



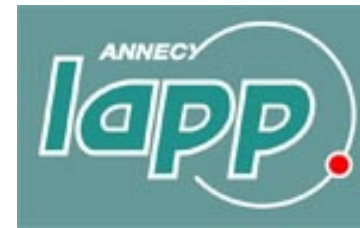
# CKM Fits: What the Data Say

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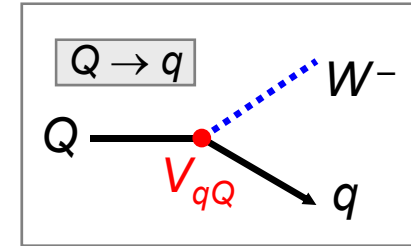
On behalf of the CKMfitter group

<http://ckmfitter.in2p3.fr>



# The CKM Matrix: Four Unknowns

$$V_{\text{CKM}} = \begin{pmatrix} V_{ud} & V_{us} & V_{ub} \\ V_{cd} & V_{cs} & V_{cb} \\ V_{td} & V_{ts} & V_{tb} \end{pmatrix}$$



unitarity-exact and convention-independent version of the Wolfenstein parametrization

$$\lambda^2 \equiv \frac{|V_{us}|^2}{|V_{ud}|^2 + |V_{us}|^2}$$

$$A^2 \lambda^4 \equiv \frac{|V_{cb}|^2}{|V_{ud}|^2 + |V_{us}|^2}$$

$$\bar{\rho} + i\bar{\eta} \equiv -\frac{V_{ud}V_{ub}^*}{V_{cd}V_{cb}^*}$$

Measurement  
of Wolfenstein  
parameters:

- ☀  $\lambda$  from  $|V_{ud}|$  (nuclear transitions) and  $|V_{us}|$  (semileptonic  $K$  decays)  
→ combined precision: 0.5%
- ☀  $A$  from  $|V_{cb}|$  (inclusive and exclusive semileptonic  $B$  decays)  
→ combined precision: 2%
- ☀  $\bar{\rho}, \bar{\eta}$  from (mainly) CKM angle measurements:  
→ combined precision: 20% ( $\rho$ ), 7% ( $\eta$ )

# Inputs: $\Delta m_s$



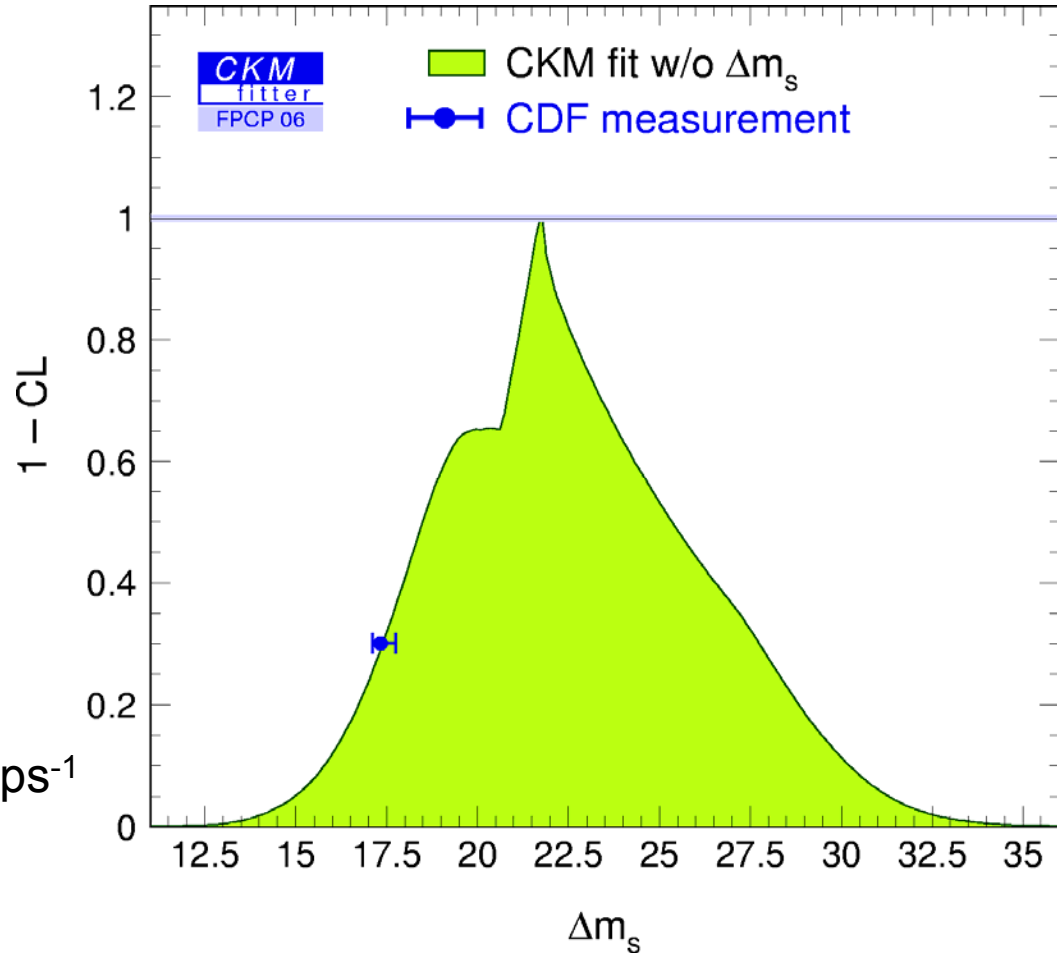
hep-ex/0603029

$17 < \Delta m_s < 21 \text{ ps}^{-1}$  @90 C.L.



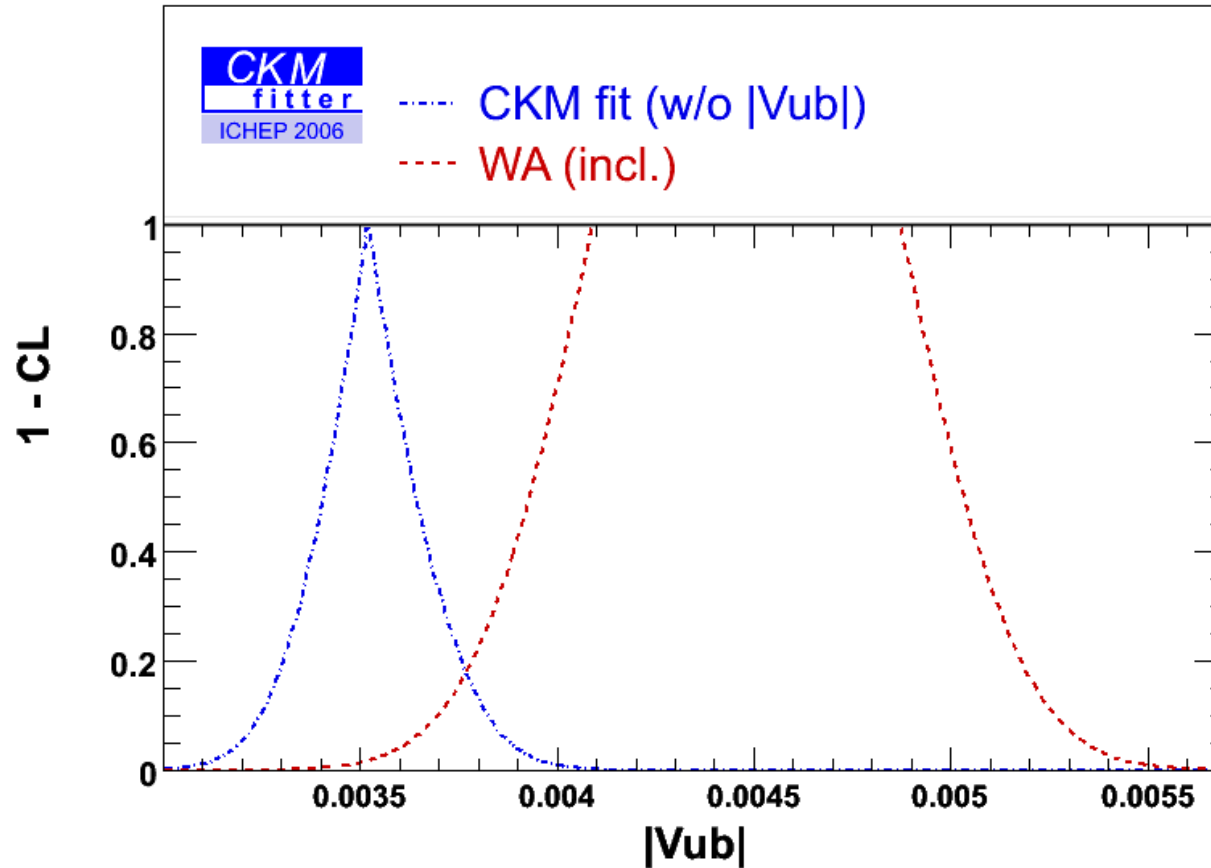
hep-ex/0606027

$\Delta m_s : 17.31^{+0.33}_{-0.18}$  (stat.)  $\pm 0.07$  (syst.)  $\text{ps}^{-1}$



