

DETERMINING PARAMETERS OF THE LITTLEST HIGGS MODEL AT THE ILC

OUTLINE:

1.) THE LITTLEST HIGGS MODEL

- Introduction and Parameters of the model
- Spectrum + Couplings

2.) ANALYSIS

$$\cdot e^+e^- \rightarrow f\bar{f}$$

- LEP II constraints

- ILC search reach

- ILC Parameter Fits

$$\cdot e^+e^- \rightarrow Z h$$

- ILC search reach

3.) CONCLUSIONS

THE LITTLEST HIGGS (LH) MODEL

- Arkani-Hamed, et. al. [hep-ph/0206021](#)

LIKE OTHER LITTLE HIGGS MODELS:

- Based on a spontaneously broken non-linear σ model
- UV cut-off ~ 10 TeV
- Higgs is pseudo-Goldstone boson of global symmetry
- Higgs mass quadratic divergences come in only at two loops *

FOR THE LITTLEST HIGGS:

* This protects m_h^2 from one-loop div.

$$4\pi f \hookrightarrow [SU(2) \times U(1)]^2 \subset SU(5) \xrightarrow{\text{Global Symmetry}} \begin{matrix} \text{Gauged} \\ \downarrow \quad vev \sim f \quad \downarrow \\ [SU(2)_L \times U(1)_Y] \subset SO(5) \\ \downarrow \quad v \sim f/4\pi \\ \text{EM} \end{matrix}$$

THE LITTLEST HIGGS MODEL CONTD.

PARAMETERS

- f
 - since $v \sim f/4\pi$ and UV cutoff $\sim 4\pi f$, then $f \sim \text{few TeV}$
 - If $f \gtrsim 10 \text{ TeV}$ it no longer addresses naturalness

- s, s'
 - mixing angles for mass e' states after $SU(5) \rightarrow SO(5)$

$$W = sW_1 + cW_2 \quad W' = -cW_1 + sW_2$$

$$B = s'B_1 + c'B_2 \quad B' = -c'B_1 + s'B_2$$

- functions of coupling constants of $[SU(2) \times U(1)]^2$

$$s = \frac{g_2}{\sqrt{g_1^2 + g_2^2}} \quad , \quad s' = \frac{g'_2}{\sqrt{g'^1_1 + g'^2_2}}$$

- The model has an interesting top sector with some more params. For our analysis ($e^+e^- \rightarrow f\bar{f}$ and $e^+e^- \rightarrow Z h$), at tree level, these three are only relevant params.

NEUTRAL GAUGE BOSON SPECTRUM

($\exists N_H^{\pm}$, but not important for our analysis)

Then after EWSB, mass e's states:

A_L → massless (phew!)

Z_L → tiny corrections

$$M_{Z_L}^2 = M_{Z_{SM}}^2 \left[1 + \mathcal{O}(v^2/f^2) \right]$$

A_H → heavy "photon"

$$M_{A_H}^2 = M_{Z_{SM}}^2 S_W^2 \left(\frac{f^2}{5 s'^2 c'^2 v^2} - 1 + \frac{x_H C_W^2}{4 s^2 c^2 S_W^2} \right)$$

Z_H → heavy "Z"

$$M_{Z_H}^2 = M_{Z_{SM}}^2 \left(\underbrace{\frac{f^2}{s^2 c^2 v^2} - 1}_{\text{Dominant term - } S \text{ vs. } S'} - \frac{x_H S_W^2}{4 s'^2 c'^2 C_W^2} \right)$$

Note:

Dominant term - S vs. S' dependence

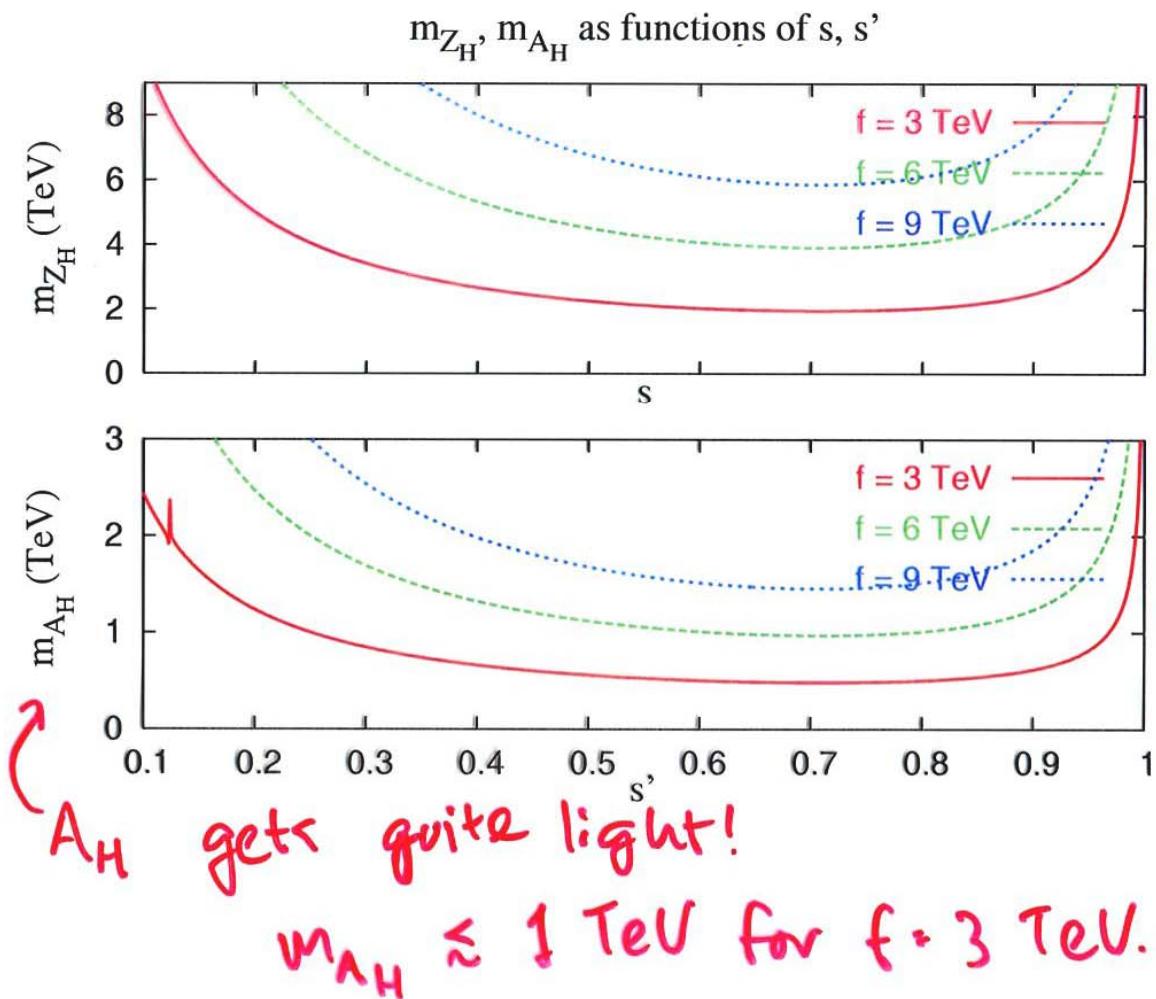
- M_Z^2 corrections ~.01 shift ρ

- Hewett, et. al. hep-ph/0211218

- x_H is $\mathcal{O}(1)$ mixing parameter,
fn. of s, s'

- Han, et. al. hep-ph/0301040

• Masses of Neutral Gauge Bosons



EXISTING CONSTRAINTS

- Strong constraints from Tevatron and precision EW ($f \gtrsim 4$ TeV at best)
 - Hewett, et. al. hep-ph/0211218
 - Csaki, et. al. hep-ph/0211124
- Contributions from light A_H dominate in generic cases
- Many models have been invented that decouple A_H in one way or another
 - e.g. Csaki, et. al. hep-ph/0303236
 - Gregoire, et. al. hep-ph/0305275
 - Kilic + Mahbubani hep-ph/0312053
- We work with Littlest Higgs model b/c it is well studied.

