

Large β_0 corrections to the energy levels and wave function at N³LO

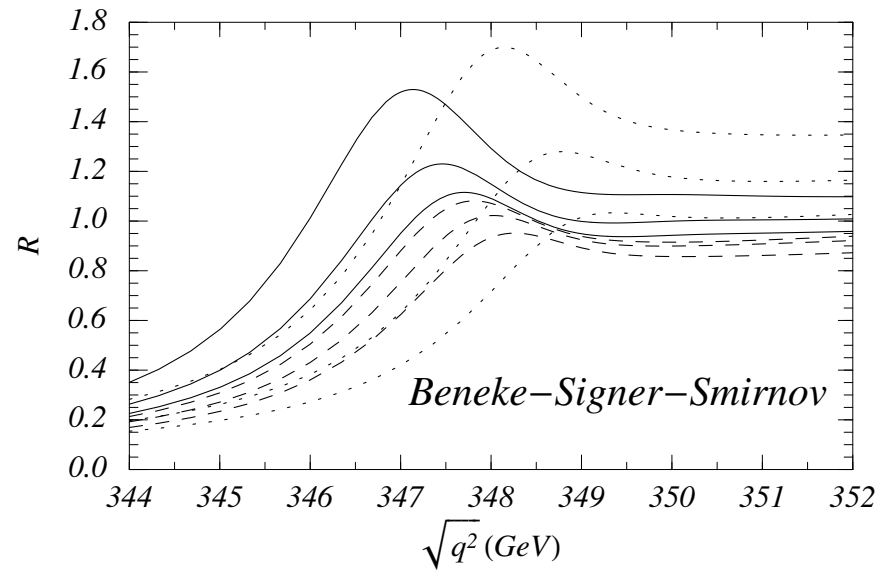
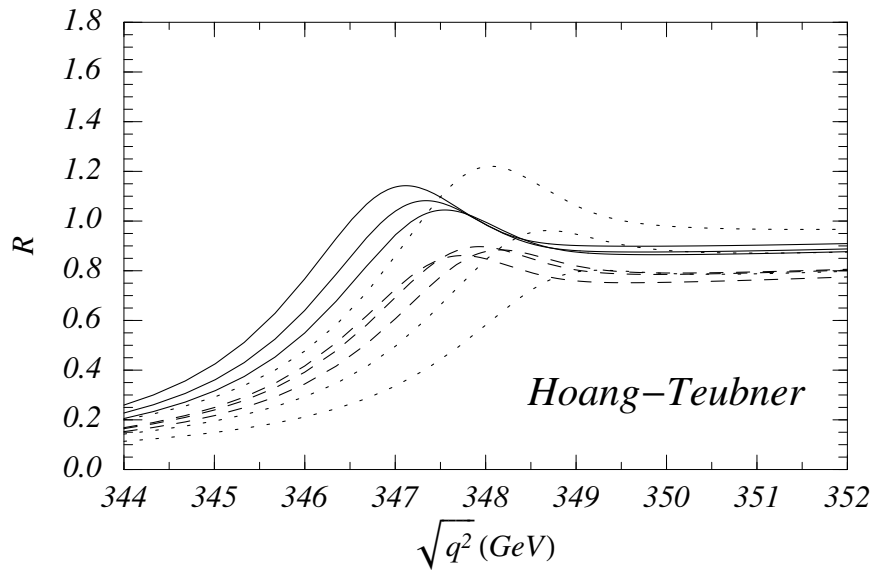
Matthias Steinhauser, University of Karlsruhe

LCWS05, March 2005

[In collaboration with A. Penin and V. Smirnov]

Motivation I — top

$$\sigma(e^+e^- \rightarrow t\bar{t}) \text{ @ threshold}$$



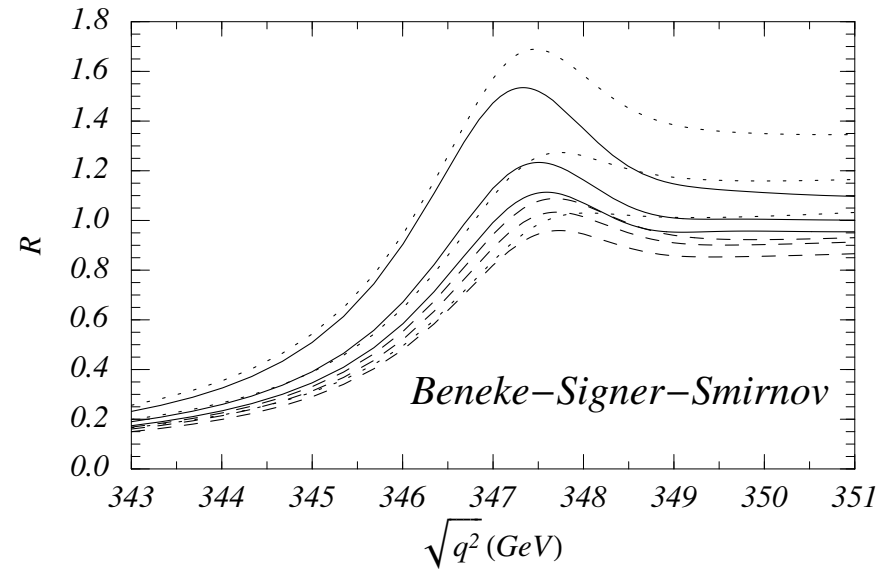
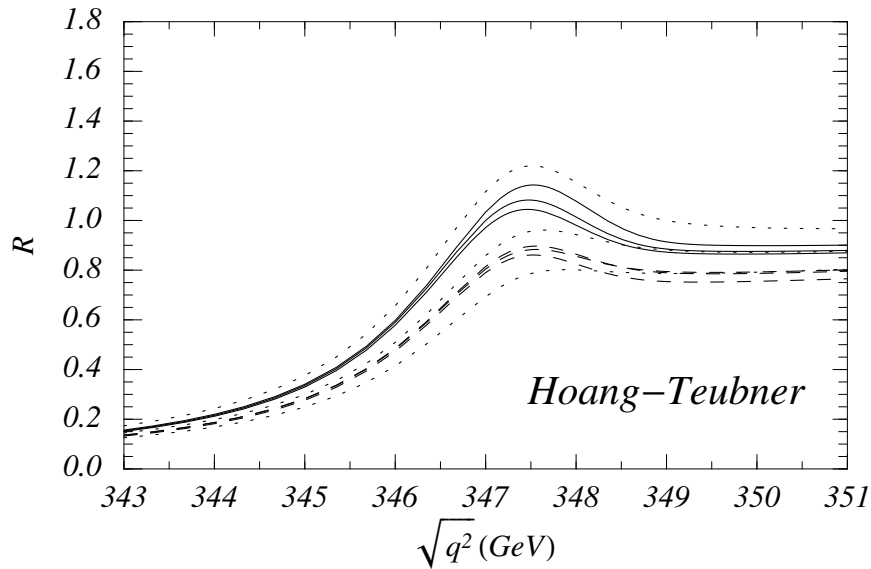
pole mass, NNLO:

- 4 different groups
- no stability in **position** of peak
- no stability in **normalization** of peak

[A.H. Hoang, M. Beneke, K. Melnikov, T. Nagano, A. Ota, A.A. Penin, A.A. Pivovarov, A. Signer, V.A. Smirnov, Y. Sumino, T. Teubner, O. Yakovlev, A. Yelkhovsky'00]