# Inputs (major changes): $|V_{ub}|$ (incl.)

$|V_{ub}|$  (incl.)  $[10^{-3}] = 4.48 \pm 0.24 \pm 0.39$  (our average)



# Inputs (major changes): $\sin 2\beta$

☀ “The” *raison d'être* of the  $B$  factories:

$$\sin 2\beta = \begin{cases} 0.710 & 0.034 & 0.019 & [\text{BABAR}(348\text{m})] \\ 0.642 & 0.031 & 0.017 & [\text{Belle} (532\text{m})] \\ 0.674 & 0.026 & (0.023 \text{ stat-only}) & \end{cases}$$

$$C(\text{J}/\psi\text{K}^0) = \begin{cases} 0.070 & 0.028 & 0.018 & [\text{BABAR}] \\ -0.018 & 0.021 & 0.014 & [\text{Belle}] \\ 0.012 & 0.022 & (0.017 \text{ stat-only}) & \end{cases}$$

ICHEP '06

☀ Conflict with  $\sin 2\beta_{\text{eff}}$  from s-penguin

modes ? (New Physics (NP)?)

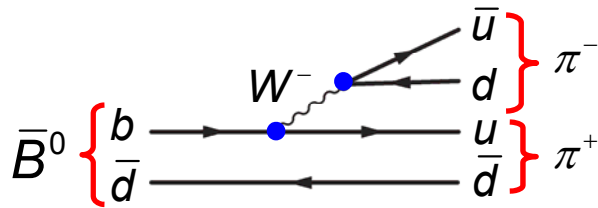
$$\sin 2\beta (\eta' \text{K}^0) = \begin{cases} 0.55 & 0.11 & 0.02 & [\text{BABAR}(347\text{m})] \\ 0.64 & 0.10 & 0.02 & [\text{Belle} (532\text{m})] \\ 0.60 & 0.08 & & \end{cases}$$

$$\sin 2\beta (\phi \text{K}^0) = \begin{cases} 0.12 & 0.34 & & [\text{BABAR}(347\text{m})] \\ 0.44 & 0.27 & 0.05 & [\text{Belle} (386\text{m})] \\ 0.31 & 0.21 & & \end{cases}$$

NP can contribute differently among the various s-penguin modes  
Meaning of the average?

NB: a disagreement would falsify the SM. The interference NP/SM amplitudes introduces hadronic uncertainties  
→ Cannot determine the NP parameters cleanly

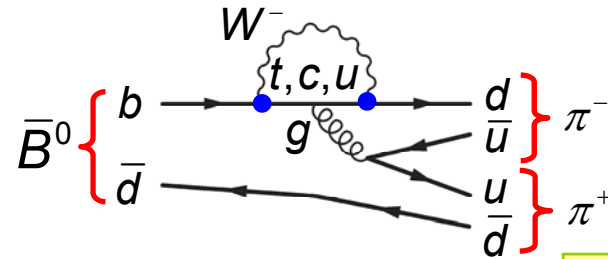
# Inputs (major changes): angle $\alpha$



Tree : dominant

$$\propto V_{ub} V_{ud}^*$$

$$\propto \lambda^3$$



Penguin : competitive ?

$$\propto V_{tb} V_{td}^*$$

$$\propto \lambda^3$$

☀ Time-dependent CP observable :

realistic scenario

$$A_{h^+h^-}(t) = S_{h^+h^-} \sin(\Delta m_d t) - C_{h^+h^-} \cos(\Delta m_d t)$$

$$= \sqrt{1 - C_{h^+h^-}^2} \sin(2\alpha_{\text{eff}}) \cdot \sin(\Delta m_d t) - C_{h^+h^-} \cos(\Delta m_d t)$$

Time-dependent CP analysis of  $B^0 \rightarrow \pi^+\pi^-$  alone determines  $\alpha_{\text{eff}}$  : but, we need  $\alpha$  !



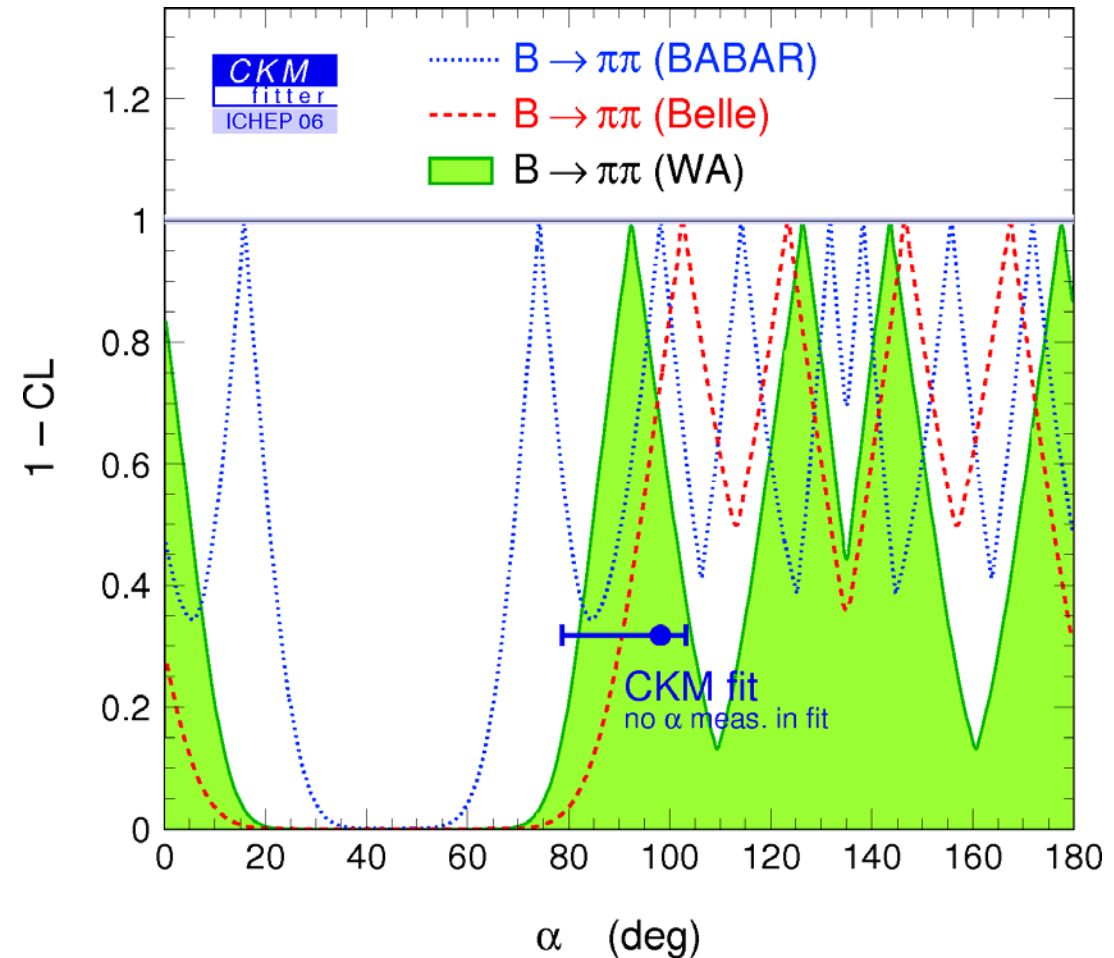
Isospin analysis ( $\alpha$  can be resolved up to an 8-fold ambiguity within  $[0, \pi]$ )

# Isospin Analysis: $B \rightarrow \pi\pi$

	BABAR (347m)	Belle (532m)	Average
$S_{\pi\pi}$	$-0.53 \pm 0.14 \pm 0.02$	$-0.61 \pm 0.10 \pm 0.04$	$-0.58 \pm 0.09$
$C_{\pi\pi}$	$-0.16 \pm 0.11 \pm 0.03$	$-0.55 \pm 0.08 \pm 0.05$	$-0.39 \pm 0.07$

