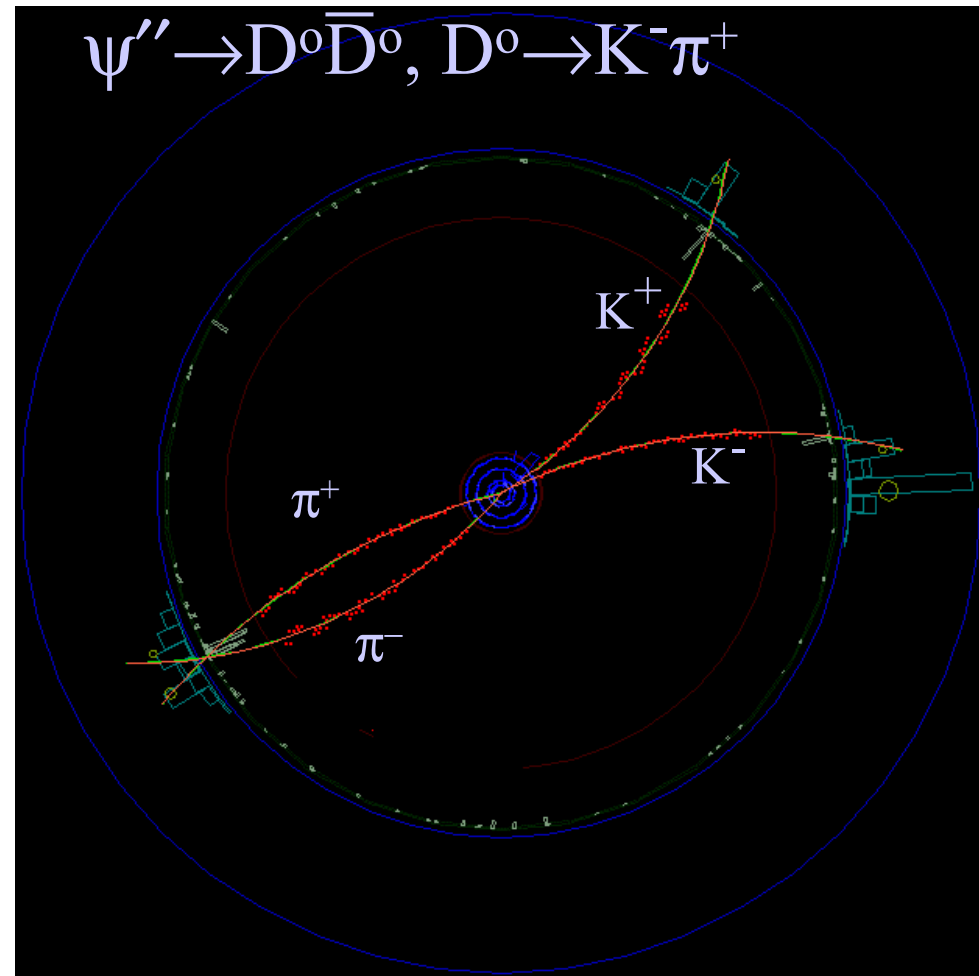


CKM Studies and New Physics Searches with Charm

Two themes:

- 1) Why Charm Physics allows B Physics to reach its full potential
- 2) Charm physics as a probe of physics beyond the Standard Model

Ian Shipsey,
Purdue University





- I am completely deaf
- I communicate by lip reading
- BUT lip reading obeys an inverse square law, and the audience is too far away
- Please write down your questions
- Pass them up to me
- I will read out your question before answering it

Outline of the Lectures

Overview: How Charm Physics Helps B Physics
→ Precision Quark Flavor Physics

Experiments That Contribute To Charm Physics

Precision CKM Physics:

Lifetimes

Hadronic Decays

Leptonic Decays and Decay constants

Semileptonic Decays and CKM matrix elements

Tests of Unitarity

Spectroscopy

Charm as a Probe of New Physics:

Mixing

CP Violation & Rare Decays

Summary & Outlook

Charm Physics: the context

*This
Decade*

Flavor Physics: is in "the $\sin^2\beta$ era" akin to precision Z. Over constrain CKM matrix with precision measurements. Limiting factor: non-pert. QCD.

*The
Future*

LHC may uncover strongly coupled sectors in the **physics that lies beyond the Standard Model**. The LC will study them. Strongly-coupled field theories are an outstanding challenge to theoretical physics. Critical need for reliable theoretical techniques & detailed data to calibrate them.

*Example:
The
Lattice*

Complete definition of pert & non. Pert. QCD. Matured over last decade, can calculate to 1-5% B, D, Y, Ψ ...

Charm can provide the data to calibrate QCD techniques

(See Peter Lepage's lectures for details of Lattice QCD)

Charm Physics: What do we need to measure?

- **flavor physics:** overcome the non pert. QCD roadblock

Precision charm lifetimes ← exist

do not exist

- precision charm abs. branching ratio measurements

Leptonic decays
 : decay constants

Semileptonic decays:
 V_{cs}, V_{cd} , unitarity
 form factors

**Abs D hadronic
 Br's normalize
 B physics**

Tests QCD techniques in
 c sector, apply to b sector

→ Improved $V_{ub}, V_{cb}, V_{td} & V_{ts}$

- **strong coupling in Physics beyond the Standard Model**

- precise measurements of quarkonia spectroscopy & decay provide essential data to calibrate theory.

← Important
 Input for the lattice

Physics beyond the Standard Model:

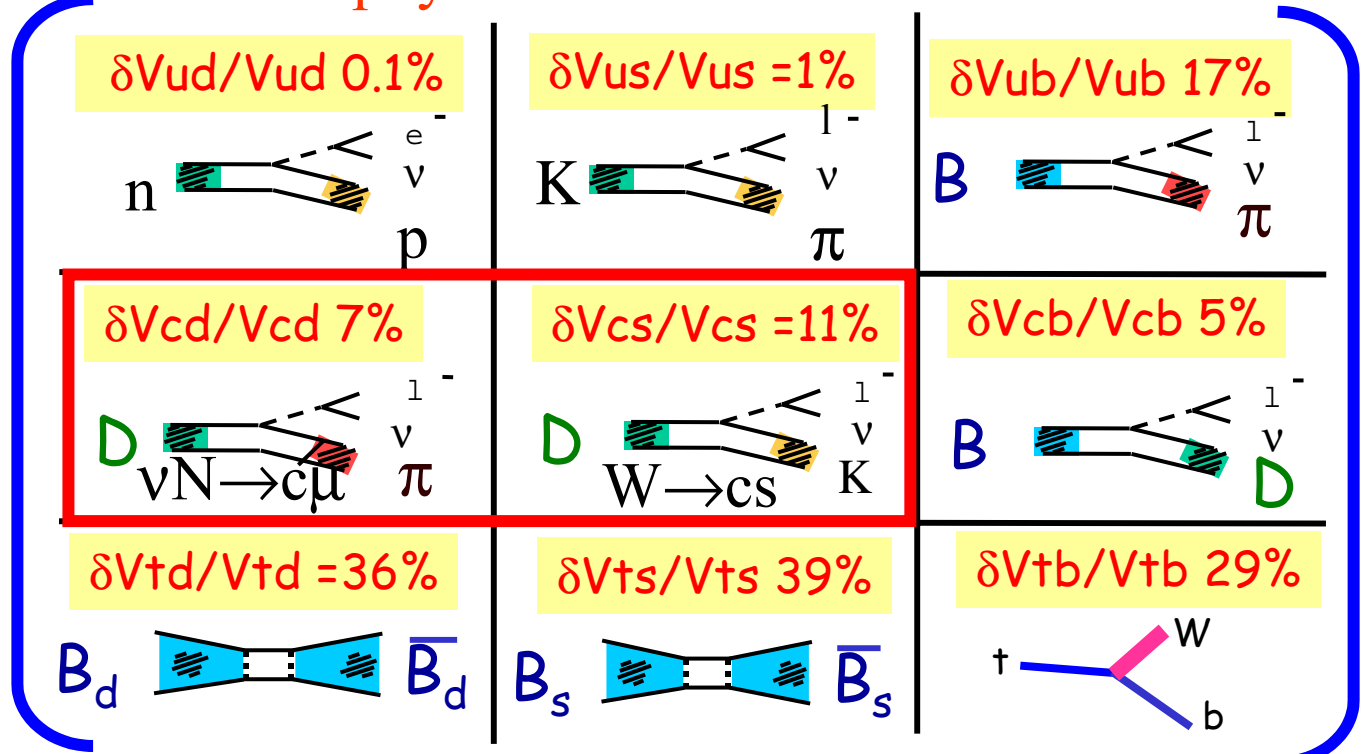
- D-mixing, CPV, rare decays. + measure strong phases

Charm physics builds the tools to enable this decade's
 flavor physics and the next decade's new physics.

Precision Quark Flavor Physics

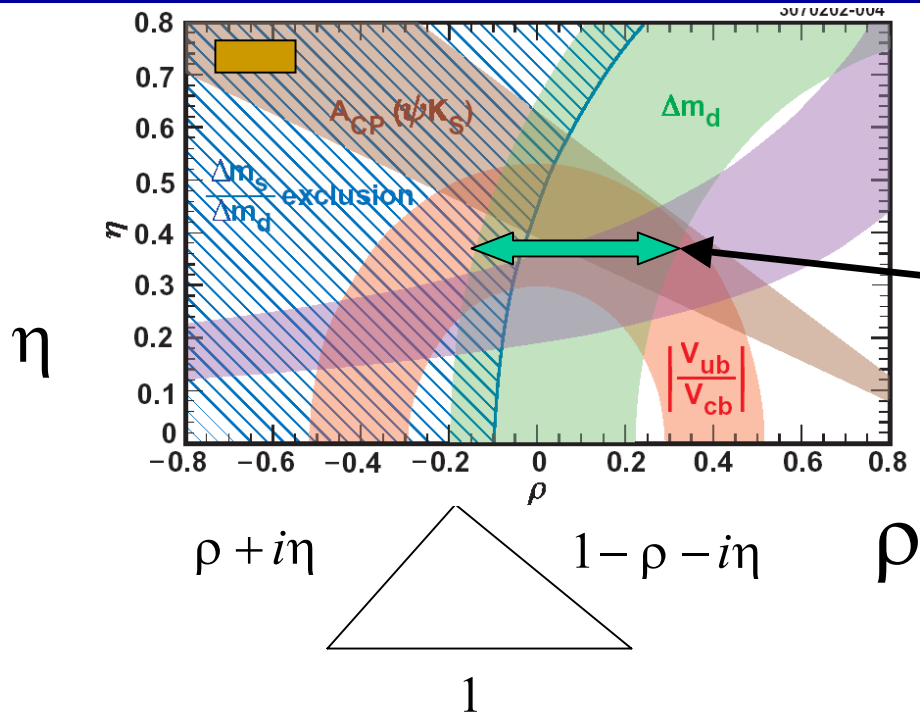
Goal for the decade: high precision measurements of V_{ub} , V_{cb} , V_{ts} , V_{td} , V_{cs} , V_{cd} , & associated phases. Over-constrain the “Unitarity Triangles”
 - Inconsistencies → New physics !

CKM
 Matrix
 Current
 Status:



Many experiments will contribute. Measurement of absolute charm branching ratios will enable precise new measurements at Bfactories/Tevatron to be translated into greatly improved CKM precision.

Importance of measuring absolute charm leptonic branching ratios: f_D & f_{D_s} : V_{td} & V_{ts}



$$\Delta M_d = 0.50 ps^{-1} \left[\frac{\sqrt{B_{B_d}} f_{B_d}}{200 MeV} \right]^2 \left[\frac{|V_{td}|}{8.8 \times 10^{-3}} \right]^2$$

$$\frac{\sigma(\rho)}{\rho} = 0.5 \frac{\sigma(\Delta M_d)}{\Delta M_d} \oplus \frac{\sigma(f_B \sqrt{B_{B_d}})}{f_B \sqrt{B_{B_d}}}$$

(ICHEP02) 1.2%

~15% (LQCD)

$$\frac{\Delta M_d}{\Delta M_s} \propto \left[\frac{\sqrt{B_{B_d}} f_{B_d}}{\sqrt{B_{B_s}} f_{B_s}} \right]^2 \left[\frac{|V_{td}|}{|V_{ts}|} \right]^2$$

~5-7%

$$\frac{\delta f_{D_c}}{f_{D_c}} \sim 14\%$$

$$\frac{\delta f_{D_c}}{f_{D_c}} \sim 100\%$$

Lattice predicts f_B/f_D & f_{B_s}/f_{D_s} with small errors

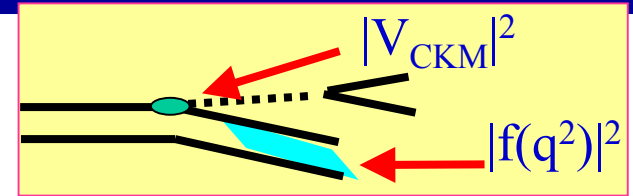
if precision measurements of f_D & f_{D_s} existed (they do not)

We could obtain precision estimates of f_B & f_{B_s} and hence precision determinations of V_{td} and V_{ts}

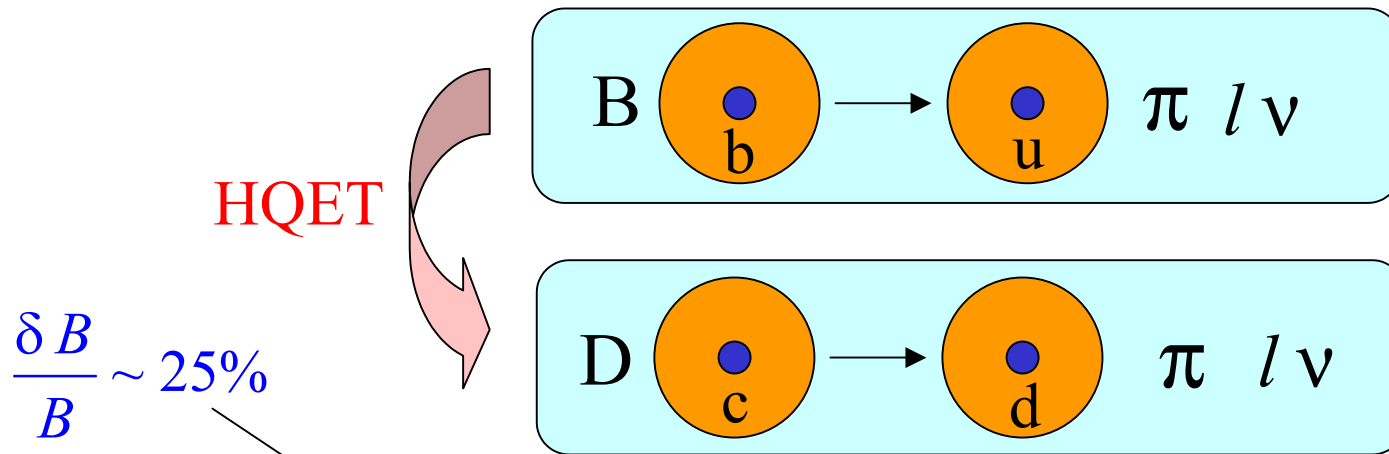
Similarly f_D/f_{D_s} checks f_B/f_{B_s}

Importance of absolute charm semileptonic decay rates.

$$\frac{d\Gamma}{dq^2} = \frac{G_F^2}{24\pi^3} |V_{cs}|^2 p_K^3 |f_+(q^2)|^2$$



- I. Absolute magnitude & shape of form factors are a stringent test of theory.
- II. Absolute charm semileptonic rate gives direct measurements of V_{cd} and V_{cs} .
- III Key input to precise V_{ub} vital CKM cross check of $\sin 2\beta$



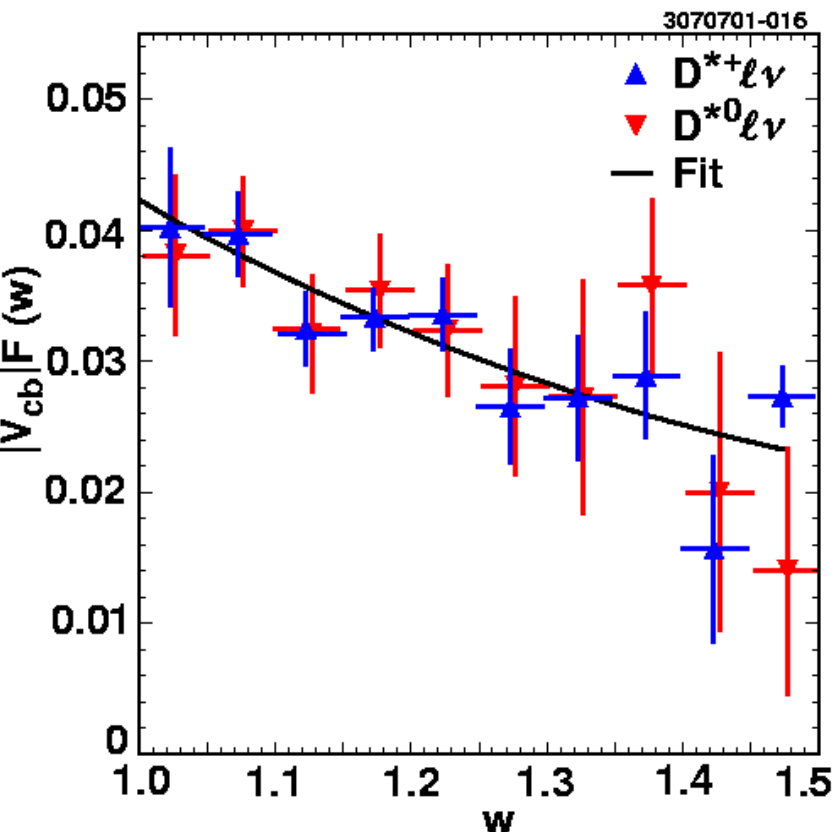
- 1) Measure $D \rightarrow \pi$ form factor in $D \rightarrow \pi l \nu$. Calibrate LQCD uncertainties .
- 2) Extract V_{ub} at BaBar/Belle using *calibrated* LQCD calc. of $B \rightarrow \pi$ form factor.
- 3) But: need absolute $\text{Br}(D \rightarrow \pi l \nu)$ and high quality $d\Gamma(D \rightarrow \pi l \nu)/dE_\pi$ neither exist

The Importance of Precision Charm Absolute Branching Ratios I

V_{cb} from zero recoil in $B \rightarrow D^* l^+ \nu$

CLEO hep-ex/0203032

Accepted for publication in PRL



$$|V_{cb}| = (46.9 \pm 1.4 \pm 2.0 \pm 1.8) \times 10^{-3}$$

CLEO has single most precise V_{cb} by this technique

Stat: 3.0% Sys 4.3% theory 3.8%

Dominant Sys: ϵ_{π} slow, form factors

As B Factory data sets grow, and theory improves

$$dB(D \rightarrow K\pi) / dB(D \rightarrow K\pi)$$

$$\rightarrow dV_{cb} / V_{cb} = 1.3\%$$

The Importance of Precision

Charm Absolute Branching Ratios II

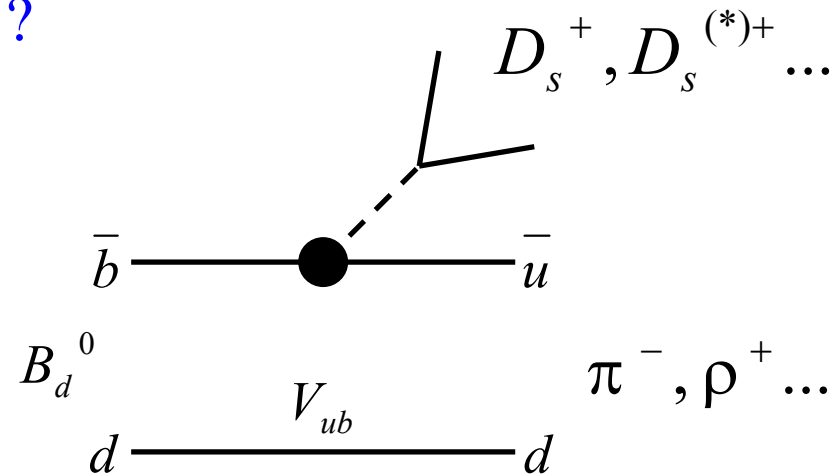
$B^0 \rightarrow D_s^{(*)+} \pi^-$ Extraction of V_{ub} ?

Dominated by $b \rightarrow u$ transition

BABAR/Belle have signals

Theory error probably large

Experimental error dominated
 by $B(D_s \rightarrow \phi \pi)$ which is known
 to 25%



V_{ub}/V_{cb} from $\frac{\Gamma(\Lambda_b \rightarrow p \ell \bar{\nu})}{\Gamma(\Lambda_b \rightarrow \Lambda_c \ell \bar{\nu})}$ at hadron machines requires:

$B(\Lambda_c \rightarrow p K \pi)$ poorly known:
 $9.7\% > B > 3.0\%$ at 90% C.L

The importance of precision absolute Charm BRs III

HQET spin symmetry test

$$\frac{\Gamma(\bar{B}^0 \rightarrow D^{*+} h^-)}{\Gamma(\bar{B}^0 \rightarrow D^+ h^-)} = 1$$

since $D^{*+} \rightarrow \pi^+ D^0$ is most useful mode, \nearrow
 this compares D^0/D^+ absolute rates

Compare $B^0 \rightarrow D^{(*)+} h^-$ and $B^+ \rightarrow D^{(*)0} h^-$ rates to extract color suppressed amplitudes

Test factorization with $B \rightarrow DD_s$

Need Abs Br
 D_s

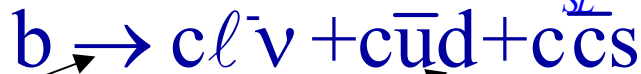
The importance of precision absolute Charm BRs IV

$BR_{SL} = B(b \rightarrow c\ell\nu)$, is low compared to theory A possible explanation is that

the c quark effective mass is low \rightarrow large decay rate for $b \rightarrow c\bar{c}s(d)$

$\rightarrow n_c = (n_c + n_{\bar{c}})$ is negatively correlated to BR_{SL}

The simultaneous measurement of BR_{SL} and n_c can clarify the theoretical picture



BR_{SL} n_c are anti-correlated

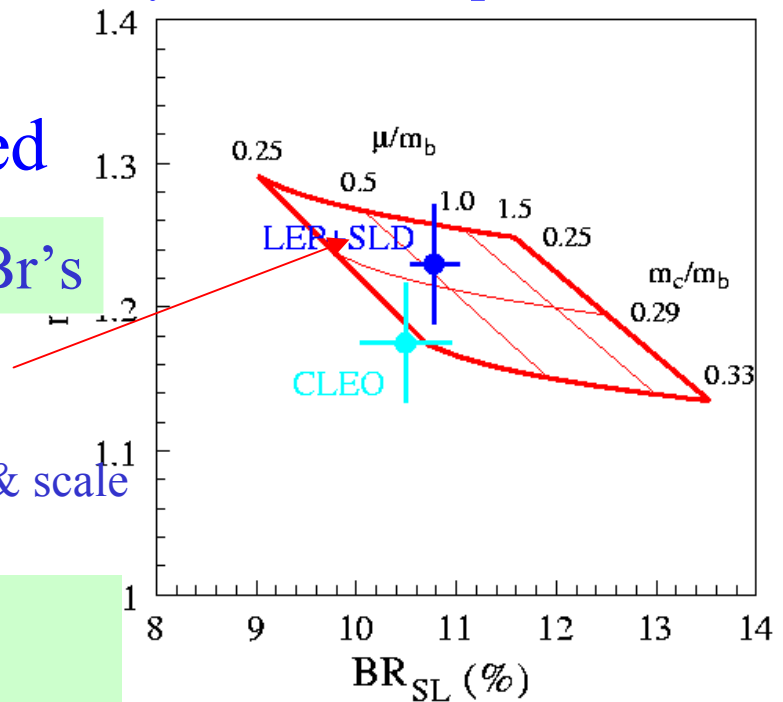
To measure n_c need absolute charm Br's

◆ QCD prediction:

(Neubert & Sachrajda, Nucl. Phys. B483,339 1997)

As a function of the effective charm quark mass & scale at which QCD corrections are evaluated

Theory can accommodate present values but experimental errors are large \rightarrow test becomes more incisive if abs charm BR's were known precisely



(Plot from Parodi at HF9)

The importance of precision absolute Charm BRs V

Test of the Standard Model.