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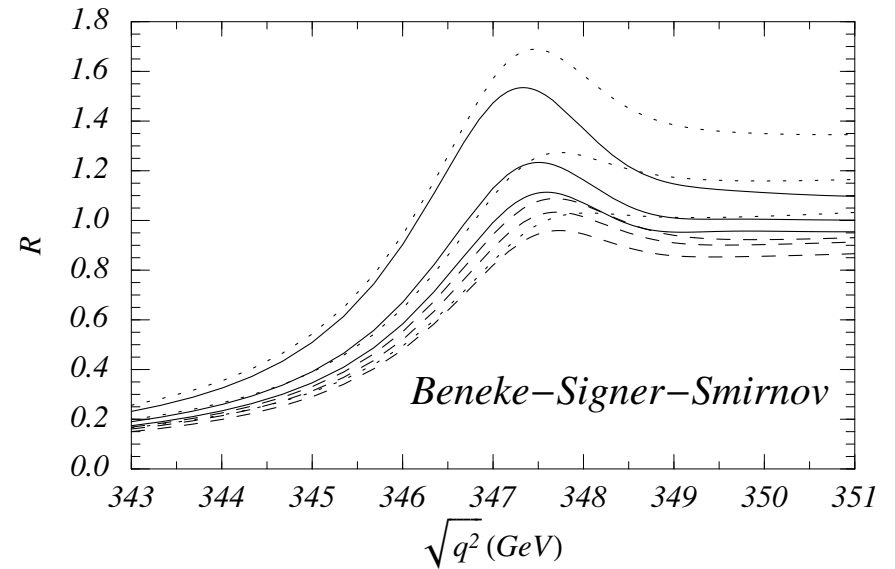
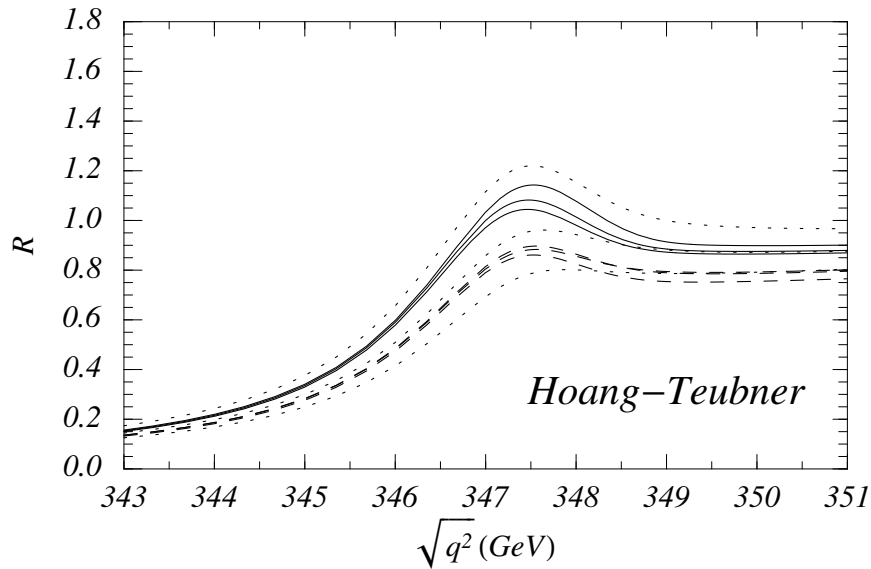
“Threshold” mass, NNLO:

- stability in **position** of peak
- no stability in **normalization** of peak

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Resummation of logarithms (to NNLL):

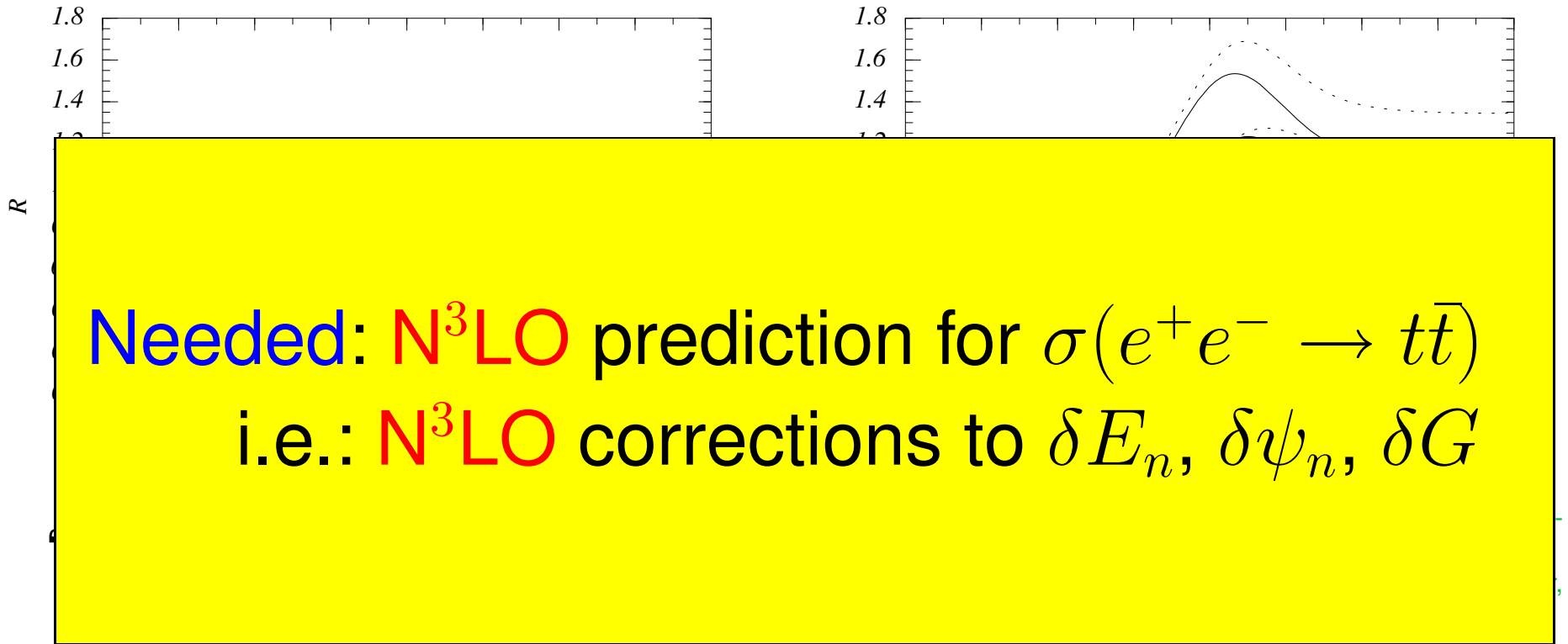
promising; however: **not complete**

[Hoang, Manohar, Stewart, Teubner'01, Hoang'04]

[Penin, Pineda, Smirnov, MS'04]

Motivation I — top

$\sigma(e^+e^- \rightarrow t\bar{t})$ @ threshold



V.A. Smirnov, Y. Sumino, T. Teubner, O. Yakovlev,

A. Yelkhovsky'00]

- stability in position of peak
- no stability in normalization of peak

Resummation of logarithms (to NNLL):

[Hoang,Manohar,Stewart,Teubner'01,Hoang'04]

promising; however: not complete

[Penin,Pineda,Smirnov,MS'04]

Motivation II — bottom

Determination of m_b from bottom system:

- low-moment SR:

NNLO: $\delta m_b = 50 \text{ MeV}$

[Kühn,Steinhauser'01; Corcella,Hoang'03]

N³LO: still missing

- $\Upsilon(1S)$ -system: $M_{\Upsilon(1S)} = 2m_b + E_1^{\text{p.t.}} + \delta^{\text{n.p.}} E_1$

needed: $E_1^{\text{p.t.}}$ to **N³LO**

[Penin,MS'02]

$\Leftrightarrow \delta m_b = 70 \text{ MeV}$

- Υ SR:

NNLO: $\delta m_b = 50 - 100 \text{ MeV}$

[Penin,Pivovarov'98; Beneke,Signer'99; Melnikov,Yelkhovsky'00; Hoang'00]

N³LO: ???