“agreement”:  $2.6\sigma$

BABAR & Belle



# Isospin Analysis: $B \rightarrow \rho\rho$

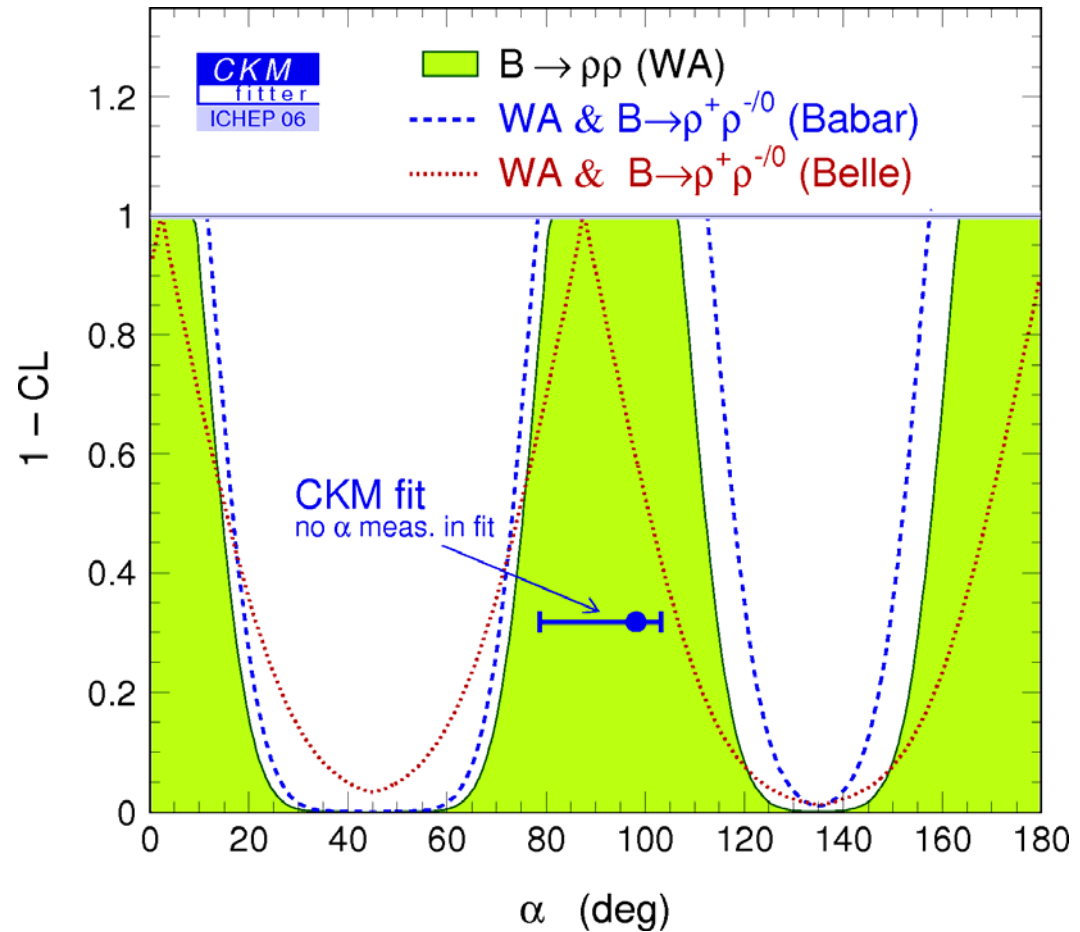
	BABAR (347m)			Belle (275m)			Average	
$S_{\rho\rho}$	-0.19	0.21	$^{+0.05}_{-0.07}$	0.08	0.41	0.09	-0.13	0.19
$C_{\rho\rho}$	-0.07	0.15	0.06	0.0	0.3	0.09	-0.06	0.14

BABAR & Belle

BABAR (347m)			
$f_L^{00}$	0.86	$^{+0.11}_{-0.13}$	0.06
$BR^{00}$	$(1.2 \ 0.4 \ 0.3) \times 10^{-6}$		

☀ Isospin analysis :

$$\alpha = [ 94 \quad 21 ]^\circ$$



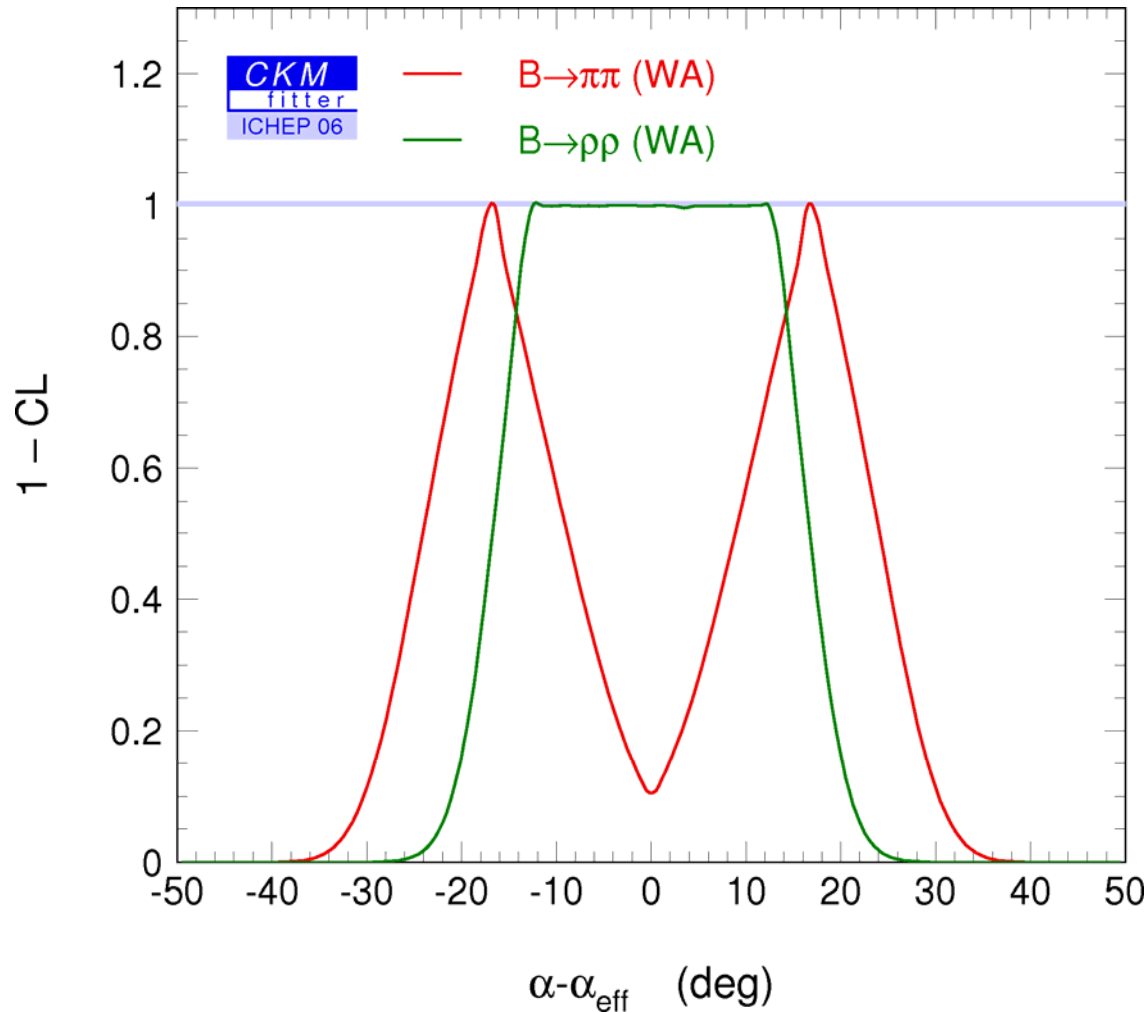
# Isospin Analysis: angle $\alpha_{\text{eff}}$ [ $B \rightarrow \pi\pi/\rho\rho$ ]

