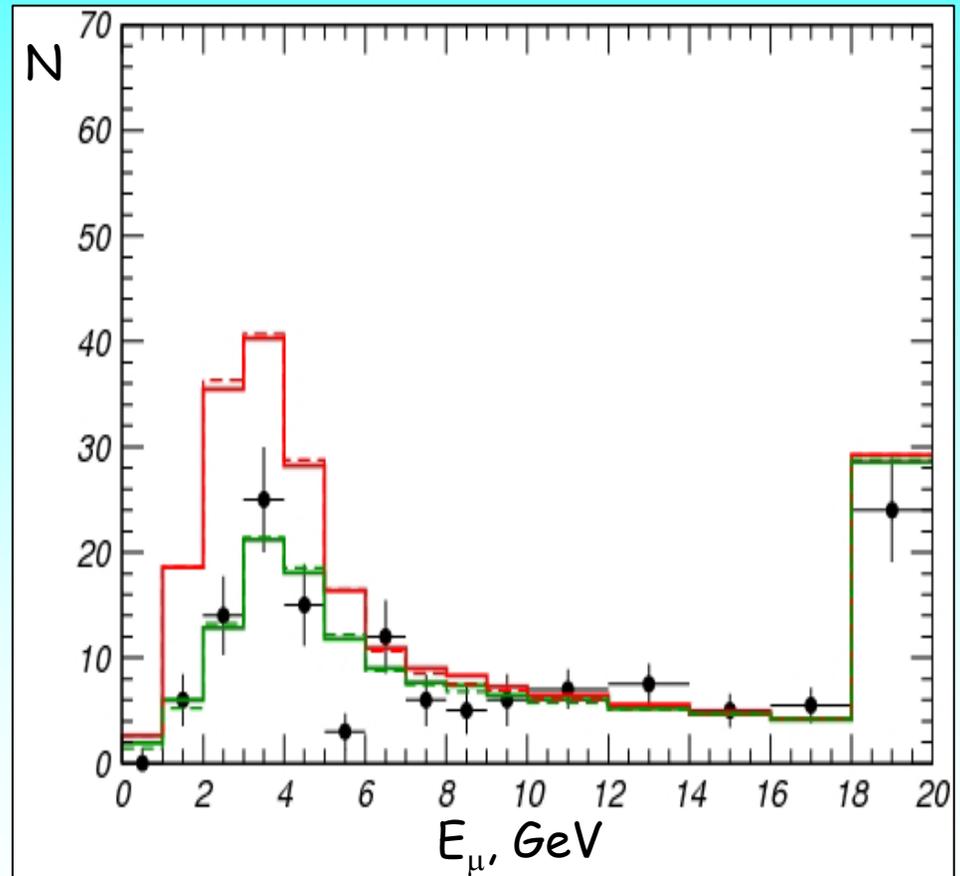
CC events induced by ν_μ

b.f.:

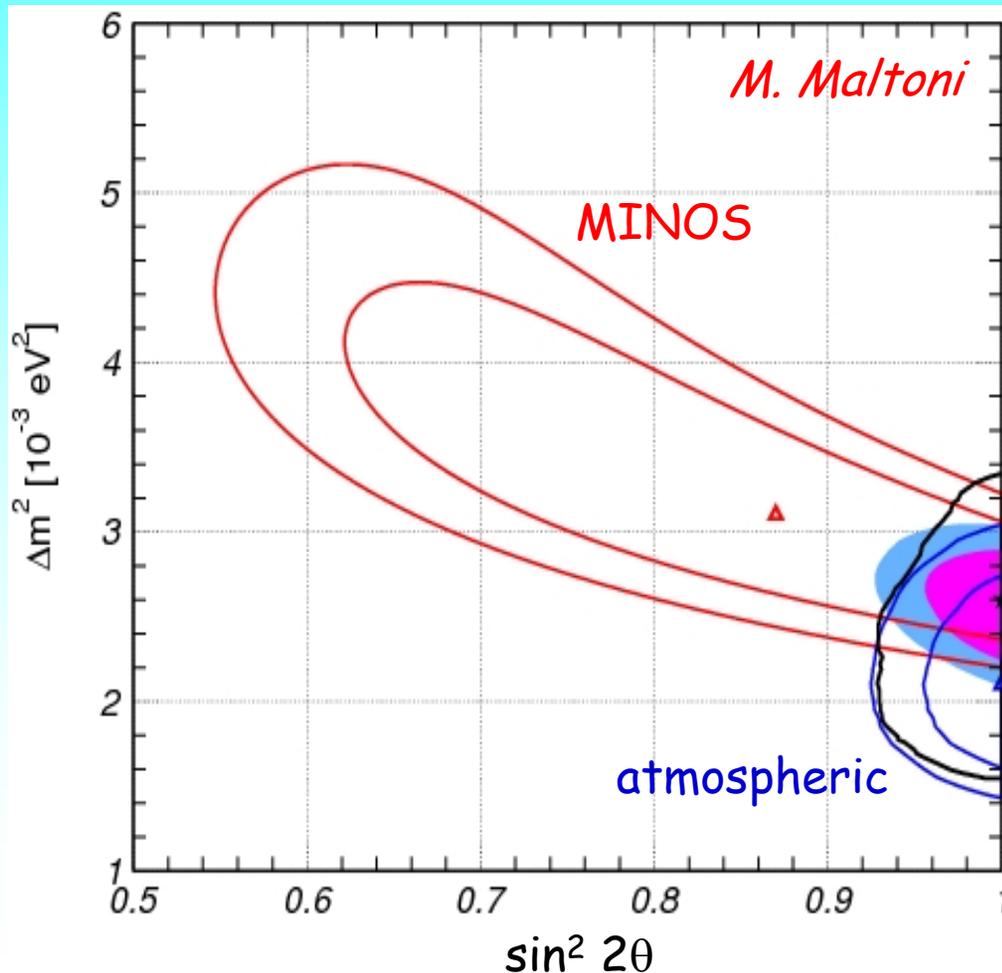
$$\Delta m^2 = 3.1 \cdot 10^{-3} \text{ eV}^2$$

The first oscillation minimum:
 $E = 1.8 \text{ GeV}$

Second minimum: $E = 5.4 \text{ GeV}$



Combined analysis



Allowed regions at
68% and 90% C.L.

b.f.:

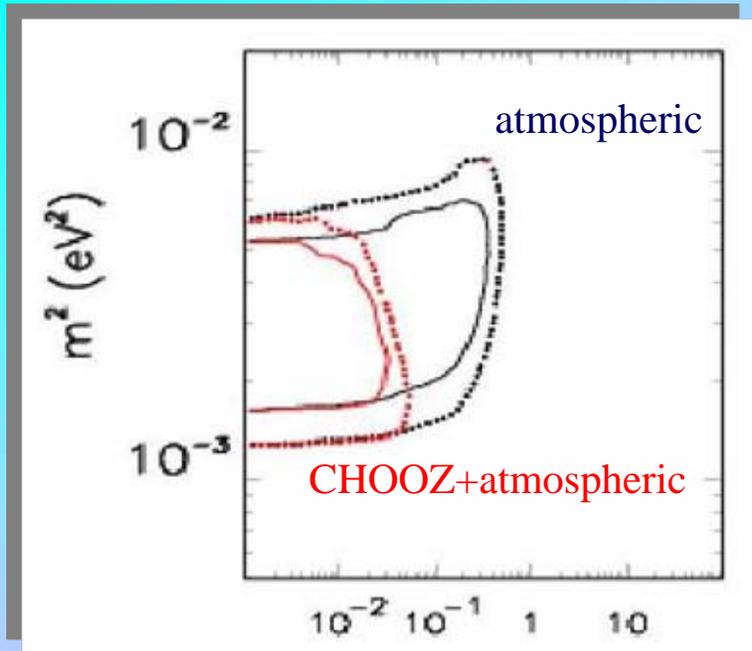
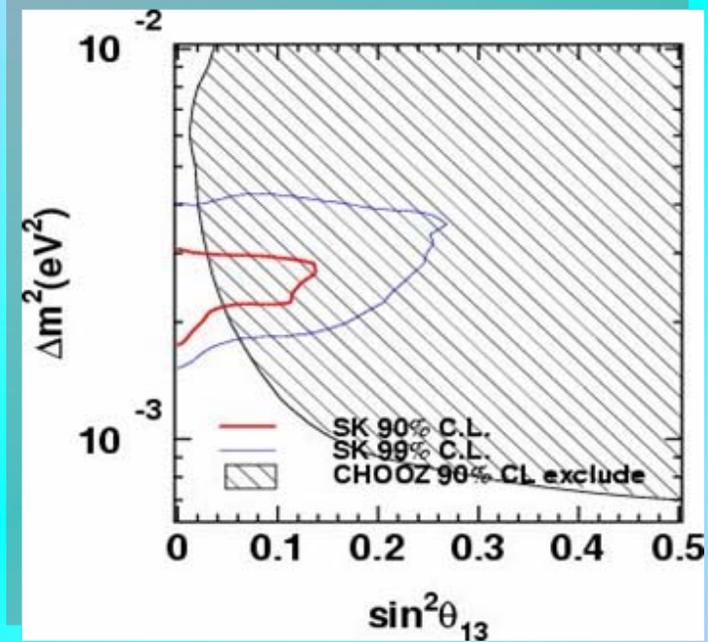
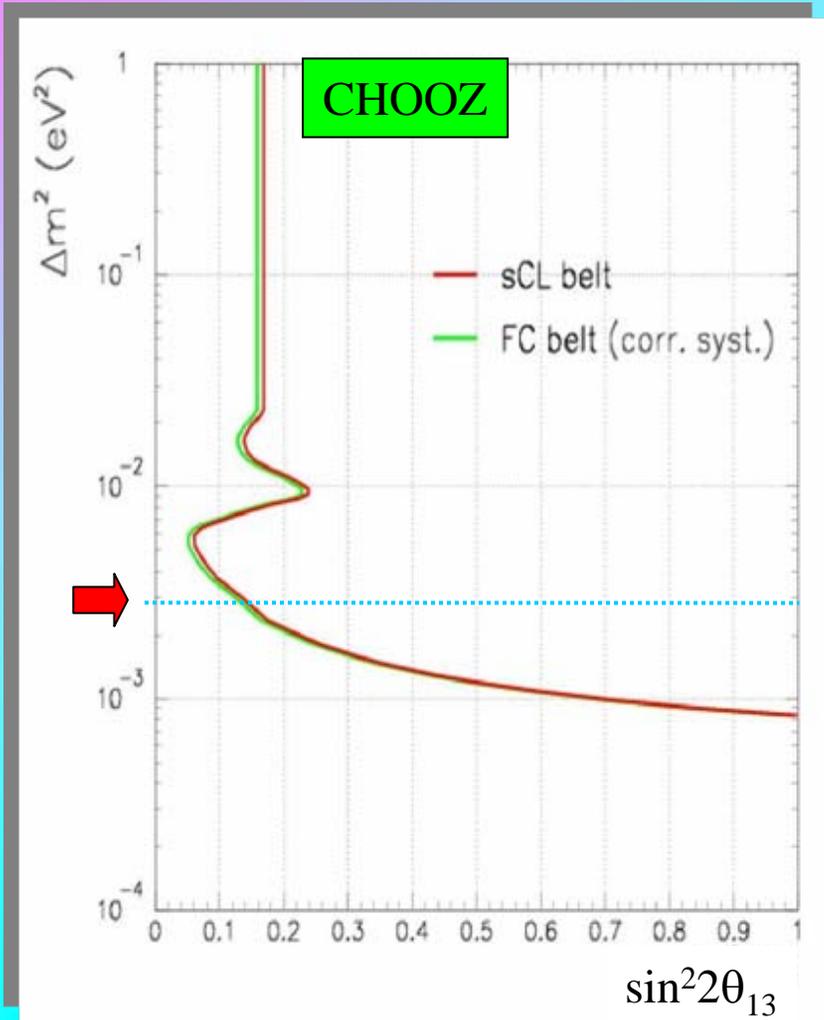
$$\Delta m^2 = 2.6 \cdot 10^{-3} \text{ eV}^2$$

