

Physics Impact of Detector Performance

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SLAC

March 18, 2005

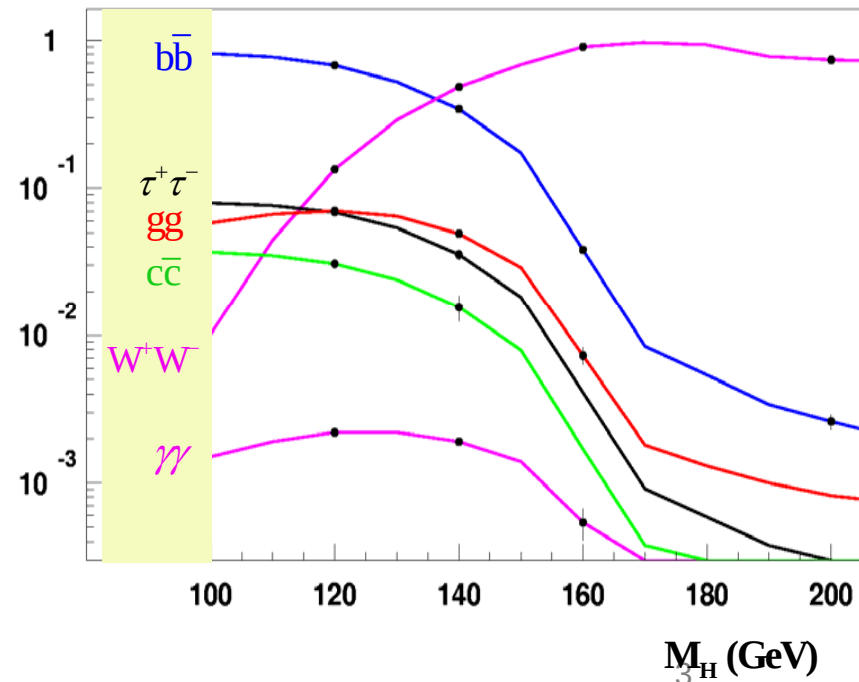
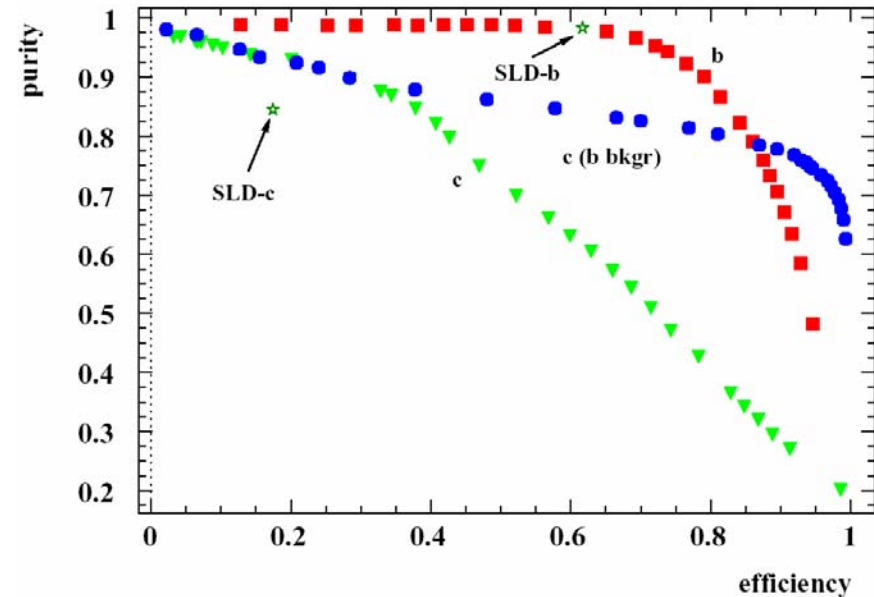
Outline

- General Considerations
 - Vertex Detector
 - Tracker
 - Calorimeter
 - Far forward detector
- Examples of parametric physics studies
 - Calorimeter ΔE_{jet}
 - Tracker Δp_t
- Summary

Vertex Detector

- Classic application of b,c tagging to Higgs branching ratios.
- But there's more:
 - vertex charge
 - top, W helicity
 - $q\bar{q}$ asymmetries
 - τ tagging
 - stau analyses
 - Higgs tau BR
 - b jets with several ν 's

*Talk by Chris Damerell 21Mar2005



Vertex Detector – tau tagging example

$e^+e^- \rightarrow \nu\bar{\nu}t\bar{t}$ is an important strong symmetry breaking signal ($WW \rightarrow t\bar{t}$).

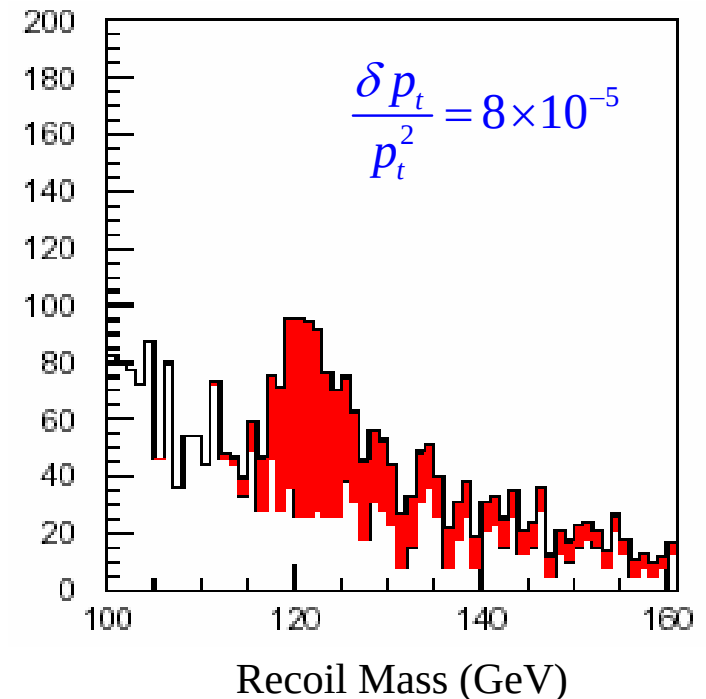
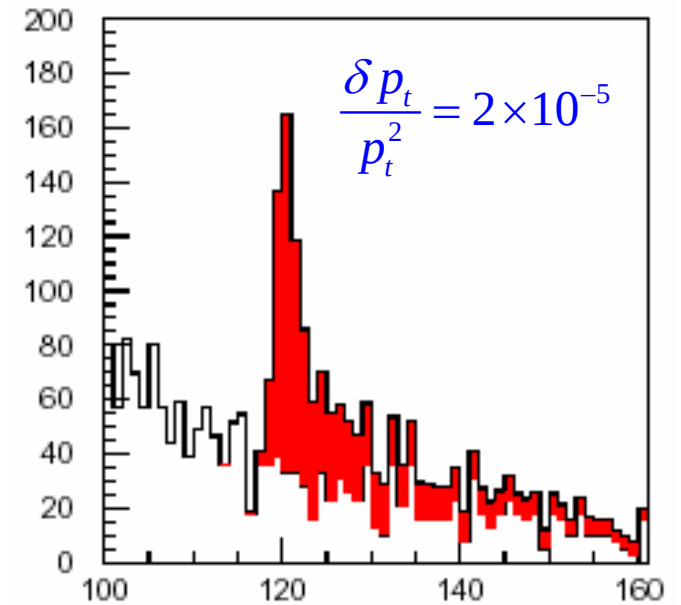
The large $e^+e^- \rightarrow e^+e^-t\bar{t}$ background can be mostly suppressed by vetoing the forward $e^\pm$ and requiring unbalanced p_t . But there remains a seemingly irreducible background from $e^+e^- \rightarrow e^+e^-t\bar{t} \rightarrow e^+e^-b\bar{b}W^+W^-$ where one of the b quarks undergoes the decay

$b \rightarrow c\tau^-\bar{\nu}_\tau \rightarrow c\rho^-\nu_\tau\bar{\nu}_\tau$, $ce^-\bar{\nu}_e\nu_\tau\bar{\nu}_\tau$, etc. $\Rightarrow$ **$b \rightarrow \tau$ decays have at least 2 ν 's**

Tau tagging could reduce this background (and help b jet energy flow analysis in general).

Tracker

- Momentum resolution set by recoil mass analysis of $ZH \rightarrow l^+l^-X$
- K_S^0 , Λ^0 reconstruction and long-lived new particles (GMSB SUSY)
- Multiple scattering effects
- Forward tracking
- Measurement of E_{cm} , differential luminosity and polarization using physics events



Calorimeter

- Separate hadronically decaying W's from Z's in reactions where kinematic fits won't work:

$$e^+e^- \rightarrow \nu\nu W^+W^- , \nu\nu ZZ$$

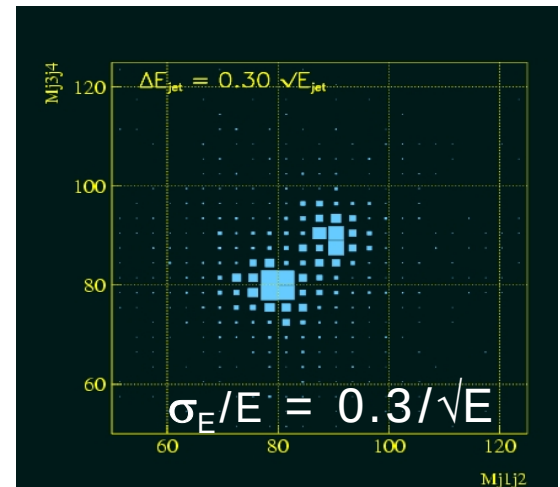
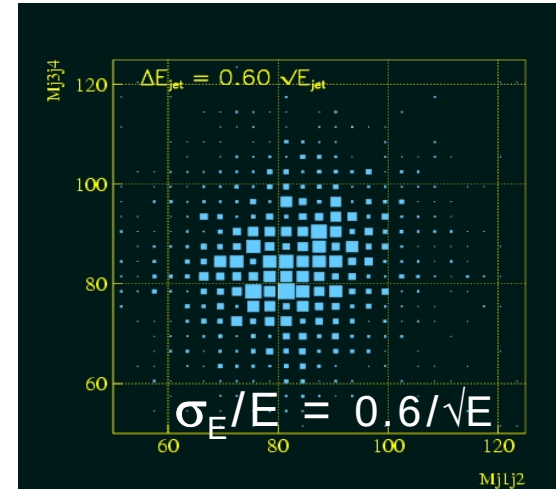
$$e^+e^- \rightarrow \chi_1^+ \chi_1^- \rightarrow \chi_1^0 \chi_1^0 W^+W^-$$

$$e^+e^- \rightarrow \chi_2^0 \chi_2^0 \rightarrow \chi_1^0 \chi_1^0 ZZ$$

- Help solve combinatoric problem in reactions with 4 or more jets

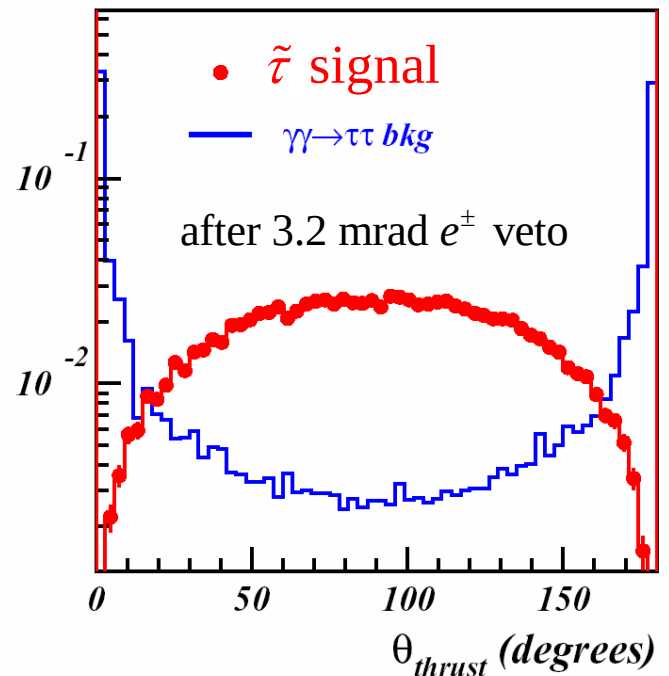
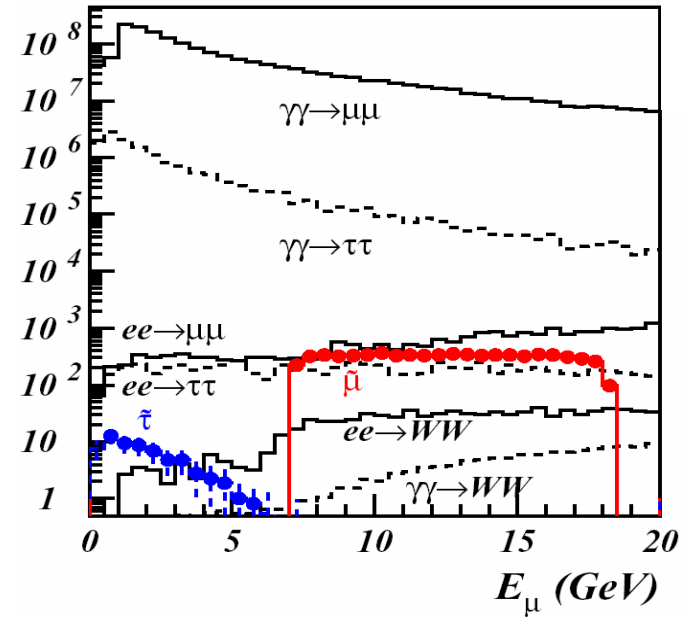
$$e^+e^- \rightarrow ZH \rightarrow q\bar{q}WW^* \rightarrow q\bar{q}q\bar{q}l\nu$$