needed: δE_n and $\delta \psi_n$ to N³LO

Framework

$$G(0, 0, E) = \sum_{n=0}^{\infty} \frac{|\psi_n(0)|^2}{E - E_n}$$

$$|\psi_n(0)|^2 = |\psi_n^C(0)|^2 \left(1 + \delta\psi_n^{(1)} + \delta\psi_n^{(2)} + \delta\psi_n^{(3)} + \dots \right)$$

$$E_n = E_n^C \left(1 + \delta E_n^{(1)} + \delta E_n^{(2)} + \delta E_n^{(3)} + \dots \right)$$

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1. insert

2. expand for $E \rightarrow E_n^C$

single pole: $\Delta^{(n)}\psi$

double pole: $\Delta^{(n)}E$

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$$\delta G(0, 0, E) \sim \langle G_C | \mathcal{H} | G_C \rangle + \langle G_C | \mathcal{H} | G_C | \mathcal{H} | G_C \rangle + \dots$$

Framework

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$$\delta G(0, 0, E) \sim \langle G_C | \mathcal{H} | G_C \rangle + \langle G_C | \mathcal{H} | G_C | \mathcal{H} | G_C \rangle + \dots$$

$$G_C(\mathbf{x}, \mathbf{y}, k) \sim \sum_{l=0}^{\infty} (2l+1) \sum_{m=0}^{\infty} \frac{L_m^{2l+1}(2kx) L_m^{2l+1}(2ky) m!}{(m+l+1-\alpha_s C_F m_q/(2k)) (m+2l+1)!} P_l((\mathbf{x}\mathbf{y})/xy)$$

⇒ single, double and triple sums

⇒ analytical solution in terms of ζ function possible

Framework II

- potential NRQCD [Pineda,Soto'98,Brambilla,Pineda,Soto,Vairo'00]
- 2 expansion parameters: $\alpha_s, 1/m_q$
- dynamical degrees of freedom:
potential quarks, ultra-soft gluons
- effective Hamiltonian known to N³LO [Kniehl,Penin,Smirnov,MS'02]

Effective Hamiltonian

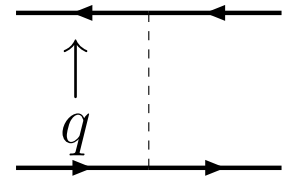
$$\mathcal{H} = (2\pi)^3 \delta(\vec{q}) \left(\frac{\vec{p}^2}{m_q} - \frac{\vec{p}^4}{4m_q^3} \right) + C_c(\alpha_s) V_C(|\vec{q}|) + C_{1/m}(\alpha_s) V_{1/m}(|\vec{q}|) \\ + \frac{\pi C_F \alpha_s(\mu)}{m_q^2} \left[C_\delta(\alpha_s) + C_p(\alpha_s) \frac{\vec{p}^2 + \vec{p}'^2}{2\vec{q}^2} + C_s(\alpha_s) \vec{S}^2 \right]$$

Static potential: $V_C(|\vec{q}|) = -\frac{4\pi C_F \alpha_s(|\vec{q}|)}{\vec{q}^2}$

$1/m_q$ potential: $V_{1/m}(|\vec{q}|) = \frac{\pi^2 C_F \alpha_s^2(|\vec{q}|)}{m_q |\vec{q}|}$

“Breit” potential: $\propto 1/m_q^2$

$$\vec{q} = \vec{p}' - \vec{p}$$



$$C_C(\alpha_s) = 1 + \frac{\alpha_s(|\mathbf{q}|)}{4\pi} a_1 + \left(\frac{\alpha_s(|\mathbf{q}|)}{4\pi} \right)^2 a_2 + \left(\frac{\alpha_s(|\mathbf{q}|)}{4\pi} \right)^3 \left(a_3 + 8\pi^2 C_A^3 \ln \frac{\mu^2}{q^2} \right) + \dots$$

$\vec{S}$: spin; $L = 0$; no $1/m_q^3$ (tree-level) potential

Perturbation theory for E_n

Energy level: $E_n = E_n^C + \delta E_n^{(1)} + \delta E_n^{(2)} + \delta E_n^{(3)} + \dots$ $E_n^C = -\frac{C_F^2 \alpha_s^2 m_q}{4n^2}$

(a) $\delta E_n^{(1)} : \langle G_C | \delta \mathcal{H}^{\text{NLO}} | G_C \rangle$

(b) $\delta E_n^{(2)} : \langle G_C | \delta \mathcal{H}^{\text{NNLO}} | G_C \rangle$, 2nd iteration of $\delta \mathcal{H}^{\text{NLO}}$

(c) $\delta E_n^{(3)} :$

- $\langle G_C | \delta \mathcal{H}^{\text{N}^3\text{LO}} | G_C \rangle$

- iteration of $\delta \mathcal{H}^{\text{NLO}}$ and $\delta \mathcal{H}^{\text{NNLO}}$; 3rd iteration of $\delta \mathcal{H}^{\text{NLO}}$

- ultrasoft contribution:

- $\langle G_C | \delta^{\text{us}} \mathcal{H} | G_C \rangle$

- retarded ultrasoft contribution

$$\delta E_1^{(3)} = \delta E_1^{(3)} \Big|_{\beta(\alpha_s)=0} + \delta E_1^{(3)} \Big|_{\beta(\alpha_s)}$$

$$\begin{aligned} \delta E_1^{(3)} \Big|_{\beta(\alpha_s)=0} = & -E_1^C \frac{\alpha_s^3}{\pi} \left\{ -\frac{a_1 a_2 + a_3}{32\pi^2} + \left[-\frac{C_A C_F}{2} + \left(-\frac{19}{16} + \frac{S(S+1)}{2} \right) C_F^2 \right] a_1 + \left[-\frac{1}{36} + \frac{\ln 2}{6} + \frac{L_{\alpha_s}}{6} \right] C_A^3 \right. \\ & + \left[-\frac{49}{36} + \frac{4}{3} (\ln 2 + L_{\alpha_s}) \right] C_A^2 C_F + \left[-\frac{5}{72} + \frac{10}{3} \ln 2 + \frac{37}{6} L_{\alpha_s} + \left(\frac{85}{54} - \frac{7}{6} L_{\alpha_s} \right) S(S+1) \right] C_A C_F^2 \\ & + \left[\frac{50}{9} + \frac{8}{3} \ln 2 + 3L_{\alpha_s} - \frac{S(S+1)}{3} \right] C_F^3 + \left[-\frac{32}{15} + 2 \ln 2 + (1 - \ln 2) S(S+1) \right] C_F^2 T_F \\ & \left. + \frac{49 C_A C_F T_F n_l}{36} + \left[\frac{11}{18} - \frac{10}{27} S(S+1) \right] C_F^2 T_F n_l + \frac{2}{3} C_F^3 L_1^E \right\} \end{aligned}$$

[Kniehl, Penin, Smirnov, MS'01]

$$\begin{aligned} \delta E_1^{(3)} \Big|_{\beta(\alpha_s)} = & E_n^C \left(\frac{\alpha_s}{\pi} \right)^3 \left\{ 32\beta_0^3 L_\mu^3 + [40\beta_0^3 + 12a_1\beta_0^2 + 28\beta_1\beta_0] L_\mu^2 + \left[\left(\frac{16\pi^2}{3} + 64\zeta(3) \right) \beta_0^3 + 10a_1\beta_0^2 \right. \right. \\ & + \left(40\beta_1 + \frac{a_1^2}{2} + a_2 + 8\pi^2 C_A C_F + \left(\frac{21\pi^2}{2} - \frac{16\pi^2}{3} S(S+1) \right) C_F^2 \right) \beta_0 + 3a_1\beta_1 + 4\beta_2 \Big] L_\mu \\ & + \left(-8 + 4\pi^2 + \frac{2\pi^4}{45} + 64\zeta(3) - 8\pi^2\zeta(3) + 96\zeta(5) \right) \beta_0^3 + \left(\frac{2\pi^2}{3} + 8\zeta(3) \right) a_1\beta_0^2 \\ & + \left(\left(8 + \frac{7\pi^2}{3} + 16\zeta(3) \right) \beta_1 - \frac{a_1^2}{8} + \frac{3}{4} a_2 + \left(6\pi^2 - \frac{2\pi^4}{3} \right) C_A C_F \right. \\ & \left. \left. + \left(8\pi^2 - \frac{4\pi^4}{3} + \left(-\frac{4\pi^2}{3} + \frac{4\pi^4}{9} \right) S(S+1) \right) C_F^2 \right) \beta_0 + 2a_1\beta_1 + 4\beta_2 \right\} \end{aligned}$$