MINOS has not
produced shift
from maximal mixing
in the global fit

1-3 mixing: effects and bounds

U_{e3}

$$U_{e3} = \sin \theta_{13} e^{-i\delta} = \langle \nu_e | \nu_3 \rangle$$

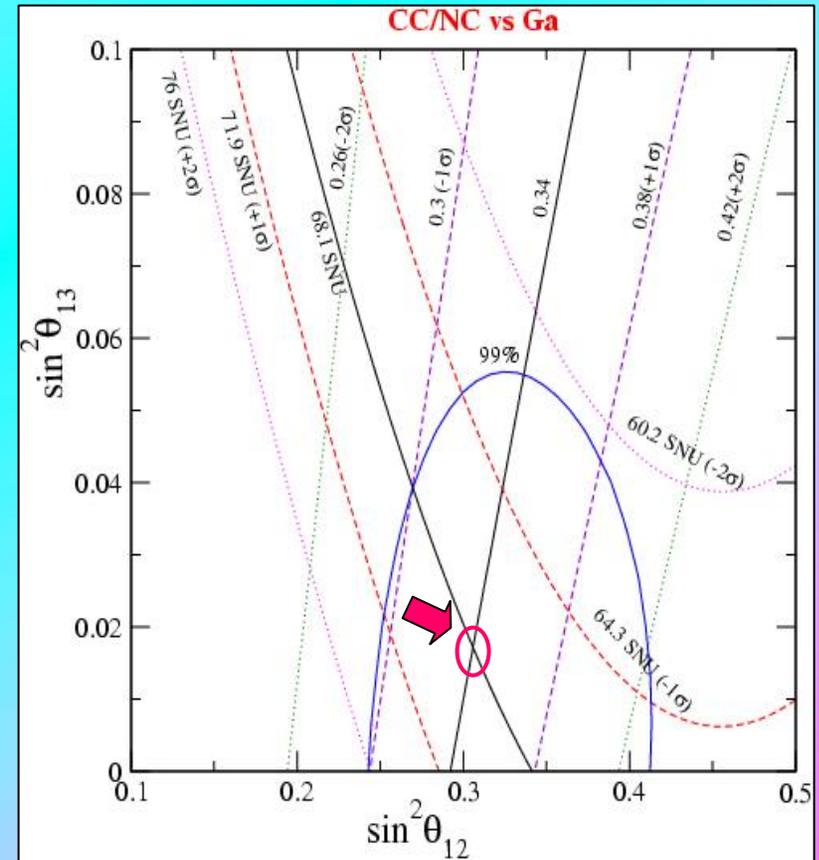
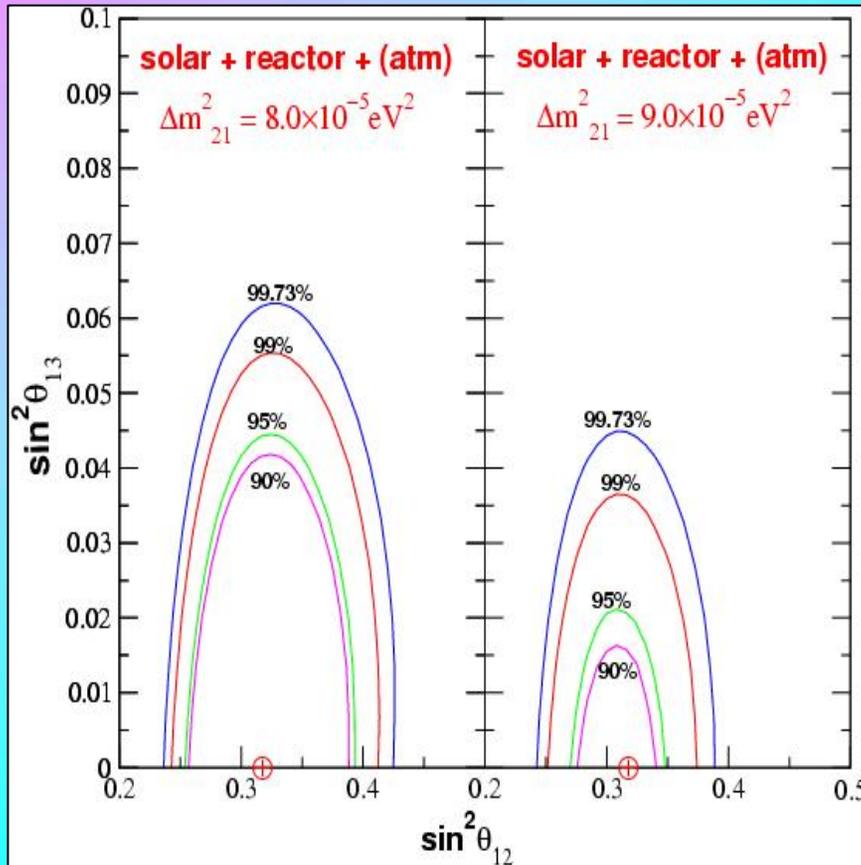


$\tan^2 \theta_{13}$

Theta 1-3

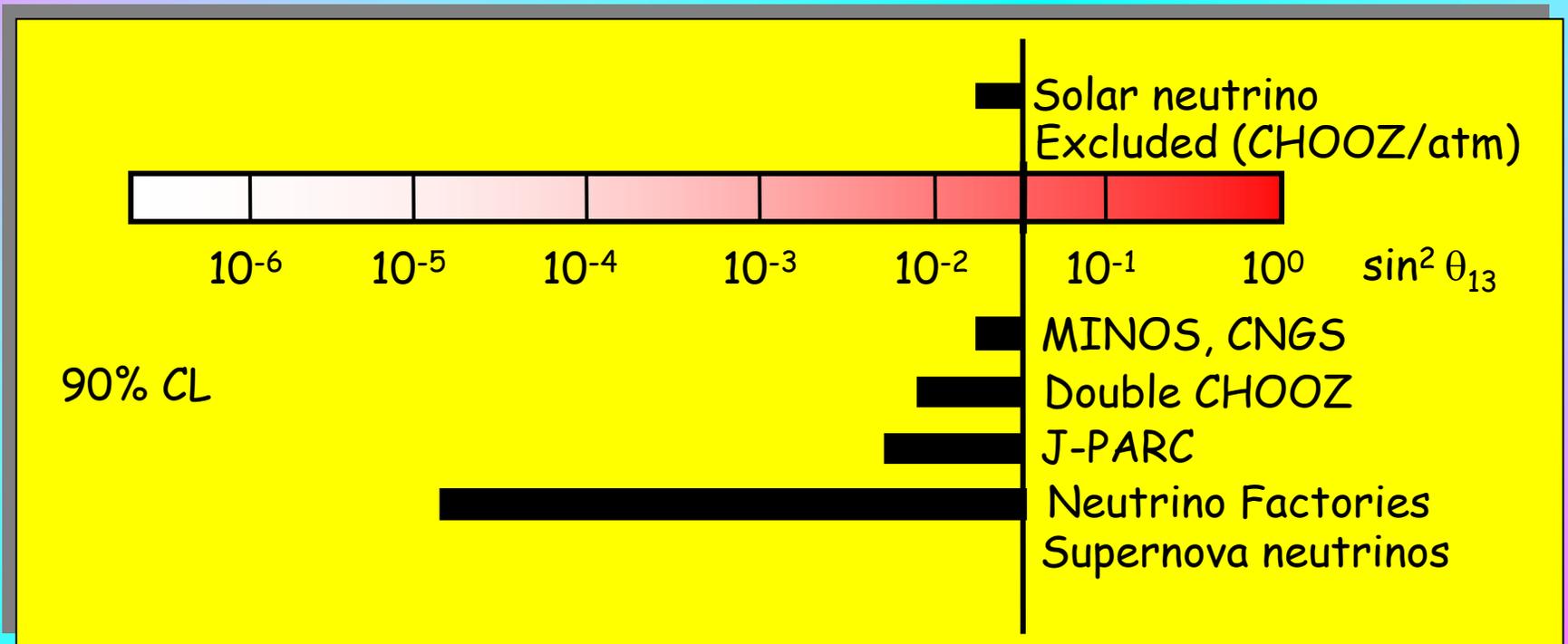
Solar neutrinos: degeneracy of 1-2 and 1-3 mixing