✿ Isospin analysis  $B \rightarrow \pi\pi$ :

$$|\alpha - \alpha_{\text{eff}}| < 32.1^\circ \text{ (95\% CL)}$$

✿ Isospin analysis  $B \rightarrow \rho\rho$ :

$$|\alpha - \alpha_{\text{eff}}| < 22.4^\circ \text{ (95\% CL)}$$



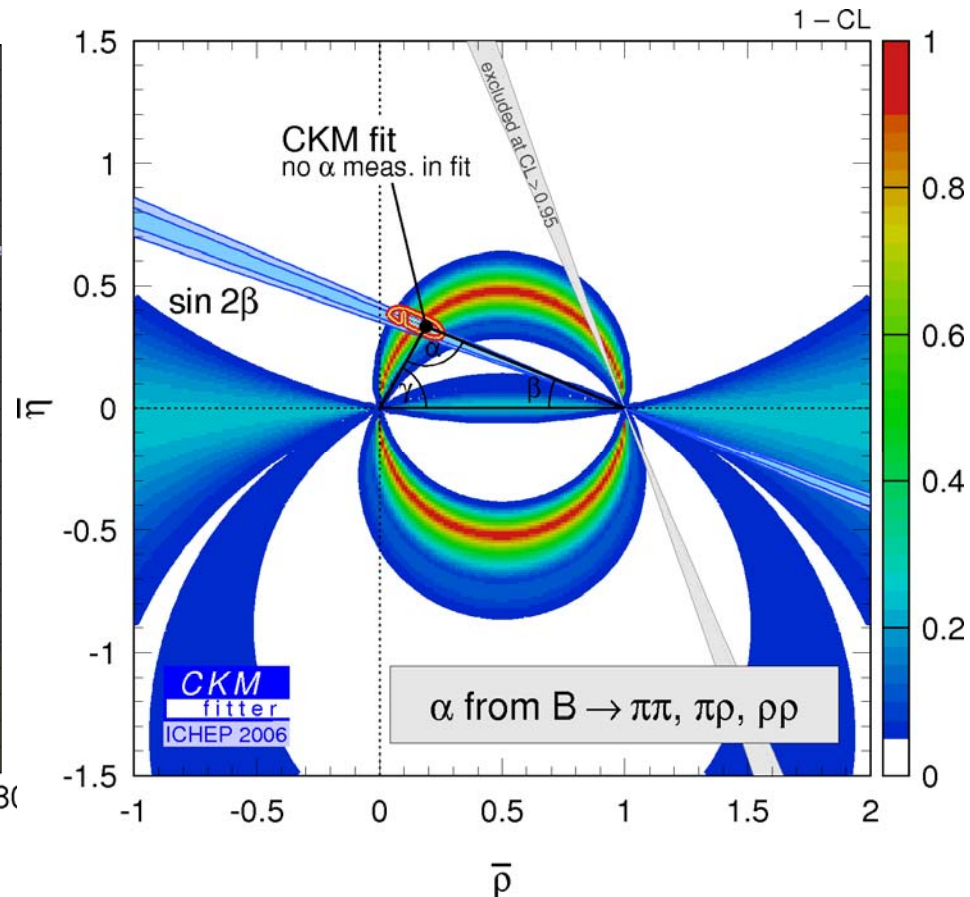
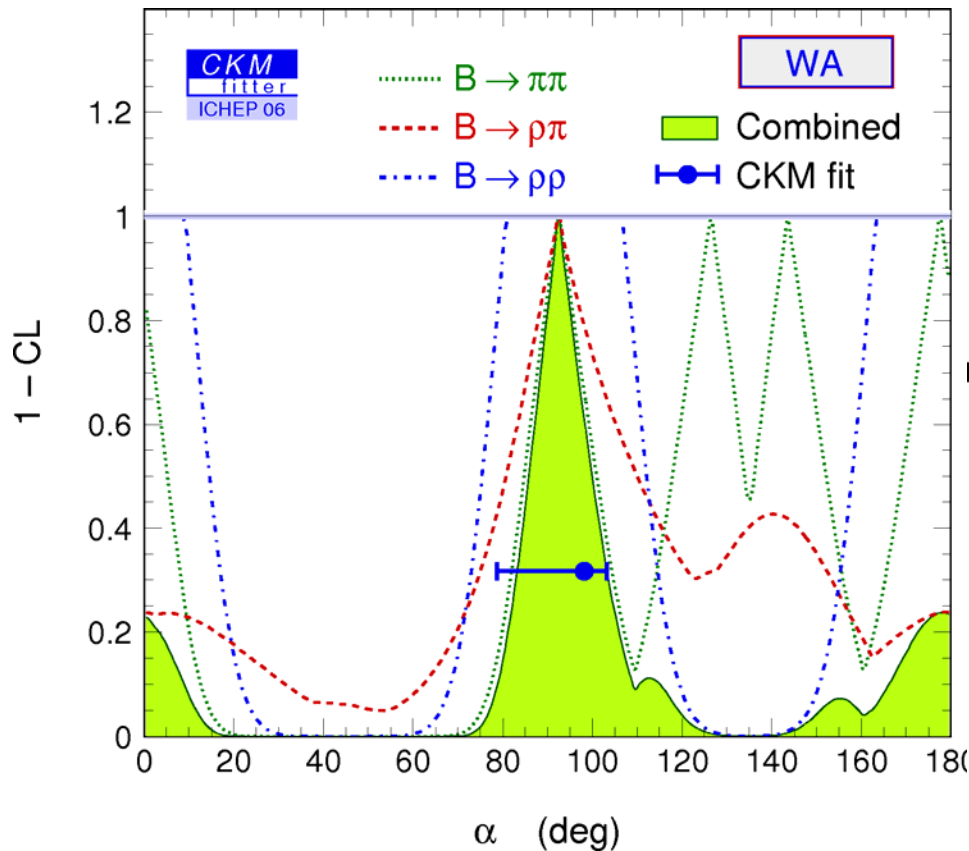
# Isospin Analysis: angle $\alpha$ [ $B \rightarrow \pi\pi / \rho\pi / \rho\rho$ ]

$$\alpha_{\text{B-Factories}} = [ 93^{+11}_{-9} ]^\circ$$



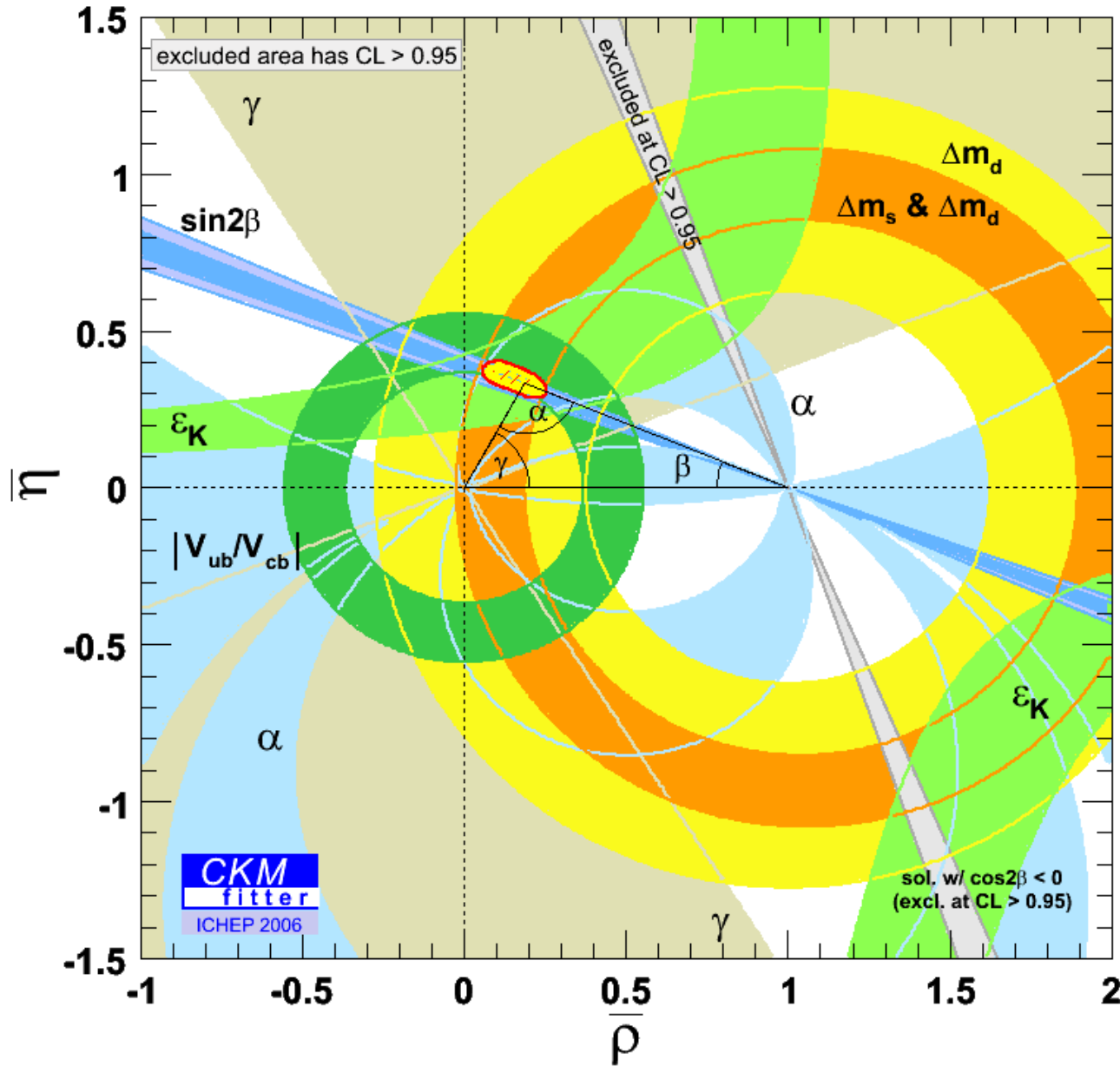
$$\alpha_{\text{Global Fit}} = [ 98^{+5}_{-19} ]^\circ$$