[Penin, MS'02]

also known: $n = 2, 3$

$$L_{\alpha_s} = -\ln(C_F \alpha_s), \quad L_\mu = \ln(\mu / (C_F \alpha_s m))$$

[Penin, Smirnov, MS'05; Beneke, Kiyo, Schuller'05]

Perturbation theory for ψ_n

Wave function:

$$|\psi_n(0)|^2 = |\psi_n^C(0)|^2 \left(1 + \delta\psi_n^{(1)} + \delta\psi_n^{(2)} + \delta\psi_n^{(3)} + \dots \right) + \dots$$

$$|\psi_n^C(0)|^2 = \frac{C_F^3 \alpha_s^3 m_q^3}{8\pi n^3}$$

$$\delta\psi_n^{(3)} = \underbrace{K_2 \ln^2 \alpha_s + K_1 \ln \alpha_s + K_0}_{\text{known}} \Big|_{\beta_0^3} + \underbrace{K_0}_{\text{unknown}} \Big|_{\text{rem}}$$

$$V_C(r) = -\frac{C_F \alpha_s}{r} \left\{ 1 + \frac{\alpha_s}{4\pi} (8\beta_0 L_r + a_1) + \left(\frac{\alpha_s}{4\pi} \right)^2 \left[64\beta_0^2 L_r^2 + (16a_1\beta_0 + 32\beta_1) L_r \right. \right. \\ \left. \left. + a_2 + \frac{16\pi^2}{3} \beta_0^2 \right] + \left(\frac{\alpha_s}{4\pi} \right)^3 \left[512\beta_0^3 L_r^3 + (192a_1\beta_0^2 + 640\beta_0\beta_1) L_r^2 \right. \right. \\ \left. \left. + (128\pi^2\beta_0^3 + 24a_2\beta_0 + 64a_1\beta_1 + 128\beta_2 + 16\pi^2 C_A^3) L_r \right. \right. \\ \left. \left. + a_3 + 16\pi^2 a_1\beta_0^2 + 1024\zeta(3)\beta_0^3 + \frac{160\pi^2}{3} \beta_0\beta_1 \right] + \mathcal{O}(\alpha_s^4) \right\}$$

$$L_r = \ln(e^{\gamma_E} \mu r)$$

$$\left. \delta\psi_n^{(3)} \right|_{\ln^2 \alpha_s} \quad \text{and} \quad \left. \delta\psi_n^{(3)} \right|_{\ln \alpha_s}$$

$$\begin{aligned} \delta_{\ln^2 \alpha_s} \psi_1^{(3)} &= \frac{\alpha_s^3}{\pi} \left\{ \left(-2C_A C_F + \left(-4 + \frac{4}{3} S(S+1) \right) C_F^2 \right) \beta_0 - \frac{2}{3} C_A^2 C_F \right. \\ &\quad \left. + \left(-\frac{41}{12} + \frac{7}{12} S(S+1) \right) C_A C_F^2 - \frac{3}{2} C_F^3 \right\} \ln^2 (C_F \alpha_s) \end{aligned}$$

[Kniehl, Penin'00; Manohar, Stewart'01]

$$\begin{aligned} \delta_{\ln \alpha_s} \psi_1^{(3)} &= \frac{\alpha_s^3}{\pi} \left\{ \left[\left(-3 + \frac{2\pi^2}{3} \right) C_A C_F + \left[\frac{4\pi^2}{3} - \left(\frac{10}{9} + \frac{4\pi^2}{9} \right) S(S+1) \right] C_F^2 \right] \beta_0 \right. \\ &\quad + \left[-\frac{3}{4} C_A C_F + \left(-\frac{9}{4} + \frac{2}{3} S(S+1) \right) C_F^2 \right] a_1 + \frac{1}{4} C_A^3 + \left(\frac{59}{36} - 4 \ln 2 \right) C_A^2 C_F \\ &\quad + \left[\frac{143}{36} - 4 \ln 2 - \frac{19}{108} S(S+1) \right] C_A C_F^2 + \left[-\frac{35}{18} + 8 \ln 2 \right. \\ &\quad \left. - \frac{1}{3} S(S+1) \right] C_F^3 + \left[-\frac{32}{15} + 2 \ln 2 + (1 - \ln 2) S(S+1) \right] C_F^2 T_F \\ &\quad \left. + \frac{49}{36} C_A C_F T_F n_l + \left[\frac{8}{9} - \frac{10}{27} S(S+1) \right] C_F^2 T_F n_l \right\} \ln (C_F \alpha_s) \end{aligned}$$

[Kniehl, Penin, Smirnov, MS'02; Hoang'04]

new: $\delta\psi_n^{(3)} \Big|_{\beta_0^3}$

$$\delta\psi_1^{(3)} \Big|_{\beta_0^3} = \left(\frac{\beta_0 \alpha_s}{\pi} \right)^3 \left[80L_1^3 + \left(52 - \frac{80\pi^2}{3} \right) L_1^2 + \left(-40 - 6\pi^2 + \frac{10\pi^4}{9} + 200\zeta(3) \right) L_1 - 20 + \frac{22\pi^2}{3} - \frac{7\pi^4}{5} + \frac{4\pi^6}{105} + 112\zeta(3) - 12\pi^2\zeta(3) - 16\zeta(3)^2 - 40\zeta(5) \right]$$

$$\delta\psi_2^{(3)} \Big|_{\beta_0^3} = \left(\frac{\beta_0 \alpha_s}{\pi} \right)^3 \left[80L_2^3 + \left(332 - \frac{160\pi^2}{3} \right) L_2^2 + \left(308 - \frac{266\pi^2}{3} + \frac{40\pi^4}{9} + 400\zeta(3) \right) L_2 - 361 + \frac{73\pi^2}{3} - \frac{26\pi^4}{45} + \frac{32\pi^6}{105} + 496\zeta(3) - 48\pi^2\zeta(3) - 128\zeta(3)^2 - 160\zeta(5) \right]$$

$$\delta\psi_3^{(3)} \Big|_{\beta_0^3} = \left(\frac{\beta_0 \alpha_s}{\pi} \right)^3 \left[80L_3^3 + \left(612 - 80\pi^2 \right) L_3^2 + \left(\frac{2893}{3} - 228\pi^2 + 10\pi^4 + 600\zeta(3) \right) L_3 - \frac{100679}{54} + \frac{183\pi^2}{2} + \frac{52\pi^4}{15} + \frac{36\pi^6}{35} + 1374\zeta(3) - 108\pi^2\zeta(3) - 432\zeta(3)^2 - 360\zeta(5) \right]$$

[Penin,Smirnov,MS'05; Beneke,Kiyo,Schuller'05]

$$L_n = \ln \frac{n\mu}{C_F \alpha_s(\mu) m_q}$$

Excited states of bottomonium

$$M_{\Upsilon(nS)} = 2m_b + E_n^{\text{p.t.}} + \delta^{\text{n.p.}} E_n$$

$$\rho_n = \frac{E_n - E_1}{2m_b + E_1} \quad (\text{does not suffer from renormalon ambiguities})$$

$$\begin{aligned} 10^2 \times \rho_2^{\text{p.t.}} &= 1.49 (1 + 0.79_{\text{NLO}} + 1.18_{\text{NNLO}} + 1.21_{\text{N}^3\text{LO}} + \dots) = 6.2^{+1.7}_{-1.2} \\ 10^2 \times \rho_2^{\text{exp}} &= 5.95 \end{aligned}$$

$$\begin{aligned} 10^2 \times \rho_3^{\text{p.t.}} &= 1.77 (1 + 0.92_{\text{NLO}} + 1.37_{\text{NNLO}} + 1.55_{\text{N}^3\text{LO}} + \dots) = 8.6^{+2.4}_{-1.8} \\ 10^2 \times \rho_3^{\text{exp}} &= 9.46 \end{aligned}$$

- ⇒ convergence not good
- ⇒ N³LO terms needed to match experimental result;
impressive agreement
- ⇒ non-perturbative contribution small !?