S. Goswami, A.S.



$\sin^2 \theta_{13} = 0.017 \pm 0.26$

Sensitivity to 1-3 mixing



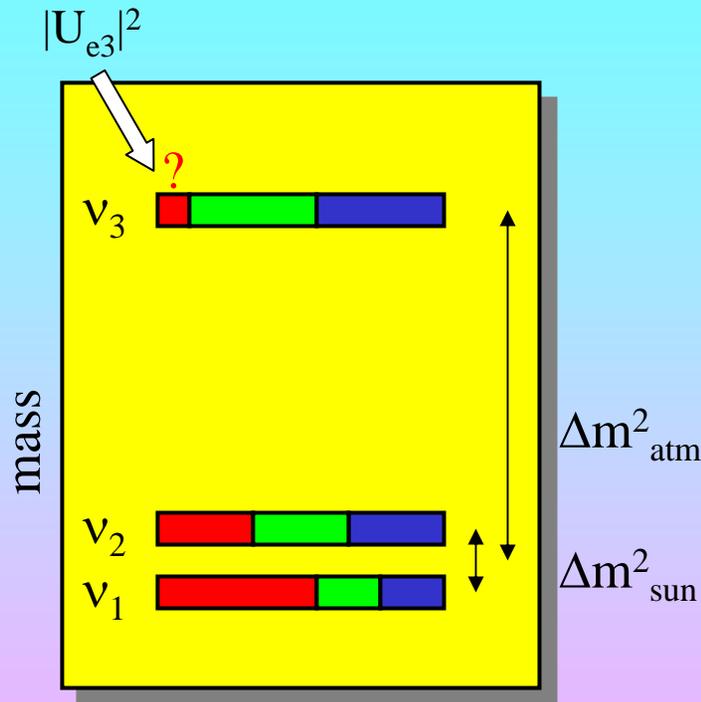
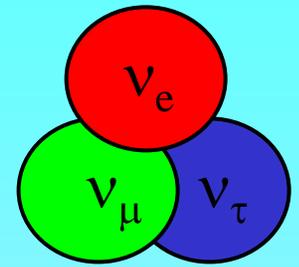
Plausible values of parameters

No new physics apart from 3 massive and mixed neutrinos

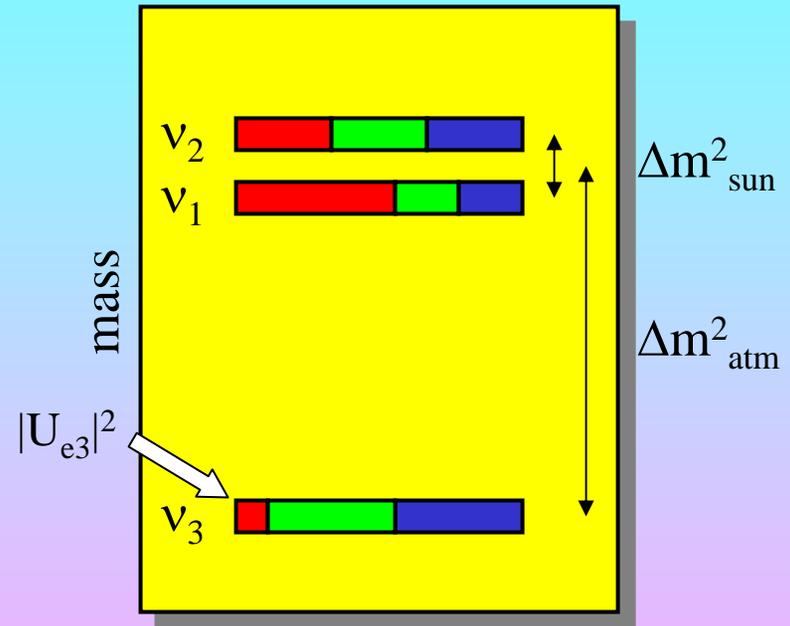
	S-V	Bari
Δm_{12}^2	$8.0 \cdot 10^{-5} \text{ eV}^2$	$7.9 \cdot 10^{-5} \text{ eV}^2$
$\sin^2 \theta_{12}$	0.310	0.314
Δm_{23}^2	$2.5 \cdot 10^{-3} \text{ eV}^2$	$2.4 \cdot 10^{-3} \text{ eV}^2$
$\sin^2 \theta_{23}$	0.50	0.44
$\sin^2 \theta_{13}$	0.00 [< 0.013]	0.009 [< 0.020]

[1σ]

Mass spectrum and mixing



Normal mass hierarchy
(ordering)