(include new  $\rho\pi$  Dalitz BABAR)



$B \rightarrow \rho\rho$ : at very large statistics, systematics and model-dependence will become an issue

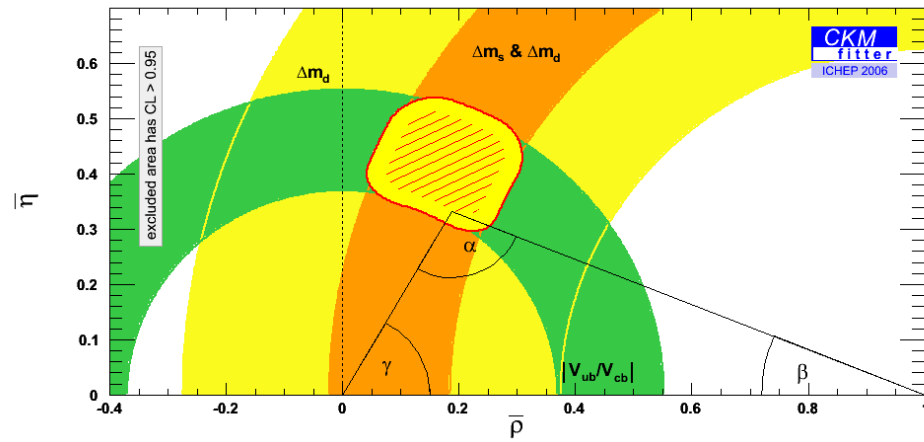
$B \rightarrow \rho\pi$  Dalitz analysis: model-dependence is an issue !



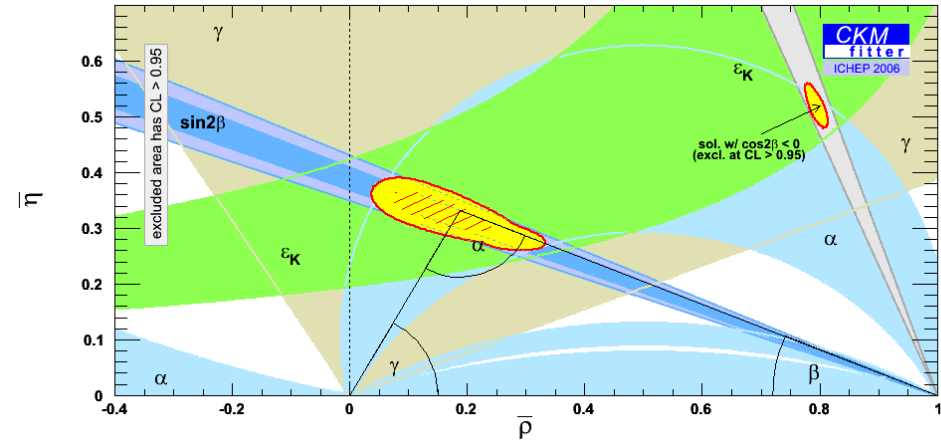
Inputs:

- $\left| \frac{V_{ub}}{V_{cb}} \right|$
- $\Delta m_d$
- $\Delta m_s$
- $B \rightarrow \tau \nu$
- $|\varepsilon_K|$
- $\sin 2\beta$
- $\alpha$
- $\gamma$

# The global CKM fit: Testing the CKM Paradigm

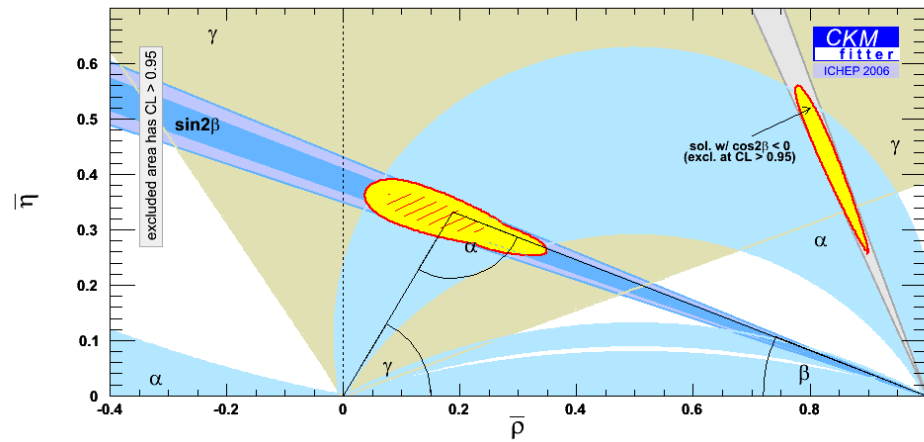


CP Conserving

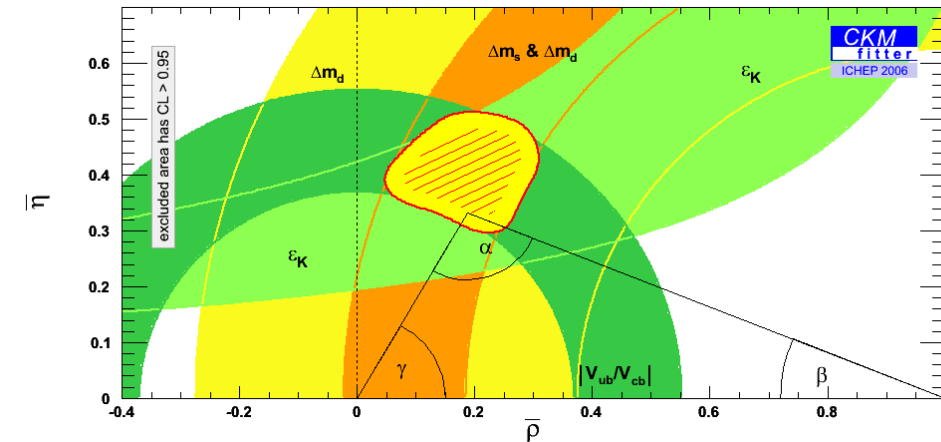


CP Violating

*CP-insensitive observables imply CP violation !*

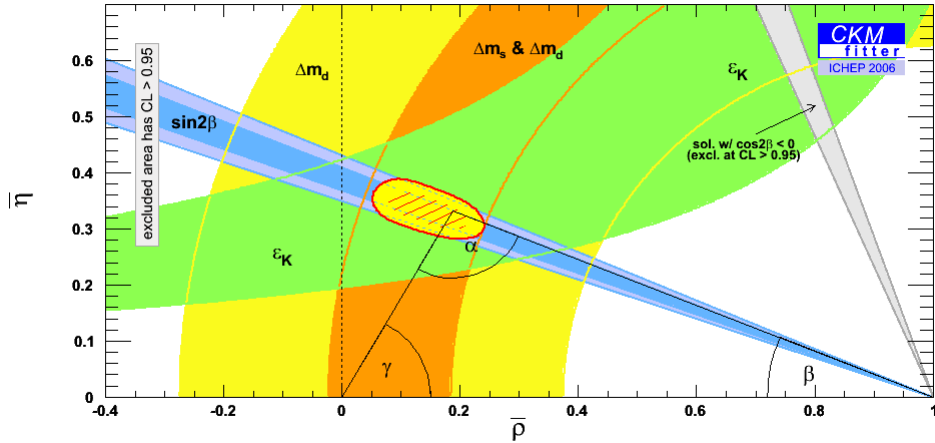
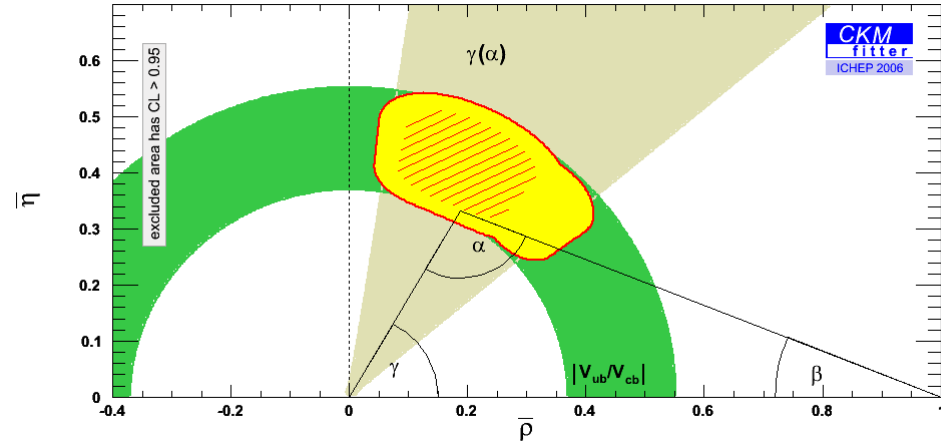