$$e^+e^- \rightarrow ZHH \rightarrow q\bar{q}b\bar{b}b\bar{b}$$

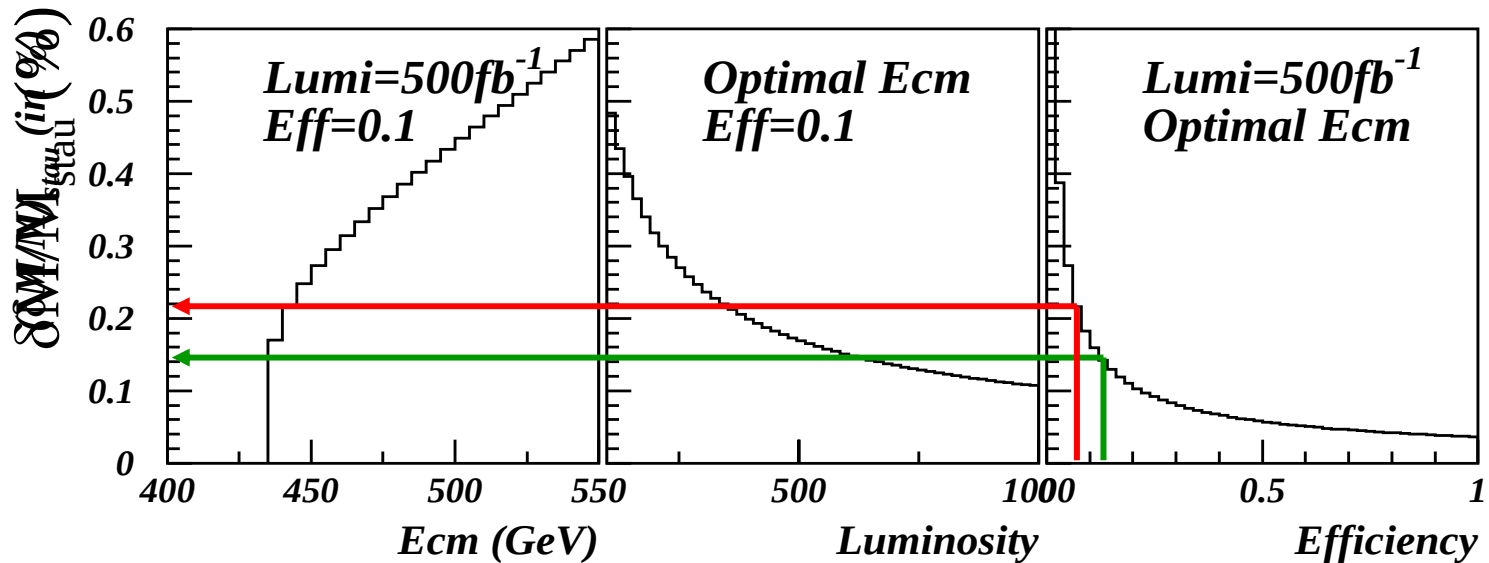
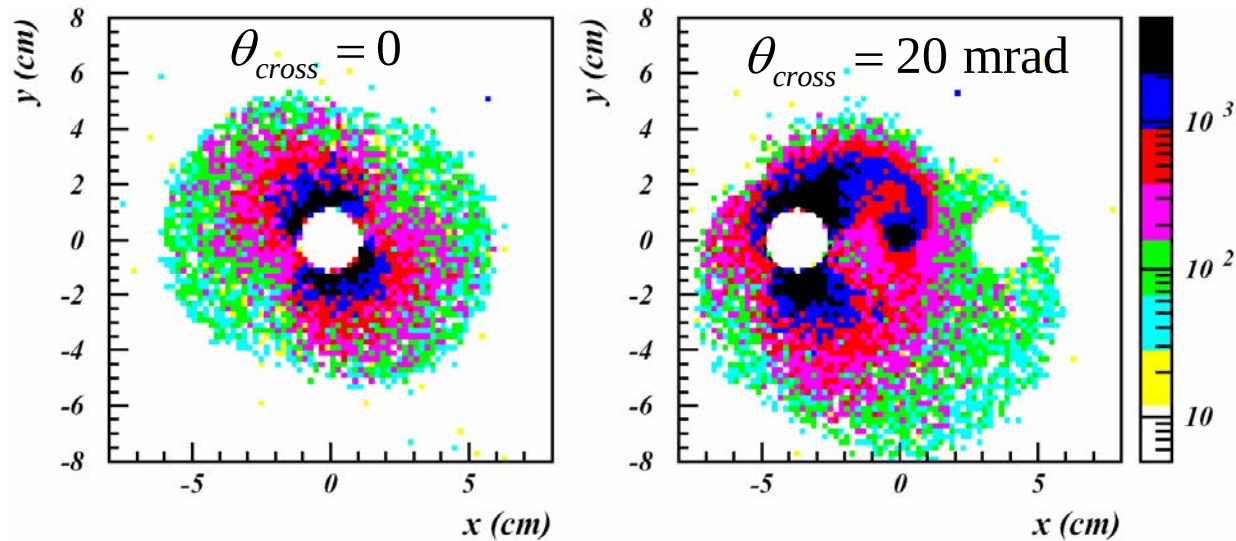


Far Forward Detector

- Electron veto down to 3.2 mrad in presence of very large e^+e^- pair background
- Useful in general to suppress $\gamma\gamma \rightarrow ff$ background. Takes on added importance given that the SUSY parameter space consistent with Dark Matter density includes region with nearly degenerate $\tilde{\chi}_1^0$, $\tilde{\tau}$
- Crossing angle implications.



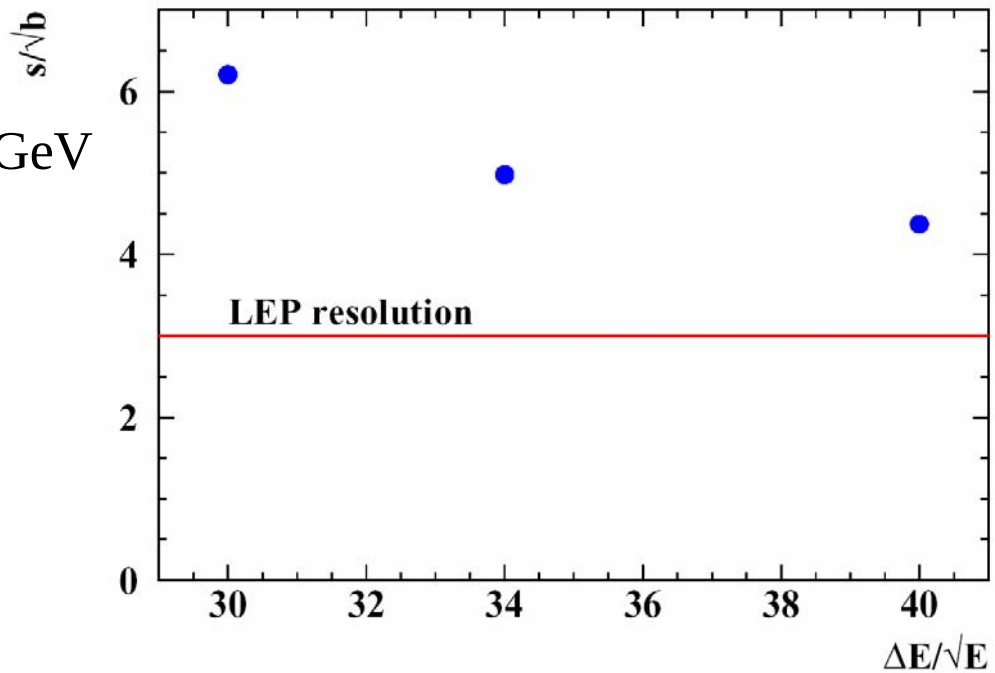
Far Forward Detector



Rel. stau mass error increases from 0.14% to 0.22% with 20 mrad cross angle

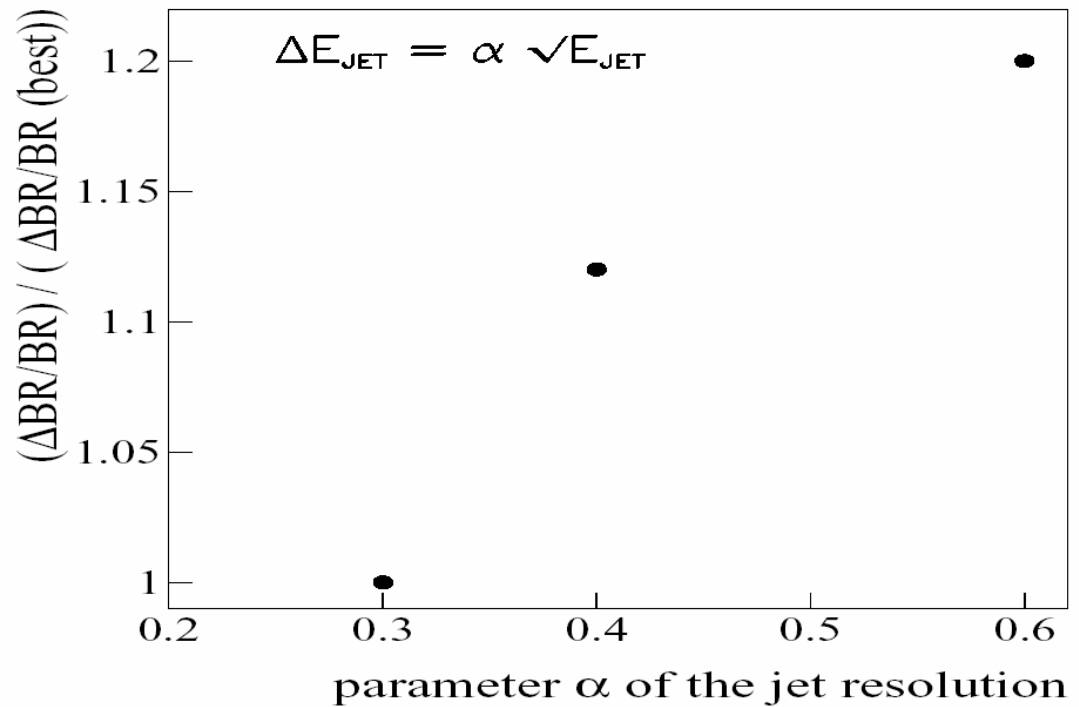
Signal significance at $\sqrt{s} = 500$ GeV
for $e^+e^- \rightarrow ZHH \rightarrow q\bar{q}b\bar{b}b\bar{b}$

C. Castanier et al. hep-ex/0101028



Error on $BR(H \rightarrow WW^*)$ from
measurement of
 $e^+e^- \rightarrow ZH \rightarrow q\bar{q}WW^* \rightarrow q\bar{q}q\bar{q}l\nu$
at $\sqrt{s} = 360$ GeV, $L=500 \text{ fb}^{-1}$