↳ To evade A_H problems, we can:

- We do both -
- 1) Choose parameters ($s' = \sqrt{3/5}$) for which A_H decouples
 - 2) Artificially send $m_{A_H} \rightarrow \infty$
 - This is not theoretically consistent, but gives a general idea of behavior of models that evade A_H problem

COUPLINGS IN LITTLEST HIGGS

- Heavy particle couplings

A_H sector
s'-dep.

$$g_{V/A}(f\bar{f}A_H) \sim \frac{g'}{s'c'} \left(\frac{c'^2}{2} - \frac{1}{5} \right)$$

\hookrightarrow vanish at $s' = \sqrt{3/5}$

Z_H sector
s-dep.

$$g_{V/A}(f\bar{f}Z_H) \sim \frac{g_C}{s}$$

- 'SM' particle couplings shifted

$$\Delta g_{V/A}(f\bar{f}Z_L) \sim v^2/f^2 \quad (\text{both } s \text{ and } s' \text{ dependent})$$

- Note for e⁺e⁻Z_L:

$$|\Delta g_V| \sim |\Delta g_A|, \text{ while} \\ |g_V| \ll |g_A|.$$

But since $\sigma \propto g_V^2 + g_A^2$,

while $A_{LR} \propto g_V g_A$

need A_{LR} to be sensitive to $A_{V/A}$

\hookrightarrow Thus in models with A_H

decoupled, only sensitive
to s' if can measure A_{LR}

\rightarrow Importance of beam polarization

ANALYSIS: $e^+e^- \rightarrow f\bar{f}$

- Looked at

$e^+e^- \rightarrow f\bar{f}$ at

LEP II (constraints)

and LC* (search reach, param fits)

- Final states + efficiencies:

$\tau^+\tau^-$, $\mu^+\mu^-$, e^+e^- (95%), \leftarrow ^{For} _{LC}

$b\bar{b}$ (60%), $c\bar{c}$ (35%).

- Observables:

LEP II: Normalized, binned
angular dist. and total σ

LC: same as above + binned ACR

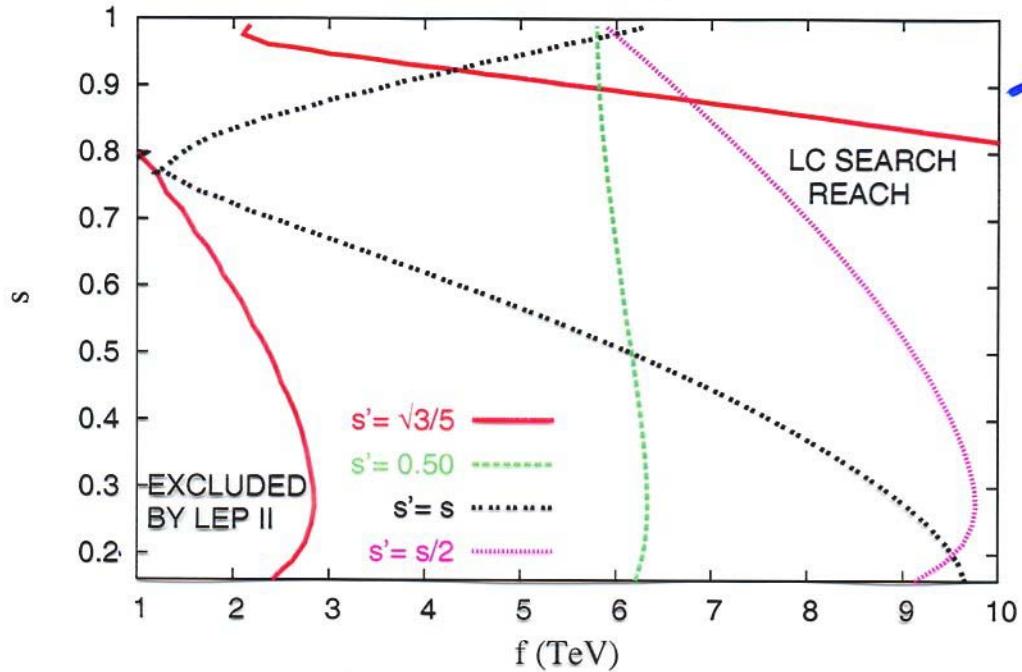
- Used $\chi^2 = \sum_i \left(\frac{\theta_{LH}^i - \theta_{SM}^i}{\delta \theta^i} \right)^2$