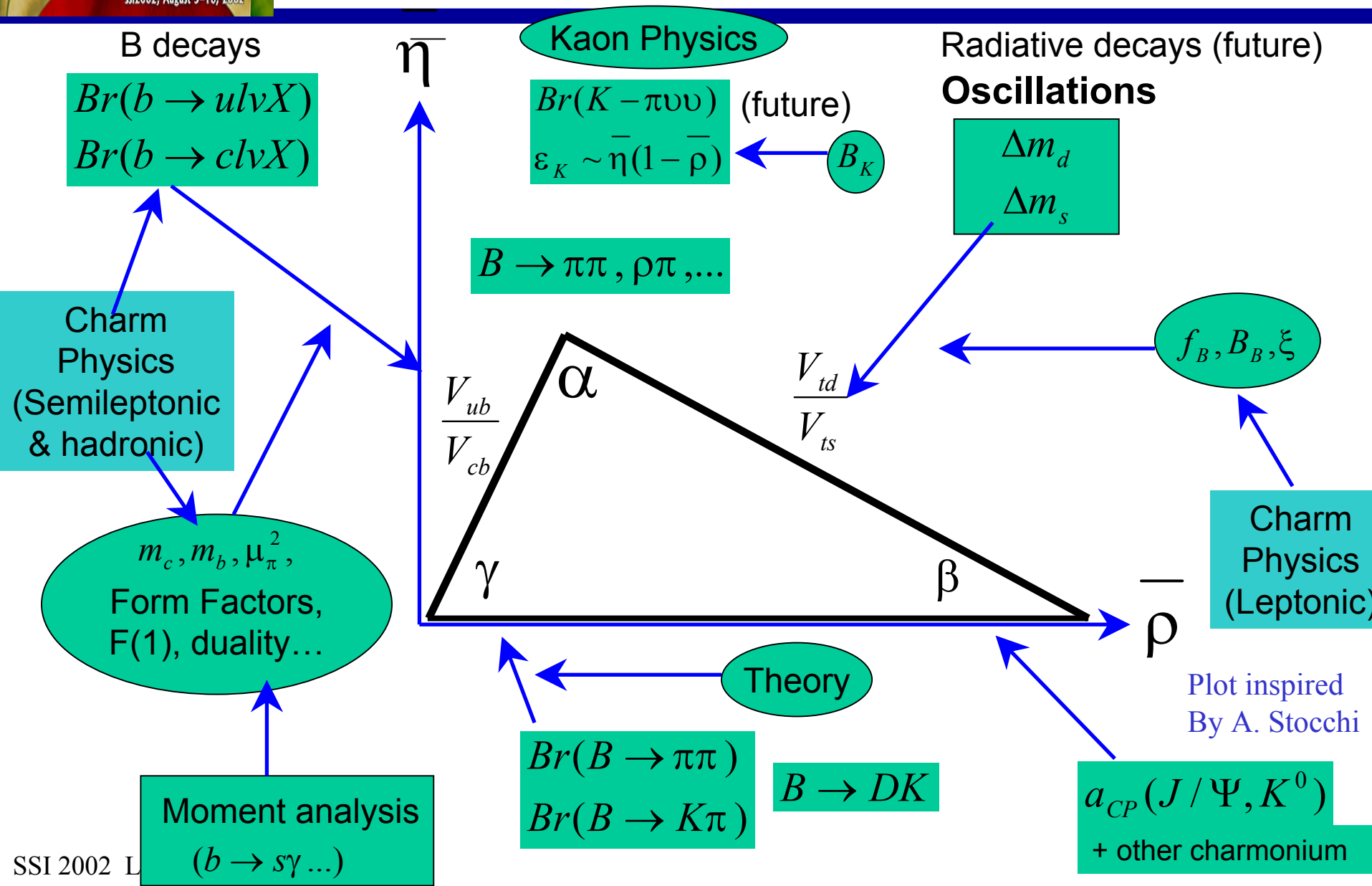
Precision: $Z \rightarrow bb$ and $Z \rightarrow cc$ (R_b & R_c)
is systematically limited by knowledge of
absolute charm branching ratios

To understand the Higgs at LHC/LC

$B(H \rightarrow bb)$ $B(H \rightarrow cc)$

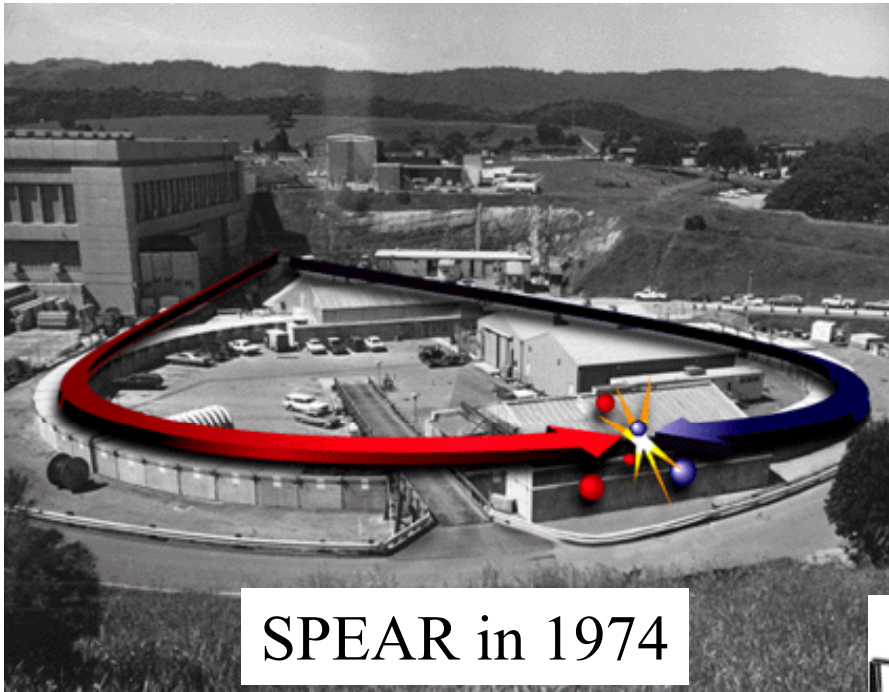
precision will depend on absolute charm
branching ratios

The Unity of Quark Flavor Physics



Charm physics: 1974

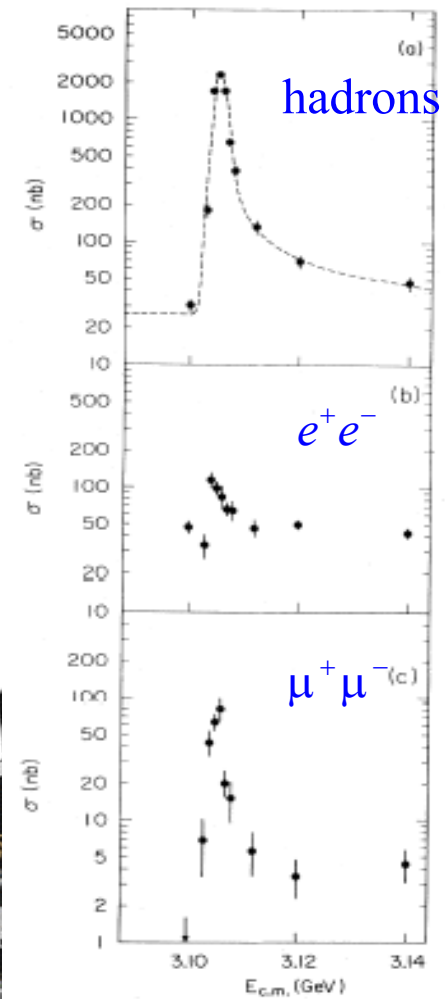
Charm discovery SLAC and BNL November 1974



$e^+ e^- \rightarrow$ multihadron enhancement
 J/Ψ width $\ll 2$ MeV (beam width)



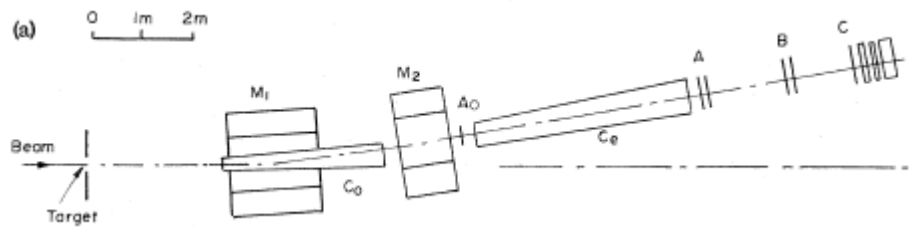
W LETTERS 2 DECEMBER 1974



Goldhaber, Perl, Richter 1974 15

Charm Physics: 1974

$$p + Be \rightarrow e^+ e^-$$



Broad band probe, clean final state

EW LETTERS

2 DECEMBER 1974

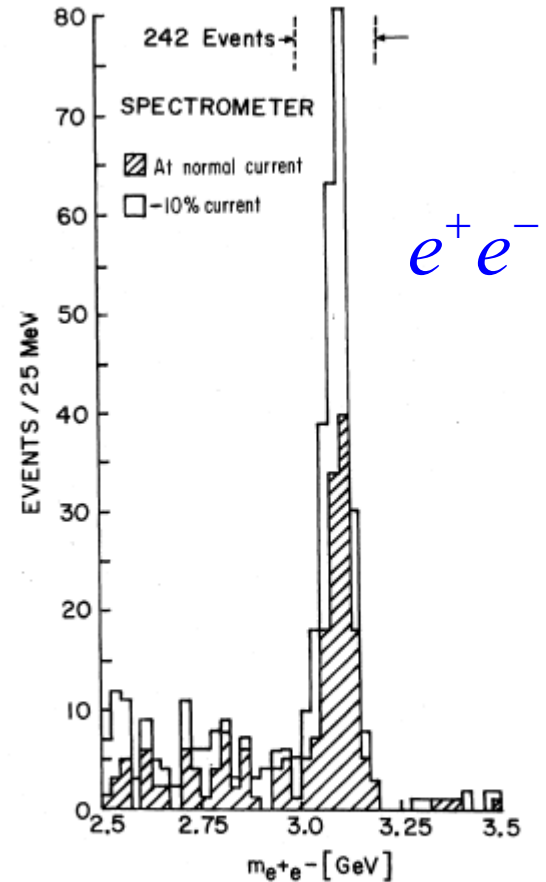


FIG. 2. Mass spectrum showing the existence of J_c . Results from two spectrometer settings are plotted showing that the peak is independent of spectrometer currents. The run at reduced current was taken two months later than the normal run.

The J / Ψ as a frontier

Ψ 's are narrow, insufficient energy to decay to

open charm, (i.e. $D\bar{D}$) or $(c\bar{u})(\bar{c}u)$

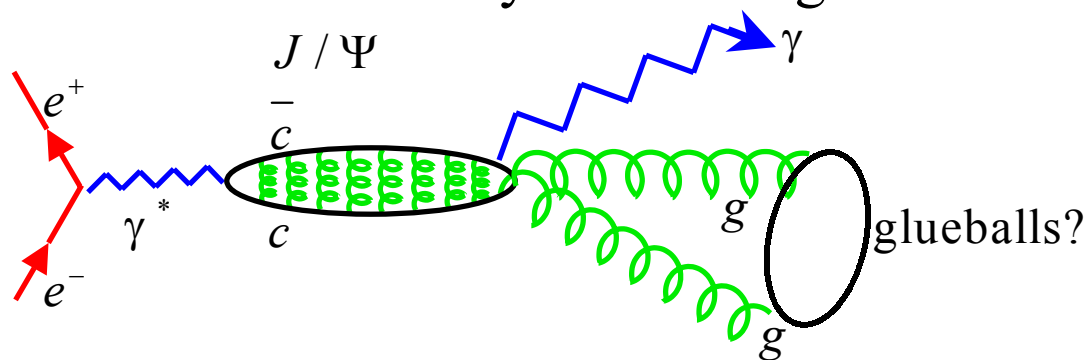
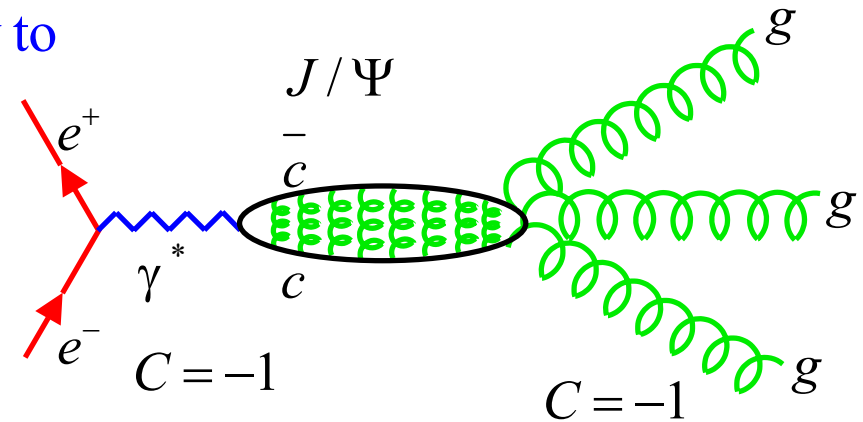
$C=-1$ easy to produce with virtual photon, but decay into three gluons suppressed

Despite this, 88% of J / Ψ decays are hadronic

(only the $\sim 12\%$ to $e^+e^- / \mu\mu$) used in $\sin 2\beta$

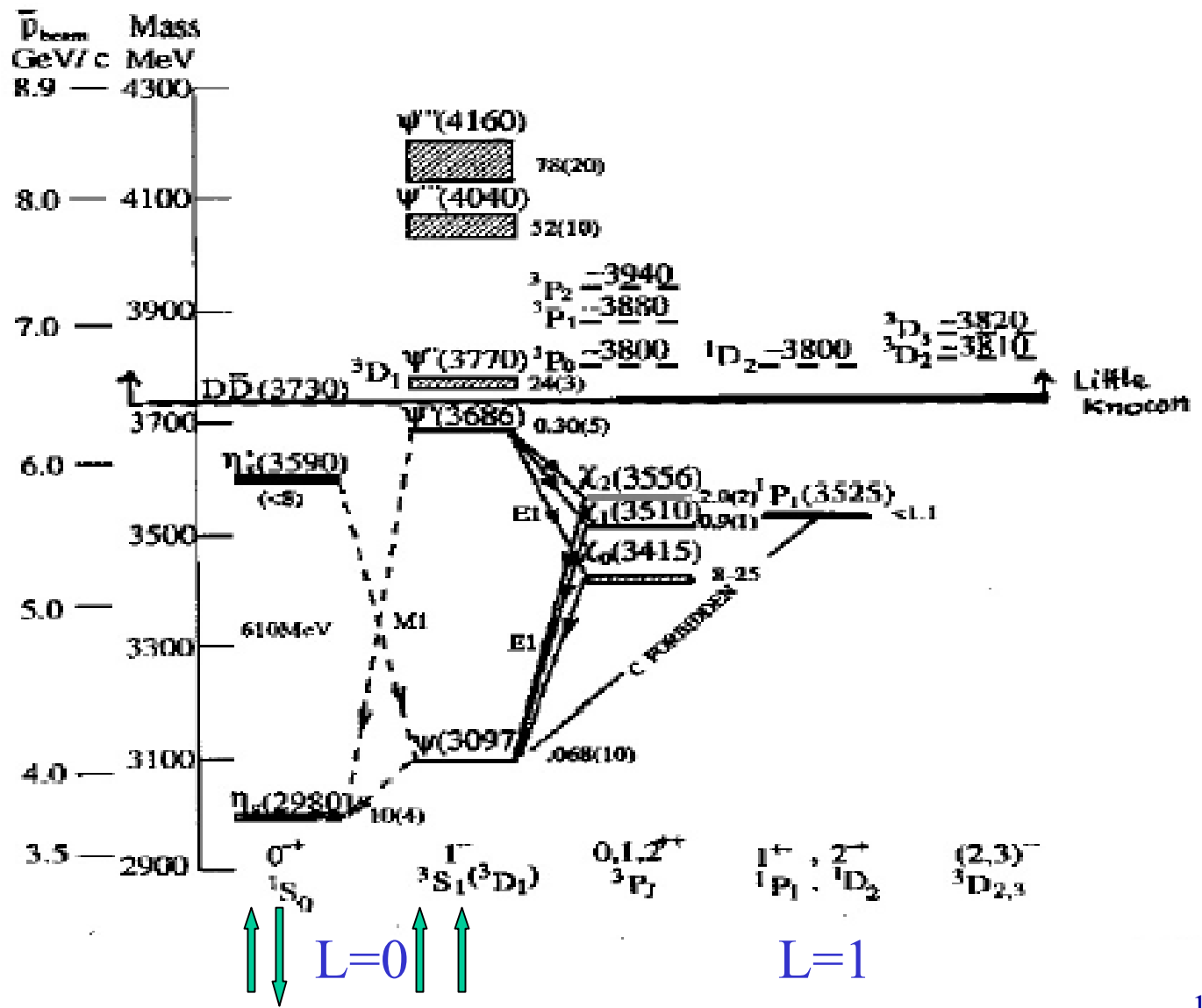
measurements

Radiative J / Ψ Br $\sim 6\%$ are very useful for glueball searches



Charmonium Spectroscopy

Open Charm

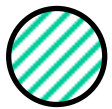


Open Charm Production at Threshold

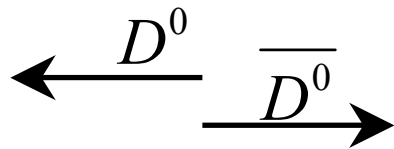
D meson discovered 1976 Goldhaber/Trilling (LBL)

At $\Psi'' = \Psi(3S) = \Psi(3770)$

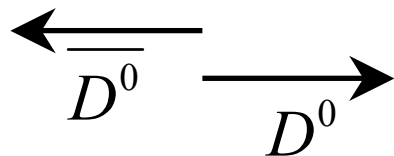
$$\Psi'' \rightarrow D^0 \bar{D}^0, D^+ D^-$$



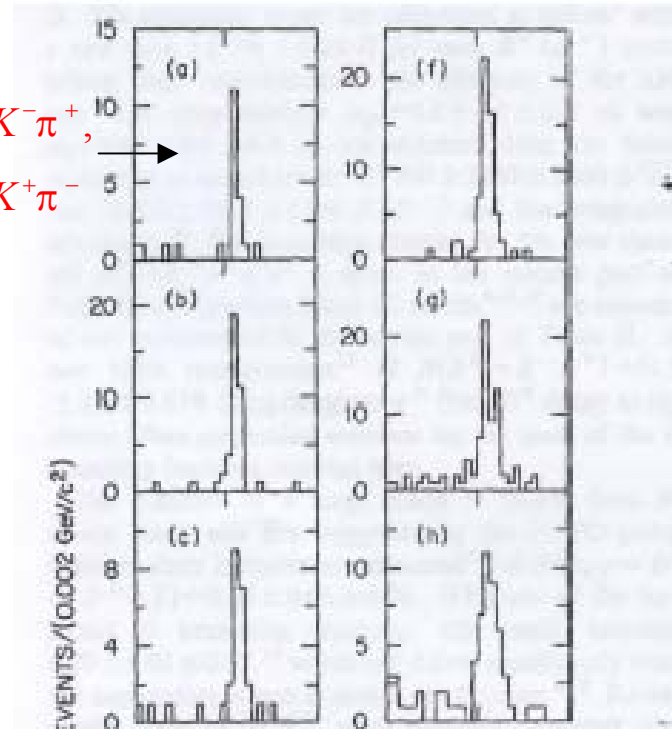
$$\Psi'' \rightarrow L=1$$



$$+(-1)^L$$



$$D^0 \rightarrow K^- \pi^+, \bar{D}^0 \rightarrow K^+ \pi^-$$



Beam constrained mass

Mark III
 Events
 where
 both D's
 are
 reconstructed

The role of the $\Psi(3770)$ in charm physics is analogous to the role of the $Y(4S)$ in B physics



Absolute Branching Ratio Measurements at the $Y(4S)$ and $\psi(3770)$

$$Br(D \rightarrow X) = \frac{\#X \text{ Observed}}{\text{efficiency} \times \#D\text{'s produced}}$$

In B decay absolute branching ratios are measured at the $Y(4S)$

$$\int L dt \cdot \sigma_{(Y(4S))} = N_{Y(4S)}$$

$$N_{Y(4S)} \cdot Br(Y(4S) \rightarrow B\bar{B}) = 2N_B$$

$$Br(Y(4S) \rightarrow B\bar{B}) = 1$$

The number of B's produced is well known

With sufficient statistics, it is possible to eliminate the $Y(4S)$ Br assumption, and complications from the fraction of $B^+ B^+$ and $B^0 B^0$ at the $Y(4S)$

by tagging (fully reconstructing one B in the event)

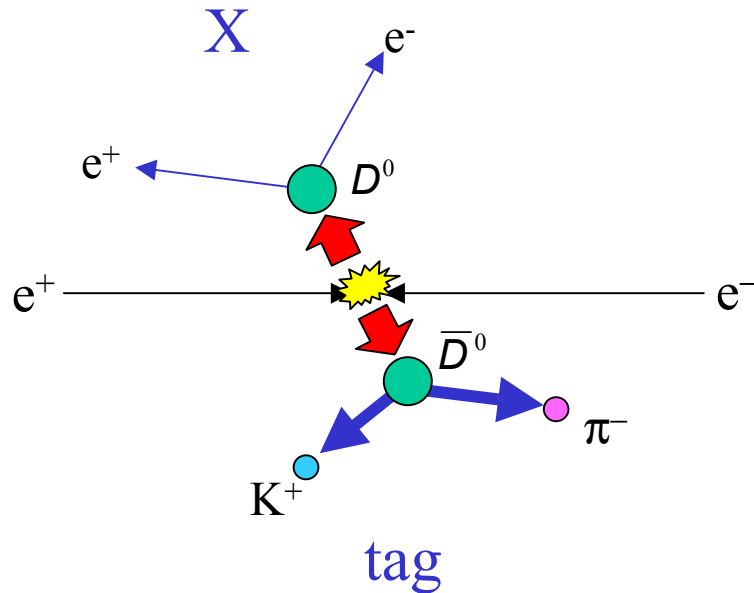
$$Br(B \rightarrow X) = \frac{\#X \text{ Observed}}{\text{efficiency} \times \#B\text{tags}}$$

Full B reconstruction has a low efficiency $\epsilon(\text{tag}) \sim 0.7\%??$, But will become a staple as B Factory data sets grow. Similarly, charm branching ratios could be measured at the $\Psi(3770)$

Absolute Charm Branching Ratios at Threshold

$$\psi(3770) \rightarrow DD$$

$$Br(D \rightarrow X) = \frac{\#X \text{ Observed}}{\text{efficiency} \times \#D\text{'s produced}}$$



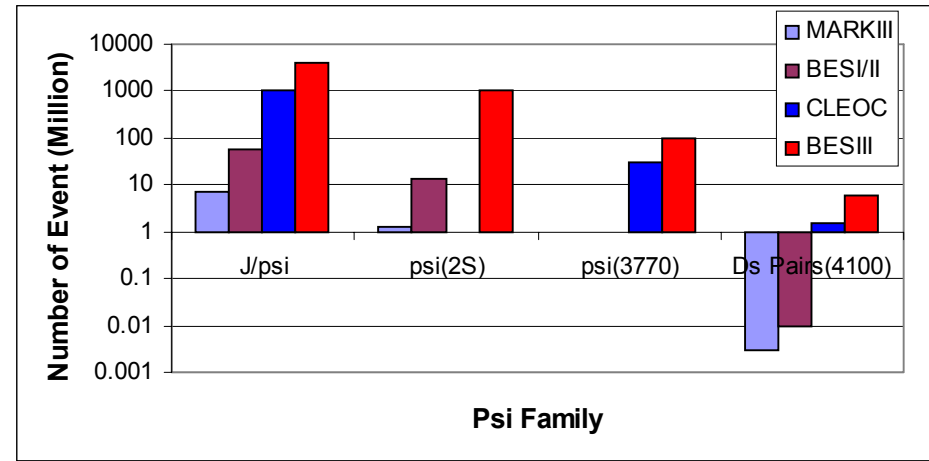
Where the # of D's produced
 Is the # of tags

The tag efficiency at the
 $\psi(3770)$ is expected
 to be about 20% as the
 D has large branching
 ratios to 2-body final states

$$\sigma \psi(3770) = 10 \text{ nb } (\sim 10 \sigma Y(4S))$$

Open Charm Production

experiments at $\Psi(3770)$		# D's
Mark III	9.6 pb^{-1}	2×10^4
CLEO-c (proposed)	3 fb^{-1}	6×10^7
BESIII (proposed)	30 fb^{-1}	6×10^8



As we will see, the $\Psi(3770)$ is by far the best place to determine absolute charm branching ratios. But nobody has operated there since 1984. There are plans to change this situation, \rightarrow

CLEOC
 phys. run \rightarrow

MARKIII \leftarrow

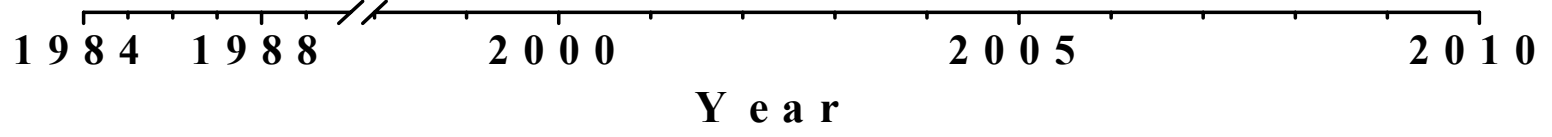
BESII

BESIII

BESIII

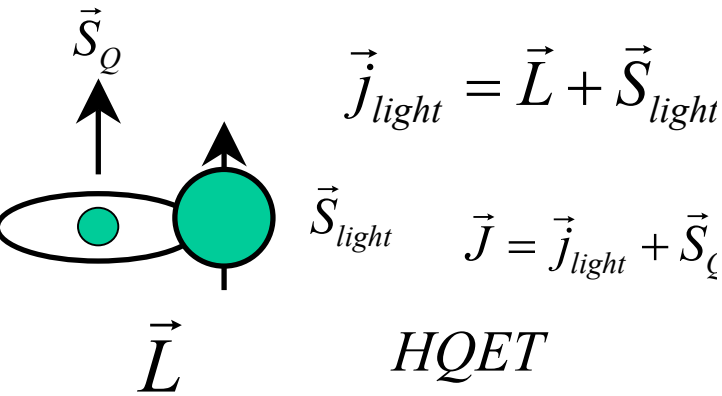
Construction

Engineer & phys. run



Charm Hadrons

Ciulli hep-ex/991104



Heavy quark (Q) hadrons: Q spin decouples $\propto 1/m_Q$. Spin of Q and total spin j of light quark are separately conserved quantum numbers.