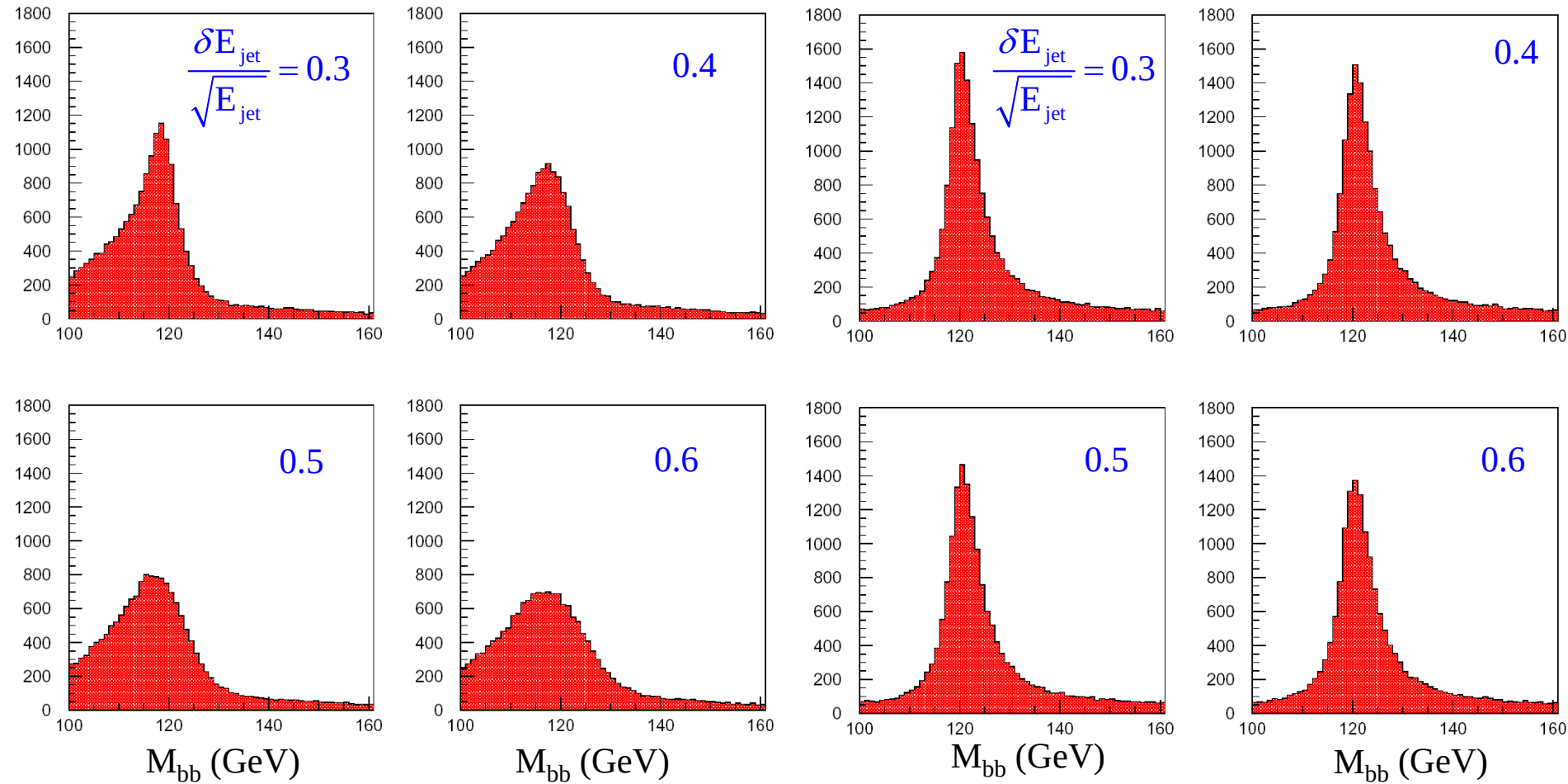
J.-C. Brient, LC-PHSM-2004-001



$$e^+e^- \rightarrow ZH \rightarrow qqbb\bar{b}$$

Reconstructed M_{bb}

4C Fitted M_{bb}

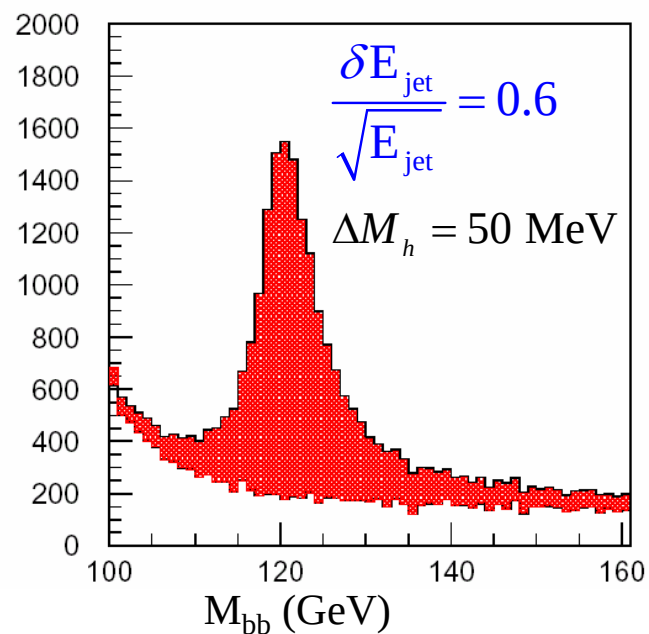
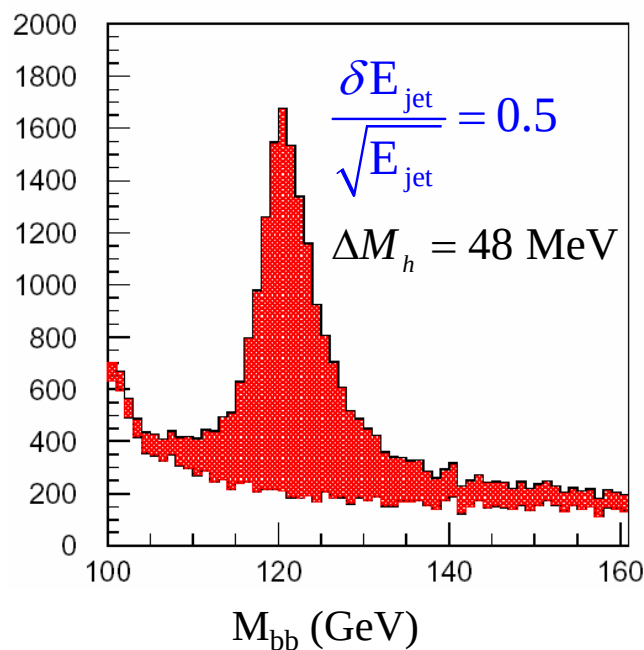
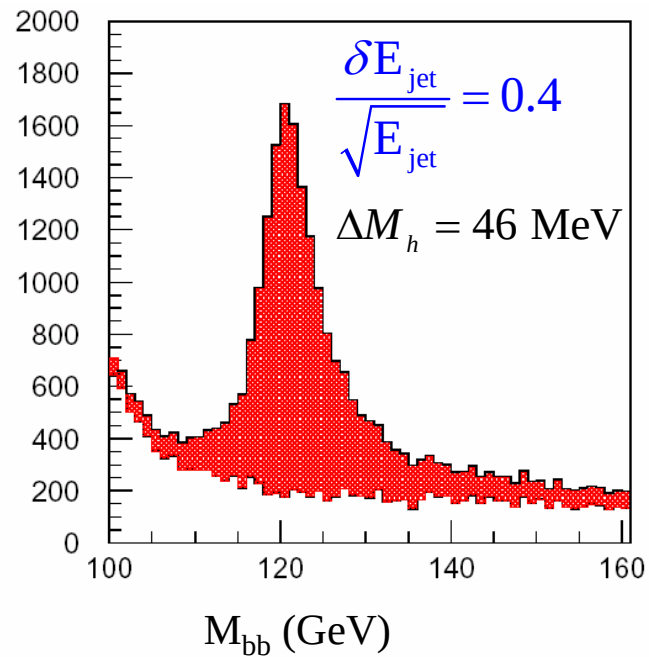
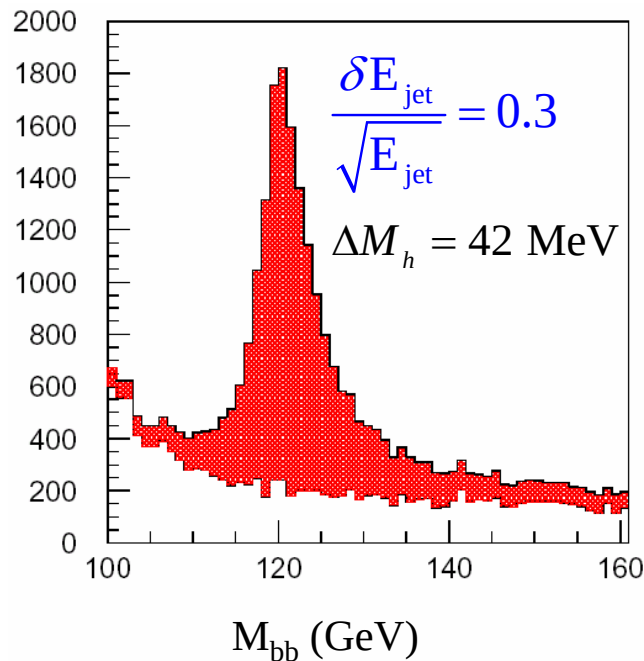