- 95% CL Regions correspond to

$\chi^2 > 5.88$ (2 DOF) always fix 1 comb.
of 3 params.

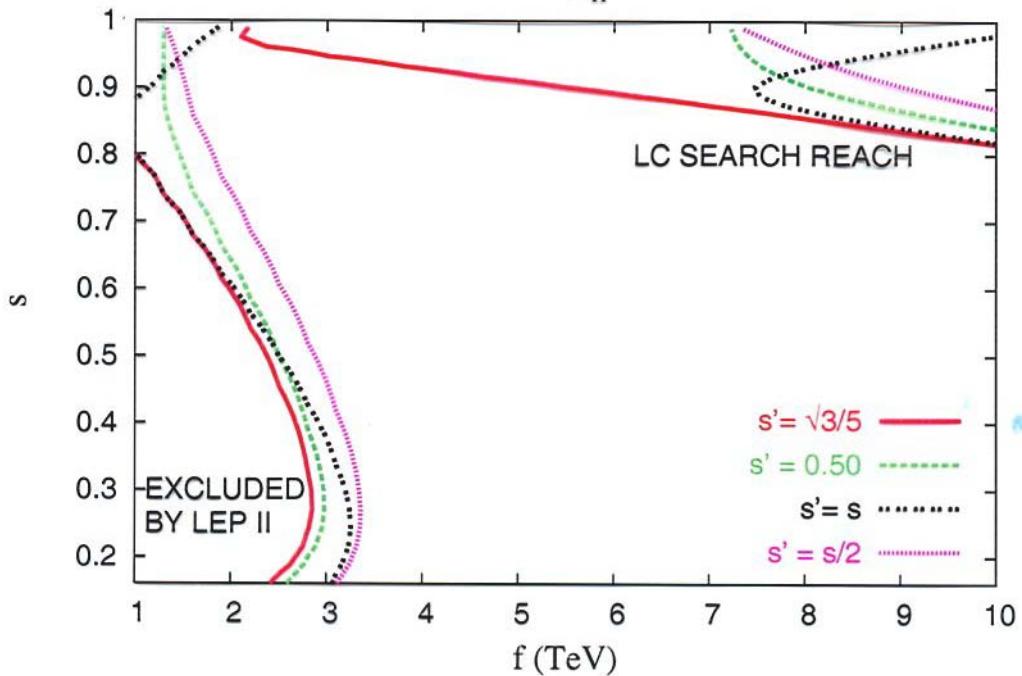
LEP II Constraints and ILC search reach using $e^+e^- \rightarrow f\bar{f}$