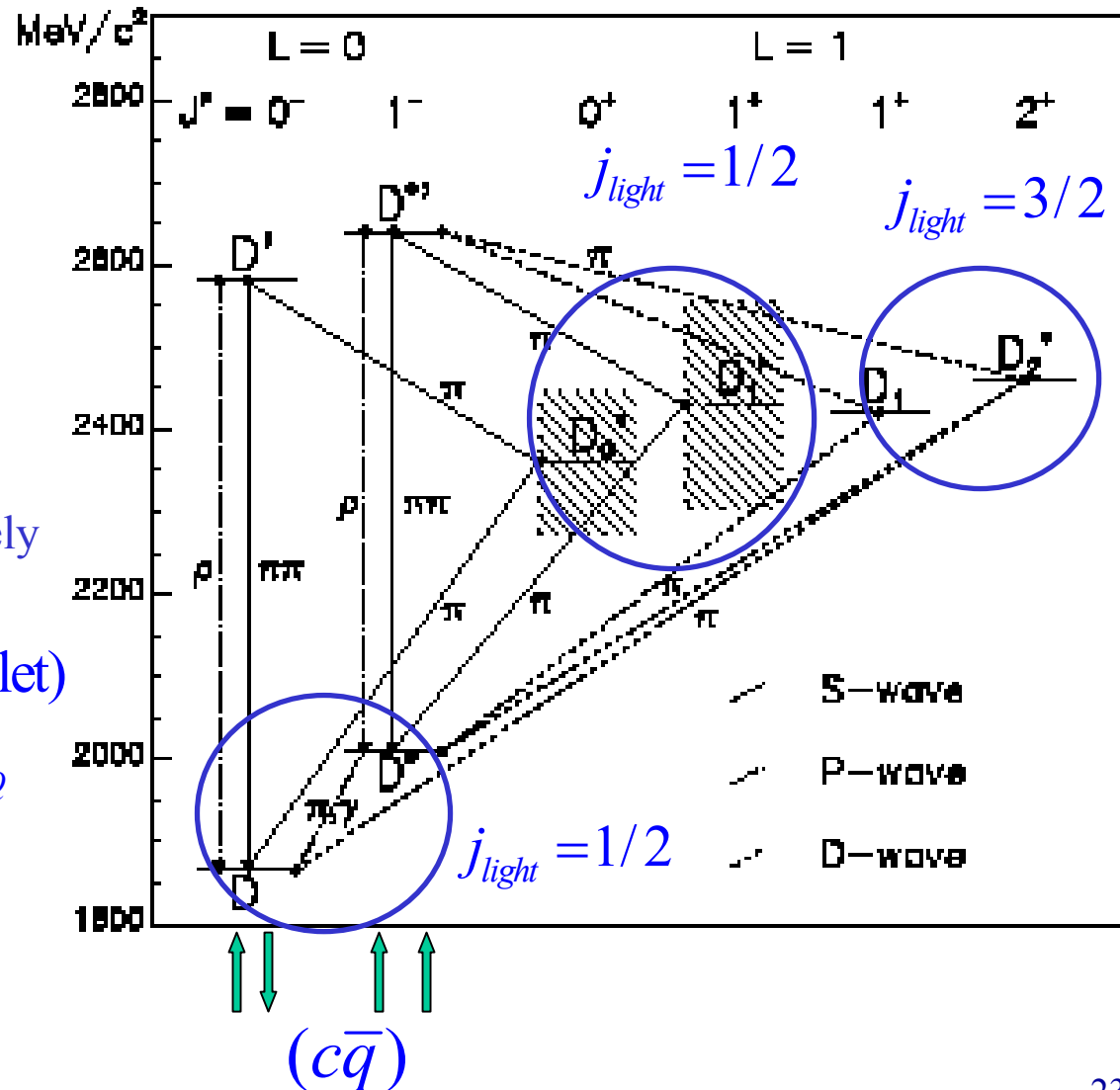
$\Rightarrow J = j_{light} \pm 1/2$ (degenerate doublet)

Corrections go like Λ_{QCD} / m_Q

$\Delta m = m(D^* - D) \sim 142 \text{ MeV}$

$\Delta m = m(B^* - B) \sim 46 \text{ MeV}$

The same description works for heavy baryons



Charm Physics Facilities

B Physics

→	ALEPH, DELPHI, L3, OPAL	at LEP
→	SLD	at SLD
→	CDF / D0	at TeVatron I / II
→	CLEO	at CESR
→	BELLE	at KEK
→	BABAR	at PEP
	FOCUS	E831 FNAL
	E731, E781 (SELEX)	at FNAL
	CLEO-C	at CESR-C
→	CMS / ATLAS / LHC-B	at LHC
→	BTeV	at TeVatron II
	BESII / III	at BEPC I / II

Charm Physics

Z^0	←
Z^0	←
pp 1.8TeV	←
Y(4S) symmetric	←
Y(4S) asymmetric	←
Y(4S) asymmetric	←
$\gamma < 300$ GeV	←
$\rho-\pi^- \Sigma^+ 600$ GeV	←
ψ	←
pp 14 TeV	←
pp 1.8 TeV	←
ψ	←

Today

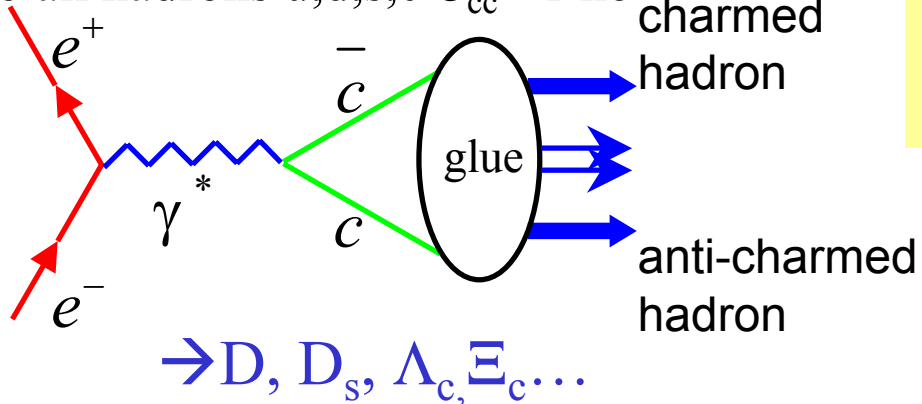
Future

Many of the charm facilities that have finished running but are still producing results, currently running facilities, and future facilities are listed above. Most charm physics facilities are also B physics facilities, exceptions are the fixed target experiments

Charm Production near/at the $Y(4S)$

Non-resonant production

to all hadrons u, d, s, c $\sigma_{cc} \approx 1 \text{ nb}$



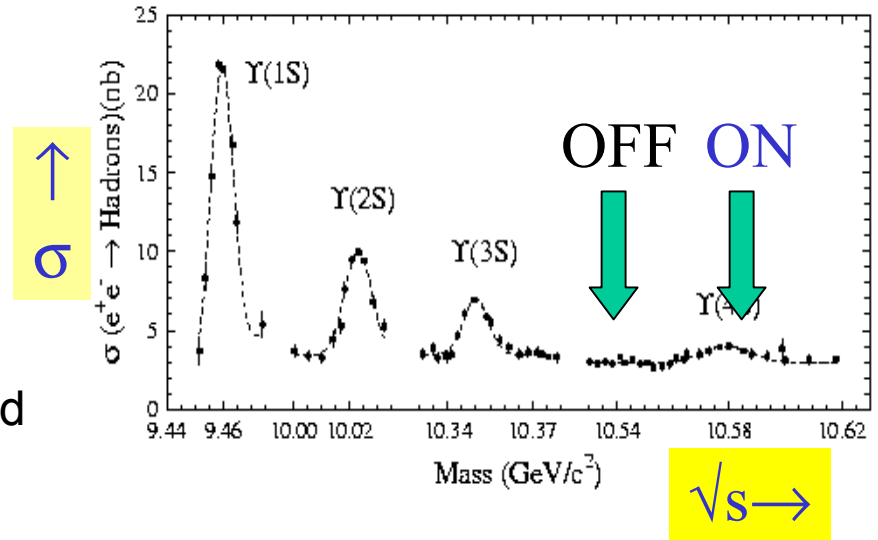
+ $\sigma(B \rightarrow c) \approx 1 \text{ nb}$ $1+1 = 2 \text{ nb}$

D 's move ~ 150 microns, but, energy of D 's is a priori unknown and the charm, and anti-charm hadron types are not strongly correlated

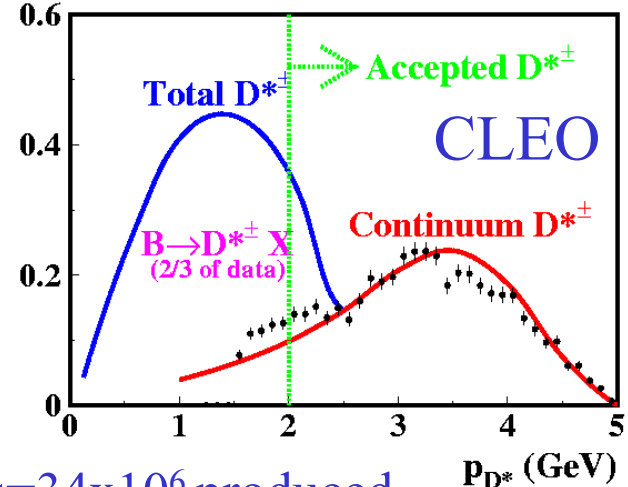
Reconstructing a \bar{D}^0 does not mean the charm hadron in the event is a D^0 , a D^+ , D_s , Λ_c , Ξ_c , Ω_c or any other ground state or excited charm hadron is also a possibility

\rightarrow absolute charm Br's difficult

3 experiments, CLEO $\#c$'s = 34×10^6 produced
BABAR, BELLE $\#c$'s = $\times(5-6)$ CLEO now, by 2005 $\times 40$ CLEO



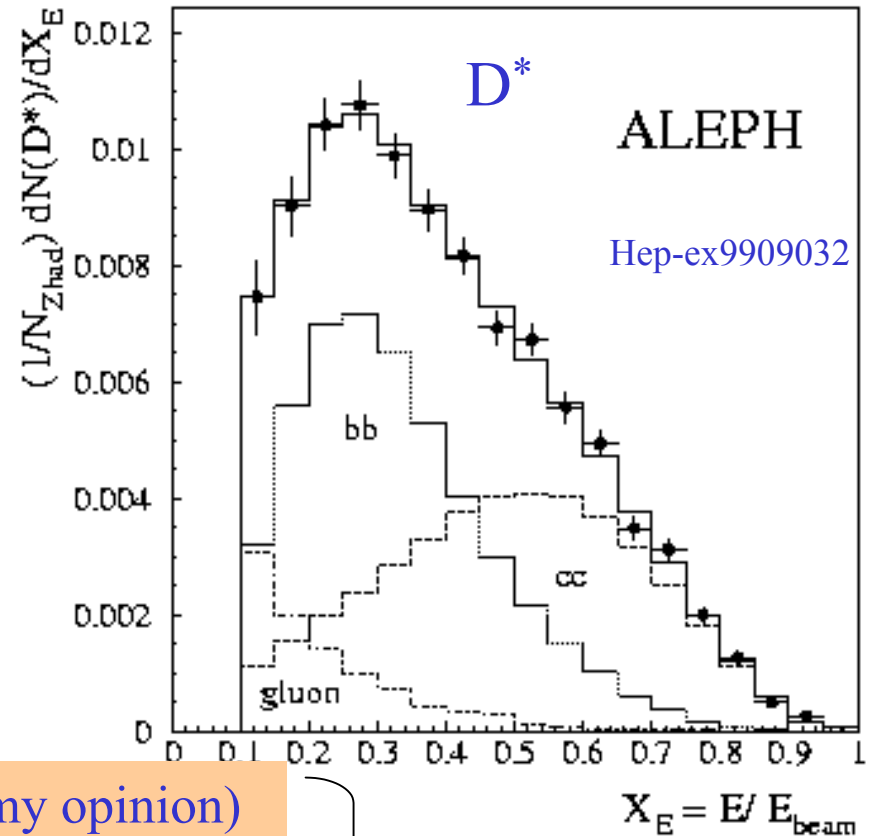
$D^{*\pm}$ Production at 10.58 GeV



Charm Production at the Z^0

$\text{Br}(Z^0 \rightarrow cc) \sim 11\%$
 $\rightarrow 3 \times 10^6 \text{ c's / LEP expt}$
 $\rightarrow D, D_s, \Lambda_c, \Xi_c \dots$
 $\rightarrow \langle p \rangle \sim 40 \text{ GeV}$
 \rightarrow Excellent reconstruction of Charm vertices

As at 10 GeV, the flavors of charm anti-charm pairs produced in Z^0 decays are not correlated \rightarrow absolute charm Br's are difficult

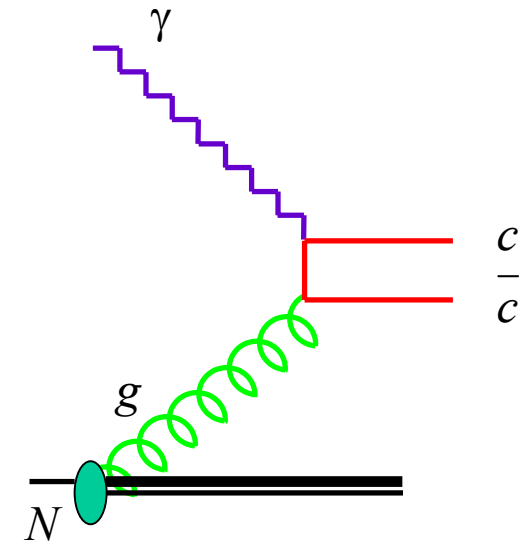
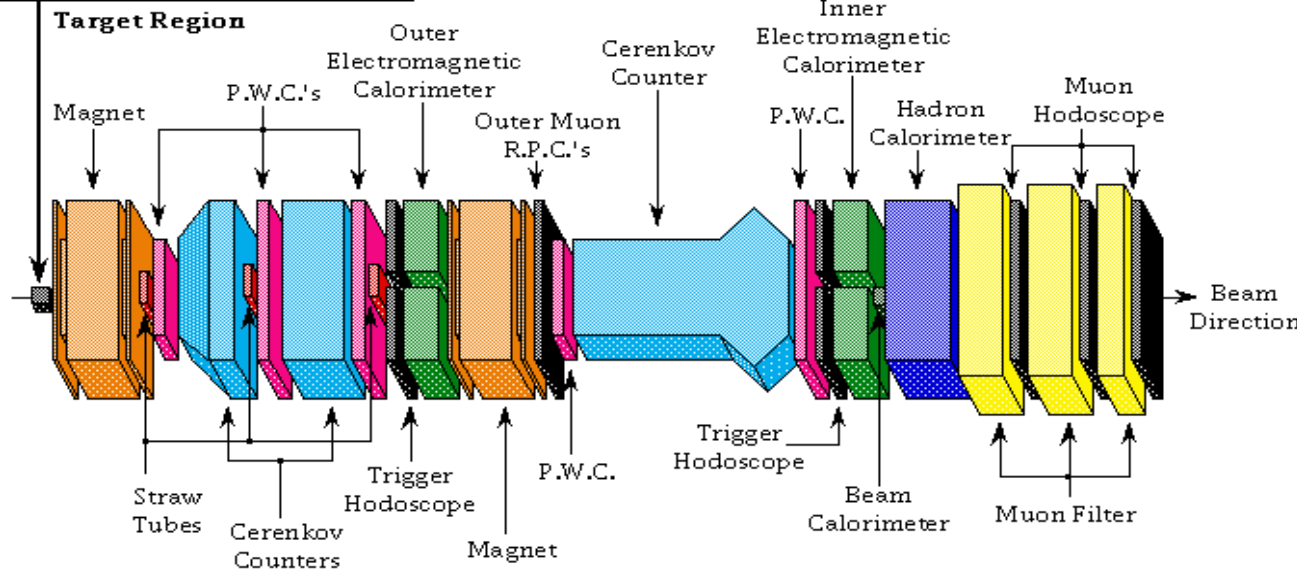
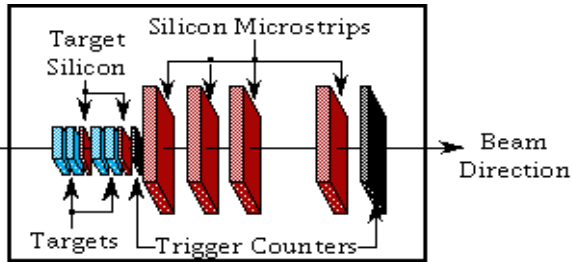
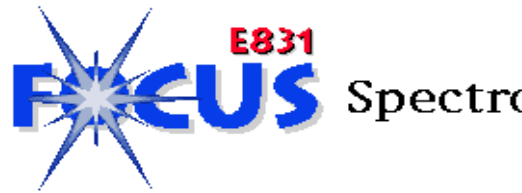
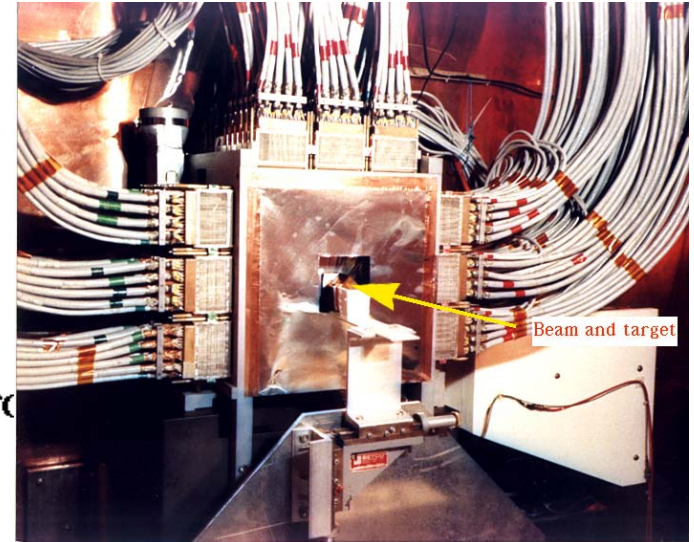


Will not discuss
 In lecture

Main LEP contributions to charm physics: (my opinion)
 - electro-weak measurements: $R_c = (Z^0 \rightarrow cc) / (Z^0 \rightarrow \text{had})$
 A_{FB}^c (charge asymmetry in $e^+ e^- \rightarrow cc$)
 See: <http://lepewwg.web.cern.ch/LEPEWWG/>
 - c-quark fragmentation function
 - D_s decay constant

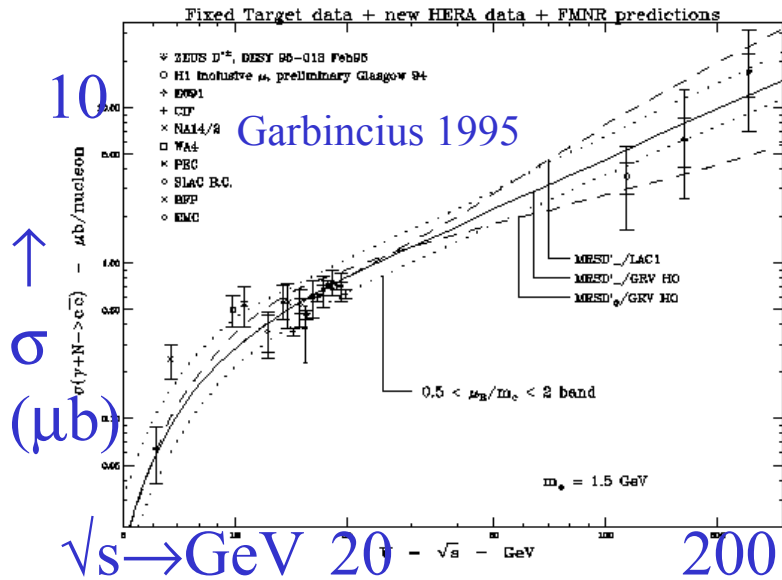
Photo production: FOCUS

Fixed target experiments have long been at the frontier of charm physics
 Detector scale typical, tiny front end

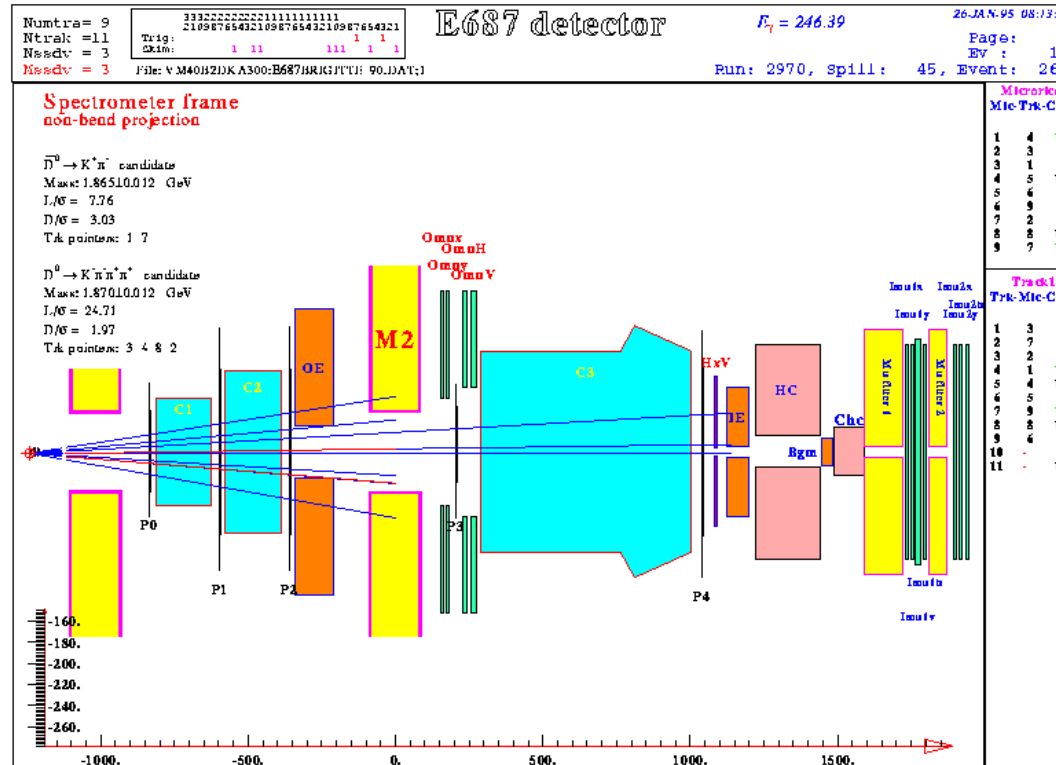


FOCUS: photoproduction

(Also HERA)