$$e^+e^- \rightarrow ZH$$

$$\rightarrow qqbb\bar{b}$$

$$\sqrt{s} = 350 \text{ GeV}$$

$$L = 500 \text{ fb}^{-1}$$



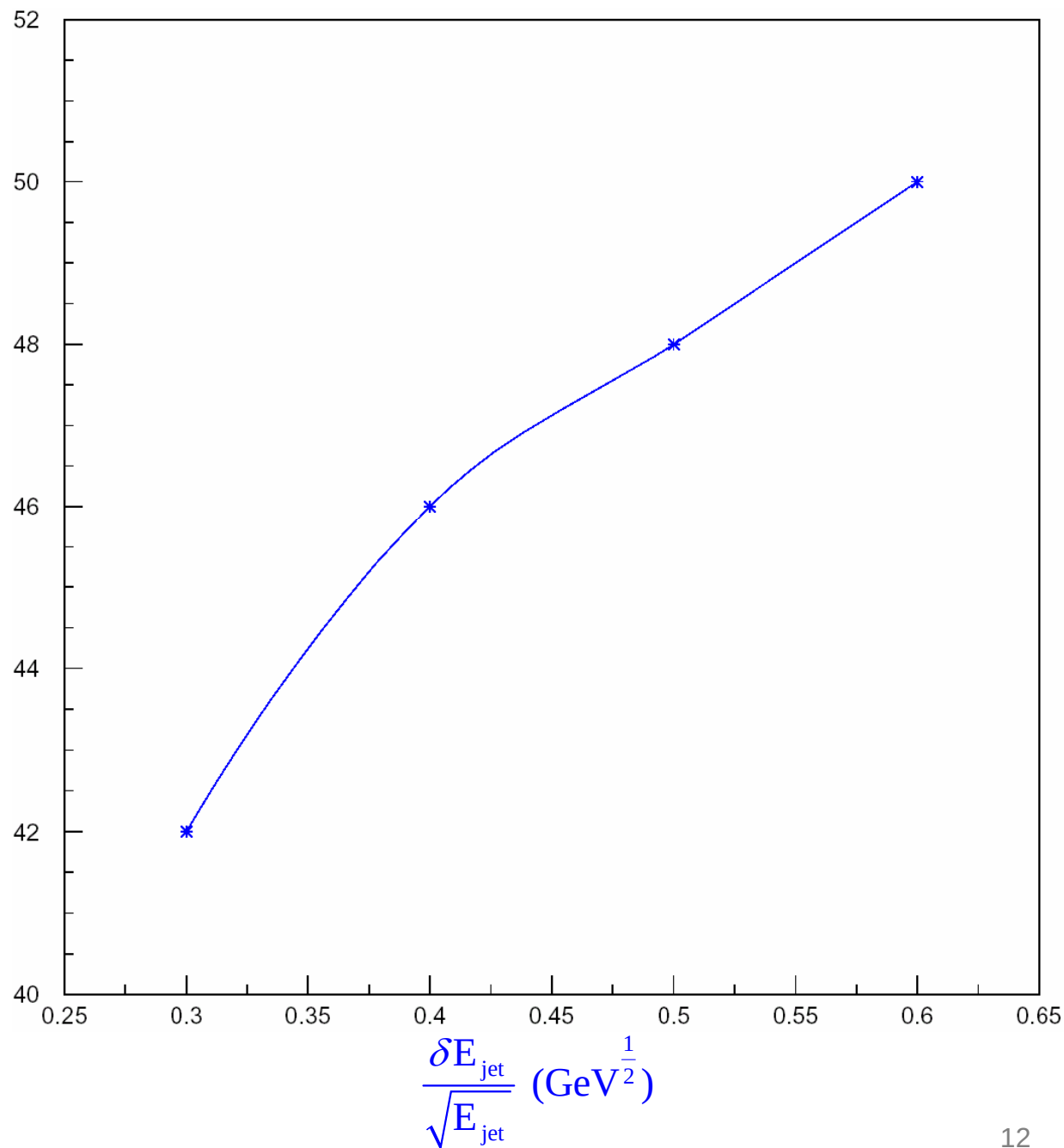
$$e^+e^- \rightarrow ZH$$

$$\rightarrow qq b \bar{b}$$

$$\sqrt{s} = 350 \text{ GeV}$$

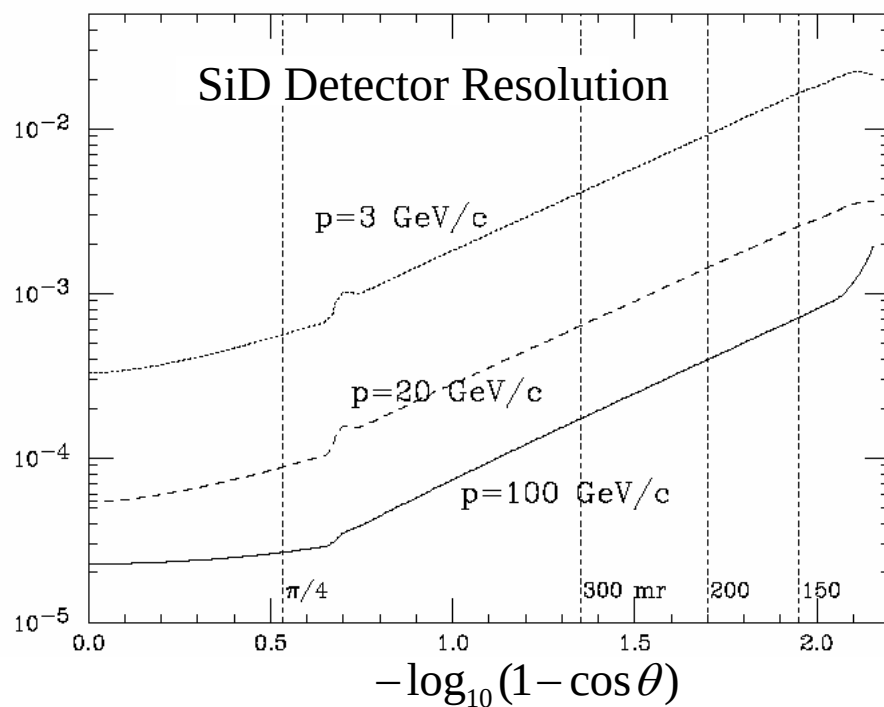
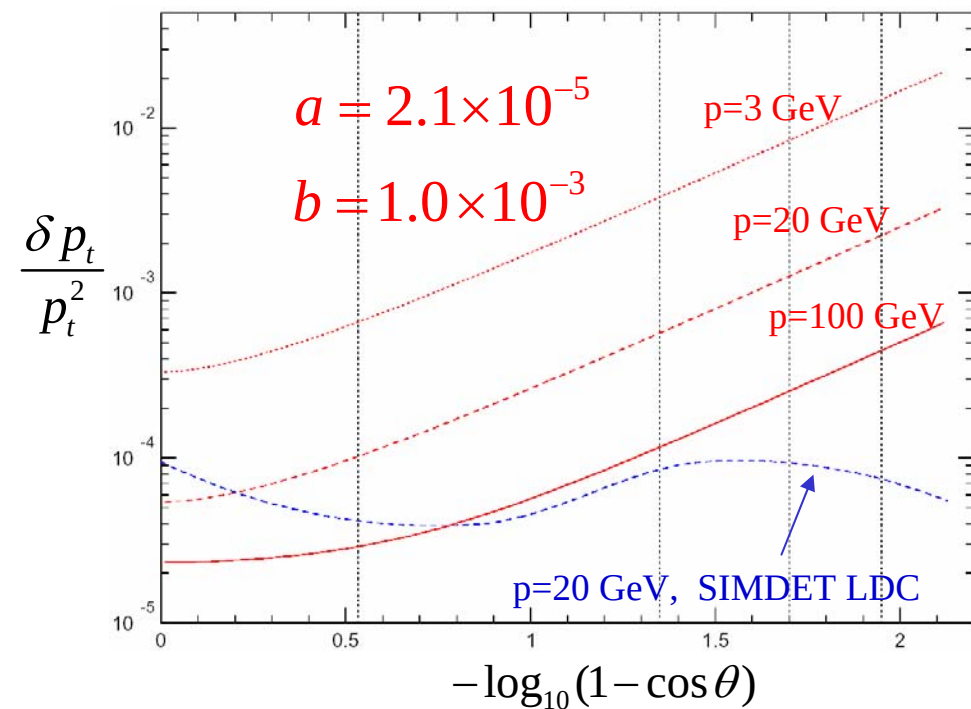
$$L = 500 \text{ fb}^{-1}$$

ΔM_h (MeV)



Simdet Fast MC with this parameterization of pt resolution in place of Simdet's emulation of LDC:

$$\frac{\delta p_t}{p_t^2} = a \oplus \frac{b}{p_t \sin \theta}$$



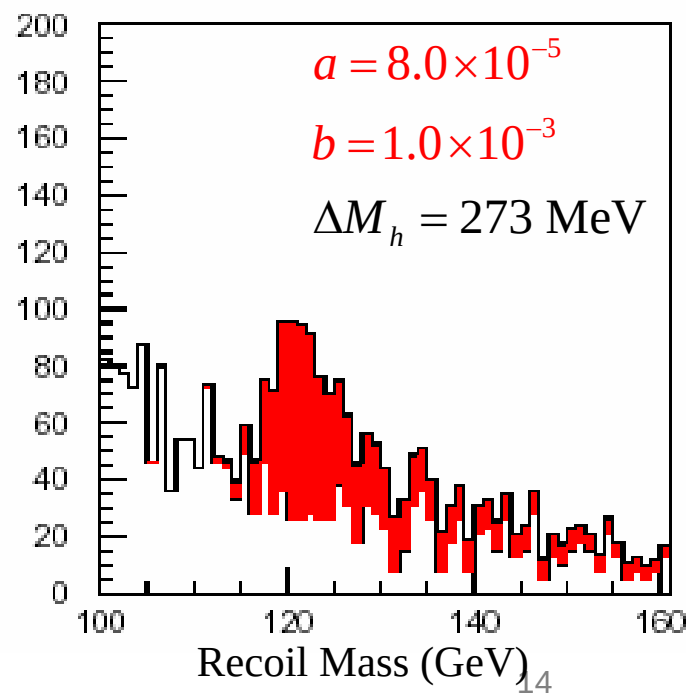
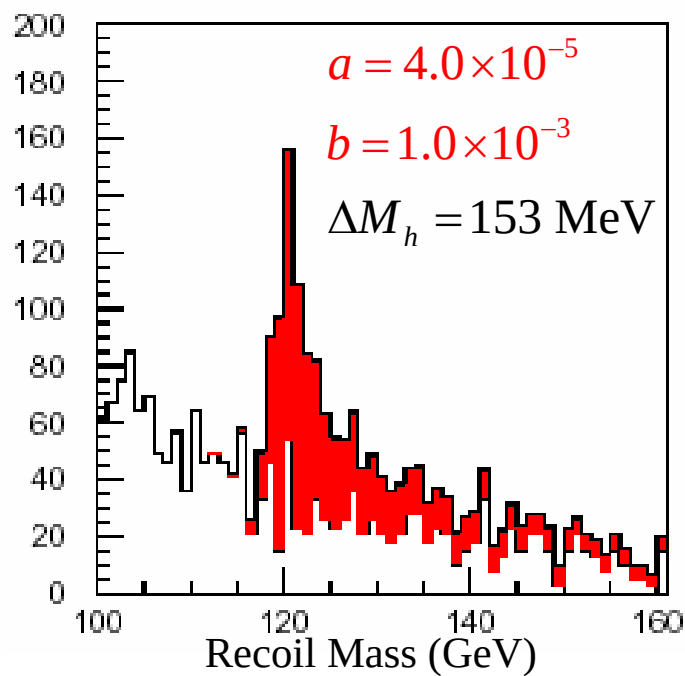
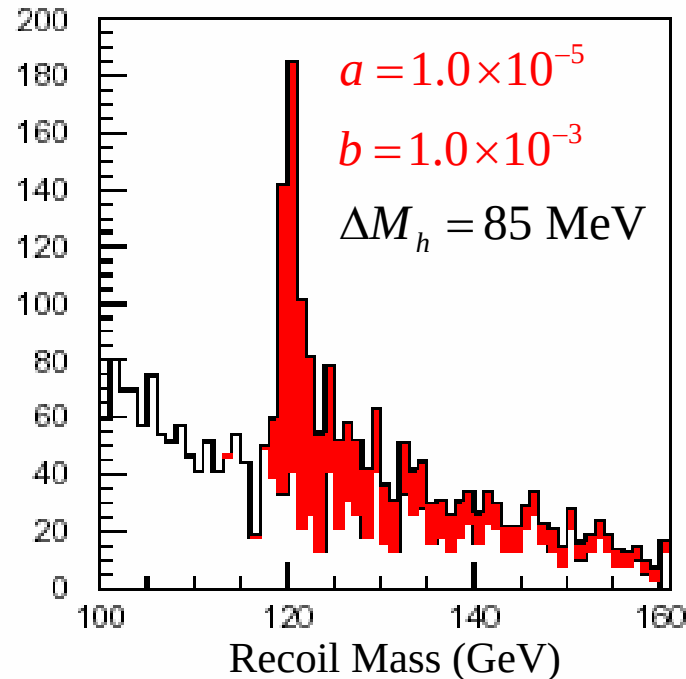
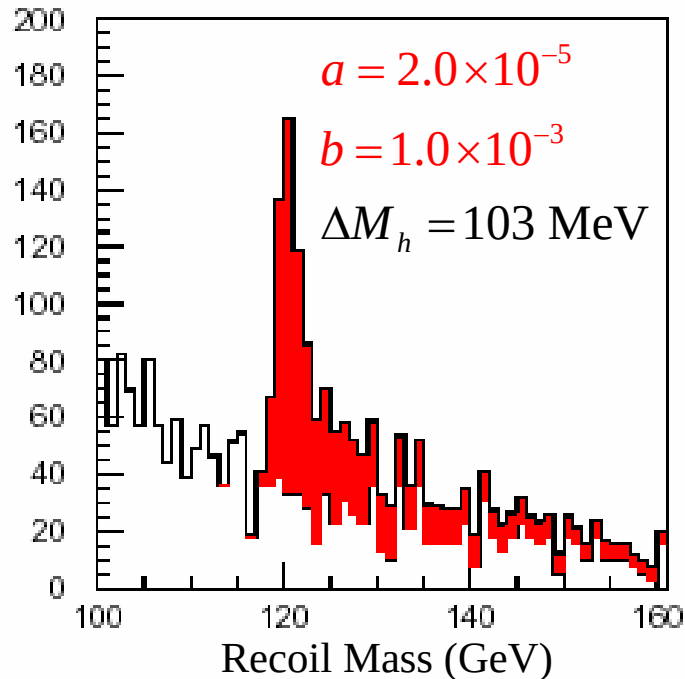
$$e^+e^- \rightarrow ZH$$