95% CL contours with light A_H at $\sqrt{s} = 500$ GeV



$$\int \mathcal{L} dt = 500 \text{ fb}^{-1}$$

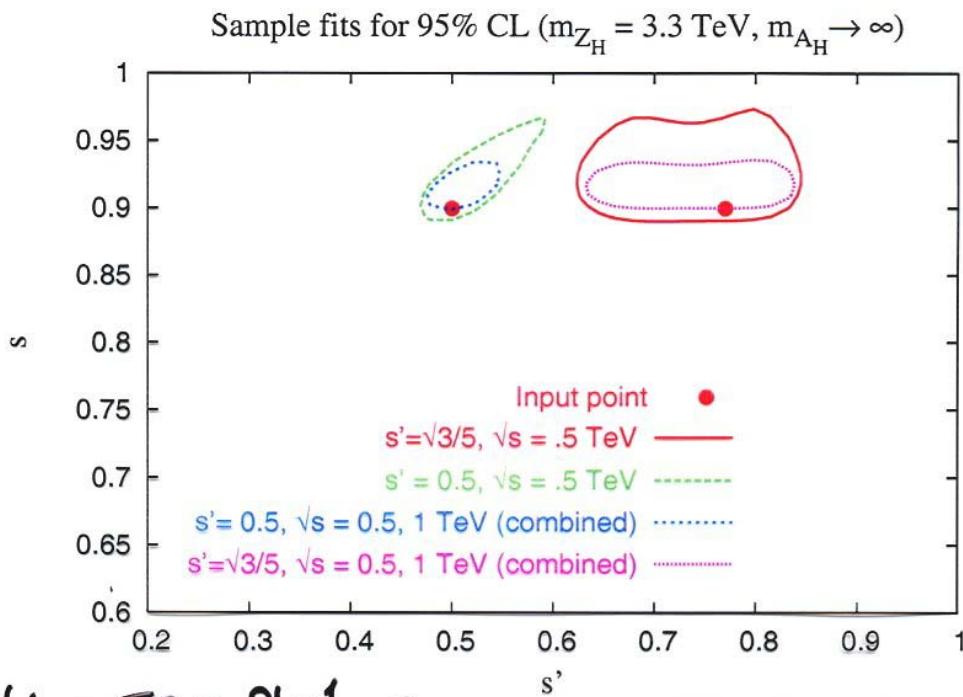
95% CL contours with $m_{A_H} \rightarrow \infty$ at $\sqrt{s} = 500$ GeV



Parameter fits at ILC using

$e^+e^- \rightarrow f\bar{f}$ (fixed m_{Z_H})

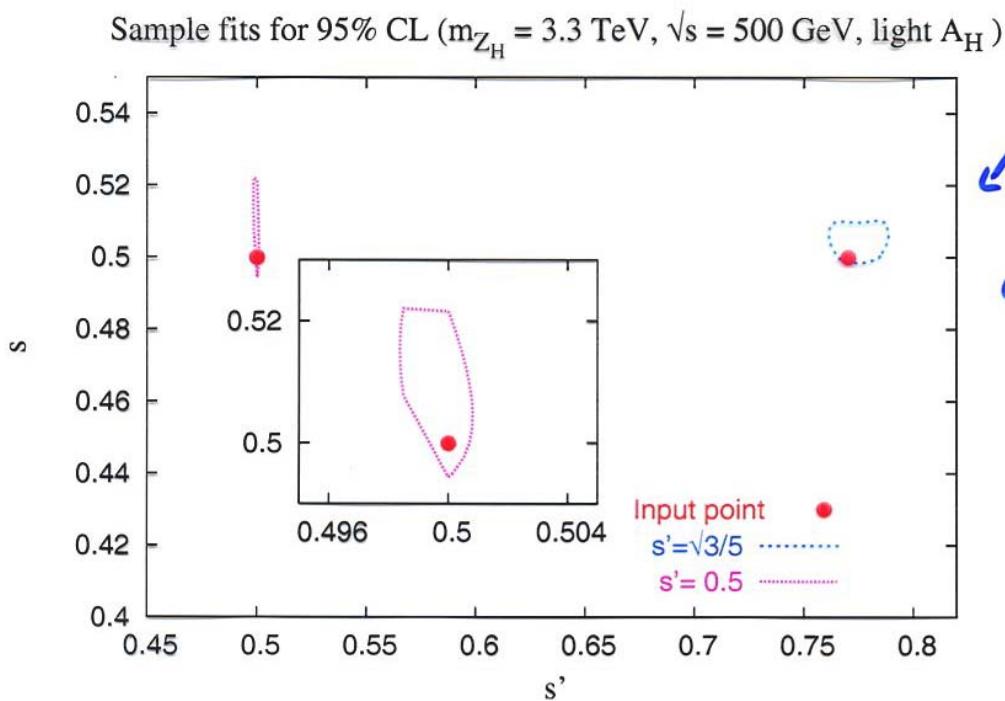
$m_{A_H} \rightarrow \infty$



Combined
 ↙ 500 GeV
 and 1 TeV
 data set
 improves fit

$\int \frac{1}{2} dt = 500 \text{ fb}^{-1}$ for each $\sqrt{s} = 500 \text{ GeV}$ and 1 TeV