$\sim 10^8$ c produced
 $> 10^6$ c reconstructed
Excellent:
-Vertex Resolution
-Particle ID & $\delta p_T/p_T$
D, D_s , Λ_c , Ξ_c
Lifetimes. Mixing
No absolute Br's



FOCUS: close-up

Numtra= 9
Ntrak= 11
Nsedv= 3
Ksedv= 3

333222222221111111111111
2109E765432109E765432109E7654321
Trig: 1 11 111 1 1 1
Skim: 1 11 111 1 1 1

File: V M40B2DK A300:E687BRK1PTE: 90.DAT:1

E687 detector

$E_{\gamma} = 246.39$

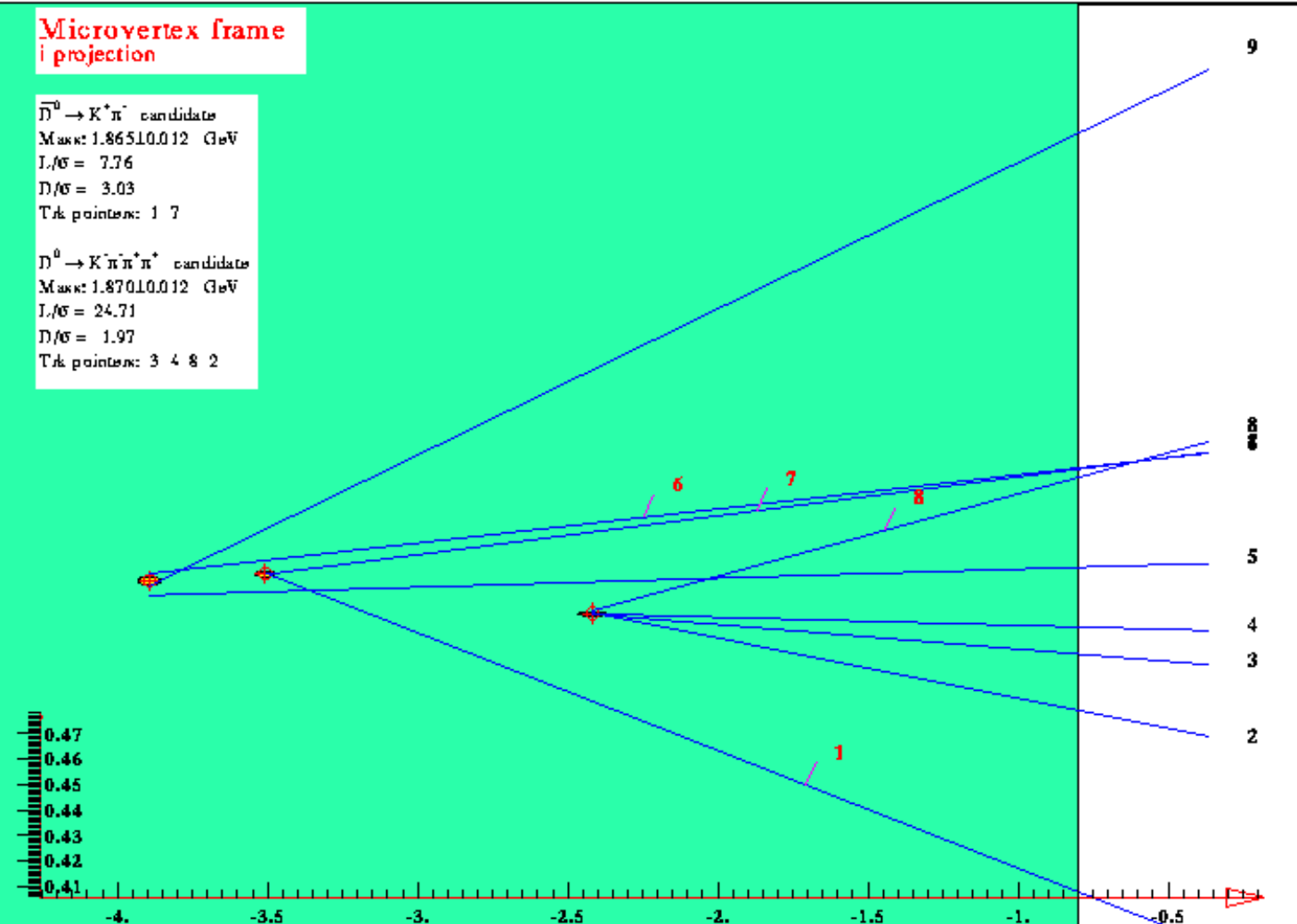
26-JAN-95 08:13:3

Page: 17
EV : 17
Run: 2970, Spill: 45, Event: 260

Microvertex frame
i projection

$\bar{D}^0 \rightarrow K^+ \pi^-$ candidate
Mass: 1.86510012 GeV
 $L/\sigma = 7.76$
 $\Pi/\sigma = 3.03$
Trk points: 1 7

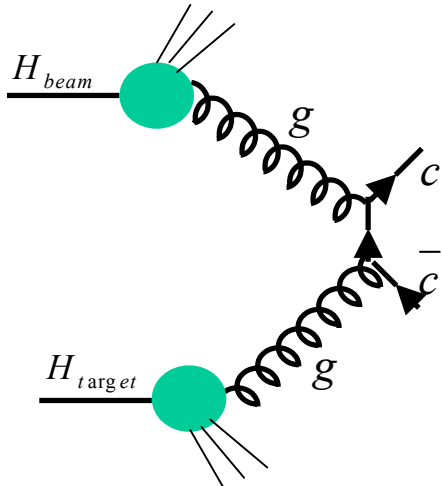
$\Pi^0 \rightarrow K^+ \pi^- \pi^+ \pi^-$ candidate
Mass: 1.87010012 GeV
 $L/\sigma = 24.71$
 $\Pi/\sigma = 1.97$
Trk points: 3 4 8 2



Microvtx		
Mic	Trk	Ce
1	4	1
2	3	3
3	1	5
4	5	1
5	6	2
6	9	1
7	2	5
8	8	1
9	7	1

Track1		
Trk	Mic	Ce
1	3	5
2	7	5
3	2	3
4	1	1
5	4	1
6	5	2
7	9	1
8	8	1
9	6	1
10	-	3
11	-	1

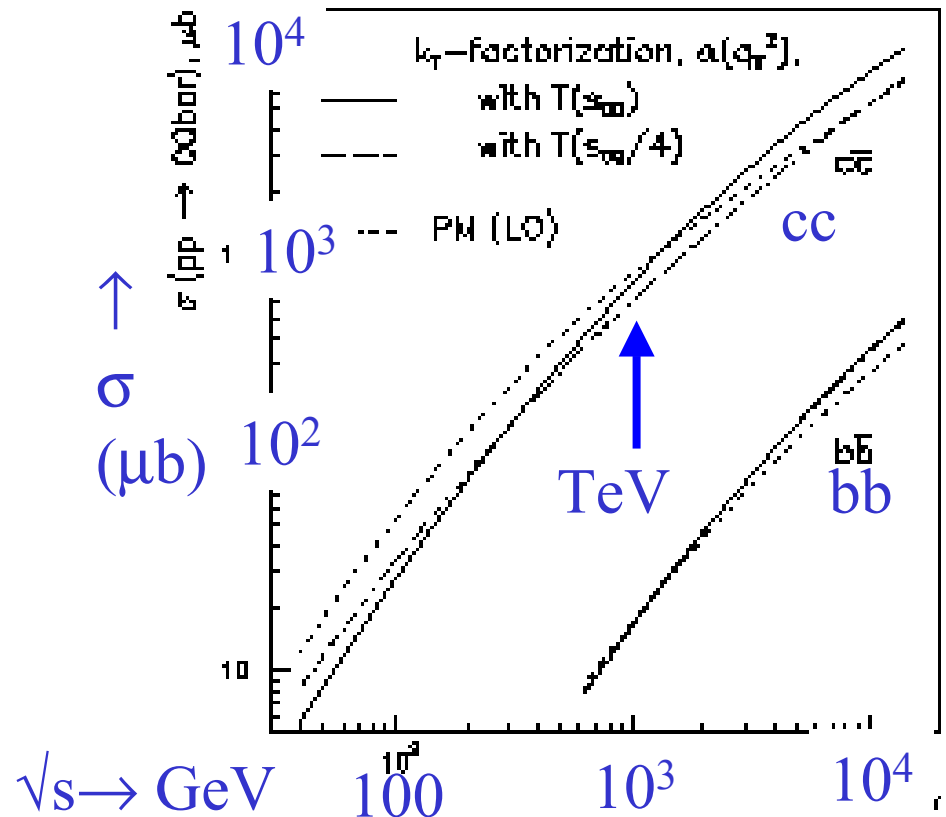
Hadroproduction



Hadronization

- SELEX (E791) at FNAL
 10^4 (10^5) c 's reconstructed
- * Millibarns at the Tevatron
 $\sim 10^{13} c$'s/year Run II
 (also BTeV)
- * X 10 at LHC
- * HERA-B

Ryskin, Shabelski, Shuvaev, 2000

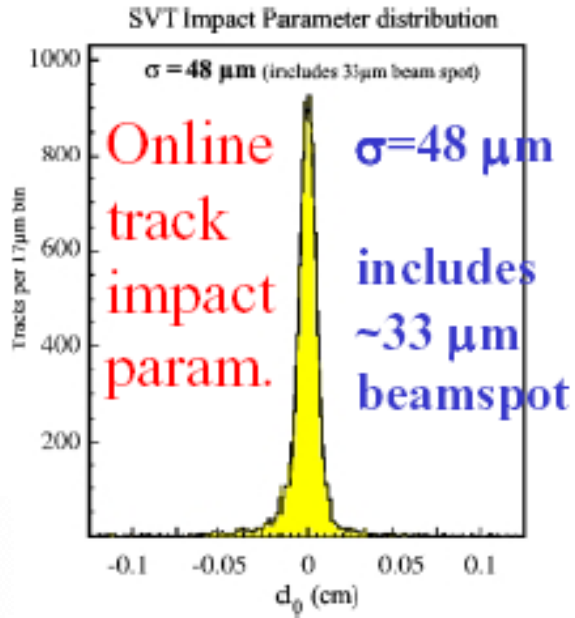


Charm at CDF

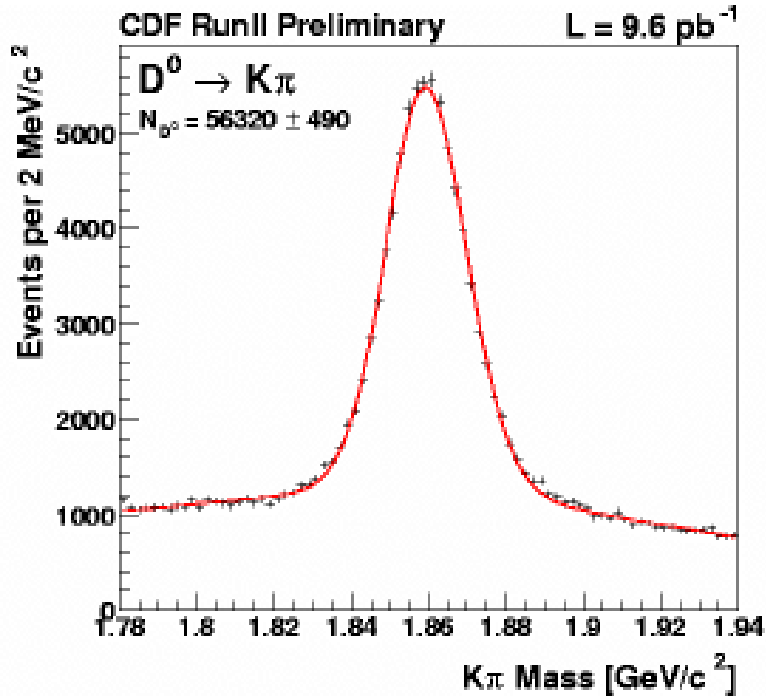
ICHEP2002



The hadronic B trigger a major milestone
 ~150 VME boards find and fit tracks in Silicon, offline accuracy in a 15μs pipeline



Secondary vertex level 2 trigger
 $|D| > 100 \mu\text{m}$ (2 body)



➤ $\Gamma(D \rightarrow KK) / \Gamma(D \rightarrow K\pi) = (11.17 \pm 0.48 \pm 0.98)\%$ (PDG: 10.83 ± 0.27)

■ Main systematic (8%): background subtraction (E687, E791, CLEO2)

Cabbibo suppressed

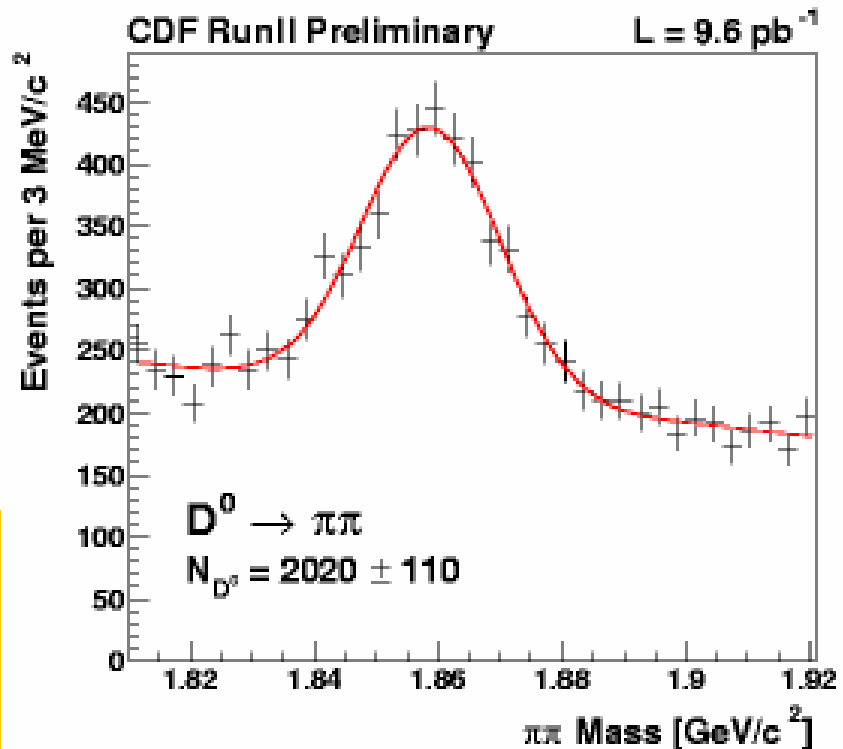
➤ $\Gamma(D \rightarrow \pi\pi) / \Gamma(D \rightarrow K\pi) = (3.37 \pm 0.20 \pm 0.16)\%$ (PDG: 3.76 ± 0.17)

■ several ~2% systematics

Already comparable!

Great potential if charm
 Stays within the trigger
 Bandwidth: expect
 $10^7 D^0 \rightarrow K\pi$ reconstructed
 in RunII 2fb^{-1}

CDF can address D mixing,
 DCPV and rare decays
 Not absolute branching ratios



Summary of current & future charm particle data sources

~total
 # recon.
 charm

	Fixed Target			e ⁺ e ⁻ Collider		
	E791	SELEX	FOCUS (& E687)	CLEO	BABAR	BELLE
Beam	Hadronic		Photon	Off-resonance e ⁺ e ⁻		
Charm	~10 ⁵	~10 ⁴	~10 ⁶	>10 ⁶		
σ _t	~40 fs	~20 fs	~40 fs	~140 fs	~160 fs	

CDF 10⁷ D → K⁻ π⁺ reconstructed

→ X40 CLEO
 By 2005

in RunII

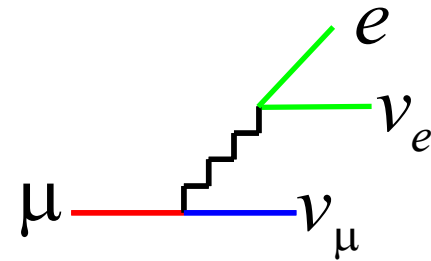
BTeV could have 10⁹ D's CLEO-c 3 x 10⁶ tagged DD

F.T. expts. Measure the c-hadron decay time very precisely this is also crucial to isolate clean event samples

e⁺e⁻ : higher relative production rate of charm compared to background, better mass resolution and great PID mean samples of comparable purity even though time resolution is X10 worse

Lifetimes

Muon decay:
$$\Gamma_o = \frac{G_F^2 m_\mu^5}{192\pi^3}$$



Naïve spectator model for charm

$$\Gamma_c = (2 + 3) \Gamma_o \quad \Gamma_o = \frac{G_F^2 m_c^5}{192\pi^3} |V_{cs}|^2$$

e, μ $u\bar{d}$ x 3 colors

Scaling from the muon:

$$\tau_c = \frac{1}{5} \left(\frac{0.105}{1.5} \right)^5 2.2 \times 10^{-6} = 7 \times 10^{-13} s$$

(700 fs)

$\tau(D^+) \sim 1,000$ fs $\tau(D^0) \sim 400$ fs. Not too bad. Including baryons lifetimes vary between ~ 100 and 1000 fs, \rightarrow non-spectator processes and higher order corrections

Charm Hadron Lifetimes

$$\frac{Br}{\tau} = \Gamma$$

Lifetime needed to compare Br(expt) to Γ (theory)

Test techniques/systematics in other areas where lifetime resolution/vertexing is important: eg B/D $\Delta m, \Delta \Gamma$

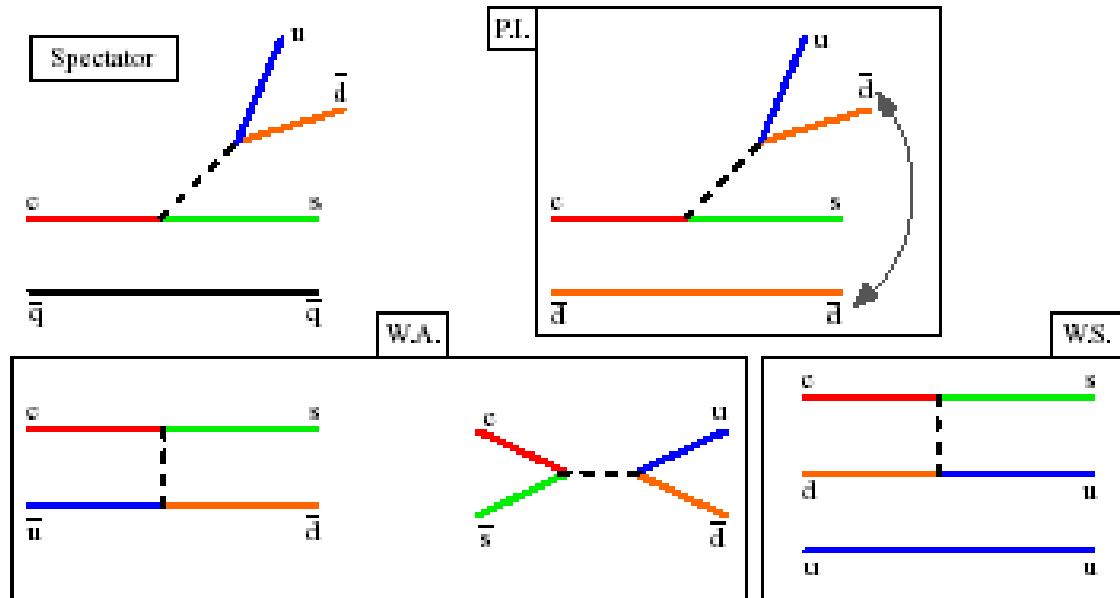
Interpreted with O.P.E.

$$\Gamma(H_c) = \Gamma_{spect} + O(1/m_c^2) + \Gamma_{PI,WA,WS}(H_c) + O(1/m_c^4)$$

Spectator effects (PI.WA,WS) are $O(1/m_c^3)$ but phase space enhanced

Note: hadrons behave more like free quarks the heavier the quark

See G. Bellini, I.I Bigi & P. Dornan
 Phys Rep. **289** (1997)

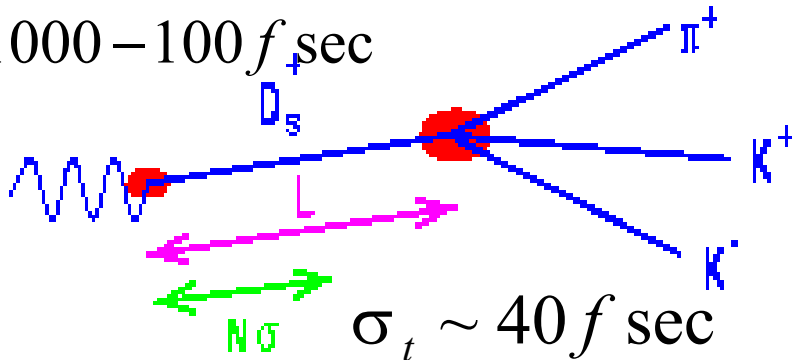


Gross features of the lifetime hierarchy can be explained

Lifetimes at Fixed Target Experiments

$t = 10^{-12} - 10^{-13}$ sec

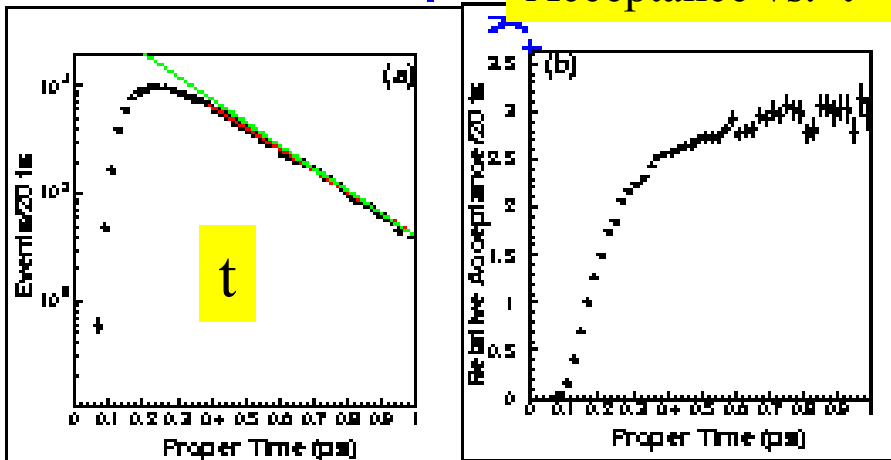
$1000 - 100 f_+ \text{ sec}$



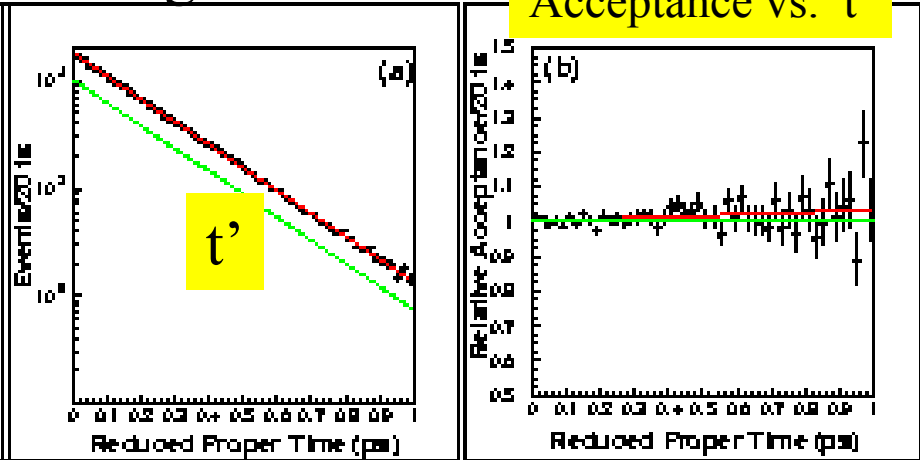
- Short flight path, need silicon
- $L > N\sigma_L$ (and outside target)
- Reduced proper time:
 $t' = L/\beta\gamma c - N\sigma_L/\beta\gamma c$
 to reduce acceptance corrections
- Acceptance checked with data (K_S)
- Systematics from acceptance &/or background