$$\rightarrow \mu^+\mu^- X$$

$$\sqrt{s} = 350 \text{ GeV}$$

$$L = 500 \text{ fb}^{-1}$$

$$\frac{\delta p_t}{p_t^2} = a \oplus \frac{b}{p_t \sin \theta}$$



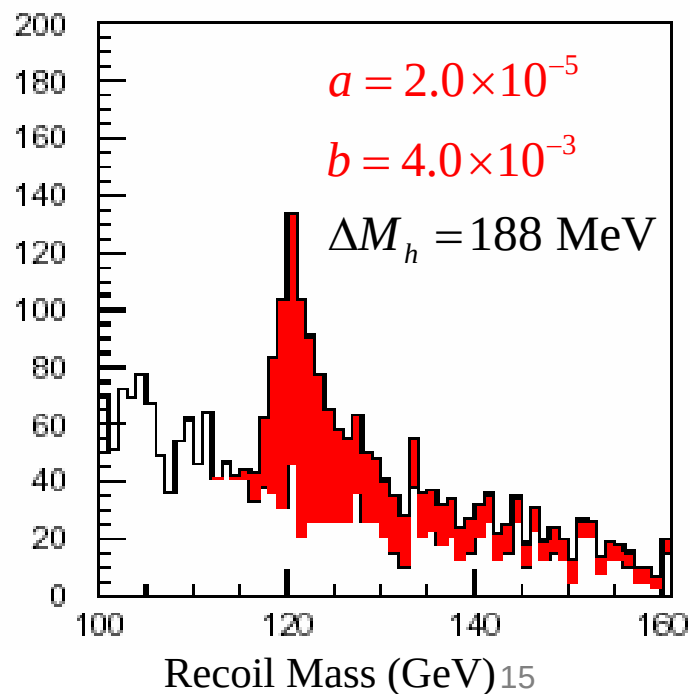
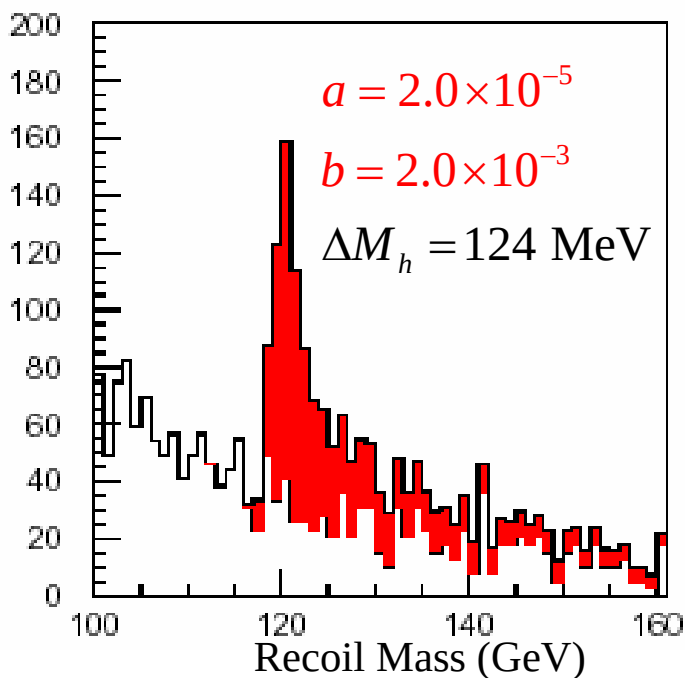
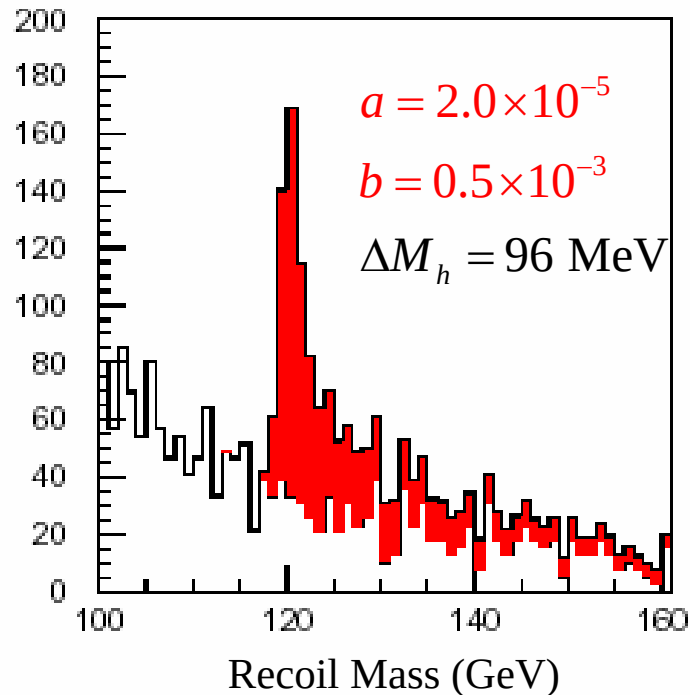
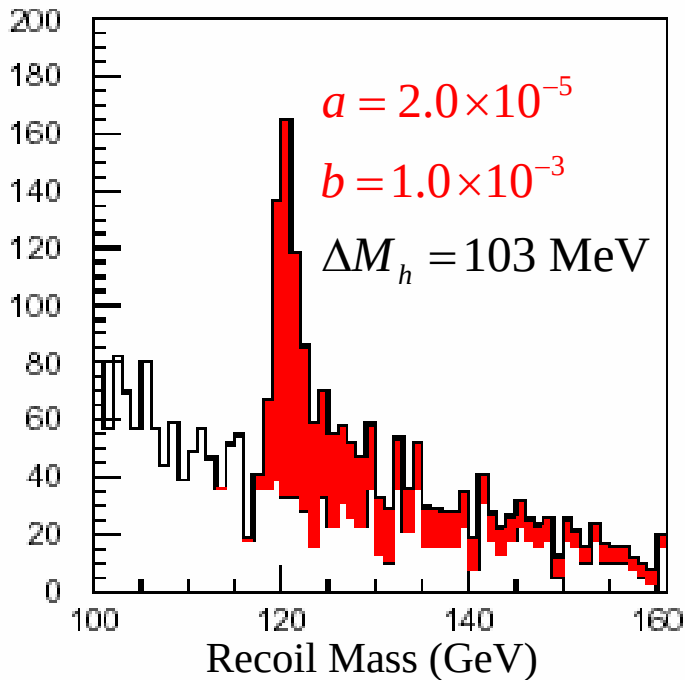
$$e^+e^- \rightarrow ZH$$

$$\rightarrow \mu^+\mu^- X$$

$$\sqrt{s} = 350 \text{ GeV}$$

$$L = 500 \text{ fb}^{-1}$$

$$\frac{\delta p_t}{p_t^2} = a \oplus \frac{b}{p_t \sin \theta}$$



$$e^+e^- \rightarrow ZH$$

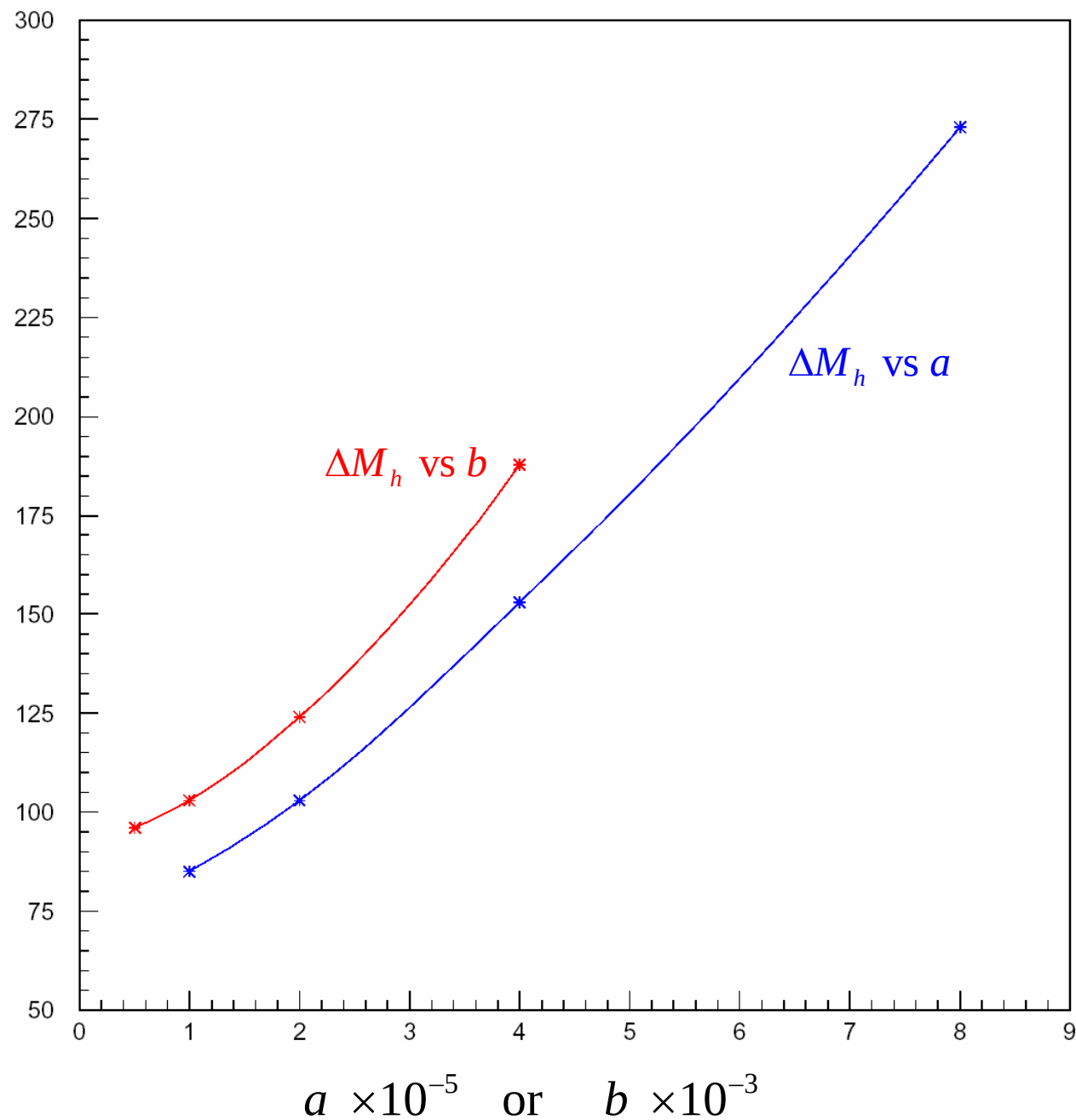
$$\rightarrow \mu^+\mu^- X$$

$$\sqrt{s} = 350 \text{ GeV}$$

$$L = 500 \text{ fb}^{-1}$$

ΔM_h (MeV)

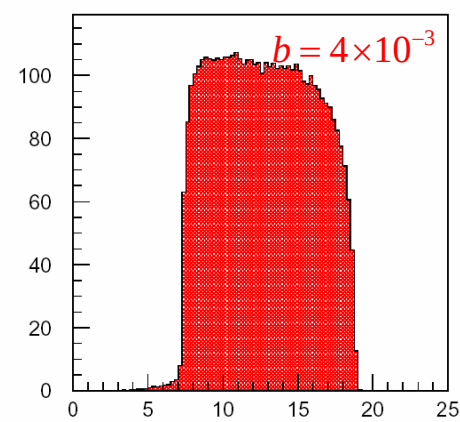
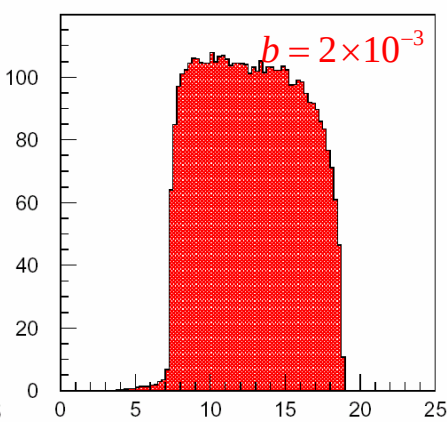
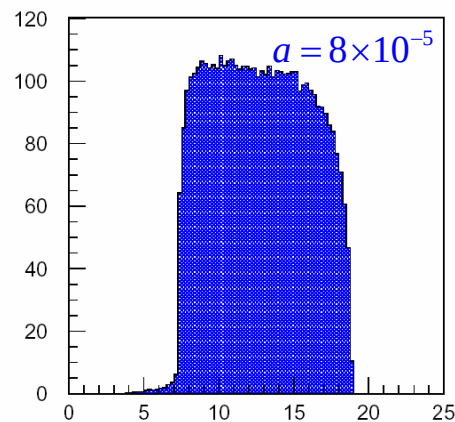
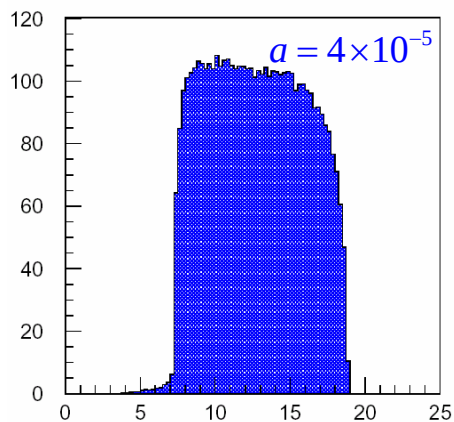
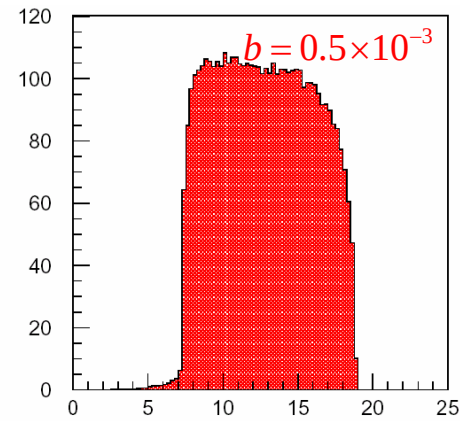
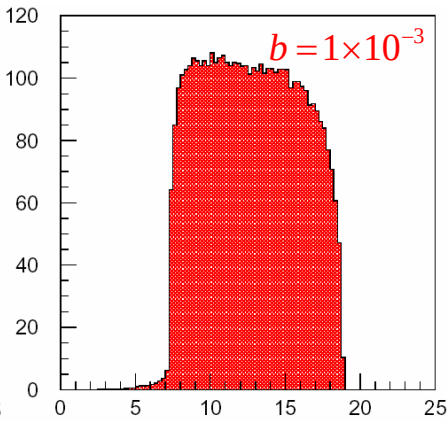
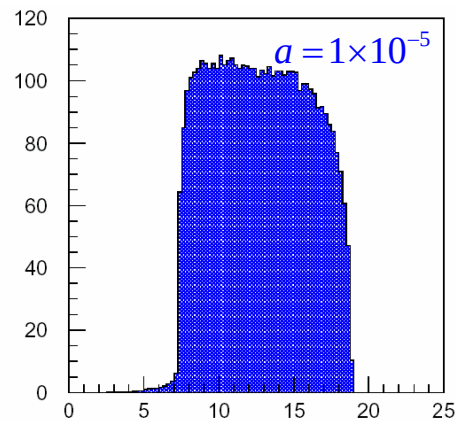
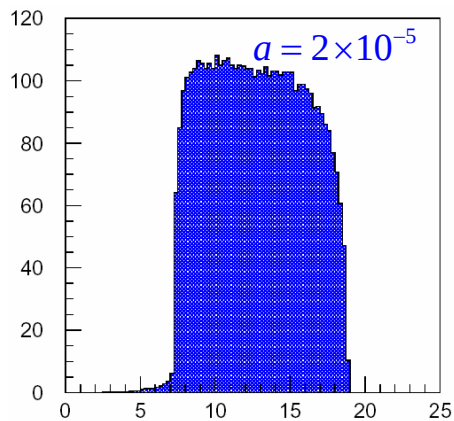
$$\frac{\delta p_t}{p_t^2} = a \oplus \frac{b}{p_t \sin \theta}$$