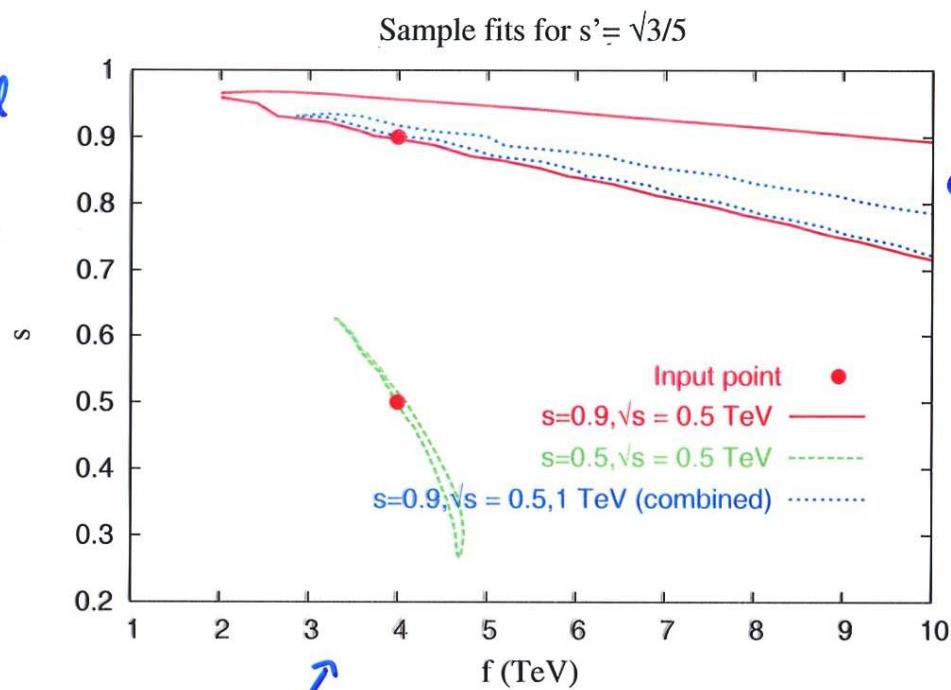
light A_H



Fit much
 better
 with light
 A_H

Parameter fit (fixed s') at ILC using $e^+e^- \rightarrow f\bar{f}$.

- This could happen if LHC doesn't determine masses



$$\int \mathcal{L} dt = 500 \text{ fb}^{-1}$$

ANALYSIS: $e^+e^- \rightarrow Z h$

- IMPORTANT: Hallmark of Little Higgs models is coupling of heavy gauge bosons to $Z h \Rightarrow$ we should measure this coupling to prove we have Little(st) Higgs