L3D decay length

Acceptance vs. t

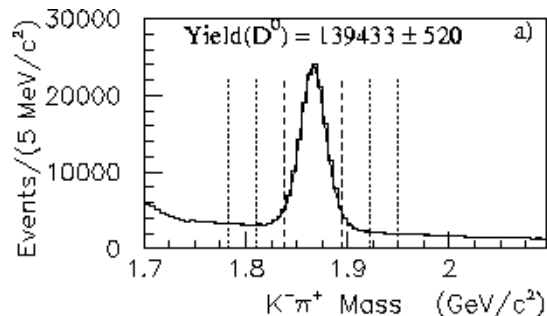


Acceptance vs. t'

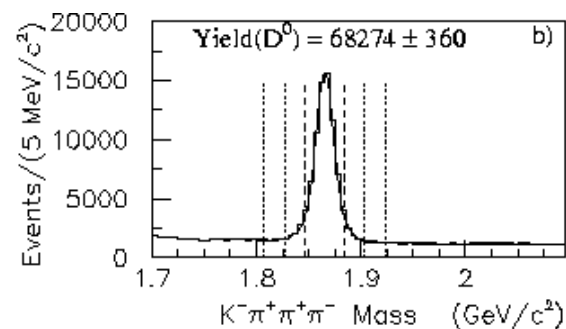
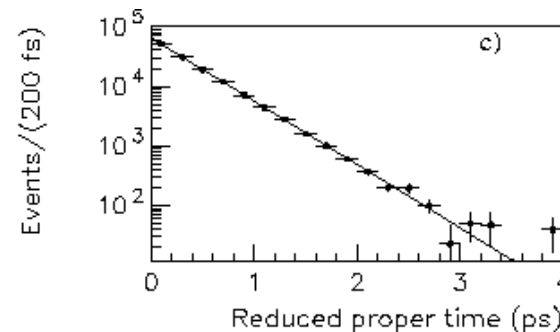


Harry W. K. Cheung - Fermilab

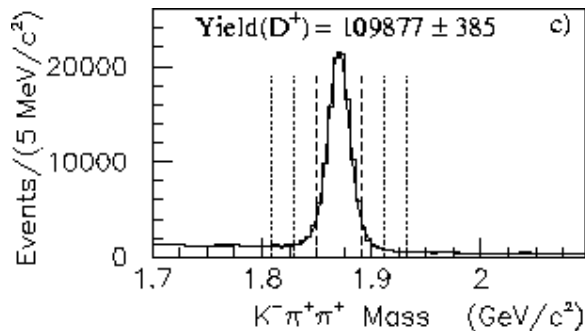
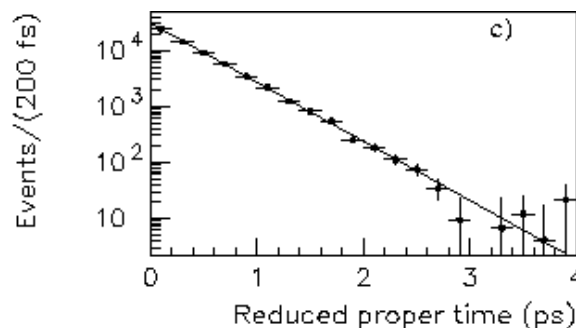
Charm Meson Lifetimes



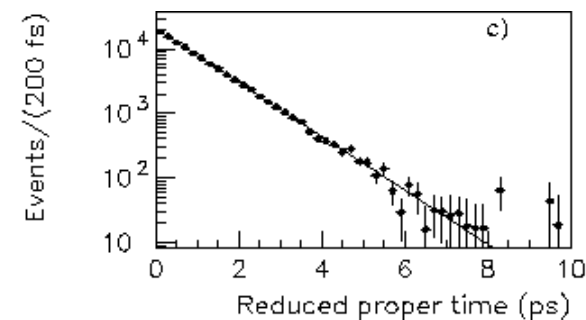
139433 ± 520 evts



68274 ± 360 evts



109877 ± 385 evts



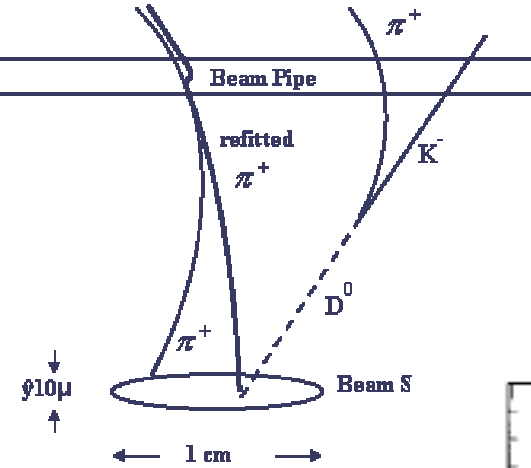
D^0, D^+ Signal

D^0, D^+ Lifetime fits

Lifetimes at e^+e^- colliders

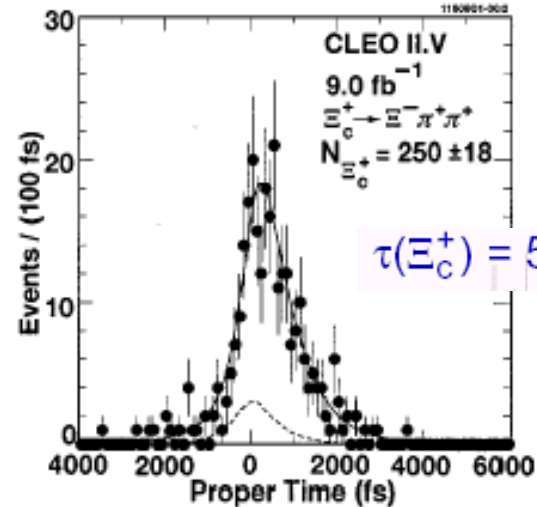
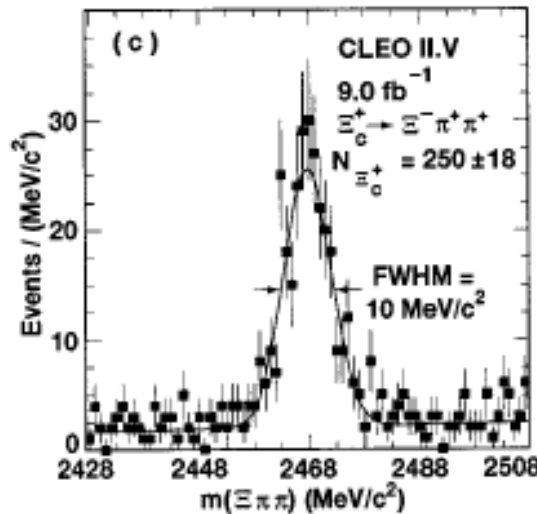
Y(4S) stationary

$L \langle D, \Xi_c^+ \rangle 150 \mu\text{m}$



PRD 650311
 2002

- * Silicon vertex detectors \rightarrow charm lifetimes
- But poorer time resolution ~ 140 fsec (CLEO)
- Needs average IP position
- Uses 2-D (or 1-D) decay length
- Needs good knowledge of mass and t resolutions
- Complicated fits using parameterized resolution and background functions
- Systematics from vertexing, resolutions and fit biases



D Meson Lifetimes

Phys.Lett.B537,192 ,2002 $\tau_{D^+} = 1042.7 \pm 6.9$

New precise measurements of $\tau(D^0)$, $\tau(D^+)$ $\tau(D_s)$ from FOCUS

$\tau(D^0) = 409.6 \pm 1.1(stat) \pm 1.5(sys)$

$\tau(D^+) = 1039.4 \pm 4.3(stat) \pm 7.0(sys)$

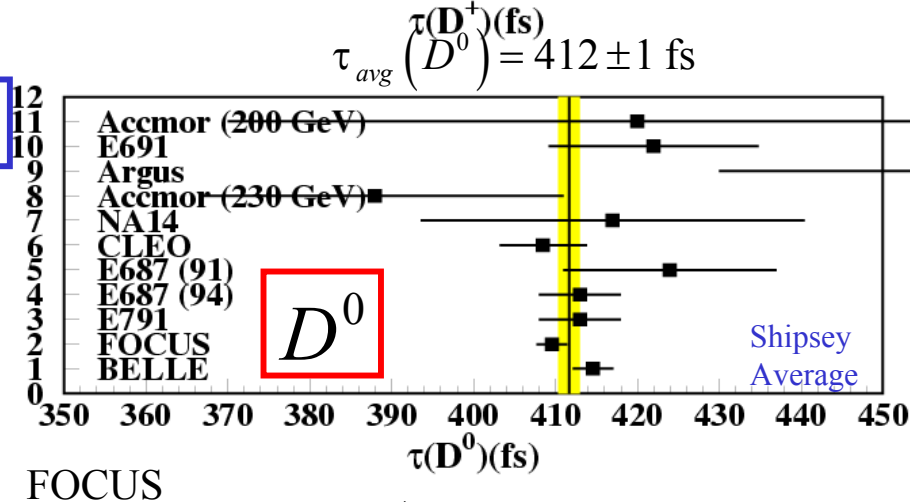
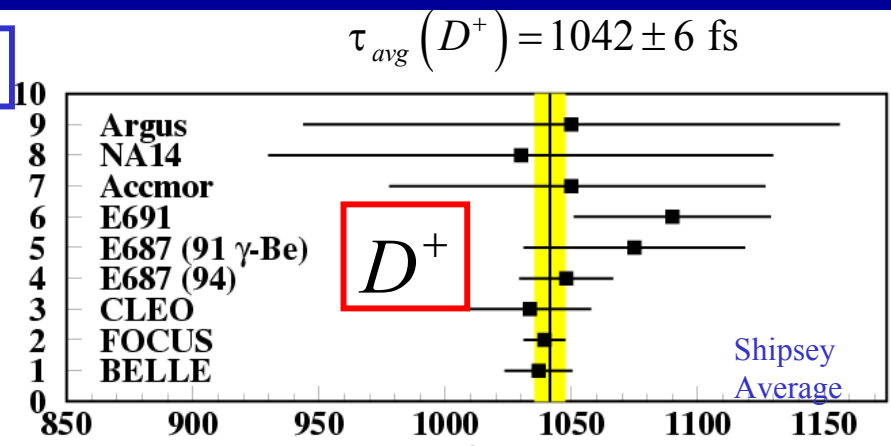
$\frac{\tau_{D^+}}{\tau_{D^0}} \approx 2.583 \pm 0.023$

$\tau_{D^0} = 410.5 \pm 1.5$

Abs. Exclusive semileptonic decays key to interpreting lifetime ratio

$\frac{\Gamma(D^0 \rightarrow eX)}{\Gamma(D^+ \rightarrow eX)} = \frac{B(D^0 \rightarrow eX)}{B(D^+ \rightarrow eX)} \times \frac{\tau(D^+)}{\tau(D^0)} = 1.01 \pm 0.13$

Large observed lifetime ratio must arise due to destructive interference in hadronic diagrams contributing only to D^+ decays (see later)



Note: lifetimes much better known than absolute BR's

$D^+ \rightarrow eX = (17.2 \pm 1.9)\%$

$D^0 \rightarrow eX = (6.87 \pm 0.28)\%$

PDG2002

Lifetime Summary I

$\tau(D^0)$	$411.7 \pm 1.3 \text{ fs}$
$\tau(D^+)$	$1041.5 \pm 6.2 \text{ fs}$
$\tau(D_s)$	$501.7 \pm 6.1 \text{ fs}$
$\tau(\Lambda_c)$	$199.7 \pm 3.3 \text{ fs}$
$\tau(\Xi^+_c)$	$422.3^{+20.1}_{-18.7} \text{ fs}$
$\tau(\Xi^0_c)$	$106.3^{+9.2}_{-7.8} \text{ fs}$
$\tau(\Omega_c)$	$73.6^{+11.8}_{-12.1} \text{ fs}$

Updated after ICHEP02

$$\frac{\tau(D^+)}{\tau(D^0)} = 2.53 \pm 0.02 \quad \text{P.I.(-)}$$

$$\frac{\tau(D_s)}{\tau(D^0)} = 1.22 \pm 0.02 \quad \text{W.A. or ??}$$

$$\frac{\tau(\Lambda_c)}{\tau(D^0)} = 0.49 \pm 0.01 \quad \text{W.S./P.I.(-)}$$

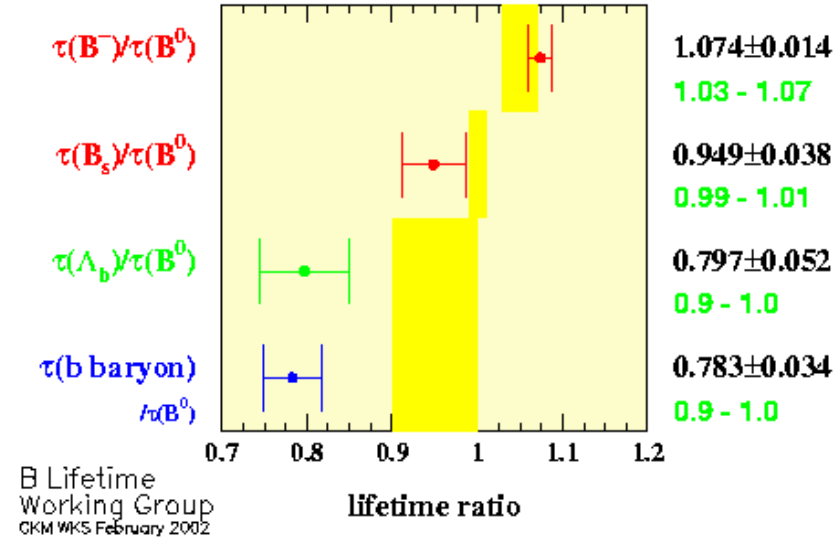
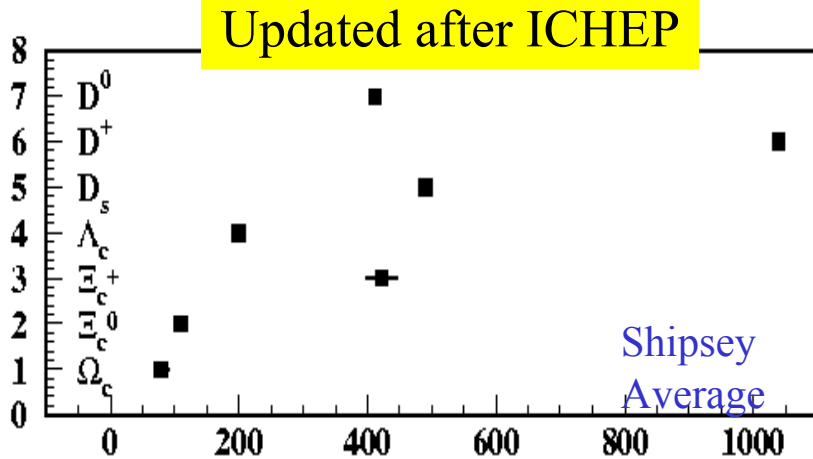
$$\frac{\tau(\Xi^+)}{\tau(\Lambda_c)} = 2.11 \pm 0.14 \quad \text{W.S.P.I.(±)}$$

To interpret this important to check
 $\Gamma(D_s \rightarrow eX)$
 $/\Gamma(D^0 \rightarrow eX)$
 But absolute
 $B(D_s \rightarrow eX)$
 is only known to 63%!

D^+ 6 ‰, D^0 3 ‰, D_s 2 ‰, Λ_c 2 ‰, Ξ^0 10 ‰, Ξ^+_c 6 ‰, Ω_c 15 ‰
 some lifetimes known as precisely as kaon lifetimes

Lifetimes span 1 order of magnitude

Compare Charm to Beauty Lifetimes



$$\frac{\Gamma_c(P.I.WA..S)}{\Gamma_c(spect)} \approx \frac{f_D^2 m_b^2}{f_B^2 m_c^2} \frac{\Gamma_b(P.I.WA..S)}{\Gamma_b(spect)} \approx 10$$

Plot from
 Stocchi ICHEP2002

Charm quarks are much more influenced by the hadronic environment than are beauty quarks

Very precisely determined lifetimes. The agreement with theory is still qualitative

Important message: errors on lifetimes are not a limiting factors in our ability to calculate absolute rates. The limiting factor is errors on absolute branching ratios.

D Nonleptonic Decays

Nonleptonic decays dominate the total rate

$$\begin{array}{l}
 D^+(c\bar{d}) : \tau_+ = 1042.7 \pm 6.9 \text{ fs} \\
 D^0(c\bar{u}) : \tau_0 = 410.5 \pm 1.5 \text{ fs}
 \end{array}
 \left. \vphantom{\begin{array}{l} D^+(c\bar{d}) \\ D^0(c\bar{u}) \end{array}} \right\} \tau_+ / \tau_0 \approx 2.5$$

Quarks or hadrons? ...in between

Compare to kaons and B-mesons:

$$\begin{array}{l}
 K^+(\bar{s}u) : \tau_+ = 12390 \pm 20 \text{ ps} \\
 K^0(\bar{s}d) : \tau_0 = 178.7 \pm 0.16 \text{ ps}
 \end{array}
 \left. \vphantom{\begin{array}{l} K^+(\bar{s}u) \\ K^0(\bar{s}d) \end{array}} \right\} \begin{array}{l} \tau_+ / \tau_0 \approx 70 \\ \text{Hadrons} \end{array}$$

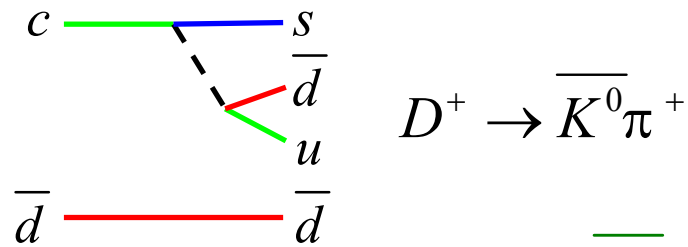
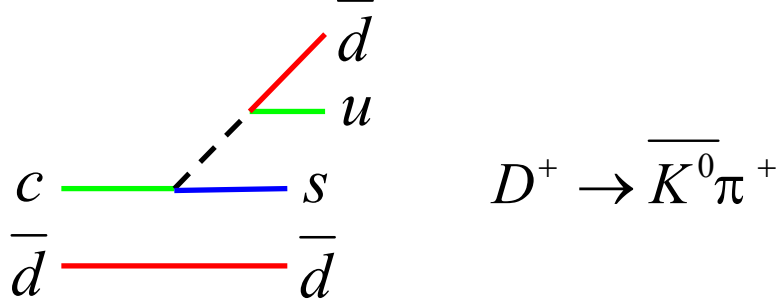
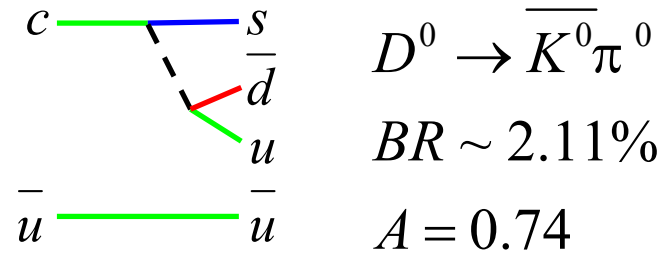
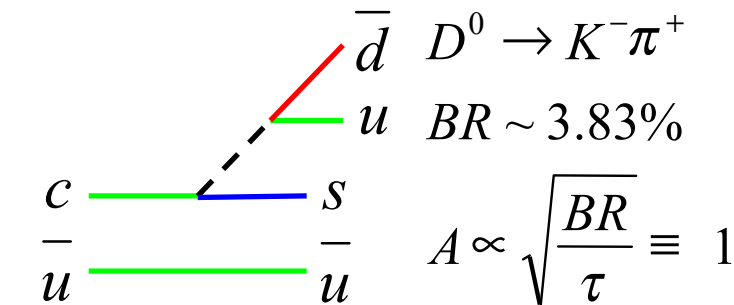
$$\begin{array}{l}
 B^+(\bar{b}u) : \tau_+ = 1655 \pm 24 \text{ fs} \\
 B^0(\bar{b}d) : \tau_0 = 1540 \pm 24 \text{ fs}
 \end{array}
 \left. \vphantom{\begin{array}{l} B^+(\bar{b}u) \\ B^0(\bar{b}d) \end{array}} \right\} \begin{array}{l} \tau_+ / \tau_0 \approx 1.07 \\ \text{quarks} \end{array}$$

The lifetime hierarchy (quark diagram level)

B decays : small BRs to 2- body final states (phase space)

2-body decays dominate D decays (multi-body decays found to be quasi 2 body) →

Is the $D^0 D^+$ lifetime hierarchy understandable in terms of 2 body hadronic decays?

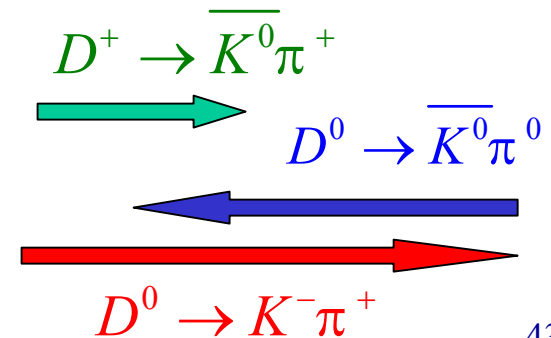


For the D^0 (D^+) the two states are distinct (identical)

100% Destructive interference predicts : $A = 1 - 0.74 = 0.26$

Measure: $A = 0.54$

Difference due to hadronic final state interactions



.....and at the hadronic level

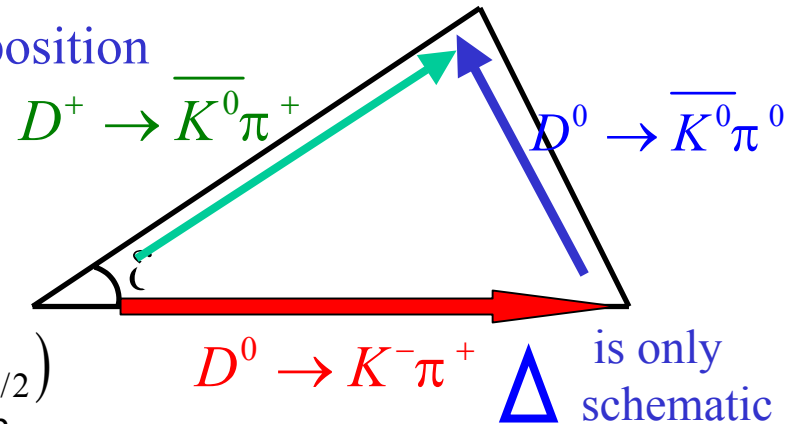
Simple factorization picture describes 2 body hadronic decays established for B's. For charm sizeable final state interactions are the norm.

$$A(D^0 \rightarrow K^- \pi^+) = \sqrt{\frac{2}{3}} A_{1/2} + \sqrt{\frac{1}{3}} A_{3/2} \quad \text{Isospin decomposition}$$

$$A(D^0 \rightarrow \overline{K}^0 \pi^0) = -\frac{1}{\sqrt{3}} A_{1/2} + \sqrt{\frac{2}{3}} A_{3/2}$$

$$A(D^+ \rightarrow \overline{K}^0 \pi^+) = \sqrt{3} A_{3/2} \quad A_I = A_I e^{i\delta}$$

$$|A_{1/2} + A_{3/2}|^2 = |A_{1/2}|^2 + |A_{3/2}|^2 + 2|A_{1/2}||A_{3/2}|\cos(\delta_{3/2} - \delta_{1/2})$$



measure: $|A(D^0 \rightarrow K^- \pi^+)|^2 + |A(D^0 \rightarrow \overline{K}^0 \pi^0)|^2 = |A_{1/2}|^2 + |A_{3/2}|^2$:extract

measure: $|A(D^+ \rightarrow \overline{K}^0 \pi^+)|^2 = 3|A_{3/2}|^2$:extract

find: $|A_{3/2}|/|A_{1/2}| = 0.37 \pm 0.03 \quad \delta = (\delta_2 - \delta_0) = 90^\circ \pm 7^\circ$ Rosner hep-ph/9903543

Many similar cases. Substantial modification of hadronic 2-body BR's due to FSI.