$$e^+e^- \rightarrow \tilde{\mu}_R^+ \tilde{\mu}_R^- \quad M_{\tilde{\mu}_R} = 224 \text{ GeV} \quad \sqrt{s} = 500 \text{ GeV} \quad L = 500 \text{ fb}^{-1}$$

$$\frac{\delta p_t}{p_t^2} = a \oplus \frac{b}{p_t \sin \theta}$$

Fit for $M_{\tilde{\mu}}$ only



Muon Energy (GeV)

Muon Energy (GeV)

Muon Energy (GeV)

Muon Energy (GeV)

No dependence on a or b in the range

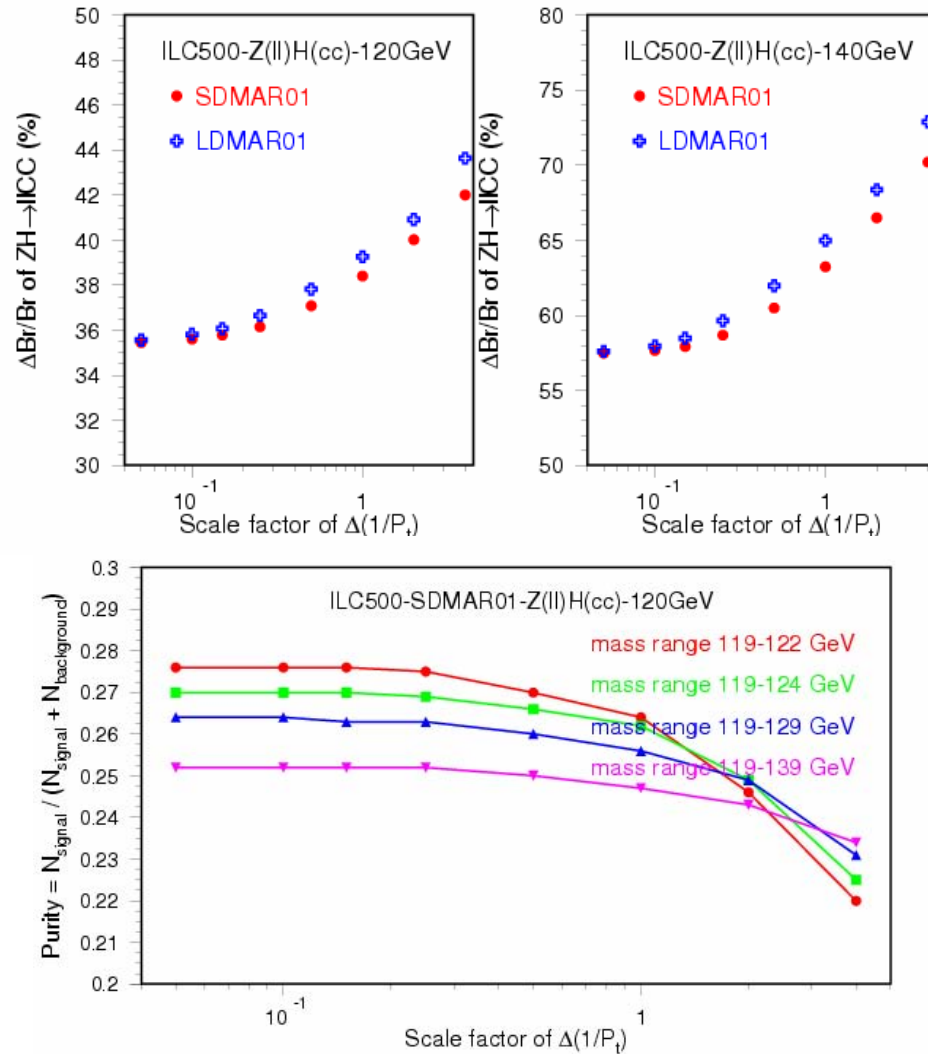
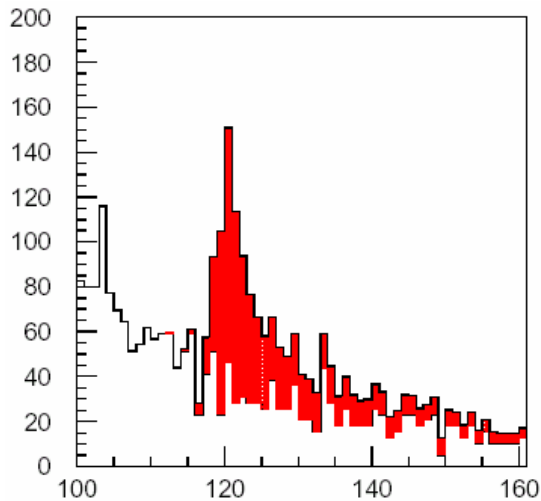
$$\Delta M_{\tilde{\mu}} = 34 \text{ MeV}$$

$$1.0 \times 10^{-5} < a < 8.0 \times 10^{-5}, \quad 0.5 \times 10^{-3} < b < 4.0 \times 10^{-3}$$

Branching Ratio of $H \rightarrow CC$

$$e^+e^- \rightarrow ZH$$

$$\rightarrow l^+l^- X$$

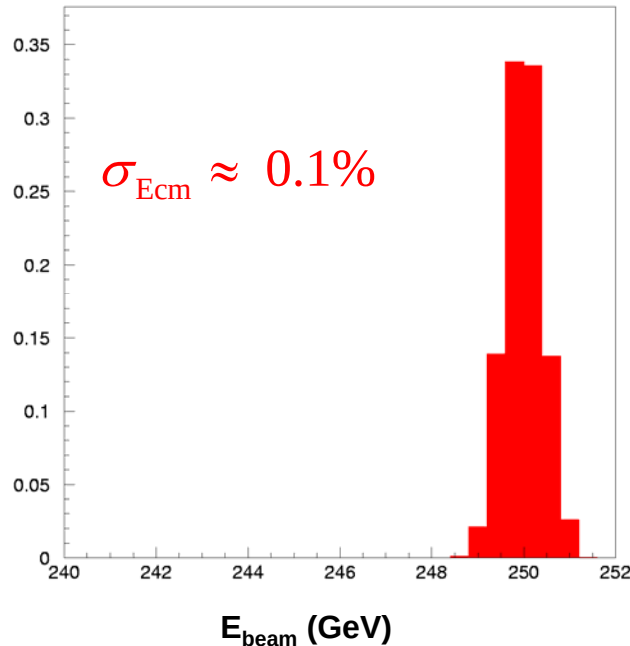


* Talk by Haijun Yang
in TRK session 20Mar2005

→ $\Delta Br/Br \sim 39\%$ (120GeV), 64% (140GeV) for $Z \rightarrow l+l^-$, 1000 fb^{-1}

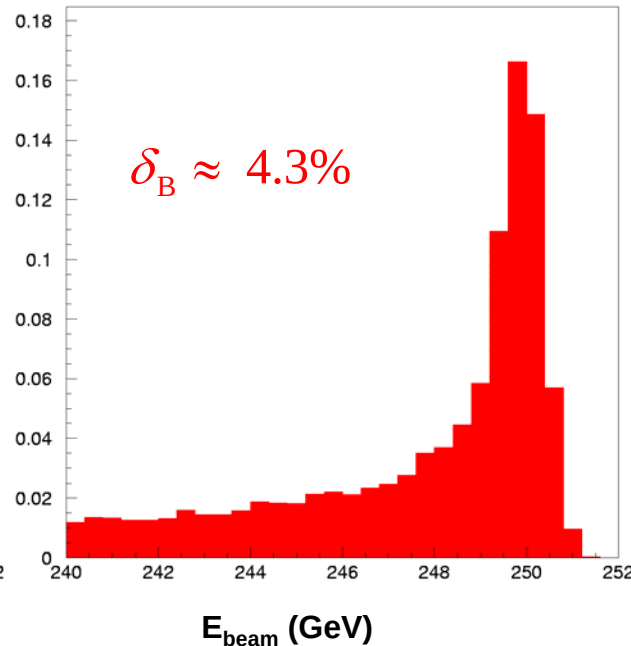
Beam Energy Profiles $\langle E_{\text{beam (incoming)}} \rangle = 250 \text{ GeV}$

Before Collision

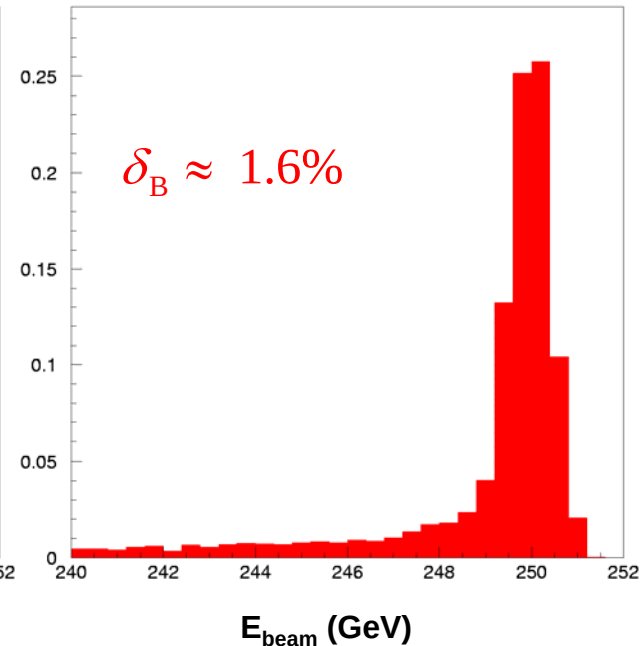


After Collision

$50 \text{ ppm} < E_{\text{CM}}^{\text{bias}} < 250 \text{ ppm}$



Lumi Weighted



Center of Mass Energy Error Requirements

- Top mass: 200 ppm (35 MeV)
- Higgs mass: 200 ppm (60 MeV for 120 GeV Higgs)
- Giga-Z program: 50 ppm