Angles (no theory)



No angles (with theory)

# The global CKM fit: Testing the CKM Paradigm (cont.)



Tree (NP-Free)

Loop

[No NP in  $\Delta I=3/2$   $b \rightarrow d$  EW penguin amplitude  
Use  $\alpha$  with  $\beta$  (charmonium) to cancel NP amplitude]

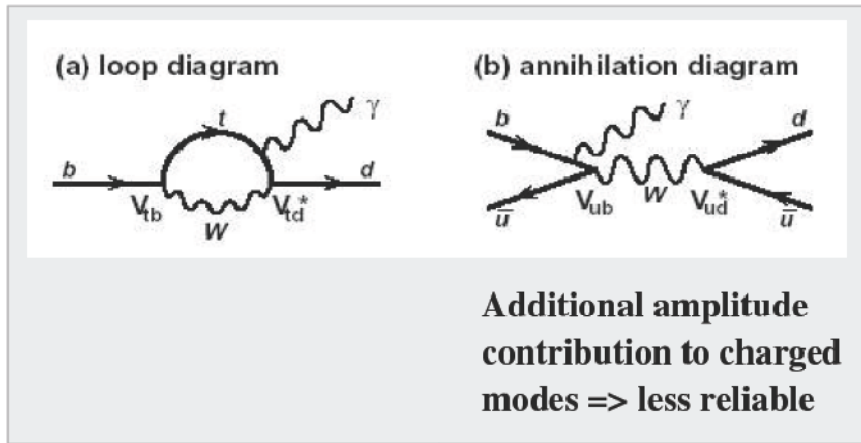
CKM mechanism: dominant source of CP violation

The global fit is not the whole story: several  $\Delta F=1$  rare decays are not yet measured

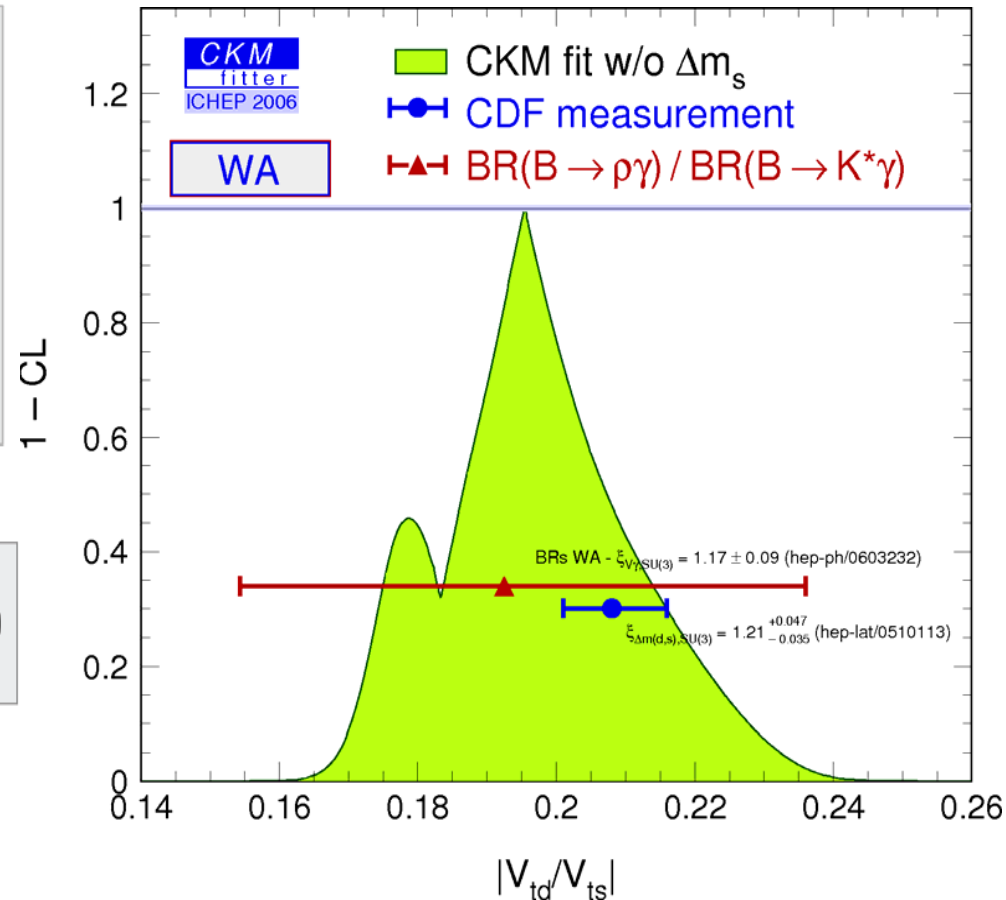
➔ Sensitive to NP

# Radiative Penguin Decays: $BR(B \rightarrow \rho \gamma) / BR(B \rightarrow K^* \gamma)$

$B \rightarrow \rho \gamma (\propto |V_{td}|^2)$  &  $B \rightarrow K^* \gamma (\propto |V_{ts}|^2)$  sensitive to New Physics



$$\frac{BF(B^0 \rightarrow \rho^0 \gamma)}{BF(B^0 \rightarrow K^{*0} \gamma)} = 1.023 \frac{1}{2} \left| \frac{V_{td}}{V_{ts}} \right|^2 \xi^{-2} (1 + \Delta)$$



	BABAR (347m)	Belle (386m)
$\rho^0 \gamma$	$0.77^{+0.21}_{-0.19} \pm 0.07$	$1.25^{+0.37}_{-0.33} \begin{smallmatrix} +0.07 \\ -0.06 \end{smallmatrix}$
$\rho^+ \gamma$	$1.06^{+0.35}_{-0.31} \pm 0.09$	$0.55^{+0.42}_{-0.36} \begin{smallmatrix} +0.09 \\ -0.08 \end{smallmatrix}$

# NP Parameterization in $B_s$ system

$$\frac{\left\langle B_s^0 \left| H_{eff}^{SM+NP} \right| \overline{B}_s^0 \right\rangle}{\left\langle B_s^0 \left| H_{eff}^{SM} \right| \overline{B}_s^0 \right\rangle} = r_s^2 e^{i2\theta_s} = 1 + h_s e^{i2\sigma_s}$$

Grossman, PL **B380**, 99 (1996)

Dunietz, Fleischer, Nierste, PRD **63**, 114015 (2001)

Hypothesis: NP in loop processes only (negligible for tree processes)

Mass difference:  $\Delta m_s = (\Delta m_s)^{SM} r_s^2$

Width difference:  $\Delta \Gamma_s^{CP} = (\Delta \Gamma_s)^{SM} \cos^2(2\chi - 2\theta_s)$

Semileptonic asymmetry:

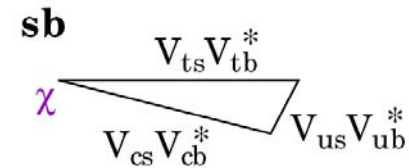
$A_{SL}^s = -\text{Re}(\Gamma_{12}/M_{12})^{SM} \sin(2\theta_s)/r_s^2$

$S_{\psi\phi} = \sin(2\chi - 2\theta_s)$

NP wrt to SM:

- reduces  $\Delta \Gamma_s$
- enhances  $\Delta m_s$

$$\chi = \arg\left(-V_{ts}V_{tb}^*/V_{cs}V_{cb}^*\right)$$

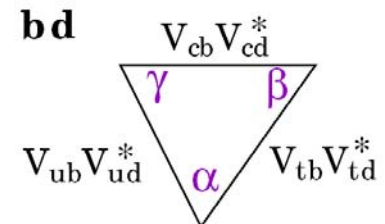


UT of  $B_d$  system: non-degenerated

→  $(h_d, \sigma_d)$  strongly correlated to the determination of  $(\rho, \eta)$