Υ SR

Compare exp. and th. moments: $\mathcal{M}_k^{\text{th}} \stackrel{!}{=} \mathcal{M}_k^{\text{exp}}$

Experimental input: masses and leptonic width of Υ resonances

Theory input: (corrections to) energy and wave function; non-perturbative effects under control

Estimate N³LO corrections:

$k \geq 20$: corrections to ground state energy dominate

$$\delta\bar{m}_b(\bar{m}_b)^{\text{N}^3\text{LO}} > 0 \text{ small}$$

$k = 4$: corrections to wave function become important

$$\delta\bar{m}_b(\bar{m}_b)^{\text{N}^3\text{LO}} \approx -100 \text{ MeV}$$

⇒ Wait for complete N³LO result

Peak for $\sigma(e^+e^- \rightarrow t\bar{t})$ at N³LO

Normalization of peak

$$R_{\text{res}}(e^+e^- \rightarrow t\bar{t}) = \frac{6\pi N_c Q_t^2}{m_t^2 \Gamma_t} c_v^2 |\psi_1(0)|^2$$

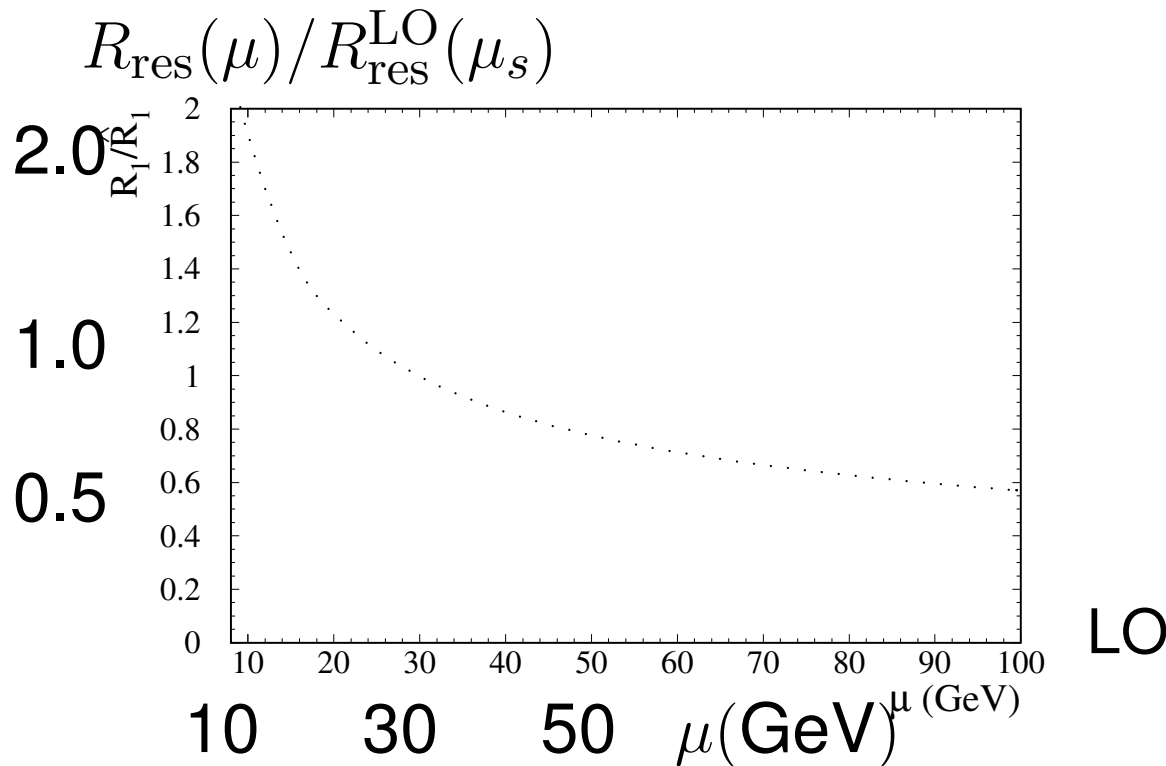
c_v : hard matching coefficient

[Czarnecki,Melnikov'97; Beneke,Signer,Smirnov'97]

$|\psi_1(0)|^2$: $\ln^{1,2} \alpha_s$, β_0^3 terms known at N³LO

[Kniehl,Penin,Smirnov,MS'02;Hoang'04]

[Penin,Smirnov,MS'05]