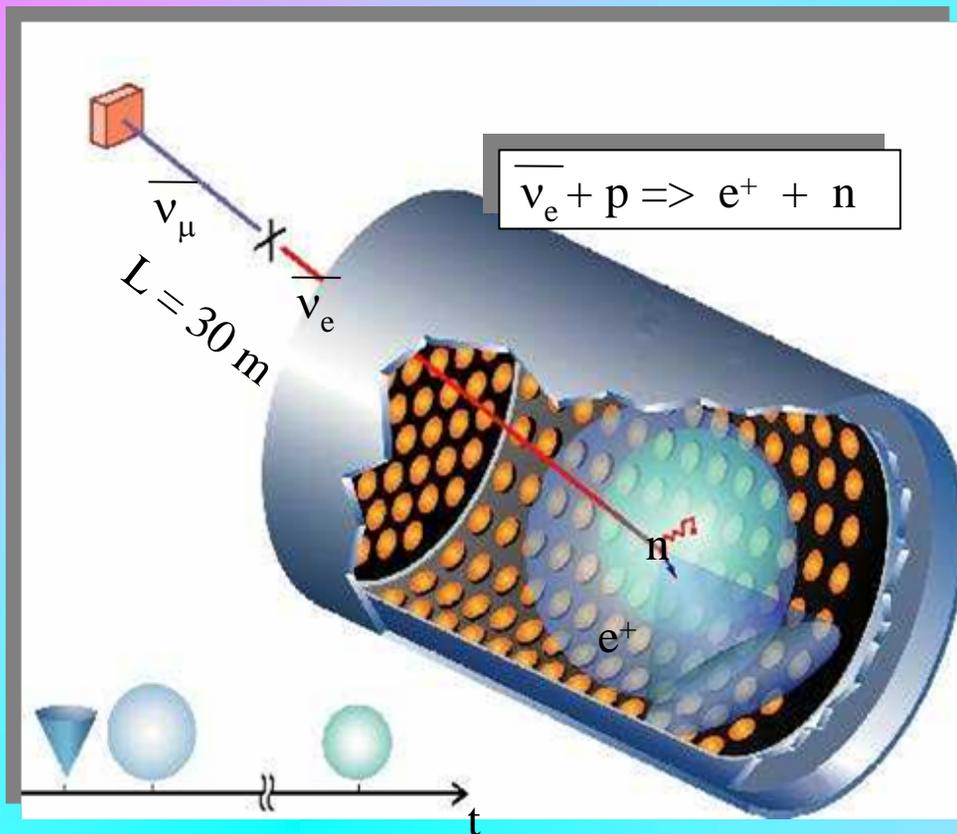
Inverted mass hierarchy
(ordering)

- Type of mass spectrum: with Hierarchy, Ordering, Degeneracy → absolute mass scale
- Type of the mass hierarchy: Normal, Inverted
- $U_{e3} = ?$

Beyond standard picture

LSND

Large Scintillator Neutrino Detector
Los Alamos Meson Physics Facility



$$\bar{\nu}_e + p \Rightarrow e^+ + n$$

Cherenkov cone + scintillations

200 t mineral oil scintillator

$$p \rightarrow \pi^+ \rightarrow \mu^+ + \nu_\mu \rightarrow e^+ + \nu_e + \bar{\nu}_\mu$$

decay at rest

$$\bar{\nu}_\mu \rightarrow \bar{\nu}_e$$

$$P = (2.64 \pm 0.67 \pm 0.45) 10^{-3}$$

Oscillations?

$$\Delta m^2 > 0.2 \text{ eV}^2$$

Beyond ``standard'' picture:

- new sector,
- new symmetry



LSND

Ultimate oscillation anomaly?

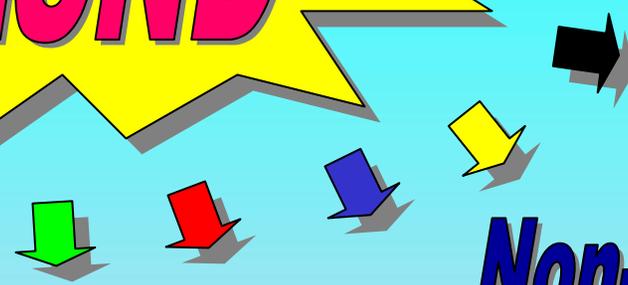
CPT-violation

G. Barenboim,
L. Borissov, J. Lykken

Disfavored by atmospheric neutrino data, no compatibility of LSND and all-but LSND data below 3σ -level

M.C. Gonzalez-Garcia,
M. Maltoni, T. Schwetz

CPT + (3+1)



Sterile neutrino (3 + 1)-scheme (3 + 2) ?

O. Peres, A.S.

M. Sorel, J. Conrad, M. Shaevitz

Decaying sterile neutrino

S. Palomares-Ruiz, S. Pascoli,
T. Schwetz

Non-standard Interactions

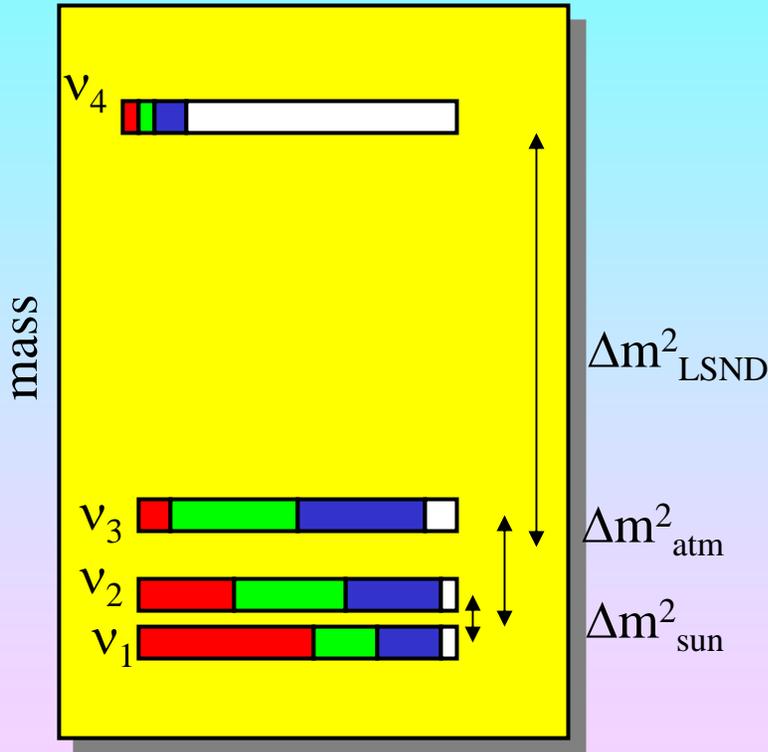
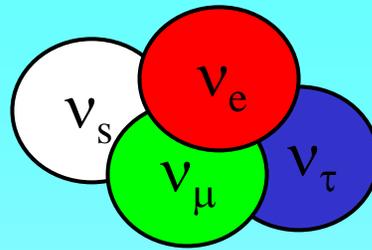
K. Babu, S. Pakvasa