The presence of strong phases between amplitudes is an important ingredient in mixing studies and in CP violation

Charm branching ratios

$$\frac{Br}{\tau} = \Gamma$$

Most branching ratios in contrast to lifetimes are not well known

We have just seen that τ is measured very precisely

Key charm decay modes used to normalize B physics

		PDG (%)	Error(%)
D^0	$K^- \pi^+$	3.83 ± 0.09	2.3
D^+	$K^- \pi^+ \pi^+$	9.0 ± 0.6	6.7
D_s	$\phi \pi^+$	3.6 ± 0.9	25
Λ_c	$p K^- \pi^+$	$9.7 > \mathcal{B} > 3.0$	@90% c.l
J/ψ	$\mu^+ \mu^-$	5.88 ± 0.10	1.7

$$Br(D \rightarrow X) = \frac{\#X \text{ Observed}}{\text{efficiency} \times \#D's \text{ produced}}$$

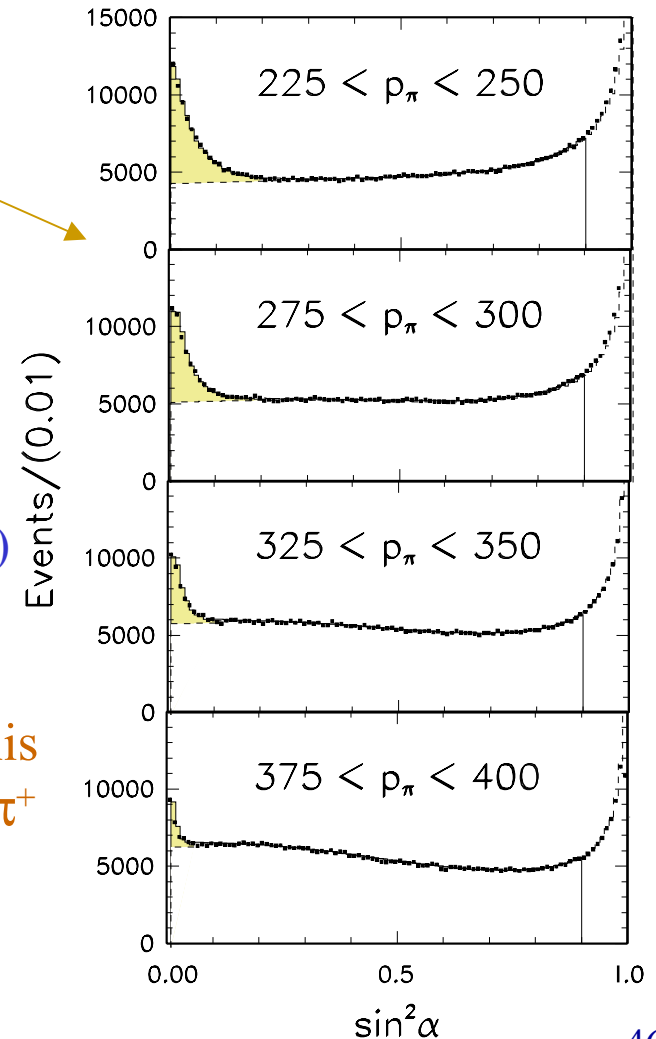
Because #D's produced is not well known

Measurement of $B(D^0 \rightarrow K^- \pi^+)$

B (%)	Error(%)	Source
$3.82 \pm 0.07 \pm 0.12$	3.6	CLEO
$3.82 \pm 0.09 \pm 0.12$	3.8	ALEPH
3.83 ± 0.09	2.3	PDG

- Method:
 - Detect $D^{*+} \rightarrow \pi^+ D^0$, $D^0 \rightarrow K^- \pi^+$ (CLEO & ALEPH Use same technique)
 - compare to: $D^{*+} \rightarrow \pi^+ D^0$, $D^0 \rightarrow$ unobserved
 - Problem: **Systematic error due to background extrapolation**
- α is \sphericalangle between thrust axis & slow π^+ (4 of 8 intervals shown)
- (Thrust is a measure of the direction of the primary quark pair in the event) shown)

$$D^{*+} \rightarrow D^0 \pi^+ \quad Q \sim 6 \text{ MeV}$$



B(D⁺ → K⁻π⁺π⁺)

B (%)	Error(%)	Source
9.3±0.6±0.8	10.8	CLEO
9.1±1.3±0.4	14.9	MKIII
9.1±0.7	7.7	PDG

From 9.6 pb⁻¹ (1984)

- Method (CLEO): Measure:

Assume this ratio is of Strong decays is given by isospin symmetry

$$\frac{B(D^{*+} \rightarrow D^0 \pi^+)}{B(D^{*+} \rightarrow D^+ \pi^0)} \frac{B(D^0 \rightarrow K^- \pi^+)}{B(D^+ \rightarrow K^- \pi^+ \pi^+)}$$

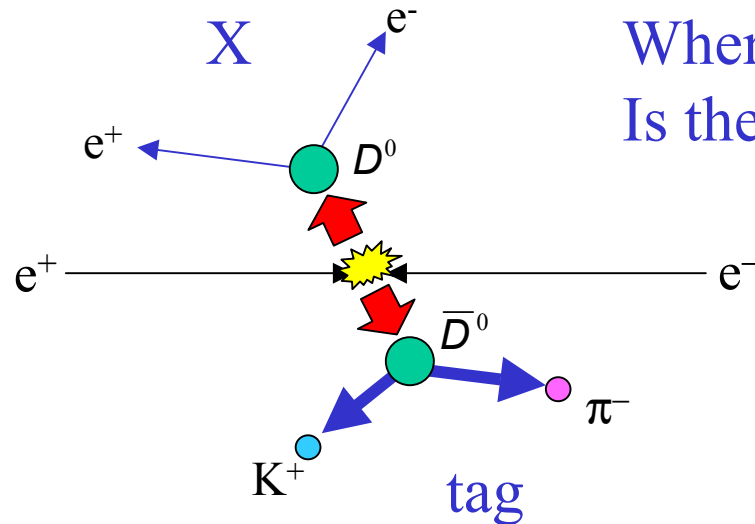
(this bootstrap method can never yield a measurement of B(D⁺ → K⁻π⁺π⁺) more accurate than B(D⁰ → K⁻π⁺)

- Method (MKIII): ψ'' → D⁺D⁻ full reconstruction, limited by size of data sample
 The determination of B(D⁺ → φπ⁺), which has a 25% error also bootstraps on B(D⁰ → K⁻π⁺)

How can we do better? Recall: Absolute Charm Branching Ratios at Threshold

$\psi(3770) \rightarrow DD$

$$Br(D \rightarrow X) = \frac{\#X \text{ Observed}}{\text{efficiency} \times \#D\text{'s produced}}$$



Where the # of D's produced
 Is the # of tags

Unique Opportunities at Charm Thresholds

- Unique event properties
 - Only $D\bar{D}$ not $D\bar{D}X$ produced
 - Can get $D^0\bar{D}^0$, D^+D^- , $D_s^+D_s^-$, $\Lambda_c^+\Lambda_c^-$
 - Probably other charmed baryons as well (not yet measured)

- Large cross sections

$$\sigma(D^0\bar{D}^0) = 5.8 \text{ nb}$$

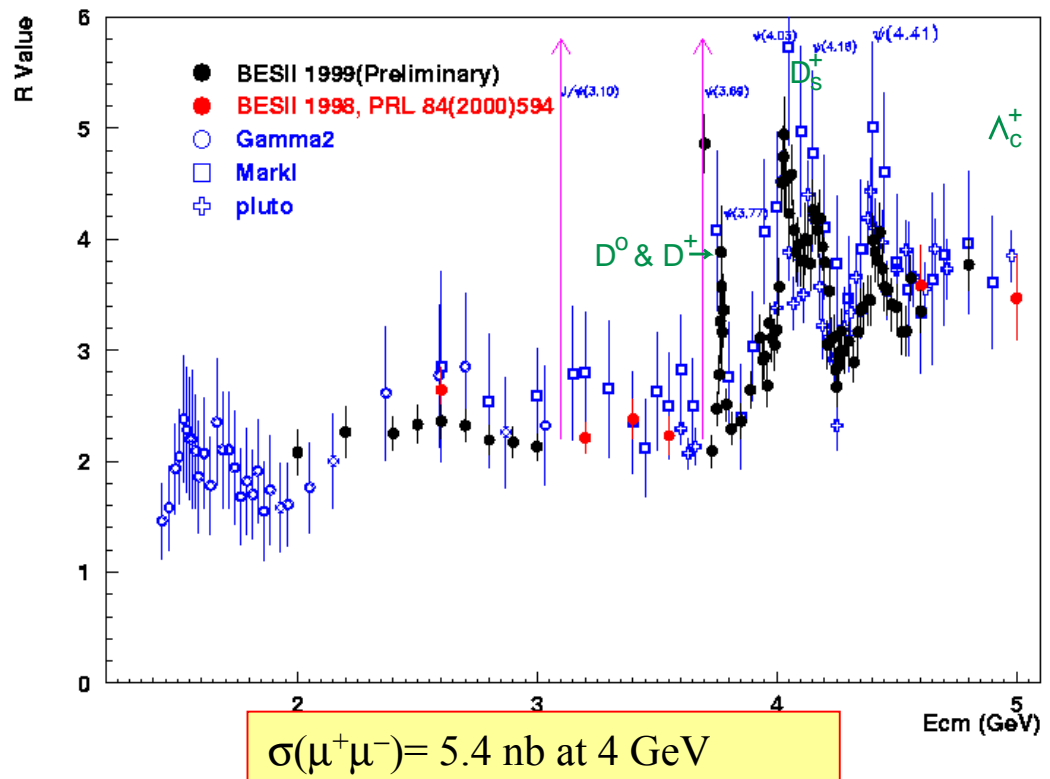
$$\sigma(D^+D^-) = 4.2 \text{ nb}$$

$$\sigma(D_s^+D_s^-) = 0.5 \text{ nb}$$

$$\psi(3770) \rightarrow D\bar{D}$$

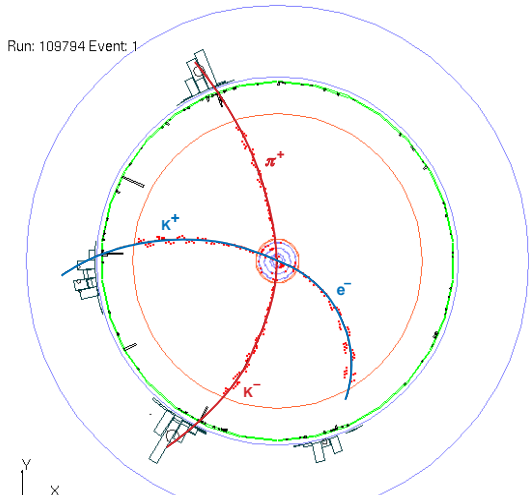
$$\sqrt{s} \sim 4140 \rightarrow D_s^+D_s^-$$

R (units of $\sigma(\mu^+\mu^-)$)



$\psi(3770)$ events are simple

$\psi(3770)$ event:



- Charm events produced at threshold are extremely clean
- Large σ , low multiplicity
- Pure initial state: no fragmentation
- Signal/Background is optimum at threshold

- Double tag events are pristine
 - These events are key to making absolute Br measurements
- Neutrino reconstruction is clean
- Quantum coherence aids D mixing & CP violation studies

precision
 flavor
 physics

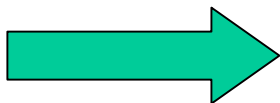
new
 physics

But: D 's don't move

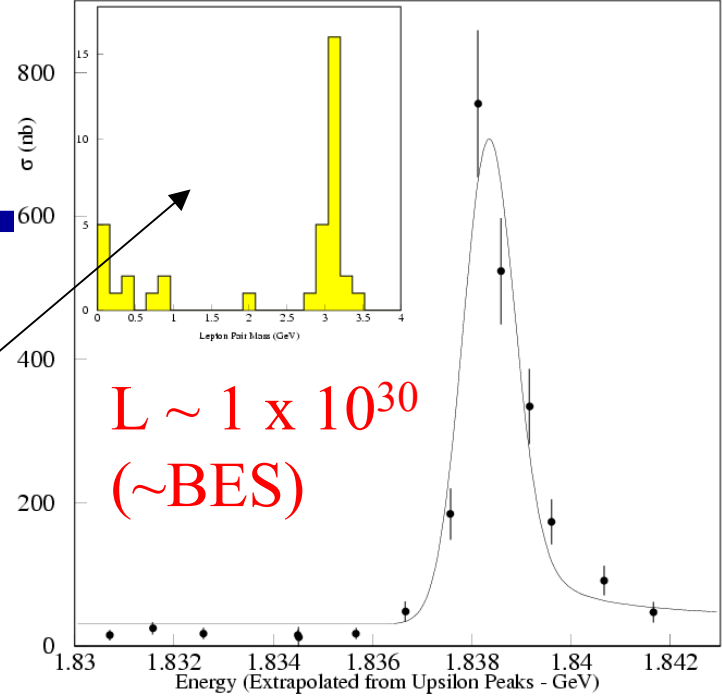
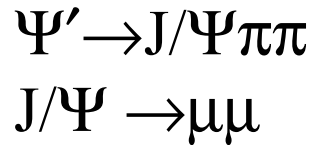
CESR

$L(@Y(4S)) = 1.3 \times 10^{33}$

- CESR-c:
 - Modify for low energy operation:
 - add wigglers for transverse cooling
- Expected machine performance:



One day scan of the Ψ' :
(1/29/02)



\sqrt{s}	$L (10^{32} \text{ cm}^{-2} \text{ s}^{-1})$
3.1 GeV	2.0
3.77 GeV	3.0
4.1 GeV	3.6

• $\Delta E_{\text{beam}} \sim 1.2 \text{ MeV at } J/\psi$

CLEO-c Proposed Run Plan

2002: Prologue: Upsilon's $\sim 1\text{-}2 \text{ fb}^{-1}$ each at $Y(1S), Y(2S), Y(3S), \dots$
 Spectroscopy, matrix element, Γ_{ee}, η_B, h_b
 10-20 times the existing world's data (Fall 2001- Fall 2002)

2003: $\psi(3770) - 3 \text{ fb}^{-1}$
 30 million DD events, 6 million *tagged* D decays
 (310 times MARK III)

2004: $\sqrt{S} \sim 4140 \text{ MeV} - 3 \text{ fb}^{-1}$
 1.5 million $D_s D_s$ events, 0.3 million *tagged* D_s decays
 (480 times MARK III, 130 times BES)

2005: $\psi(3100), 1 \text{ fb}^{-1} - 1 \text{ Billion } J/\psi \text{ decays}$
 (170 times MARK III, 20 times BES II)

C
L
E
O
-
c

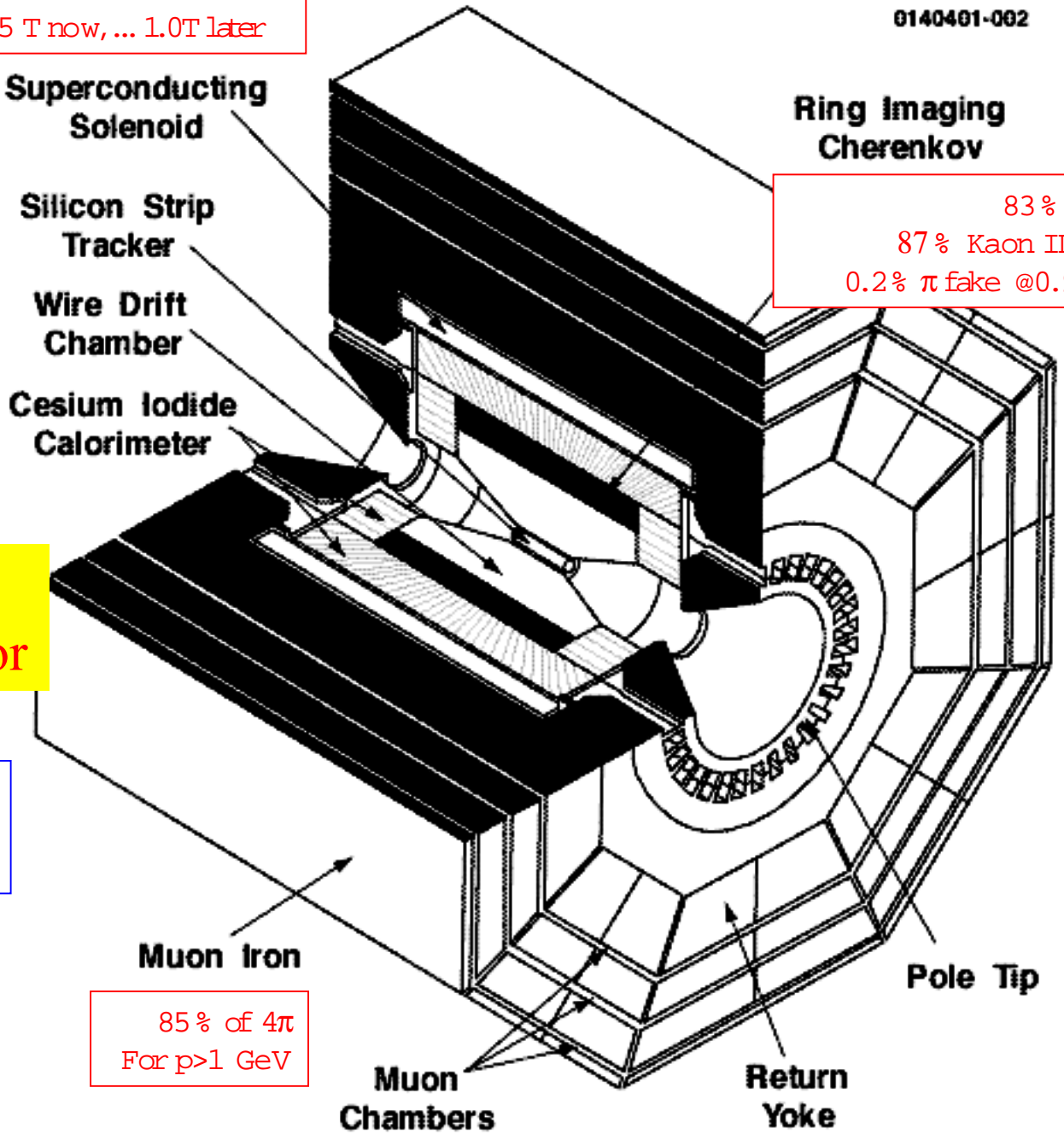
A 3 year
program

1.5 T now, ... 1.0T later

93% of 4π
 $\sigma_p/p = 0.35\%$ @1GeV
 $dE/dx: 5.7\%$ π @minI

93% of 4π
 $\sigma_E/E = 2\%$ @1GeV
 $= 4\%$ @100MeV

83% of 4π
87% Kaon ID with
0.2% π fake @0.9GeV



**CLEO III Detector
→ CLEO-c Detector**

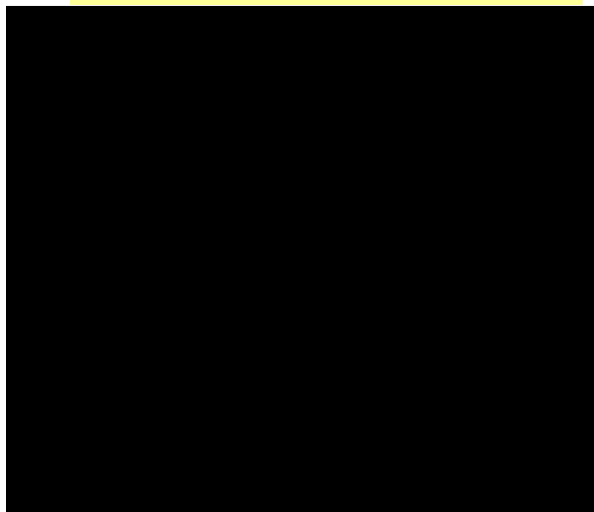
Trigger: Tracks & Showers
Pipelined
Latency = 2.5 μ s

Data Acquisition:
Event size = 25kB
Thruput < 6 MB/s

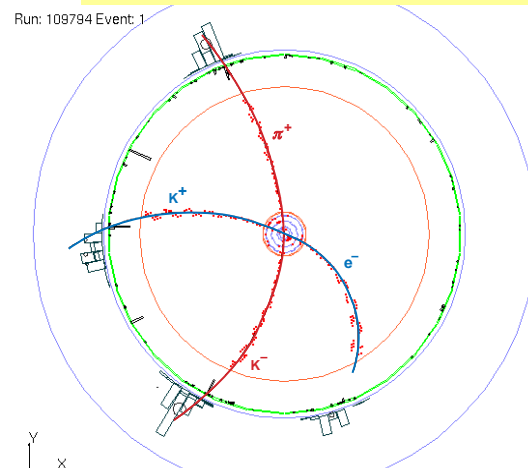
85% of 4π
For $p > 1$ GeV

$\psi(3770)$ events: simpler than $Y(4S)$ events

$Y(4S)$



$\psi(3770)$ event:



- * CLEO III state of the art detector, well understood
- * CLEO-c Replace Si \rightarrow low mass drift chamber (under construction)
- * The demands of doing physics at 3-5 GeV are easily met by the existing detector.

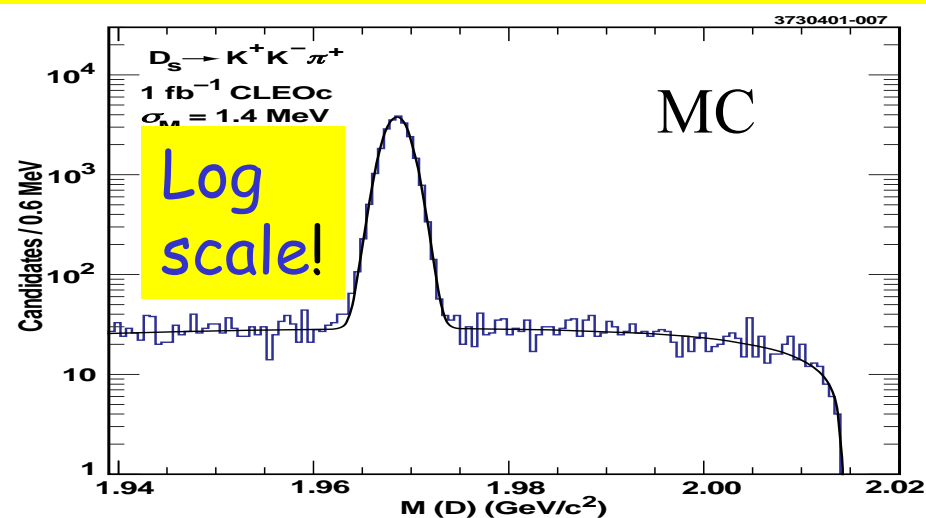
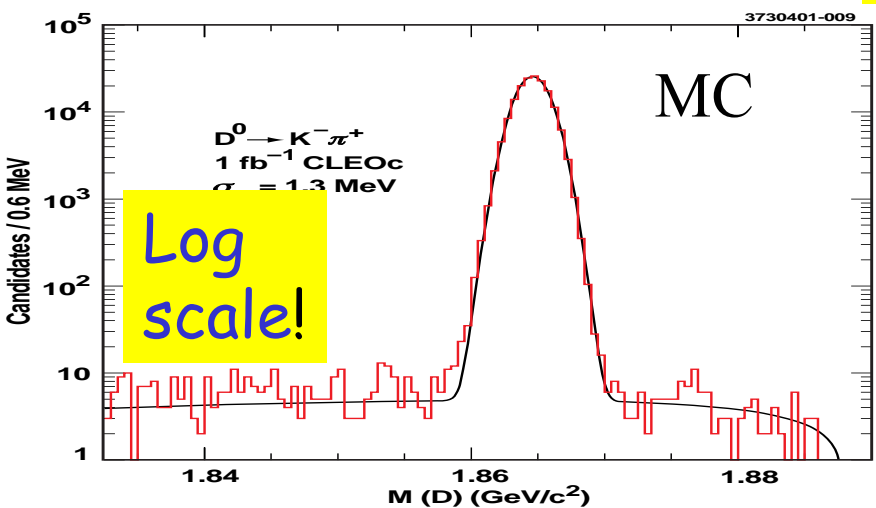
Tagging Technique, Tag Purity @ Threshold CLEO-c simulation

- $\psi(3770) \rightarrow DD$ $\sqrt{s} \sim 4140 \rightarrow D_s D_s$
- Charm mesons have many large branching ratios (~1-15%)
- low multiplicity: high reconstruction efficiency
- high net tagging efficiency ~20%!