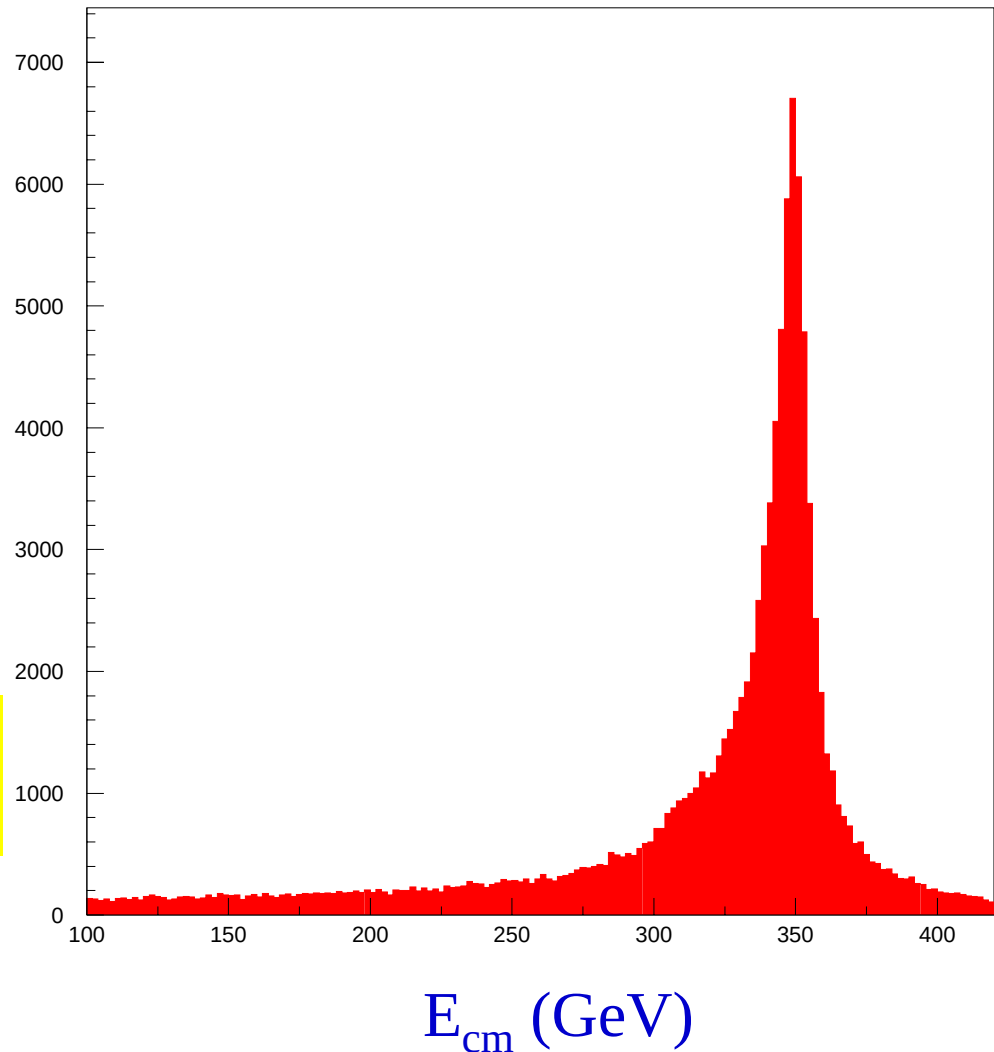
Reconstructed E_{cm} using $Z\gamma$ events and
measured angles. $Z \rightarrow \mu^+ \mu^-$

$$\sqrt{s} = 350 \text{ GeV}$$

$$L = 100 \text{ fb}^{-1}$$

$$\Delta E_{\text{cm}} = 47 \text{ MeV}$$

* Talk by Klaus Moenig
in MDI session 20Mar2005



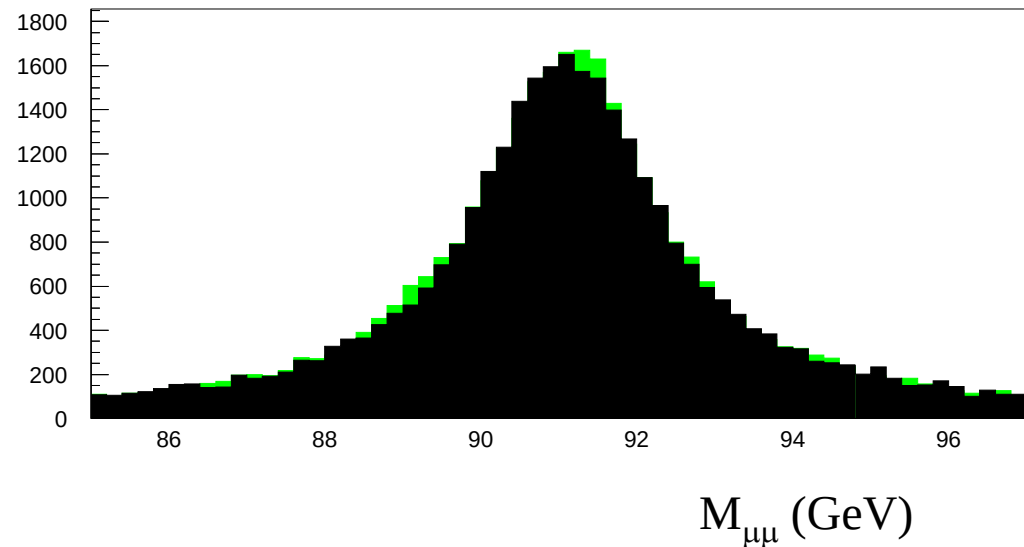
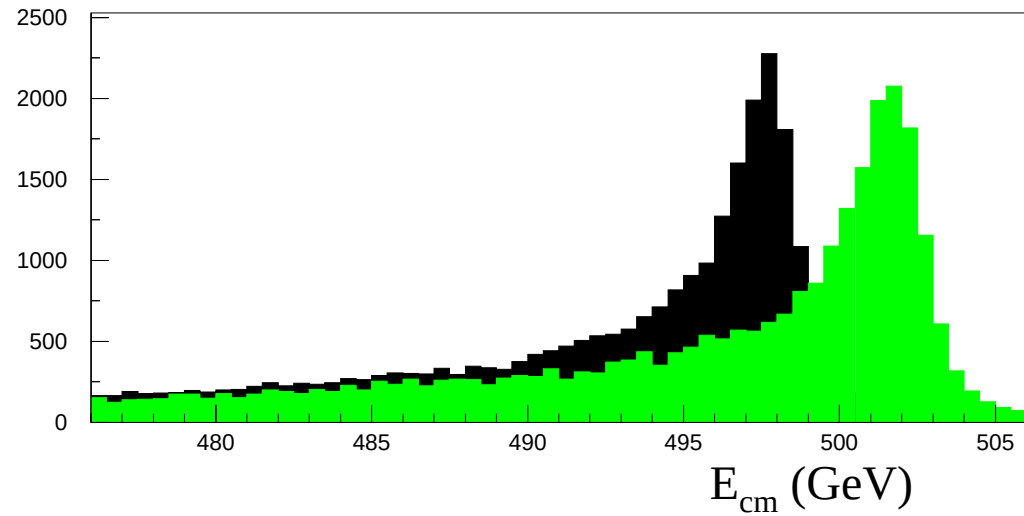
The momentum resolution is set by Higgs recoil mass measurement in $e^+e^- \rightarrow ZH \rightarrow \mu^+\mu^-H$

In the reaction $e^+e^- \rightarrow Z\gamma \rightarrow \mu^+\mu^-\gamma$ we know the mass of the photon. Why not invert the problem and use the excellent momentum resolution to solve for $\sqrt{s}$ instead of the mass of the system opposite $\mu^+\mu^-$?

Reconstructed E_{cm} & M_Z using $Z\gamma$ events and
measured momenta & angles. $Z \rightarrow \mu^+ \mu^-$

$$\Delta E_{\text{cm}} = -2 \text{ GeV}$$

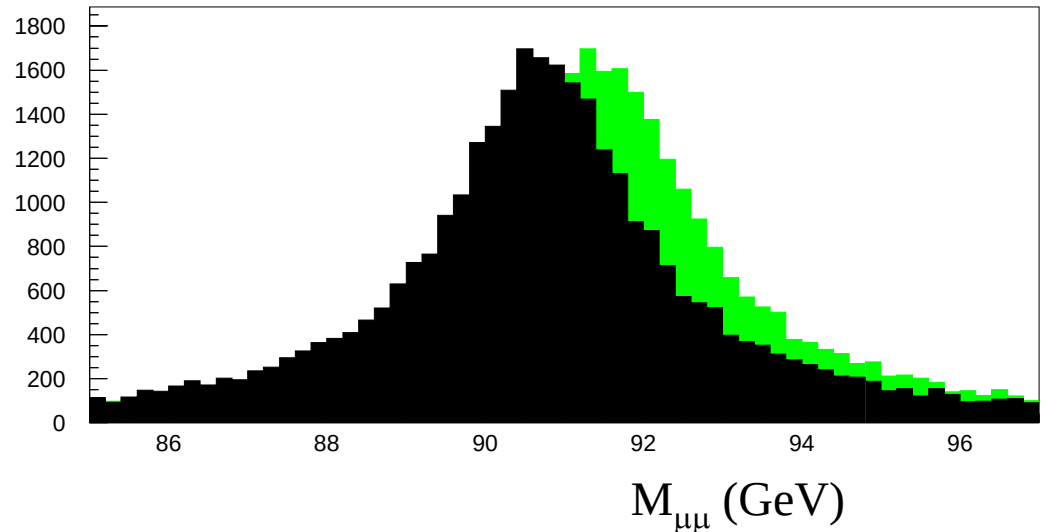
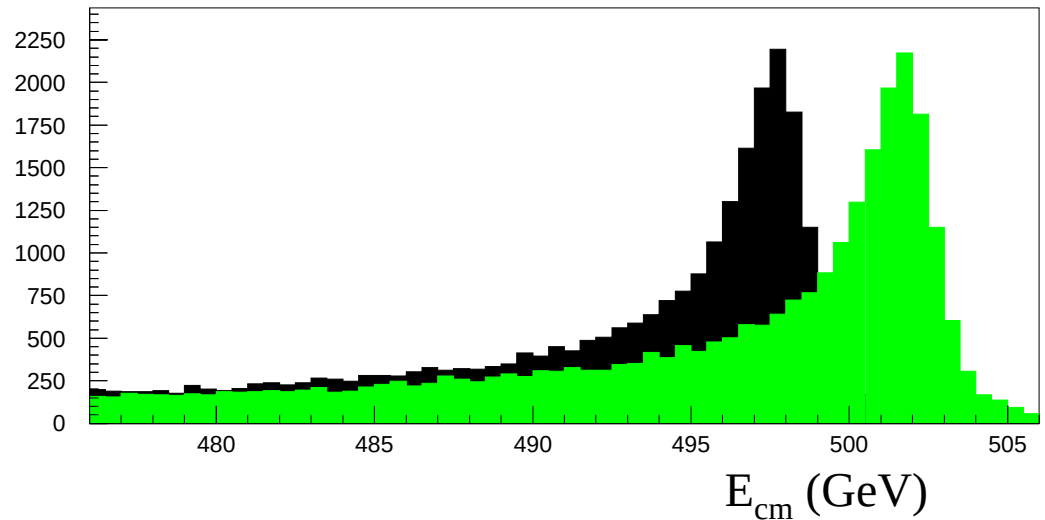
$$\Delta E_{\text{cm}} = +2 \text{ GeV}$$



Reconstructed E_{cm} & M_Z using $Z\gamma$ events and
measured momenta & angles. $Z \rightarrow \mu^+ \mu^-$