UT of  $B_s$  system: highly degenerated

→  $(h_s, \sigma_s)$  almost independent of  $(\rho, \eta)$

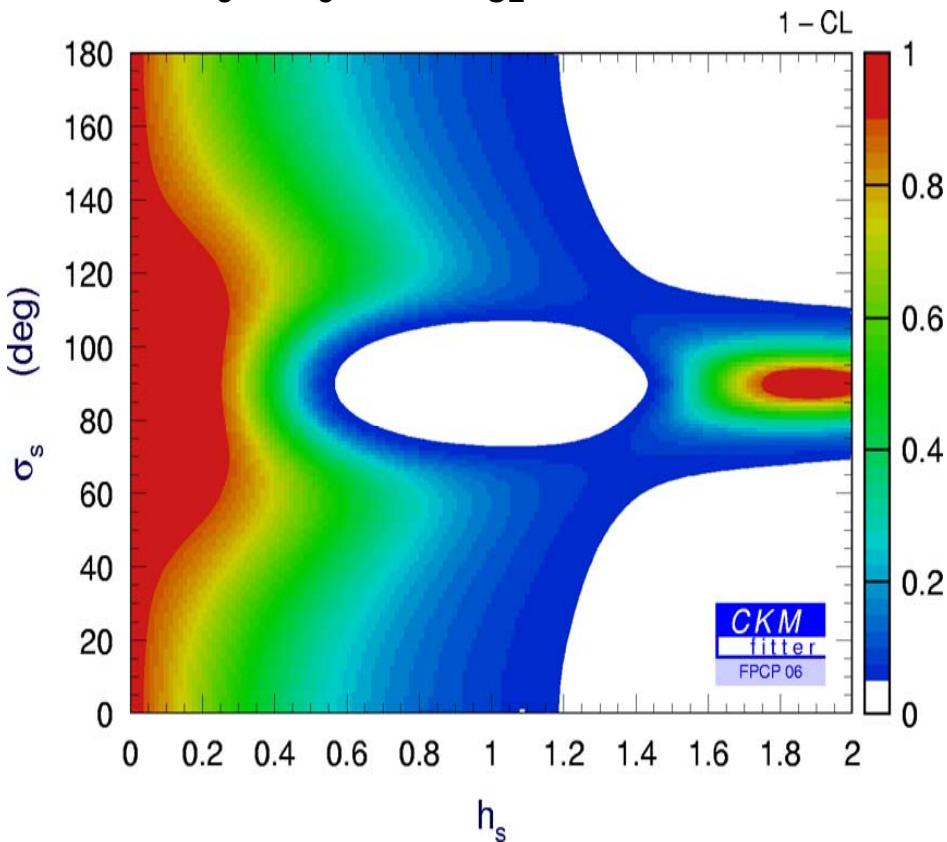


$B_s$  mixing phase very small in SM:  $\chi = 1.05 + 0.06$  (deg)

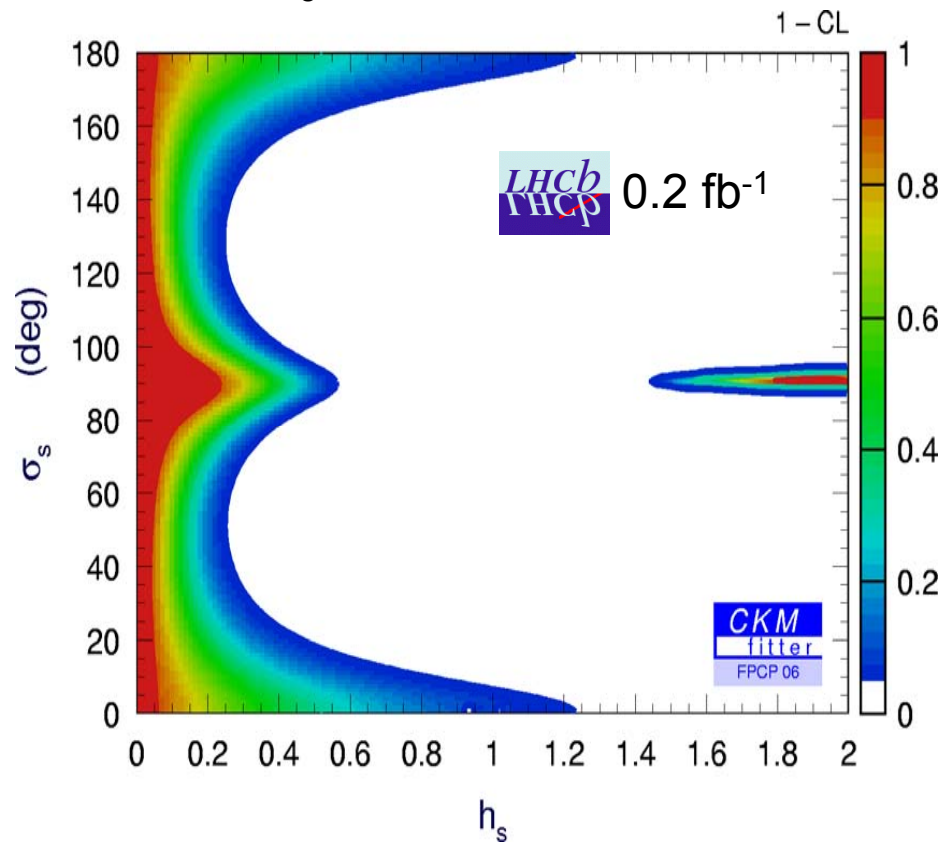
→ **Bs mixing: very sensitive probe to NP**

# NP in $B_s$ System

$\Delta m_s$ ,  $\Delta \Gamma_s$  and  $A_{SL}^s$



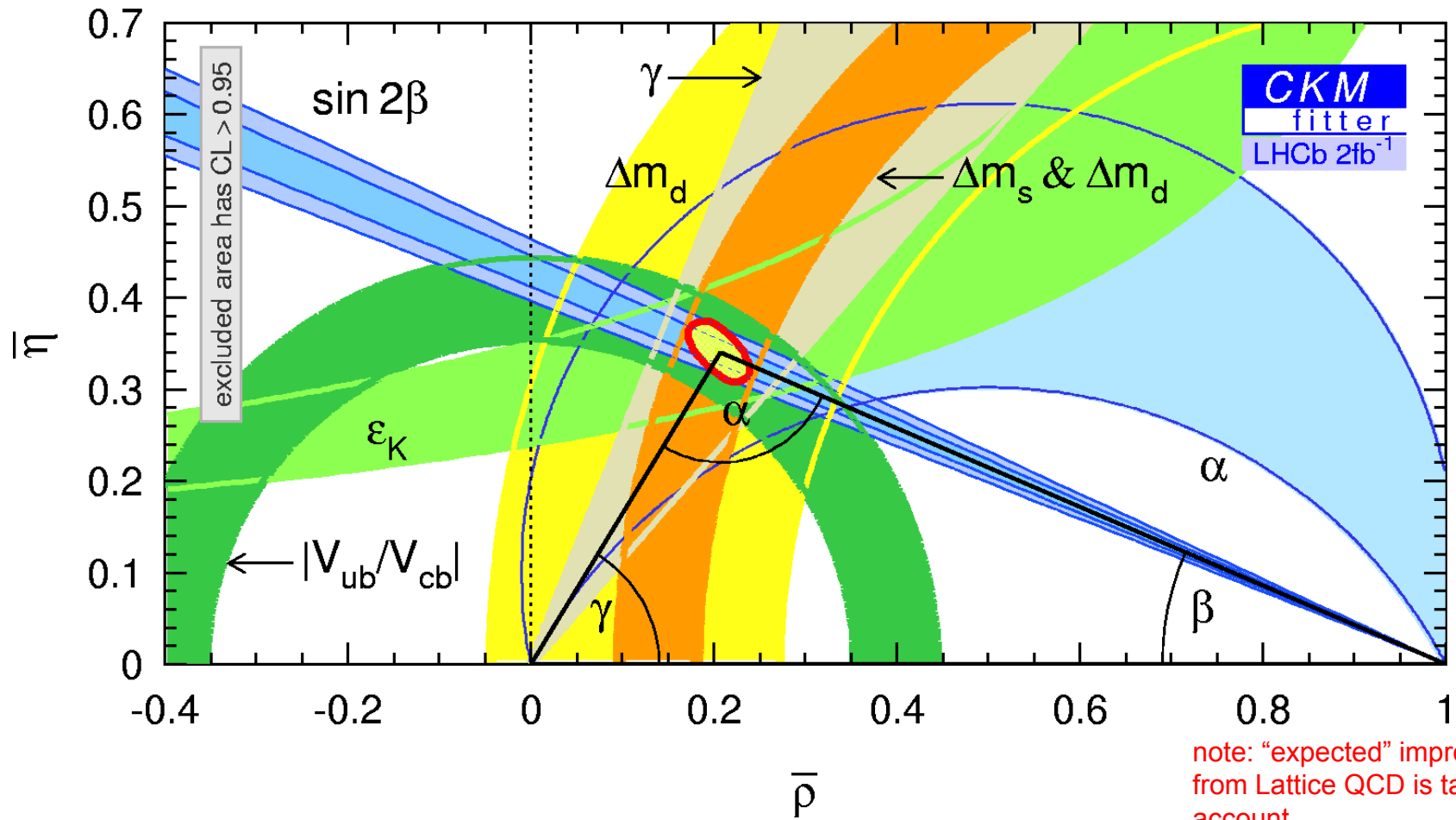
$\sigma(\Delta m_s) = 0.035$ ,  $\sigma(\sin(2\chi)) = 0.1$