Disfavored by a new analysis of KARMEN collaboration

MaVaN

R. Fardon, A. E. Nelson,
N. Weiner

(3 + 1) scheme



1-3 subsystem of levels is frozen

The problem is

$$P \sim |U_{e4}|^2 |U_{\mu 4}|^2$$

Restricted by short baseline experiments CHOOZ, CDHS, NOMAD

2 - 3 σ below the observed probability

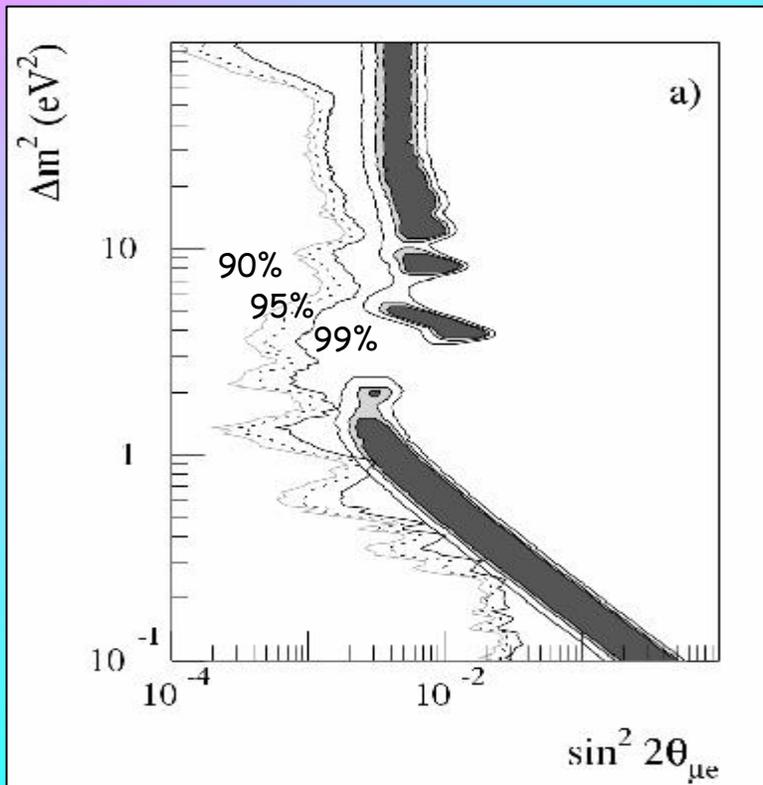
Generic possibility of interest even independently of the LSND result

Generation of large mixing of active neutrinos due to small mixing with sterile state

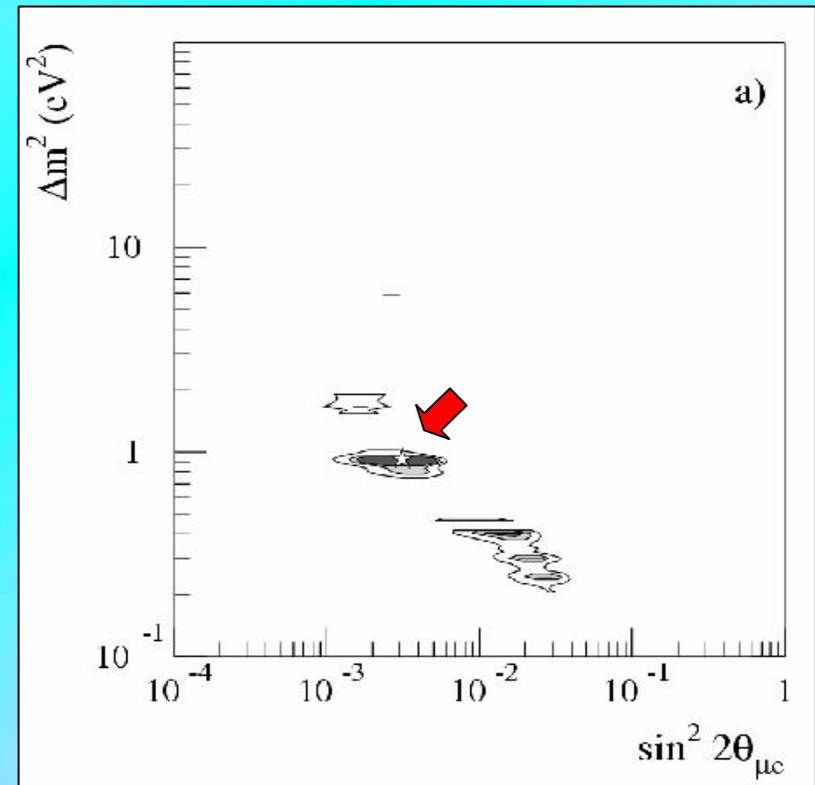
Produces uncertainty in interpretation of results