In 1 year
 At Each \sqrt{s}

Anticipate 6M D tags 300K D_s tags:

$D \rightarrow K\pi$ tag. S/B ~5000/1! $D_s \rightarrow \phi\pi$ ($\phi \rightarrow KK$) tag. S/B ~100/1

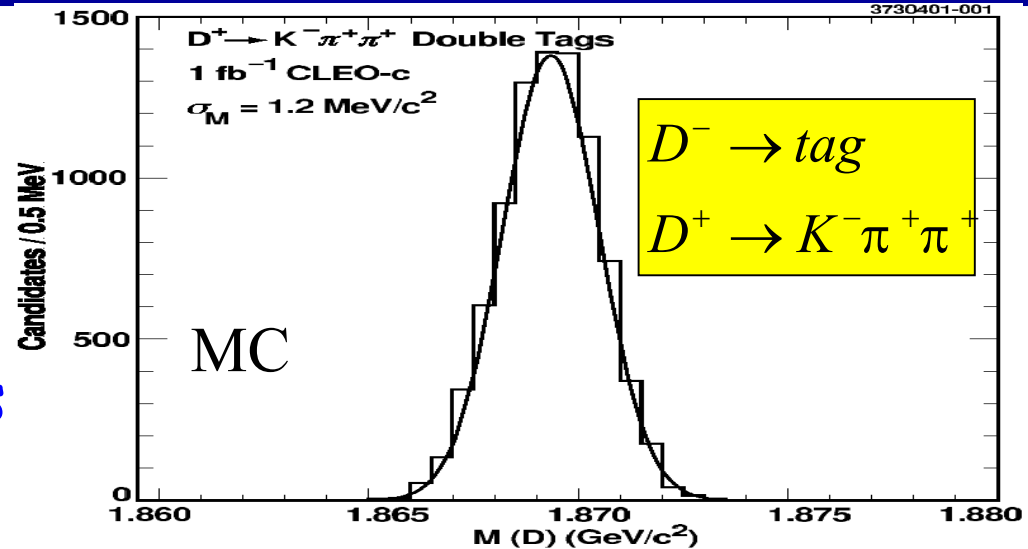


Beam constrained mass

Absolute Branching Ratios

~ Zero background in hadronic tag modes

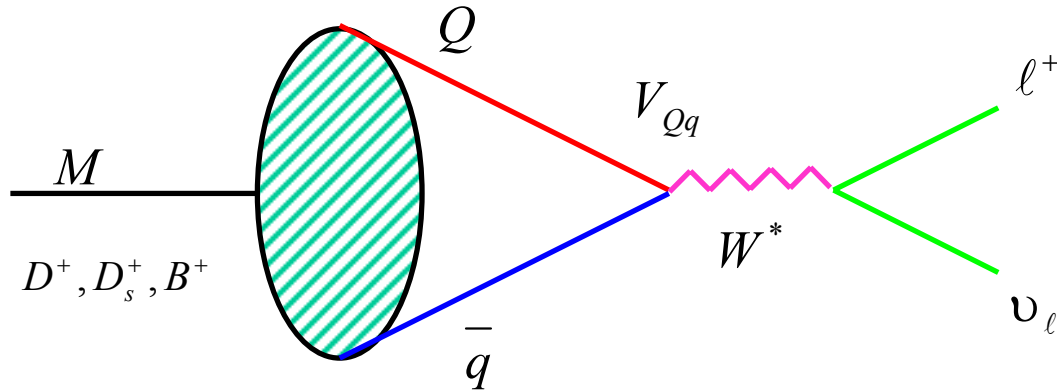
Measure absolute
 Br ($D \rightarrow X$) with double tags
 Br = # of X / # of D tags



Decay	\sqrt{s}	L fb^{-1}	Double tags	PDG ($\delta B/B$ %)	CLEOc ($\delta B/B$ %)
$D^0 \rightarrow K^- \pi^+$	3770	3	53,000	2.4	0.6
$D^+ \rightarrow K^- \pi^+ \pi^+$	3770	3	60,000	7.2	0.7
$D_s \rightarrow \phi \pi$	4140	3	6,000	25	1.9

CLEO-c potential: set the absolute scale for all heavy quark measurements

Leptonic Decays \rightarrow Decay Constants



$$q^2 = m_W^2 = M^2$$

Fixed

$$M = \frac{G_F}{\sqrt{2}} V_{Qq} \langle 0 | J_\mu | M \rangle \bar{u}(k, \sigma) \gamma^\mu (1 - \gamma_5) v(p, s) \quad (\text{Pseudoscalar Meson})$$

the meson decay constant f_M measures the probability for the Q and \bar{q} to have zero separation the annihilation probability is \propto to wave function overlap

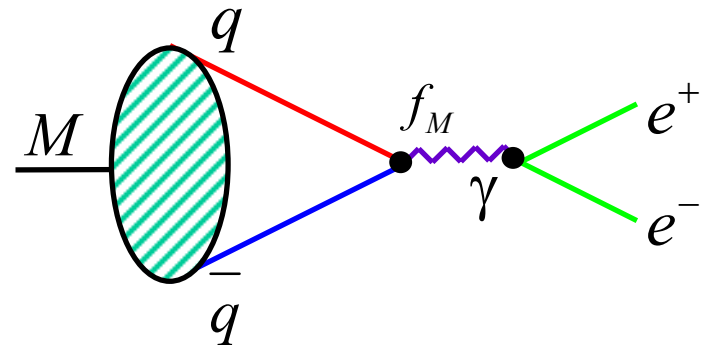
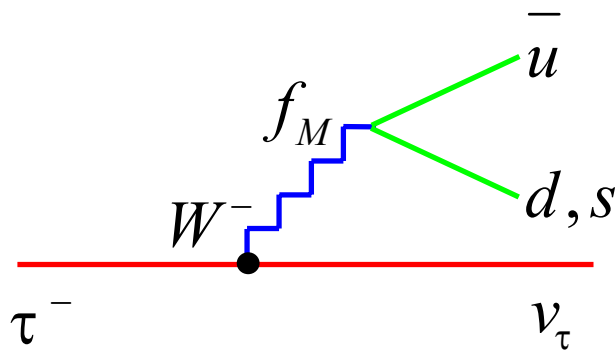
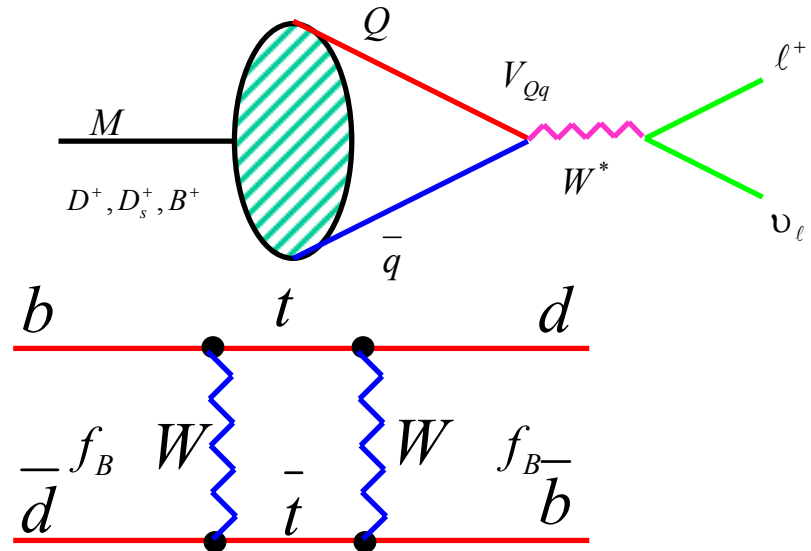
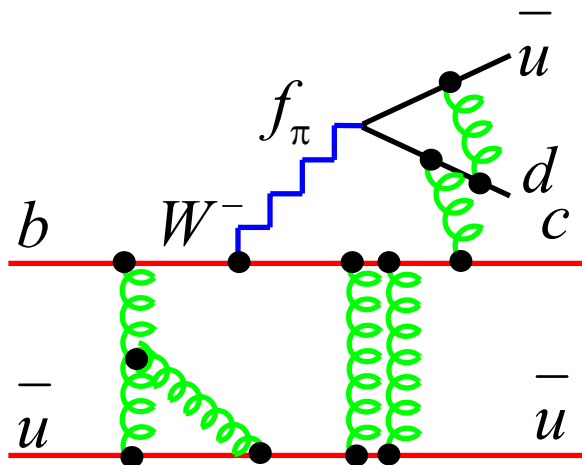
$$\langle 0 | \bar{q} \gamma_\mu \gamma_5 Q | P(p) \rangle = i f_M p_\mu$$

$$f_M^2 M = 12 |\Psi(0)|^2$$

(For a meson with two heavy quarks) (Rosner)

$$\Gamma(M_{Qq} \rightarrow \ell^- \bar{\nu}) = \frac{G_F^2}{8\pi} |V_{qQ}|^2 f_M^2 M m_\ell^2 \left(1 - \frac{m_\ell^2}{M^2} \right)^2$$

Decay constants are important in many processes



$$\Gamma(M_{Qq} \rightarrow \ell^- \bar{\nu}) = \frac{G_F^2}{8\pi} |V_{qQ}|^2 f_M^2 M m_\ell^2 \left(1 - \frac{m_\ell^2}{M^2}\right)^2$$

Helicity suppression

Decay is forbidden as $m_l \rightarrow 0$

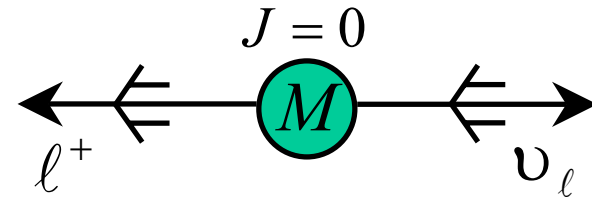
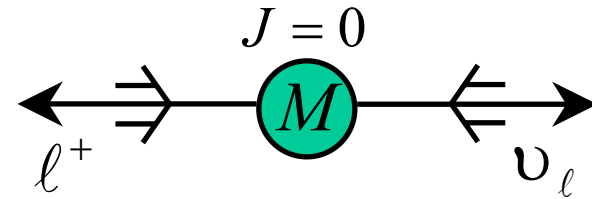
$$\frac{\Gamma(\pi^+ \rightarrow e^+ \nu_e)}{\Gamma(\pi^+ \rightarrow \mu^+ \nu_\mu)} \approx 10^{-4}$$

$$\Gamma(D_s^+ \rightarrow e^+ \nu_e) : \Gamma(D_s^+ \rightarrow \mu^+ \nu_\mu) : \Gamma(D_s^+ \rightarrow \tau^+ \nu_\tau) \approx 10^{-5} : 1 : 10$$

$$\Gamma(D^+ \rightarrow \ell^+ \nu_\ell) \propto |V_{cd}|^2 \approx (0.22)^2$$

$$\Gamma(D_s^+ \rightarrow \ell^+ \nu_\ell) \propto |V_{cs}|^2 \approx (0.97)^2$$

$$\Gamma(B^+ \rightarrow \ell^+ \nu_\ell) \propto |V_{ub}|^2 \approx (0.003)^2$$



Estimate of the leptonic Br's using $f_{B_s} = f_{B_c} = 200 \text{ MeV}$
 $f_{D_s} = 260 \text{ MeV}, f_D = 220 \text{ MeV}$

	B(ev)	B($\mu\nu$)	B($\tau\nu$)
D^+	8.2×10^{-9}	4.2×10^{-4}	1.1×10^{-3}
D_s^+	7.5×10^{-8}	5.7×10^{-3}	5.5×10^{-2}
B^+	7.5×10^{-12}	3.2×10^{-7}	7.1×10^{-5}
π^+	1.2×10^{-4}	99.99%	
K^+	1.6×10^{-5}	63.5%	

At first sight it is remarkable that : $B(D_s^+ \rightarrow \mu^+ \nu_\mu) \lll B(K^+ \rightarrow \mu^+ \nu_\mu)$

While $(f_M^2 M) \rightarrow \text{constant}$

$\Gamma(\text{total}) \propto M^5$ so leptonic branching ratio becomes smaller as $M \uparrow$

If we compare *rates*

instead of branching ratios:

The leptonic *rate* is higher for the D_s than for the K^+

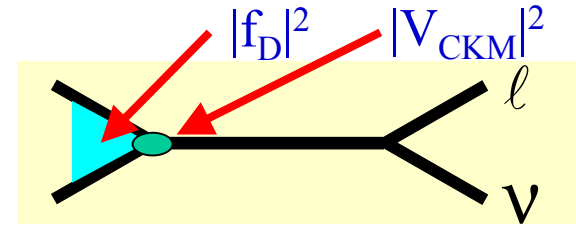
$$\Gamma(K^+ \rightarrow \mu^+ \nu_\mu) = 5.13 \times 10^7 s^{-1}$$

$$\Gamma(D_s^+ \rightarrow \mu^+ \nu_\mu) = 6.9 \times 10^9 s^{-1}$$

$$\Gamma(D_s^+ \rightarrow \tau^+ \nu_\mu) = 6.6 \times 10^{10} s^{-1}$$

D meson Decay Constants

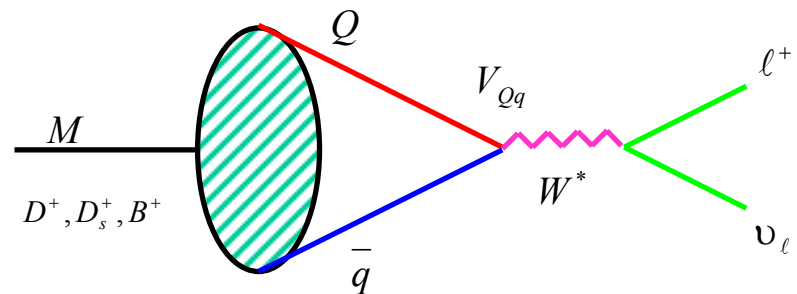
In a pseudoscalar D meson decay:
c and **q** annihilate



$$\Gamma(D_q^+ \rightarrow \ell \nu) = \frac{1}{8\pi} G_F^2 M_{D_q} m_\ell^2 \left(1 - \frac{m_\ell^2}{M_{D_q}^2}\right) f_{D_q}^2 |V_{cq}|^2$$

$$B(D^+ \rightarrow \ell \nu) / \tau_{D^+} : f_{D^+} |V_{cd}|$$

$$B(D_S \rightarrow \ell \nu) / \tau_{D_S} : f_{D_S} |V_{cs}|$$



* Charm meson lifetimes known 0.3-2%

* 3 generation unitarity

Vcs, (Vcd) known to 0.1% (1.1%) → f_{D+} f_{Ds}

Example: f_{D_s} near/at the $Y(4S)$

Signal is a single muon, or single muon + photon tag very difficult at a hadron machine

–Search for $D_s^* \rightarrow D_s \gamma$, $D_s \rightarrow \mu\nu$

–Directly detect γ , μ , Use hermeticity of detector to reconstruct ν

–Plot mass difference but Backgrounds are LARGE!

•Use $D_s \rightarrow e\nu$ (rate \sim 0) for bkgd determination but precision limited by systematics

•Compare rate to $D_s \rightarrow \phi\pi$, but $\text{Br}(D_s \rightarrow \phi\pi)$ not well known-25% error!.

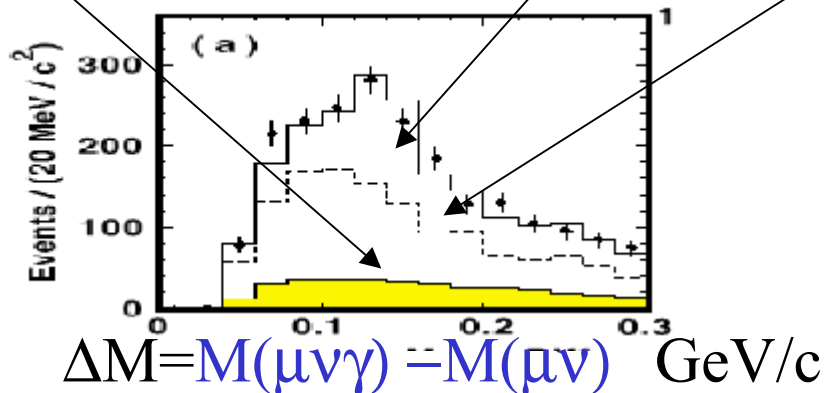
–FDs Error \sim 17% now (CLEO)

CLEO signal 4.8fb^{-1}

Excess of μ over e fakes

Background measured with electrons

$D_s \rightarrow \mu\nu$

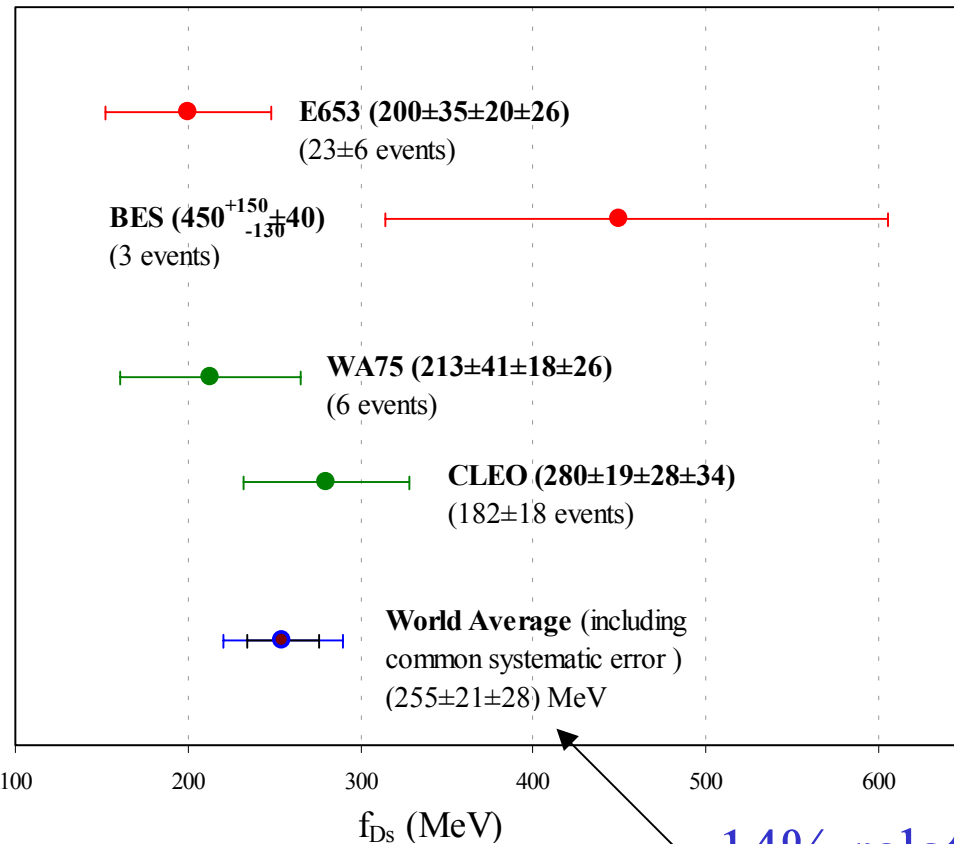


$$f_{D_s} = 280 \pm 19 \pm 28 \pm 34 \text{ (MeV)}$$

D meson Decay Constants Current Status

$f_{D^+} < 290 \text{ MeV}$ @ 90% CL
 (Mark III)

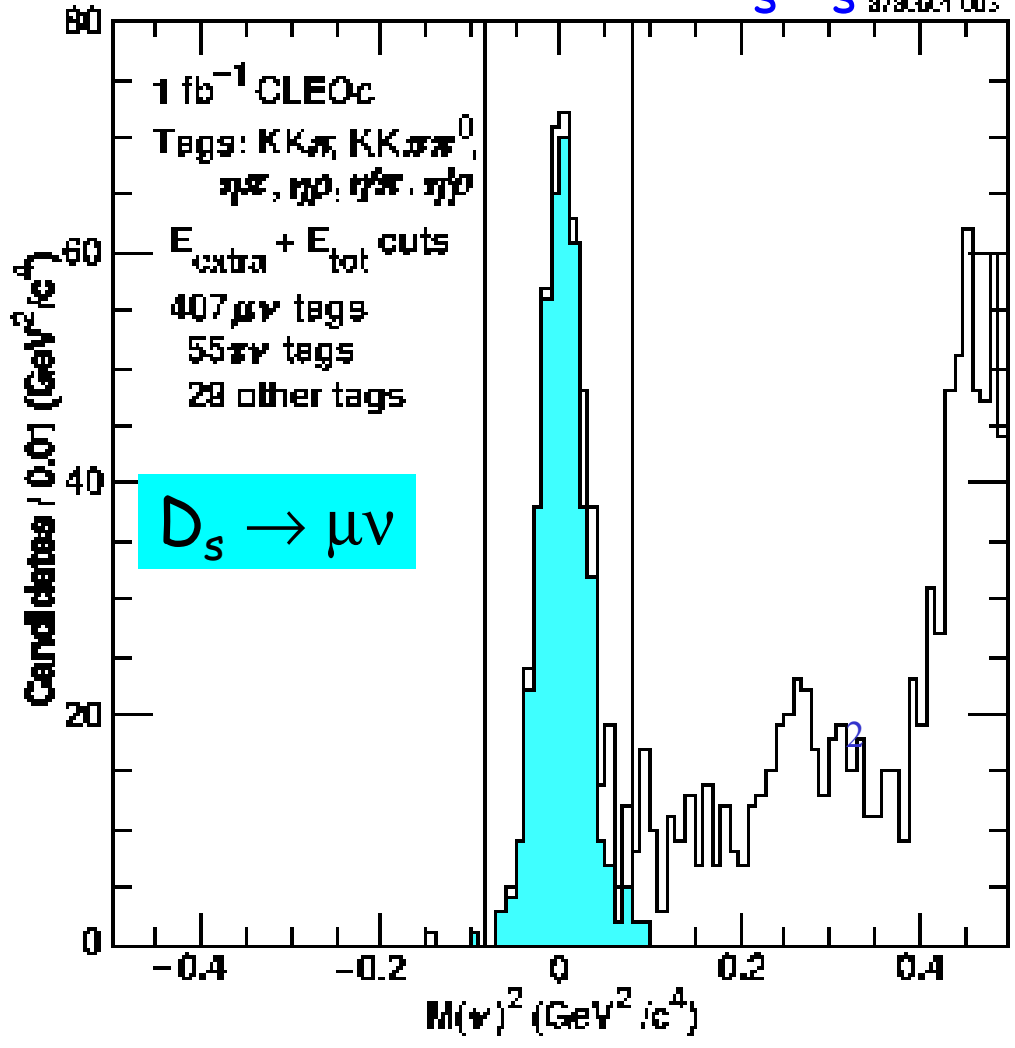
f_{D_s} has been measured by several groups, using $D_s \rightarrow \mu\nu$. There are also measurements from LEP using $D_s \rightarrow \tau^+\nu$ which I have not included in the Table or average. (Inclusion of these extra modes requires the assumption of Lepton universality, which might be interesting to test. Note large correlated common systematic error from $B(D_s \rightarrow \phi\pi)$



14% relative error

Decay Constant at Threshold (CLEO-c simulation)

$\sqrt{s} \sim 4140 \rightarrow D_s D_s$



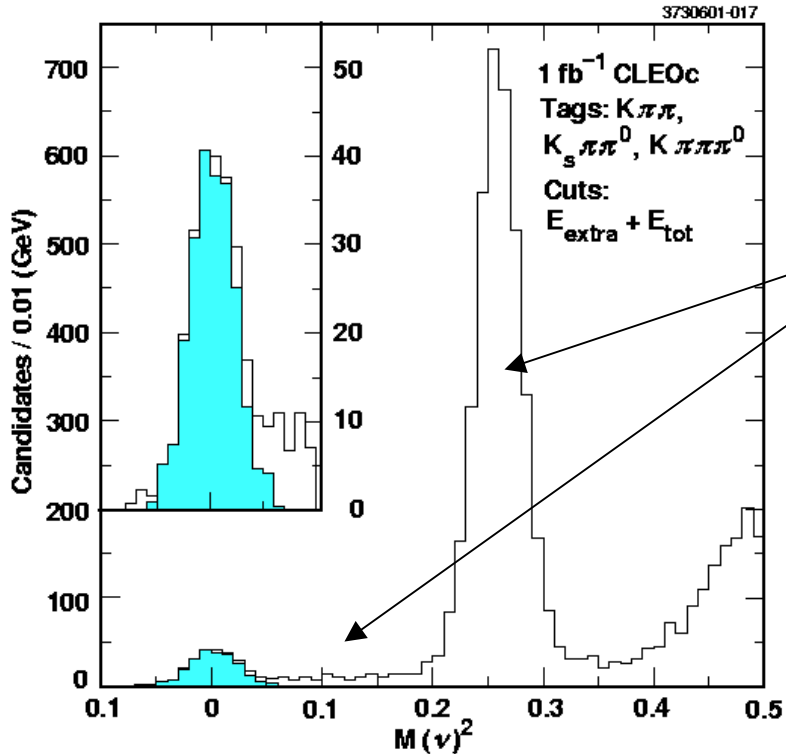
- Fully reconstruct 1 D “the tag”
- Require one additional charged track and no additional photons
- Compute MM^2 Peaks at zero for $D \rightarrow \mu + \nu$ decay.
 - No need to identify muon-helpers systematic error
 - Can identify electrons to check background level
 - Expect resolution of $\sim M_{\pi^0}$