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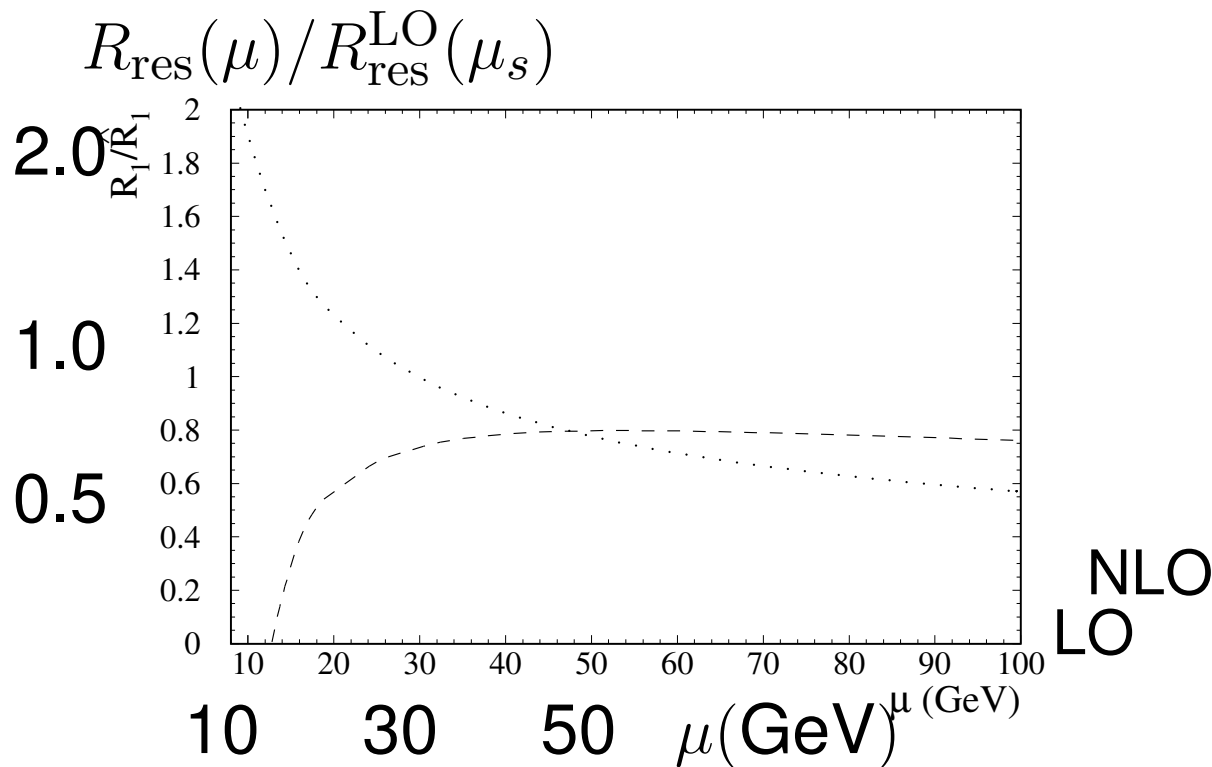
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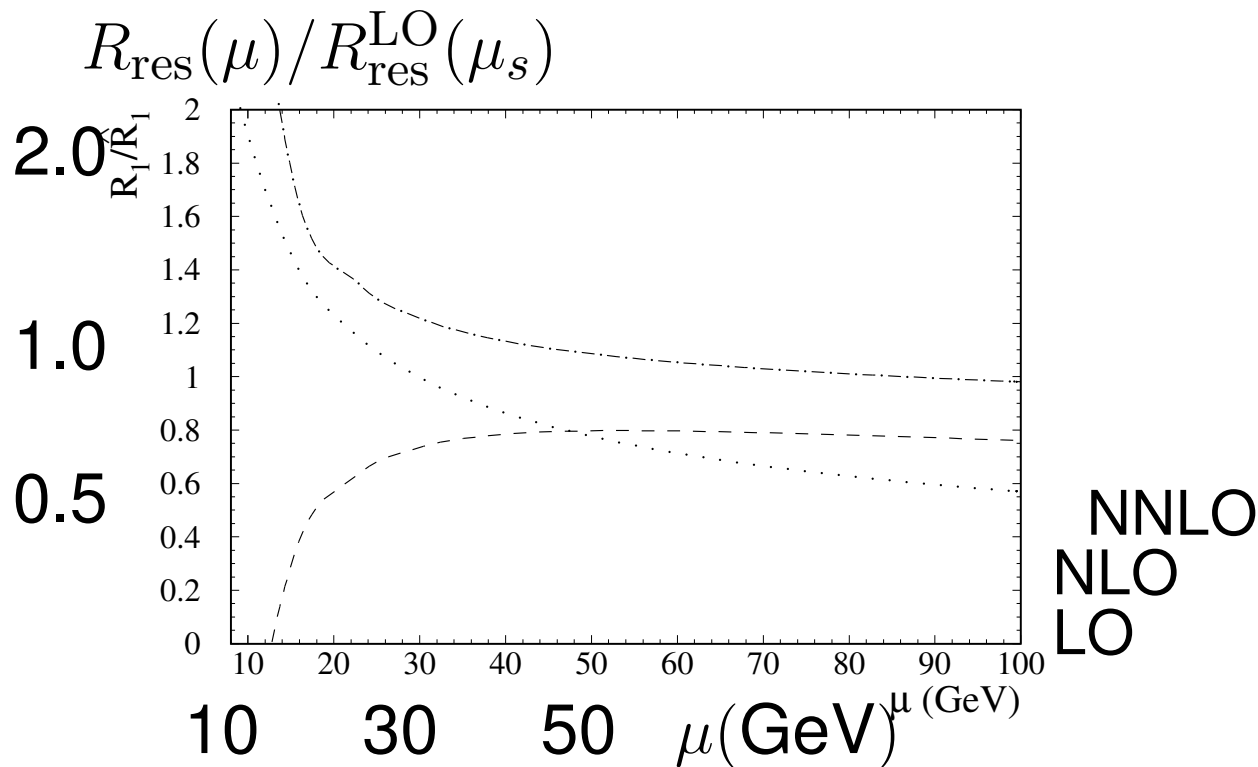
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Normalization of peak

$$R_{\text{res}}(e^+e^- \rightarrow t\bar{t}) = \frac{6\pi N_c Q_t^2}{m_t^2 \Gamma_t} c_v^2 |\psi_1(0)|^2$$

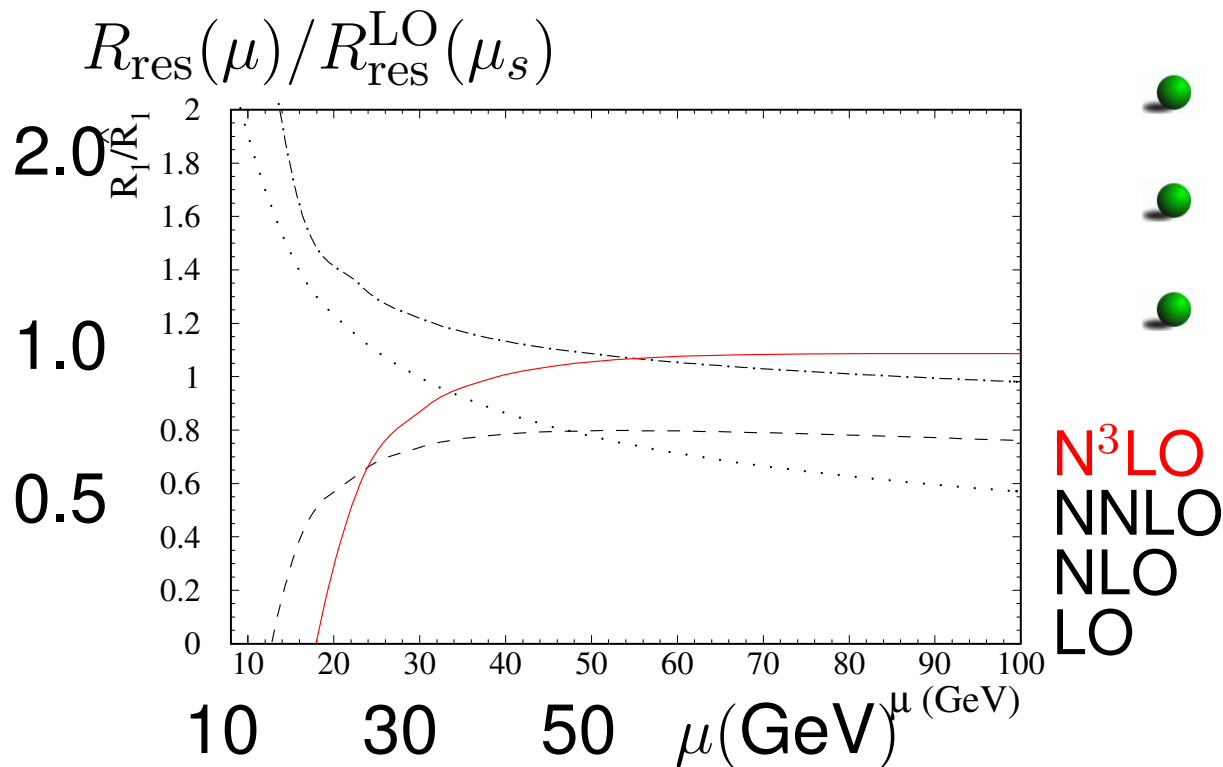
c_v : hard matching coefficient

[Czarnecki,Melnikov'97; Beneke,Signer,Smirnov'97]

$|\psi_1(0)|^2$: $\ln^{1,2} \alpha_s$, β_0^3 terms known at N³LO

[Kniehl,Penin,Smirnov,MS'02;Hoang'04]

[Penin,Smirnov,MS'05]