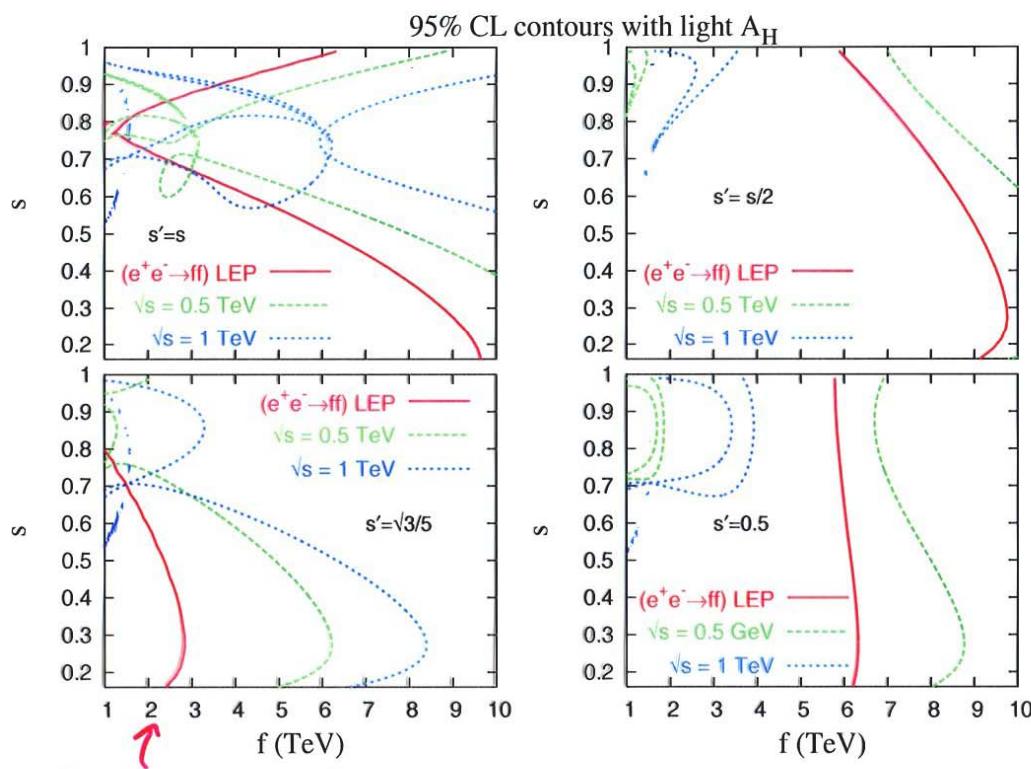


- Deviations from SM in σ_{Zh} :
 - Deviations of $Z Z h$ coupling
 - $\mathcal{O}(v^2/f^2)$, both s and s' dep.
 - Presence of Z_H, A_H in s -channel
 - $Z_H Z h$ coupling $\propto c^2 - s^2$, only s -dep.
(vanishes at $s = 1/\sqrt{2}$)
 - $A_H Z h$ coupling $\propto c'^2 - s'^2$, only s' -dep.
(vanishes at $s' = 1/\sqrt{2}$)
- We use total σ as observable
- We assume that at ILC at $\sqrt{s} = 500$ GeV, $\int \sigma dt = 500 \text{ fb}^{-1}$, we will obtain $\frac{\delta \sigma_{Zh}}{\sigma_{Zh}} \approx 1.5\%$ (TESLA TDR)