Trk mom scale factor = 0.996

Trk mom scale factor = 1.004



$$E_{\text{cm}} = 350 \text{ GeV}$$

$$\text{Lumi} = 100 \text{ fb}^{-1}$$

$E_{Z\gamma}$ = Measured E_{cm} assuming Z boson recoil against single photon

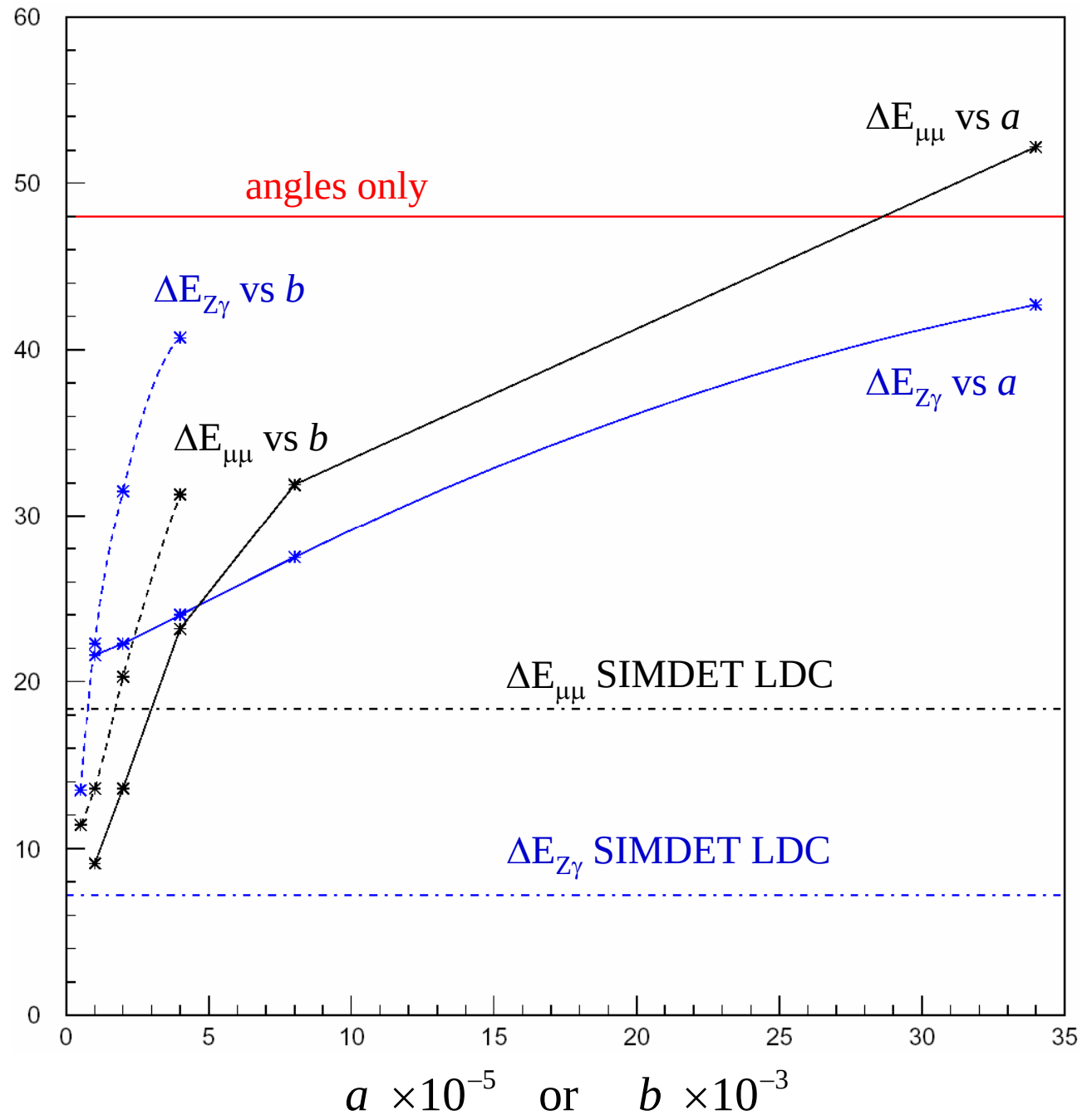
$E_{\mu\mu}$ = Measured E_{cm} using full energy $e^+e^- \rightarrow \mu^+\mu^-$

measured var	a	b	ΔE_{cm} (GeV)	ΔE_{cm} (GeV)	ΔE_{cm} (GeV)	$\frac{\Delta E_{\text{cm}}}{E_{\text{cm}}} \text{ (ppm)}$
			stat	sys(E scale)	total	total
$E_{Z\gamma}$ ang only			.0473	0	.0473	135
$E_{Z\gamma}$ $ \vec{p} $ & ang	2×10^{-5}	1×10^{-3}	.0085	.0206	.0223	64
$E_{Z\gamma}$ $ \vec{p} $ & ang	2×10^{-5}	$.5 \times 10^{-3}$	.0054	.0124	.0135	39
$E_{Z\gamma}$ $ \vec{p} $ & ang	34×10^{-5}	4×10^{-3}	.0375	.0313	.0488	139
$E_{\mu\mu}$ $ \vec{p} $ & ang	2×10^{-5}	1×10^{-3}	.0056	.0124	.0136	39

Ecm Resolution in MeV vs a or b

$$e^+e^- \rightarrow \mu^+\mu^-$$

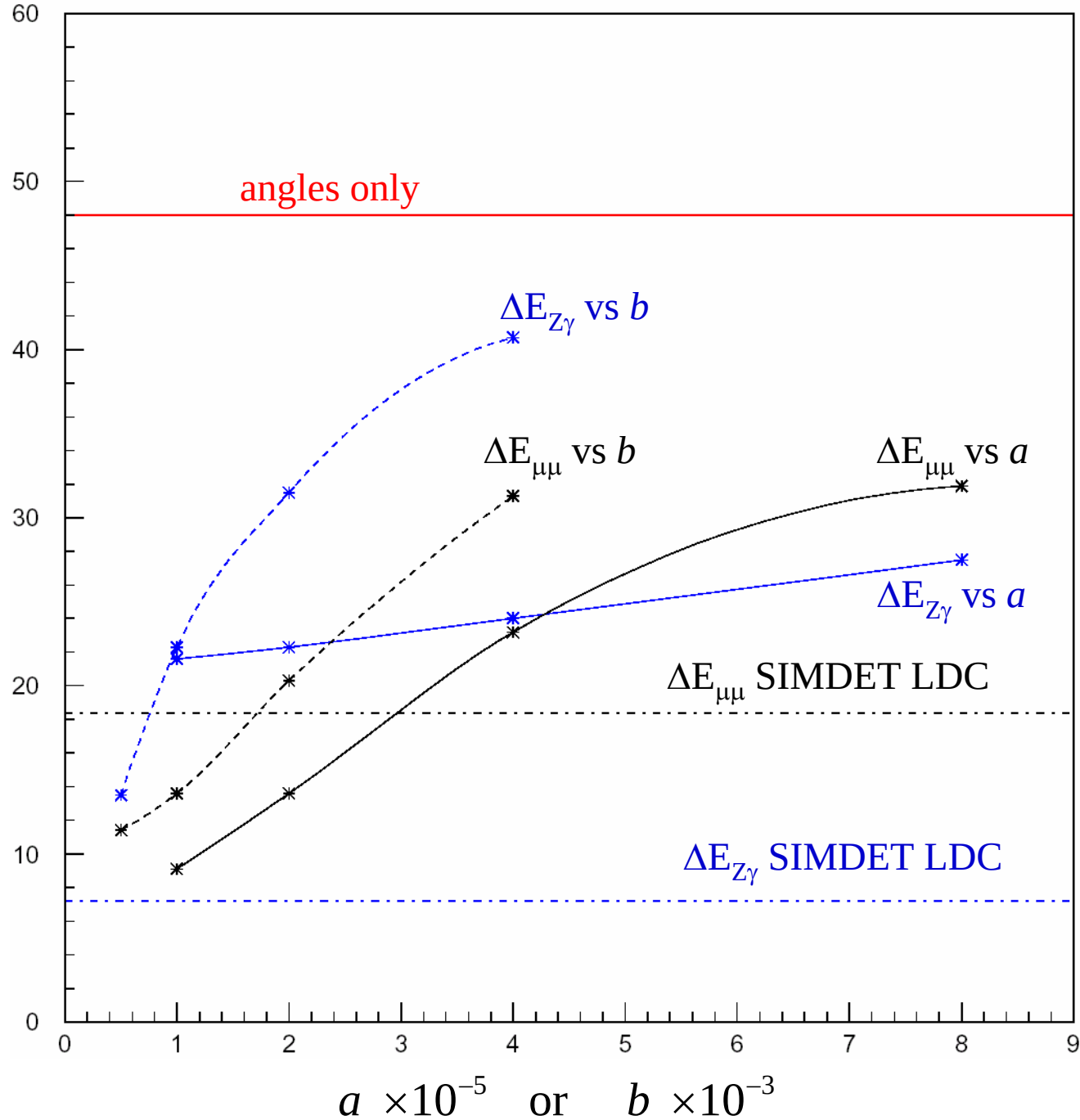
$$\frac{\delta p_t}{p_t^2} = a \oplus \frac{b}{p_t \sin \theta}$$



Ecm Resolution in MeV vs a or b

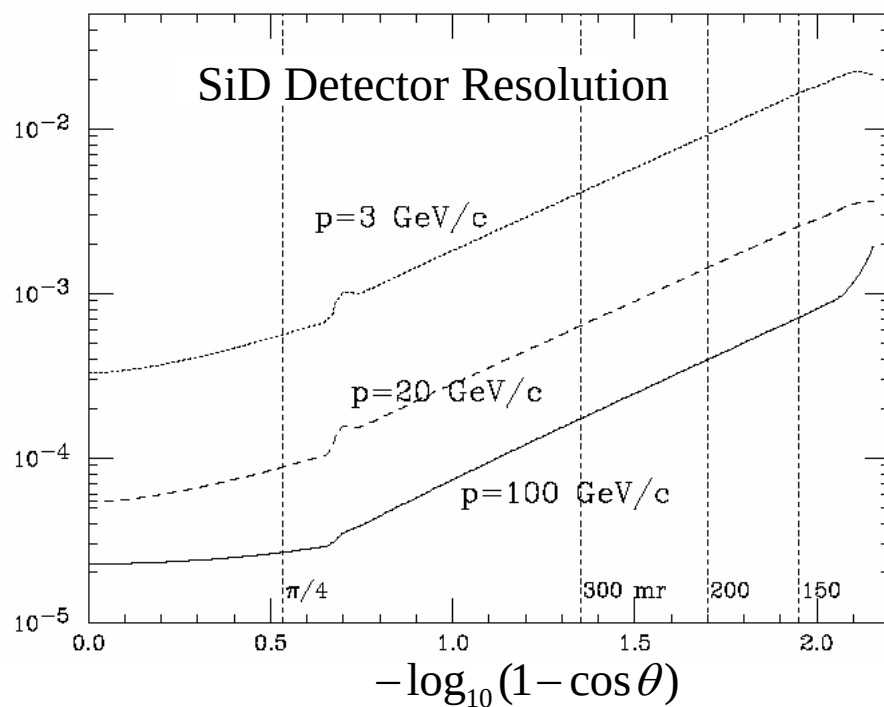
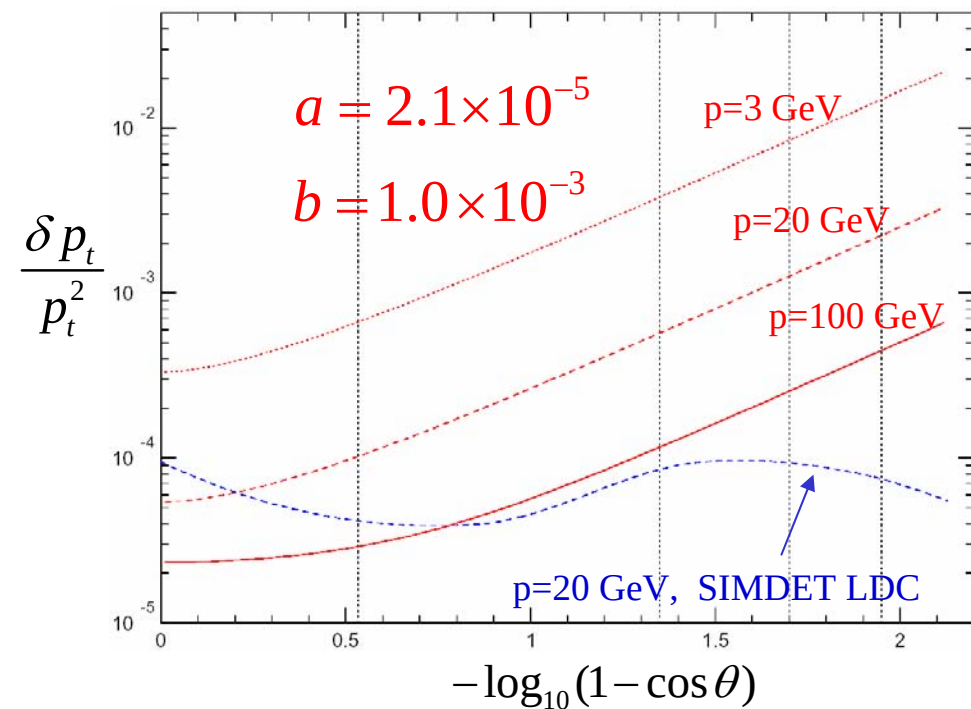
$$e^+e^- \rightarrow \mu^+\mu^-$$

$$\frac{\delta p_t}{p_t^2} = a \oplus \frac{b}{p_t \sin \theta}$$



Simdet Fast MC with this parameterization of pt resolution in place of Simdet's emulation of LDC:

$$\frac{\delta p_t}{p_t^2} = a \oplus \frac{b}{p_t \sin \theta}$$



Summary

- Current detector designs appear well matched to envisioned physics program
- Choices will have to be made as realities of detector engineering and cost are confronted – physics benchmark studies will play a crucial role. The examples of parametric studies shown here are just the beginning.