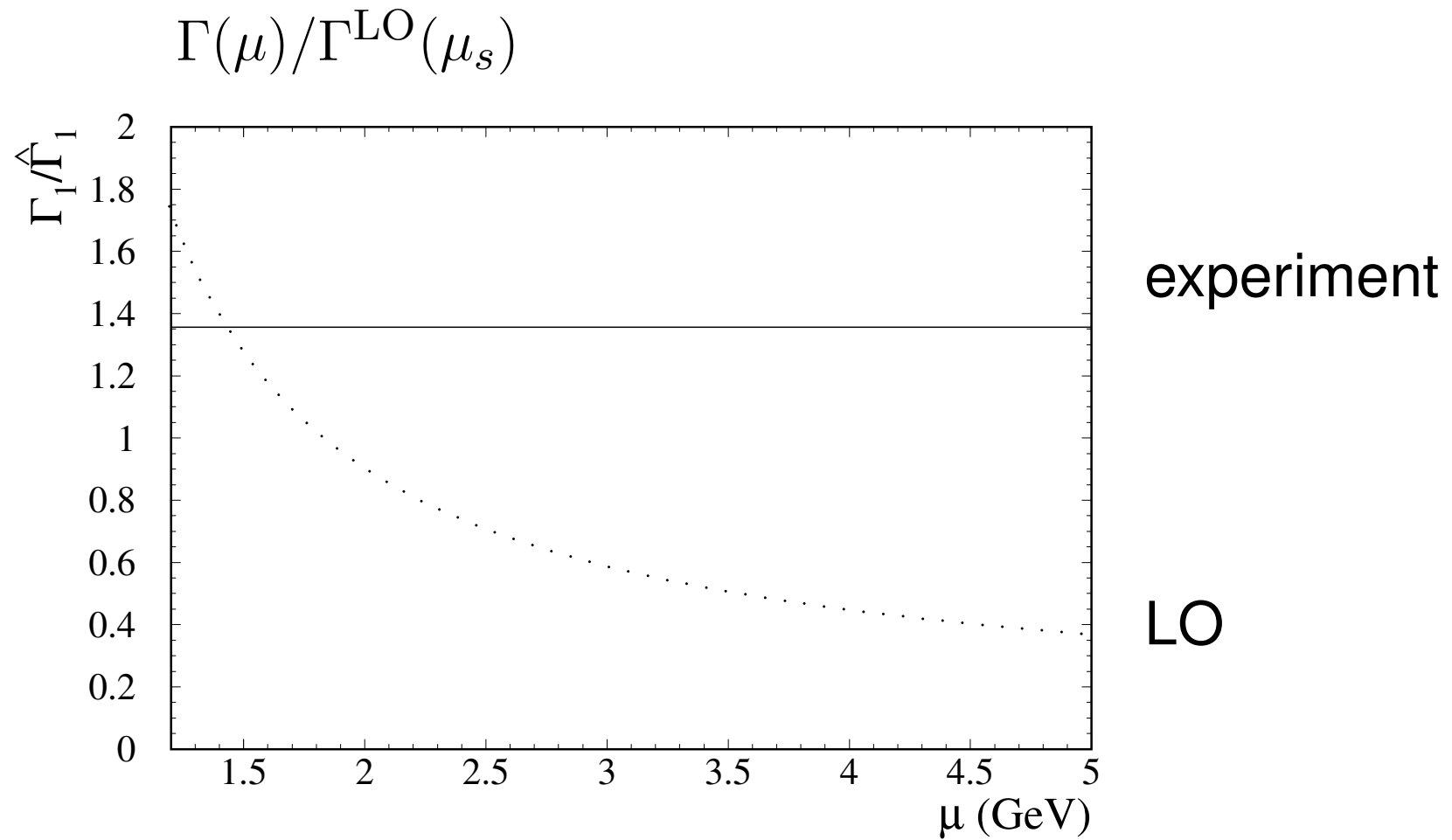


- nice μ dependence
- first sign of convergence
- complete N³LO-constant needed!

N³LO
NNLO
NLO
LO

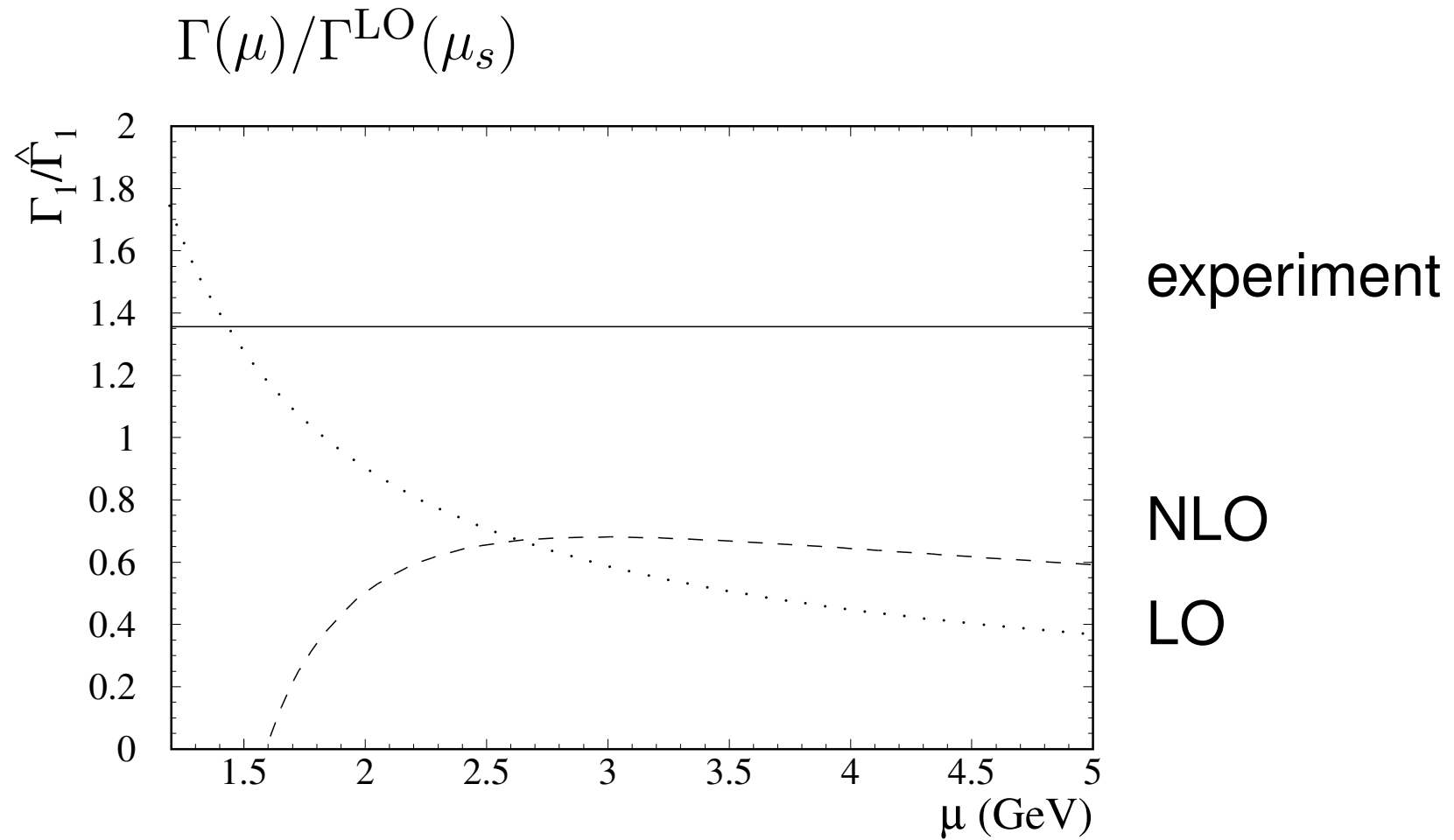
$$\Gamma(\Upsilon(1S) \rightarrow e^+e^-)$$