$\frac{\delta f_{D_s}}{f_{D_s}} \approx 1.7\%$ (Now: $\pm 14\%$)

Decay Constant at Threshold (CLEO-c simulation)

$\psi(3770) \rightarrow DD$

CLEO-c
 simulation



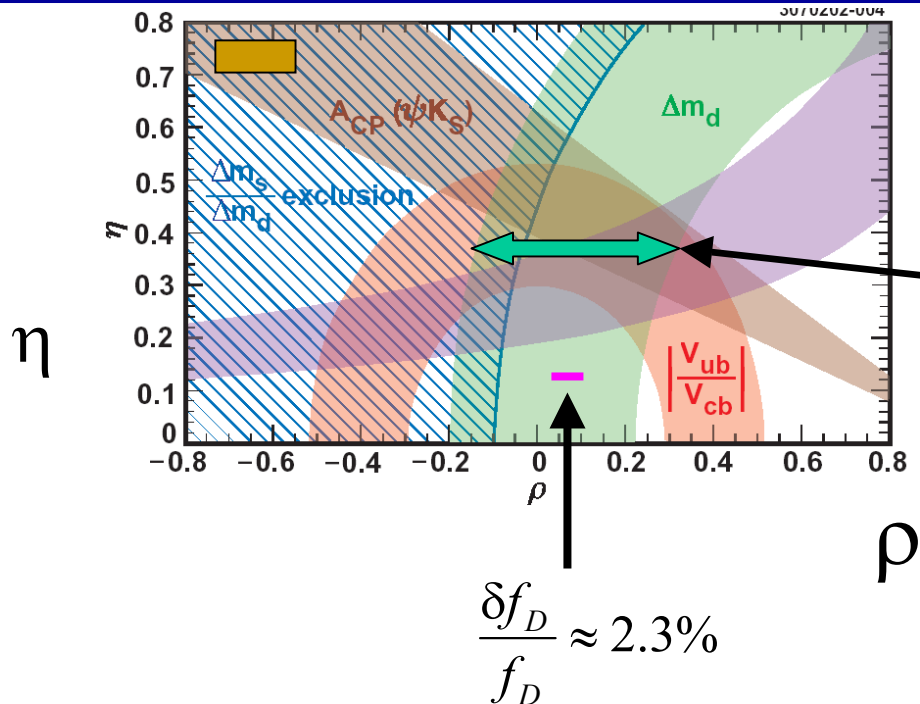
D⁺ → K_Lμν
 D⁺ → τν

$$\frac{\delta f_D}{f_D} \approx 2.3\%$$

(1 year)

Now: upper limit
 exists

Improved knowledge of the decay constant yields precision determination of V_{td}



$$\Delta M_d = 0.50 ps^{-1} \left[\frac{\sqrt{B_{B_d}} f_{B_d}}{200 MeV} \right]^2 \left[\frac{|V_{td}|}{8.8 \times 10^{-3}} \right]^2$$

$$\frac{\sigma(\rho)}{\rho} = 0.5 \frac{\sigma(\Delta M_d)}{\Delta M_d} \oplus \frac{\sigma(f_B \sqrt{B_{B_d}})}{f_B \sqrt{B_{B_d}}}$$

(ICHEP02) 1.2%

~15% (LQCD)

- Lattice predicts f_B/f_D with small errors
- precision measurement of f_D
- precision estimates of f_B
- precision determination of V_{td}



Additional Slides

Summary of Decay Constant Reach at CLEO-c

Branching
 Ratio →

Decay
 Constant ↓

Reaction	Signal	$\tau\nu/\mu\nu$	Bkgd	$\delta B/B$
$D_s^+ \rightarrow \mu\nu$	1221	165	87	3.2%
$D_s^+ \rightarrow \tau\nu$	1740	0	114	2.4%
$D^+ \rightarrow \mu\nu$	672	30	60	3.8%

	Reaction	$\frac{1}{2} \Delta B/B$	$\frac{1}{2} \Delta \tau/\tau$	$\Delta V_{cq}/V_{cq}$	CLEO-c $\delta f/f$	PDG $\delta f/f$
f_{D_s}	$D_s^+ \rightarrow \mu\nu$	1.6%	1%	0.1%	1.9%	14%
f_{D_s}	$D_s^+ \rightarrow \tau\nu$	1.2%	1%	0.1%	1.6%	33%
f_{D^+}	$D^+ \rightarrow \mu\nu$	1.9%	0.6%	1.1%	2.3%	UL

(not updated for improved lifetimes)

D_s Meson Lifetime

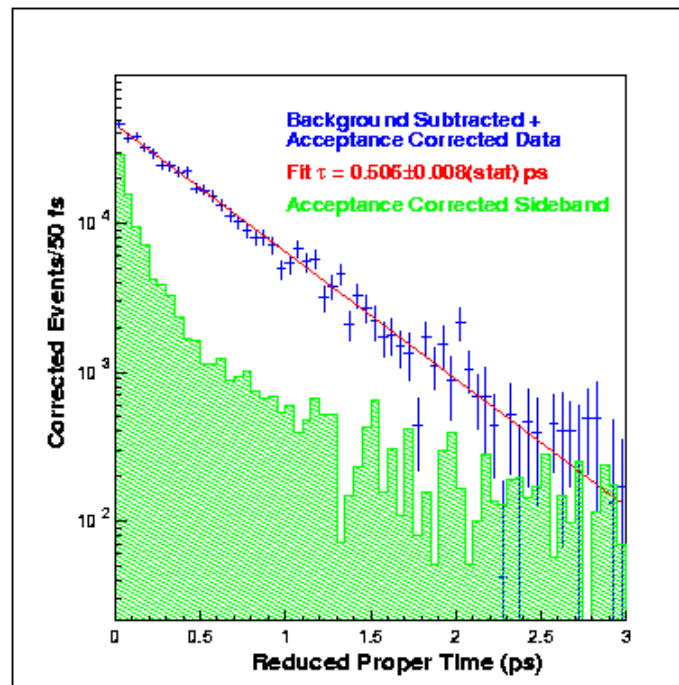
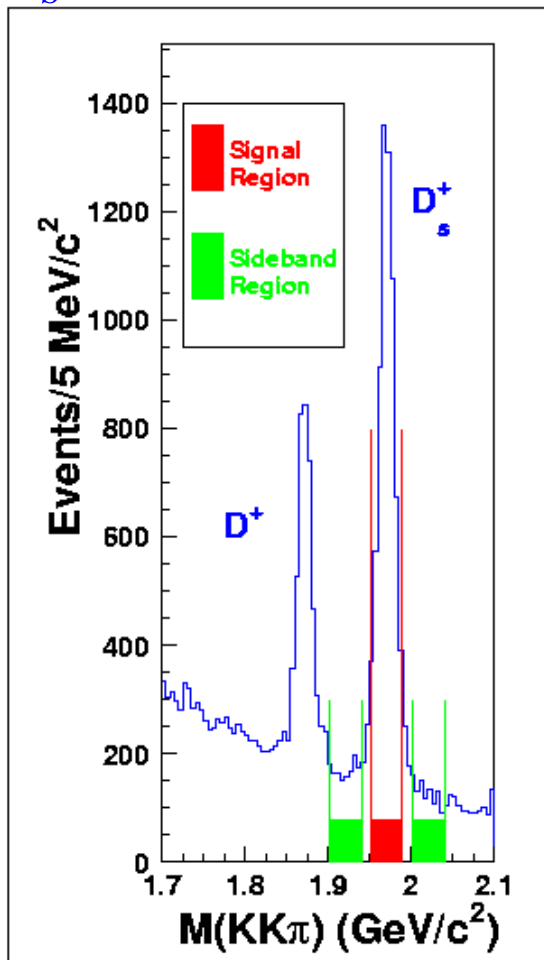
New at ICHEP

$D_s \rightarrow \phi\pi$ signal

Lifetime 506 ± 8 fs

PDG 2002 490 ± 9 fs

5668 ± 95 events (50% FOCUS data)



Preliminary

$$\frac{\tau(D_s)}{\tau(D^0)} = 1.23 \pm 0.02$$

To interpret this important to check $\Gamma(D_s \rightarrow eX)$
But absolute $dBR/BR = 63\%$!

Theoretical prediction (Bigi Uraltsev)

1.00-1.07 (no WA)

0.8-1.27 (different process interference)

4 x statistics including

Lifetime Summary

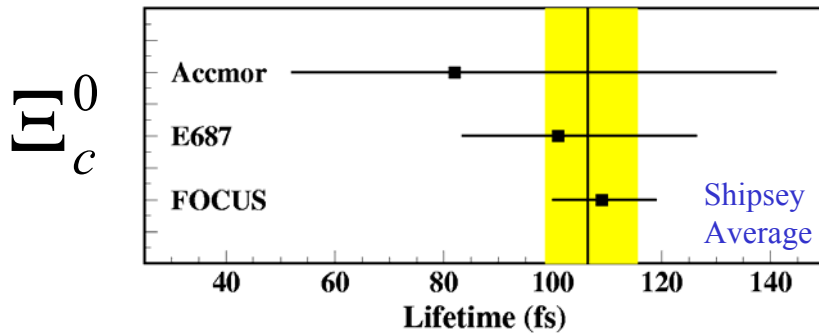
We know the charm meson lifetimes with extraordinary precision, in the best cases $<1/2\%$ (major improvement in the past year: FOCUS)
Non spectator effects are similar in size to the spectator contributions
the lifetime hierarchy is consistent with the OPE formalism
but debatable if OPE should apply to c-quark (mass)

More stringent tests of this idea would be provided if precise absolute semileptonic branching ratios of $D_s, \Lambda_c, \Xi^0, \Xi_c^+, \Omega_c$ were known

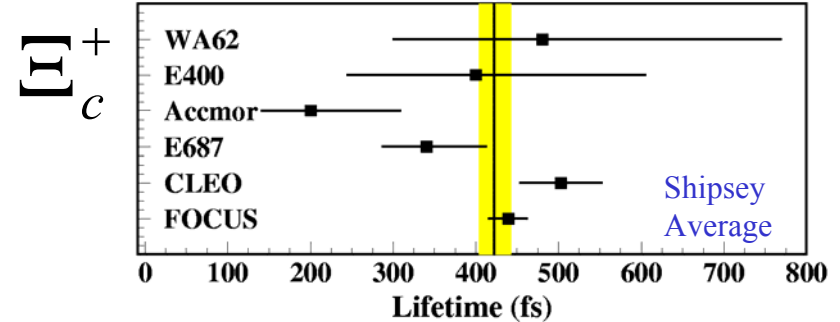
Charmed Baryon Lifetimes

- Unlike charmed mesons, decays of **charmed baryons** are not color or helicity suppressed, this results in a reduced lifetime relative to

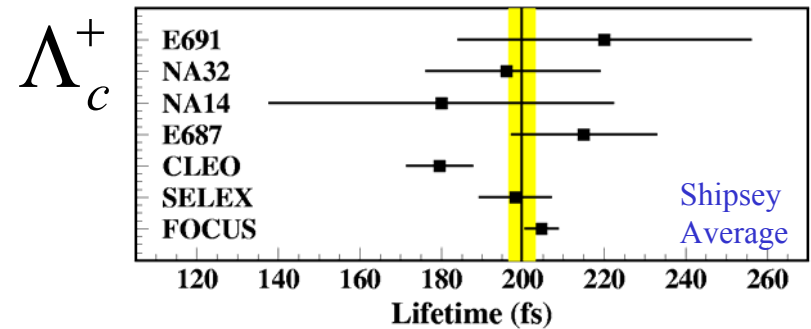
$$\tau_{\text{average}}(\Xi_c^0) = 106_{-8}^{+9} \text{ fs}$$



$$\tau_{\text{average}}(\Xi_c^+) = 422_{-19}^{+20} \text{ fs}$$



$$\tau_{\text{average}}(\Lambda_c^+) = 200 \pm 3 \text{ fs}$$

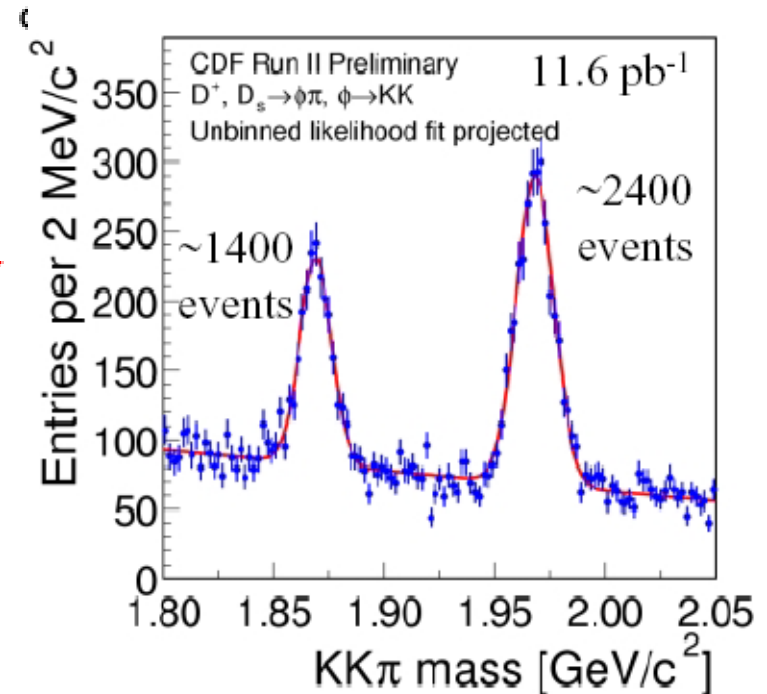


$$\Gamma(\Xi_c^+) < \Gamma(\Lambda_c^0) < \Gamma(\Xi_c^0) \sim \Gamma(\Omega_c^0)$$

P.I.(+/-) W.S.+P.I.(-) W.S.+P.I.(+) (10/3)P.I.(+)

❖ $D_s^\pm - D^\pm$ mass difference

- Both $D \rightarrow \phi\pi$ ($\phi \rightarrow KK$)
- $\Delta m = 99.28 \pm 0.43 \pm 0.27$ MeV
 - PDG: 99.2 ± 0.5 MeV
(CLEO2, E691)
- Systematics dominated by background modeling



D Hadronic decays

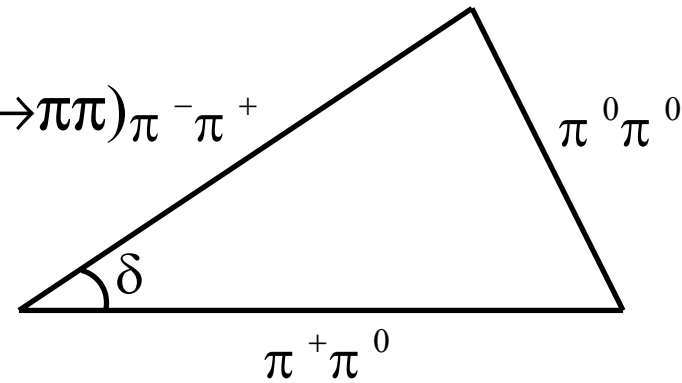
Simple factorization picture describes 2 body hadronic decays established for B's. For charm sizeable final state interactions are the norm.

Isospin decomposition (same as $B \rightarrow \pi\pi$, $K \rightarrow \pi\pi$) $\pi^- \pi^+$

$$A(D^0 \rightarrow \pi^- \pi^+) = \frac{1}{\sqrt{3}} (\sqrt{2} A_0 + A_2)$$

$$A(D^0 \rightarrow \pi^0 \pi^0) = \frac{1}{\sqrt{3}} (-A_0 + \sqrt{2} A_2)$$

$$A(D^+ \rightarrow \pi^0 \pi^+) = \sqrt{3/2} A_2$$



$$A_I = A_I e^{i\delta}$$

$$|A_0 + A_2|^2 = |A_0|^2 + |A_2|^2 + 2|A_0||A_2|\cos(\delta_2 - \delta_0)$$

measure

extract

$$|A(D^0 \rightarrow \pi^- \pi^+)|^2 + |A(D^0 \rightarrow \pi^0 \pi^0)|^2 = |A_0|^2 + |A_2|^2$$

$$|A(D^+ \rightarrow \pi^0 \pi^+)|^2 = 3/2 |A_2|^2$$

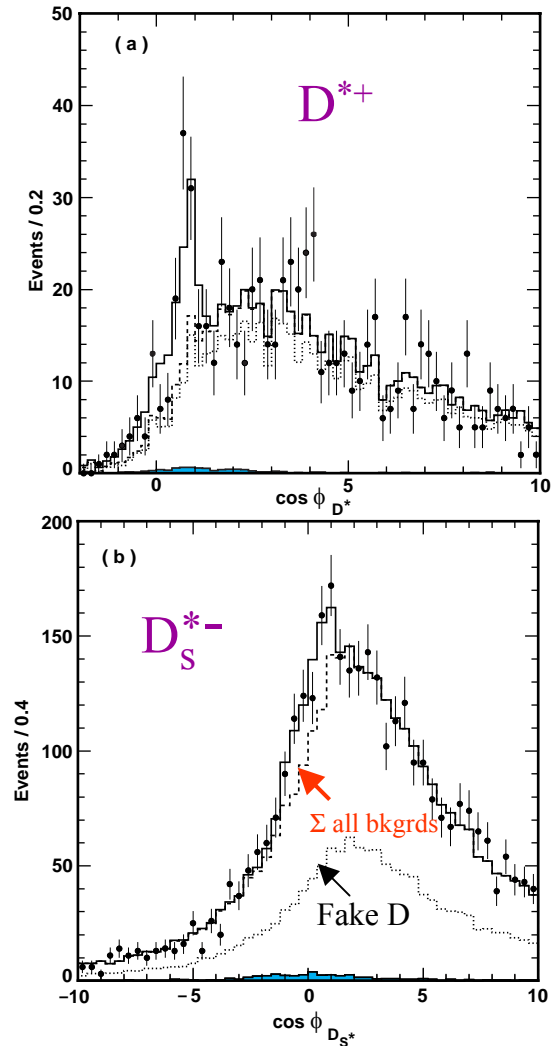
Find:

$$|A_2|/|A_0| = 0.63 \pm 0.13 \quad \delta = (\delta_2 - \delta_0) = 81^\circ \pm 10^\circ$$

B(D_s⁺ → φπ⁺)

B (%)	Error(%)	Source
3.59±0.77±0.48	25.3	CLEO
3.6±0.9	25.0	PDG

- Method: Reconstruct
 $B \rightarrow D^{*+} D_s^{*-}$, $D_s^{*-} \rightarrow \gamma D_s^-$ or
 $D^{*+} \rightarrow \pi^+ D^0$
- Observe signal both with & without explicit D_s or D⁰ reconstruction
- Measure $B(D_s \rightarrow \phi \pi^+) / B(D^0 \rightarrow K^- \pi^+)$





- Lower limit: Measure p and Λ yield in B decays and assume all such production is due to

$$\bar{B} \rightarrow \Lambda_c^+ \bar{N} X . \text{ Find } \mathcal{B} = (4.14 \pm 0.91)\%$$

- Upper limit: Measure $\Lambda_c \rightarrow \Lambda \ell \nu$, and assume that Λ saturates the rate (no Σ , for example).

$$\text{Find } \mathcal{B} = (7.7 \pm 1.5)\%$$

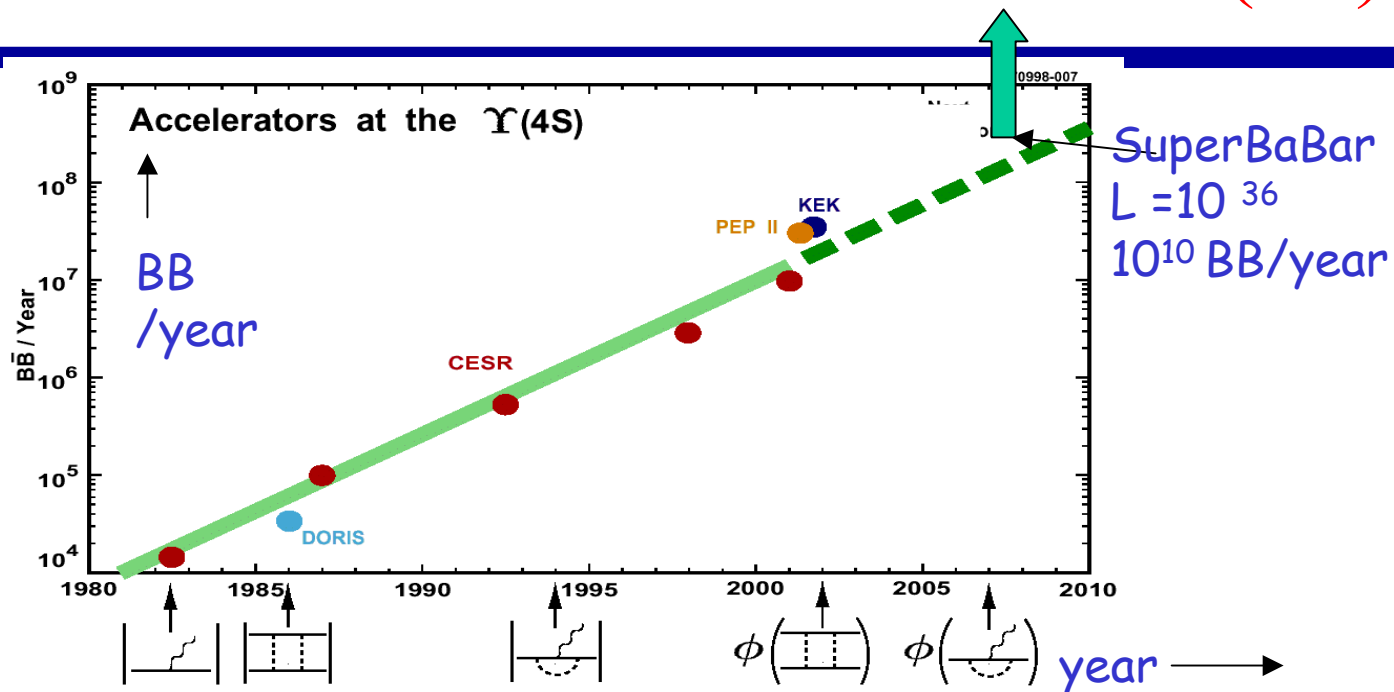
- Conclude: $9.7\% > \mathcal{B} > 3.0\% @ 90\% \text{ c. l.}$

$$J/\psi \rightarrow \mu^+ \mu^-$$

B (%)	Error(%)	Source
$5.84 \pm 0.06 \pm 0.10$	2.0	BES
6.08 ± 0.33	5.4	BES
5.88 ± 0.10	1.7	PDG

- Systematic error is the limitation. Completely correlated between the two BES measurements.
- Currently, best way to determine b yields at hadron colliders

Charm Production near/at the $\Upsilon(4S)$



	$L_{peak} \times 10^{33} \text{ cm}^{-2} \text{ s}^{-1}$	$\int L dt$	$\# B's \times 10^6$
		ON OFF	
CESR/CLEO(*)	1.3	16.0 6.7	34
KEKB/Belle	7.2	84.6 (ON+OFF)	~160
PEPII/BABAR	4.6	95.8 (ON+OFF)	~180

(*) CLEO
 No longer
 Operating at
 $\Upsilon(4S)$
 Belle/ BaBar
 ON/OFF
 June 13 '02