(3+1) - fit

hep-ex/0407027

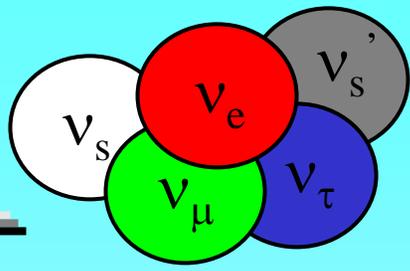


Compatibility of short baseline Experiments and LSND datasets

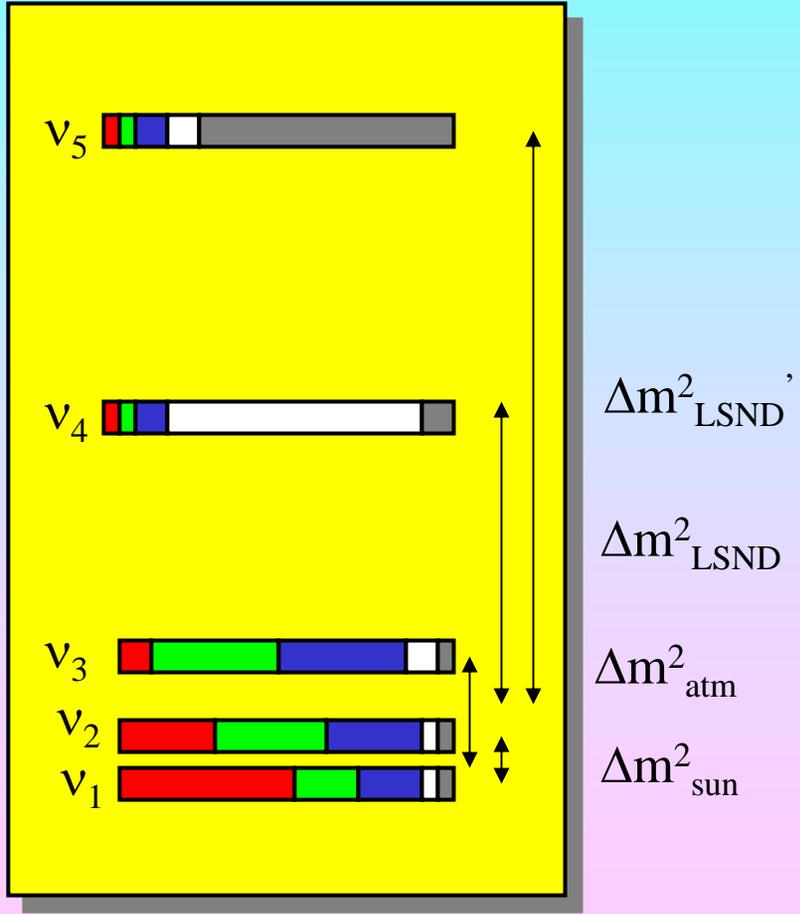


Allowed regions from combined fit of LSND and short baseline experiments

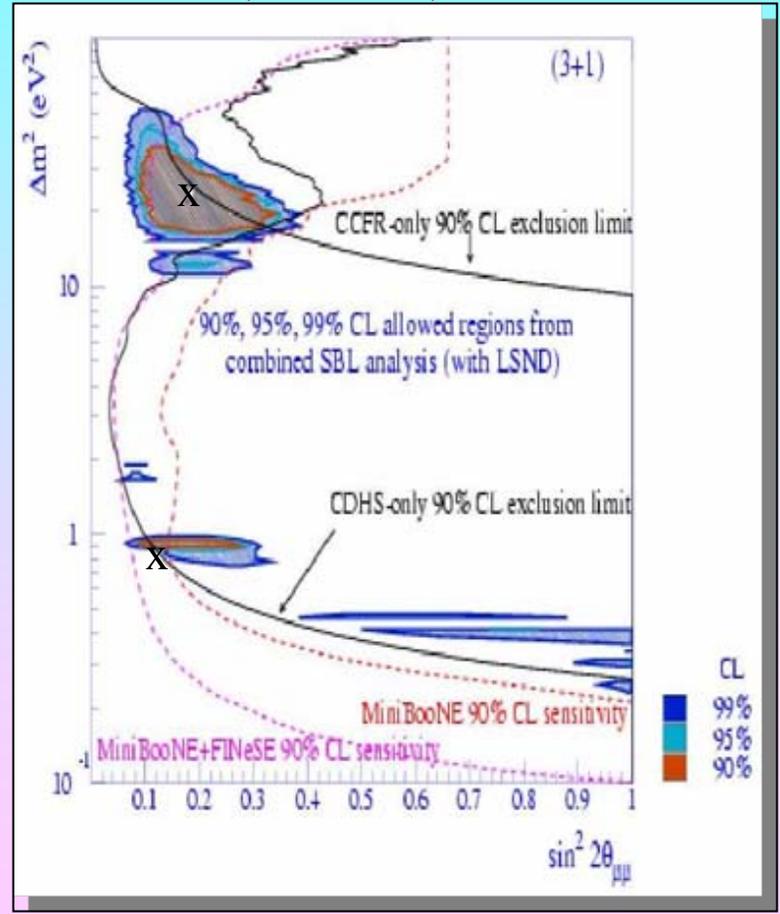
(3+2) scheme



mass



M. Sorel, J. Conrad, M. Shaevitz



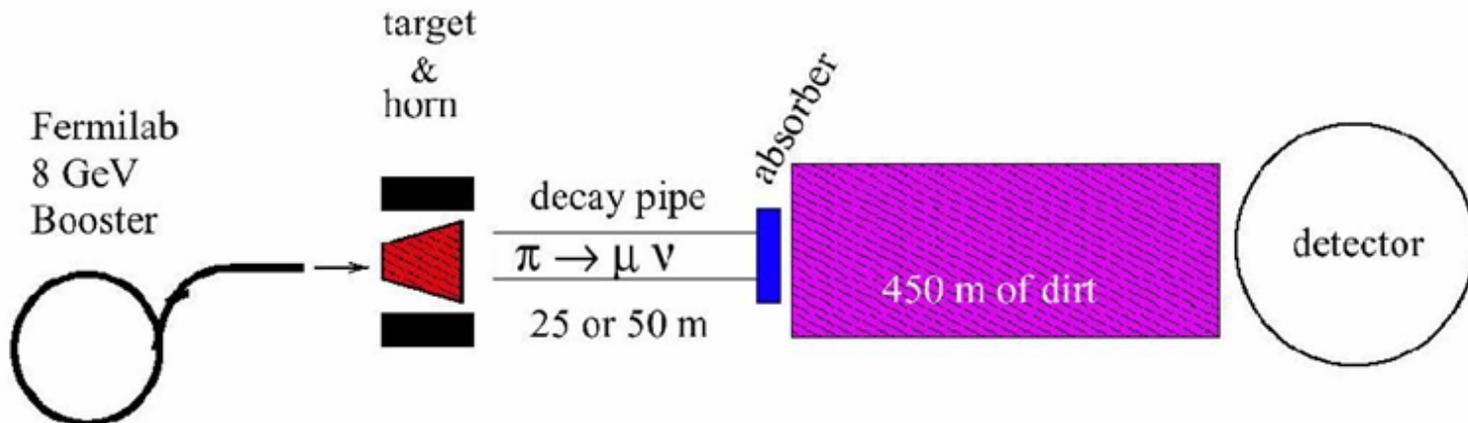
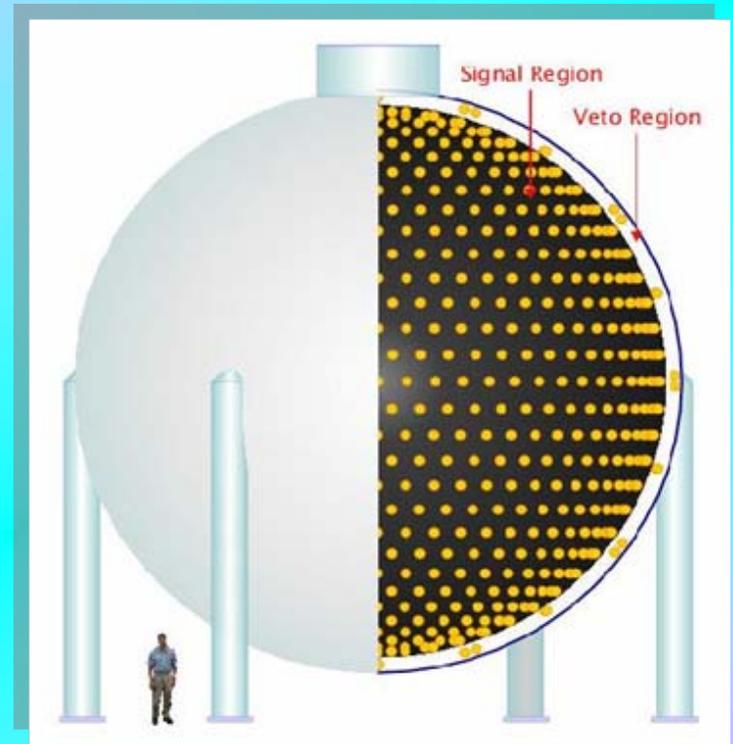
FINeSE