ILC Search Reach using $e^+e^- \rightarrow Z h$ with light A_H

$\int L dt = 500 \text{ fb}^{-1}$

for each
500 GeV
and 1 TeV

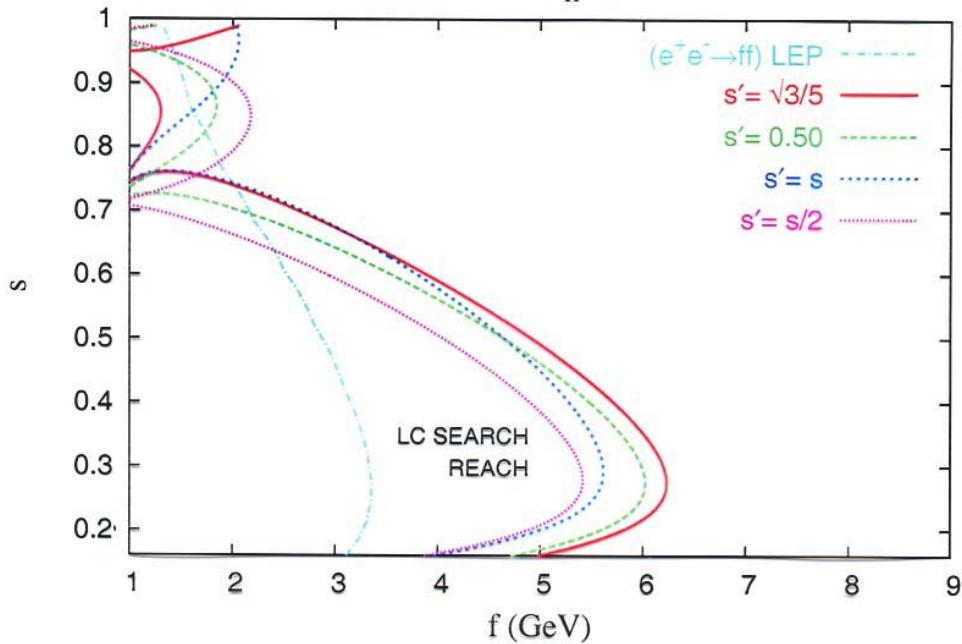


Red is LEP II constraint from $e^+e^- \rightarrow f\bar{f}$

ILC Search reach using $e^+e^- \rightarrow Z h$

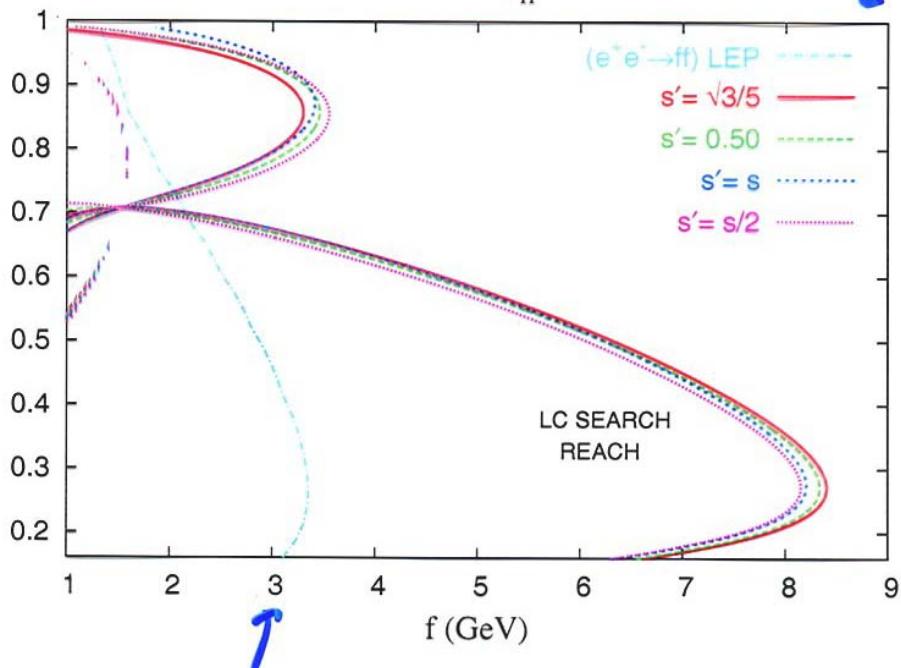
with $m_{A_H} \rightarrow \infty$

95% CL contours with $m_{A_H} \rightarrow \infty$ at $\sqrt{s} = 500$ GeV



$$\int \mathcal{L} dt = 500 \text{ fb}^{-1}$$

95% CL contours with $m_{A_H} \rightarrow \infty$ at $\sqrt{s} = 1$ TeV



$Z_L Z_H h$
coupling \rightarrow
vanishes
at $s = \sqrt{2}/2$

A lot of
param
space
not covered
(maybe other
obs. would
help)

$$\int \mathcal{L} dt = 500 \text{ fb}^{-1}$$

Blue is LEP II constraint from
 $e^+e^- \rightarrow f\bar{f}$

CONCLUSIONS:

FOR THE LITTLEST HIGGS MODEL:

- The ILC covers most of the interesting parameter space
- The ILC provides accurate determination of parameters
(helpful if we have f from LHC)
- $e^+e^- \rightarrow Z h$ gives, in some regions of parameter space, confirmation of the hallmark feature of Little Higgs.