$$\Gamma^{\text{exp}} = 1.31 \text{ keV}$$



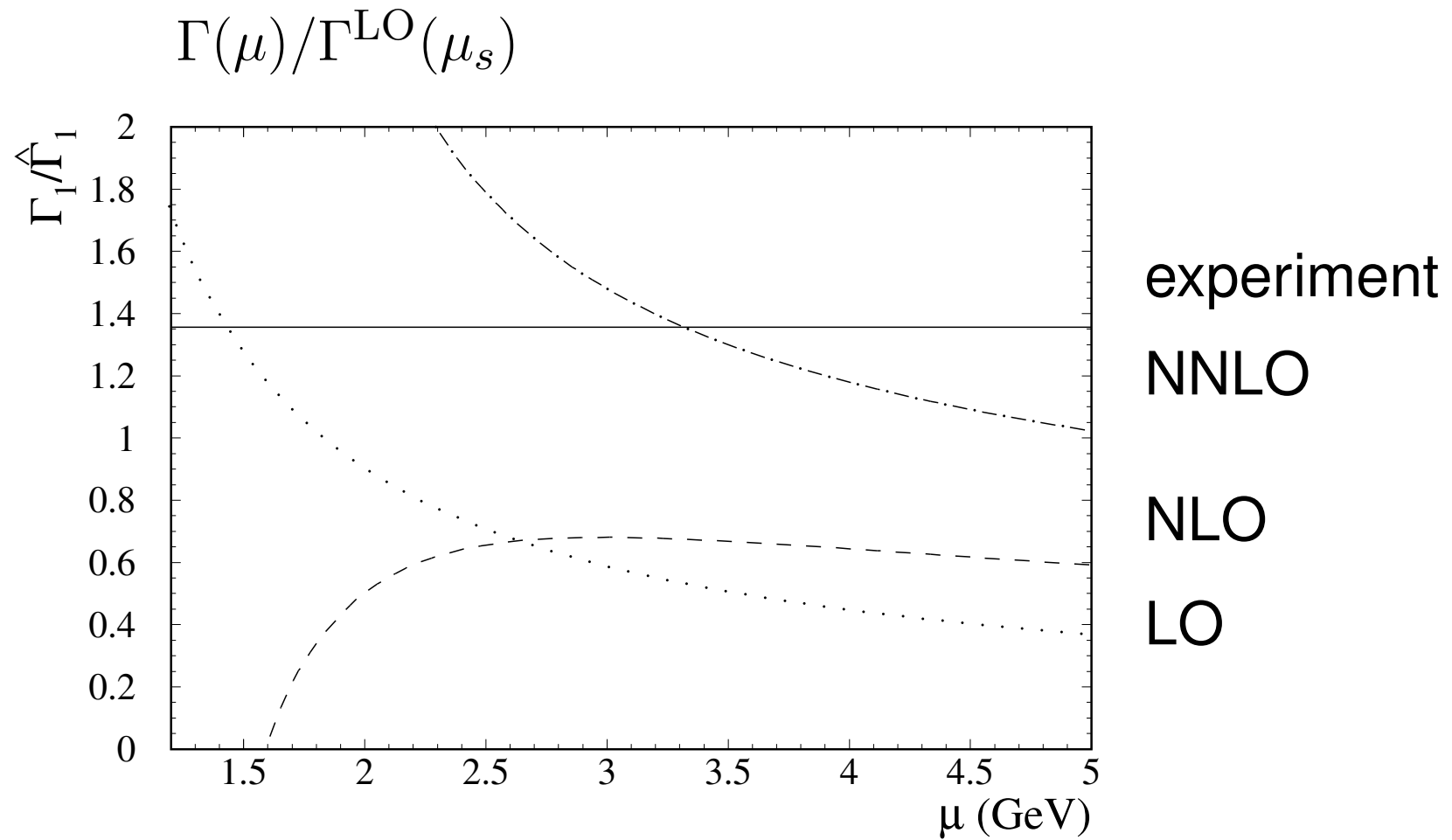
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$$\Gamma(\Upsilon(1S) \rightarrow e^+e^-)$$

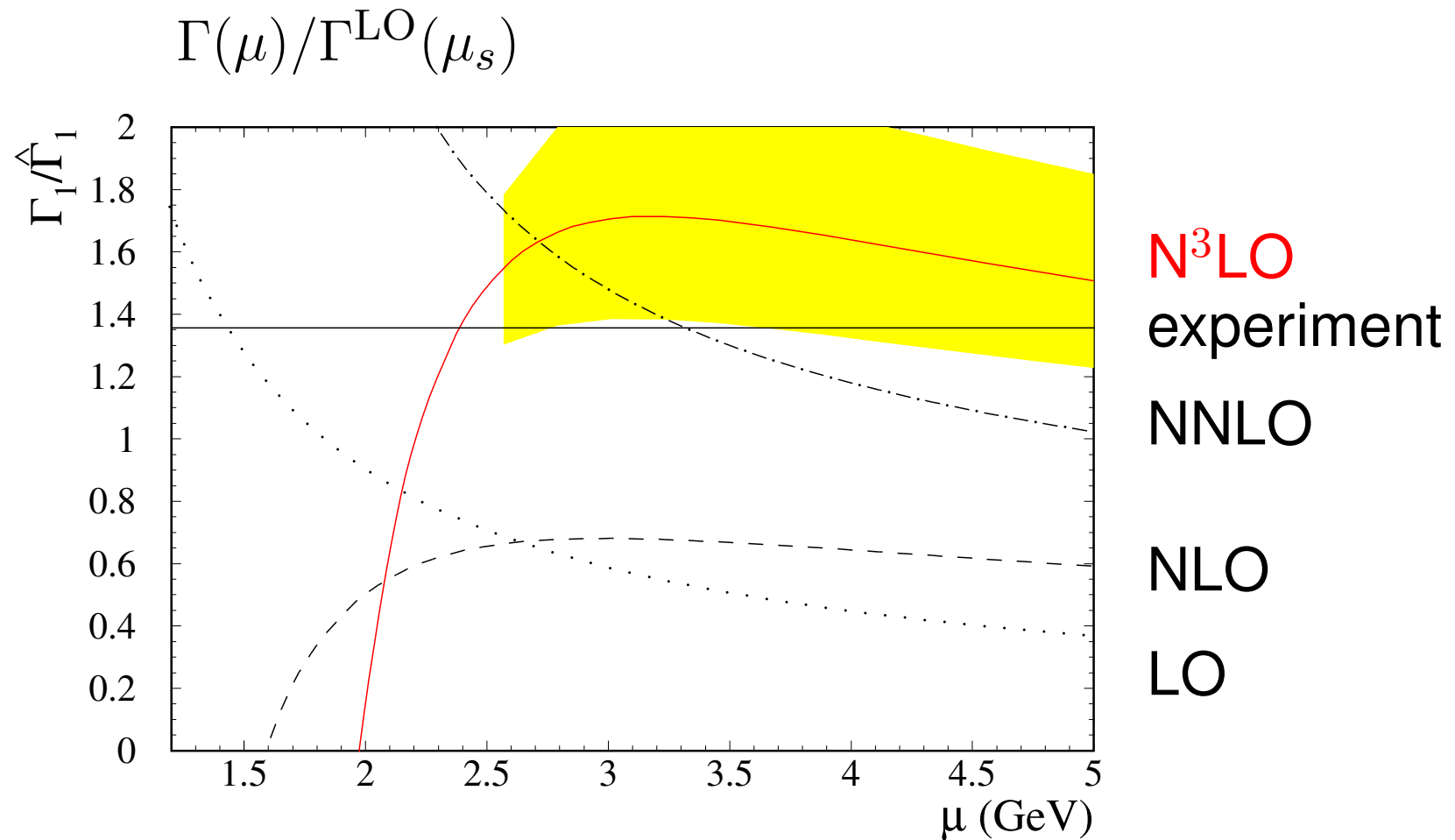
$$\Gamma^{\text{exp}} = 1.31 \text{ keV}$$



$$\Gamma(\Upsilon(1S) \rightarrow e^+e^-)$$

$$\Gamma^{\text{exp}} = 1.31 \text{ keV}$$

$$\alpha_s(M_Z) = 0.118 \pm 0.003$$



Summary

- N³LO corrections to energy levels, $\delta E_n^{(3)}$: complete
- N³LO corrections to wave function, $\delta \psi_n^{(3)}$: $\ln^2$, $\ln$, β_0^3
- stabilization of perturbative series observed
 - excited states of bottomonium
 - $\Upsilon(1S)$ leptonic width
 - Υ sum rules
 - $t\bar{t}$ threshold production
- Third-order corrections to $\delta \psi$ to be completed