MiniBooNE

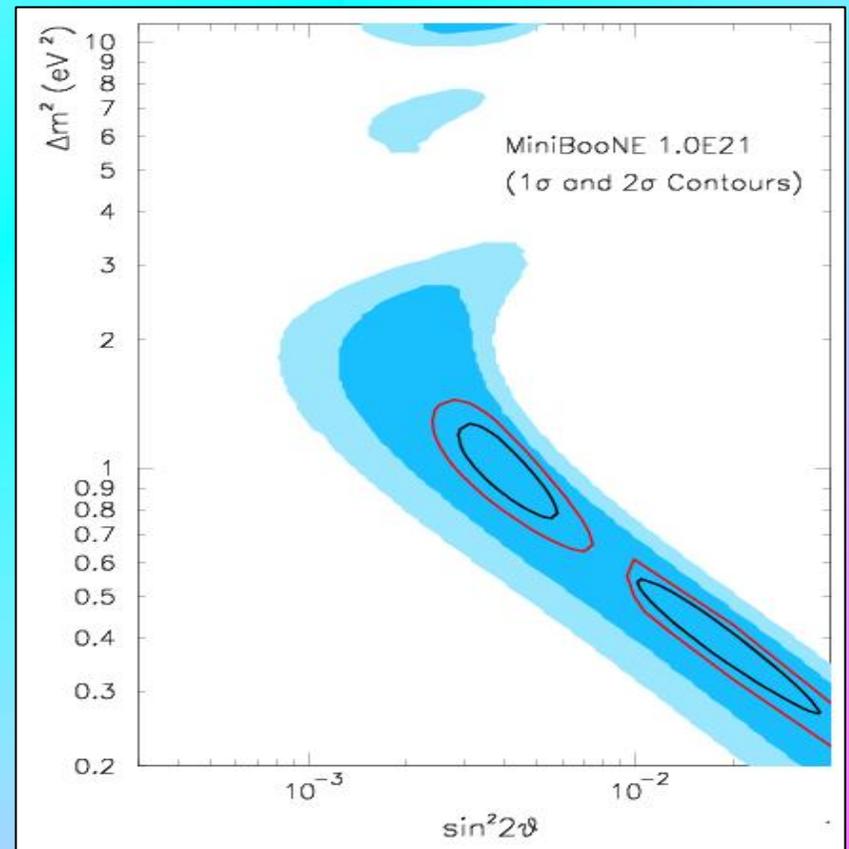
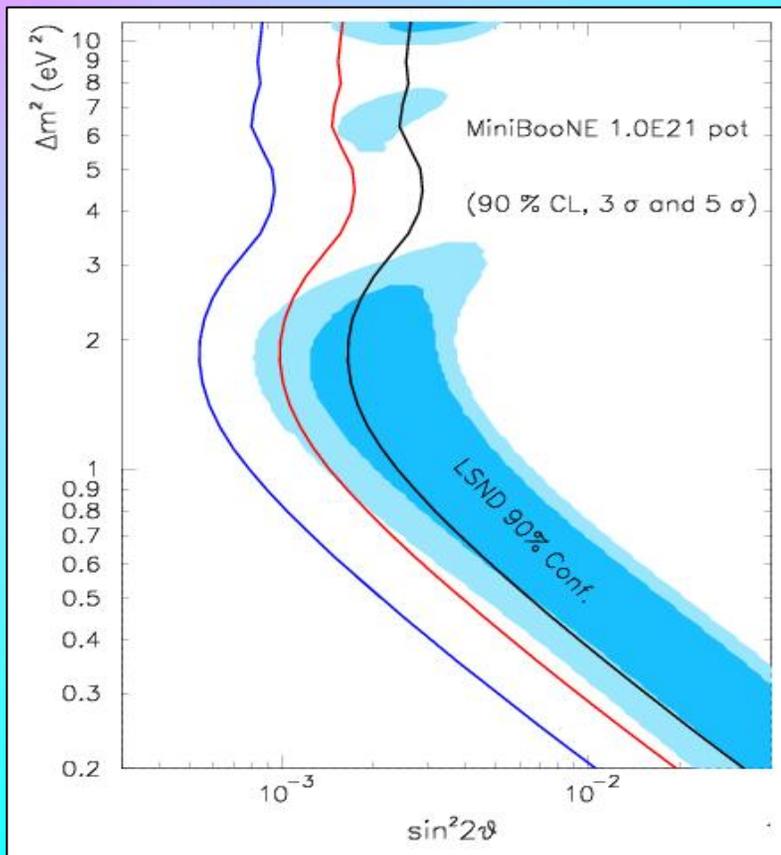
Search for ν_e appearance 12 m diameter tank

450 t (mineral oil)
1280 PMT

$L = 541 \text{ m}$, $\langle E_\nu \rangle \sim 800 \text{ MeV}$



MiniBooNE sensitivity



Empirical relations

Koide relation

$$\frac{m_e + m_\mu + m_\tau}{(\sqrt{m_e} + \sqrt{m_\mu} + \sqrt{m_\tau})^2} = 2/3$$

Y. Koide, Lett. Nuov. Cim.
34 (1982), 201

with accuracy 10^{-5}

$$\tan\theta_c = \sqrt{3} \frac{\sqrt{m_\mu} - \sqrt{m_e}}{2\sqrt{m_\tau} - \sqrt{m_\mu} - \sqrt{m_e}}$$

Both relations can be reproduced if

$$m_i = m_0(z_i + z_0)^2$$

$$\sum_i z_i = 0$$

$$z_0 = \sqrt{\sum_i z_i^2 / 3}$$

C A Brannen

Neutrinos

Another representation
which is closely connected
to circulant symmetry

Summary

Consistent picture: interpretation of all * the results in terms of vacuum mixing of three massive neutrinos

LSND ?

Two effects are important for the interpretation (at the present level of accuracy):
vacuum oscillations, adiabatic conversion (MSW)

Still unknown parameters → future
Phenomenological and experimental studies

Other effects if exist
- sub-leading contributions

Next level, sub-leading effects -
require more involved study

Strong difference of quark and lepton mixing patterns

Furthermore:

- particular neutrino symmetries
- tri-bimaximal mixing
- quark-lepton complementarity?

New puzzle?
Real or
Accidental?