First constraint for NP in the  $B_s$  sector  
 Still plenty of room for NP  
 Large theoretical uncertainties: LQCD

$$h_s \lesssim 3 \quad (h_d \lesssim 0.3, h_K \lesssim 0.6)$$

# Prospective: LHCb 2fb<sup>-1</sup> (2010 ??)?



- assumptions:
- $\sigma(\Delta m_s) = 0.01$
  - $\sigma(\sin 2\beta) = 0.02$
  - $\sigma(\alpha) = 10^\circ$
  - $\sigma(\gamma) = 5^\circ$

# Conclusion

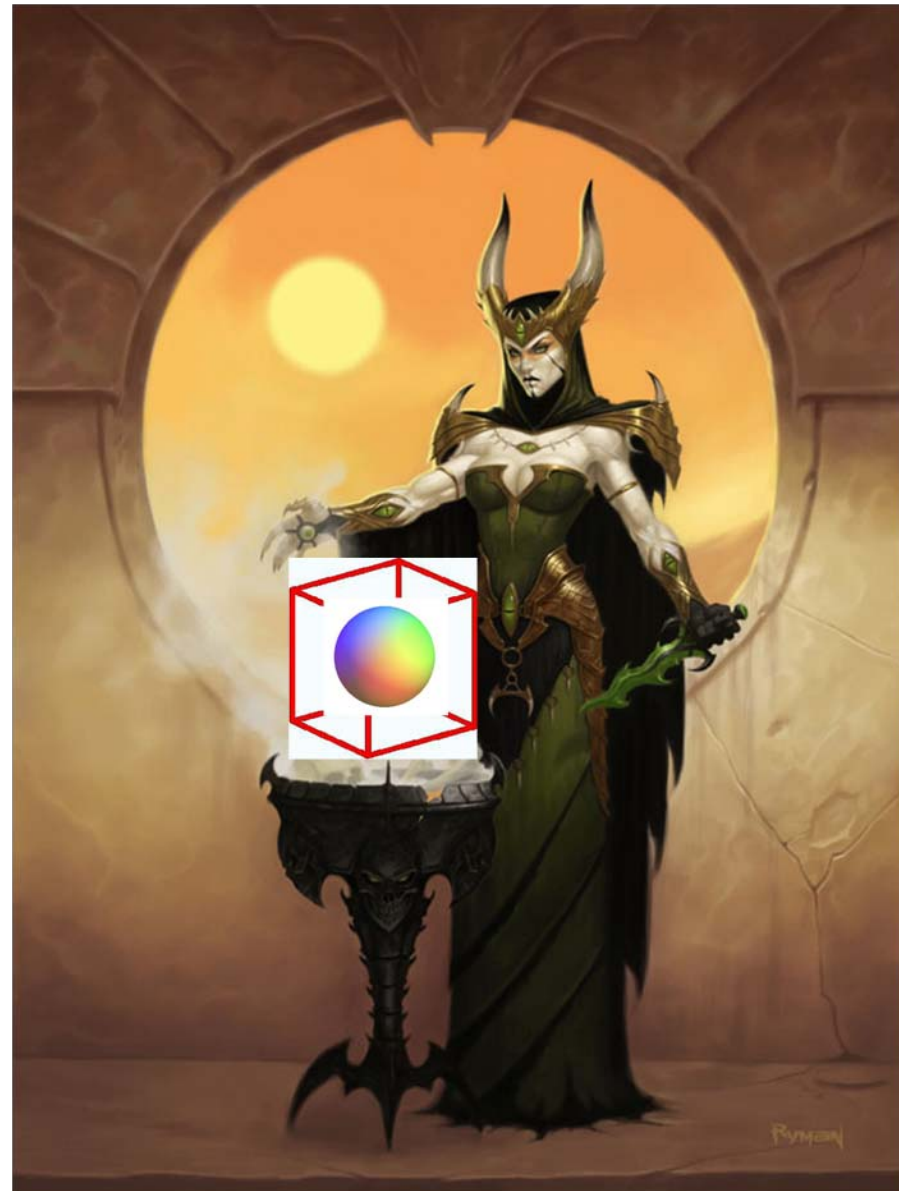
- CKM mechanism: success in describing flavor dynamics of many constraints from vastly different scales.
- With the increase of statistics, lots of assumptions will be needed to be reconsidered [e.g., extraction of  $\alpha$  from  $B \rightarrow 3\pi, 4\pi$ , etc.,  $P_{EW}$ , ...]
- $B_s$ : an independent chapter in Nature's book on fundamental dynamics
  - there is no reason why NP should have the same flavor structure as in the SM
  - $B_s$  transitions can be harnessed as powerful probes for NP ( $\chi$ : "NP model killer")
- Before claiming NP discovery, be sure that everything is "under control" (assumptions, theoretical uncertainties, etc.)
- There are still plenty of measurements yet to be done

Do Not Miss:

[hep-ph/0607246](https://arxiv.org/abs/hep-ph/0607246)

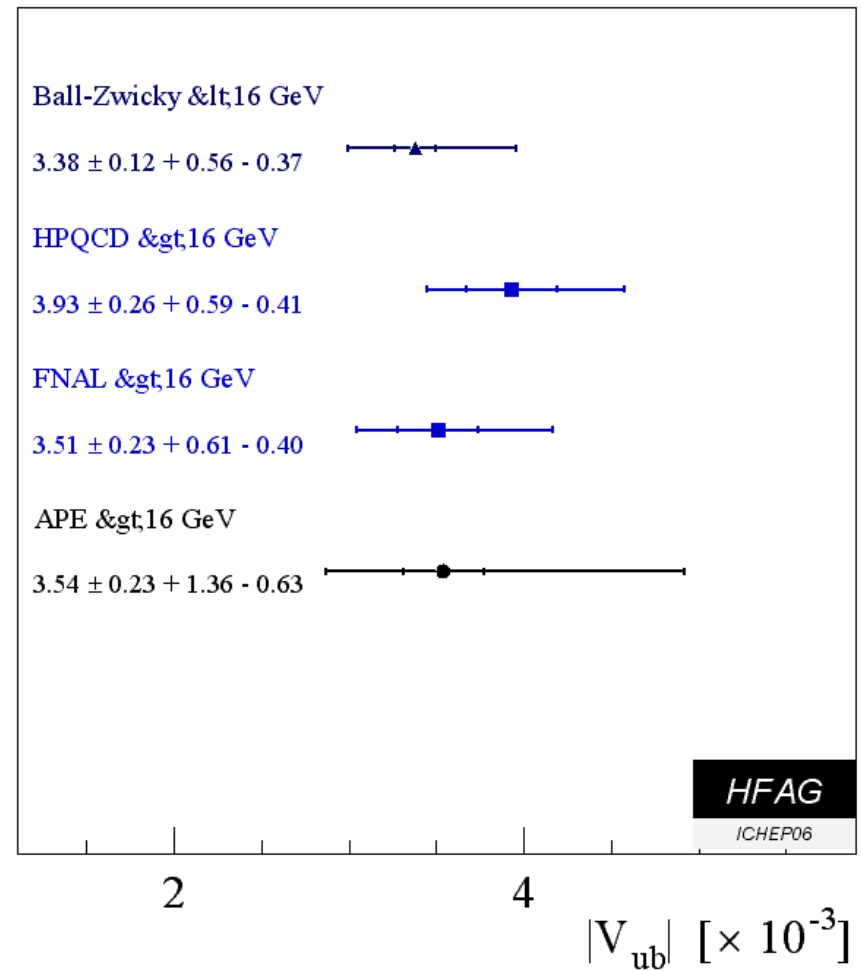
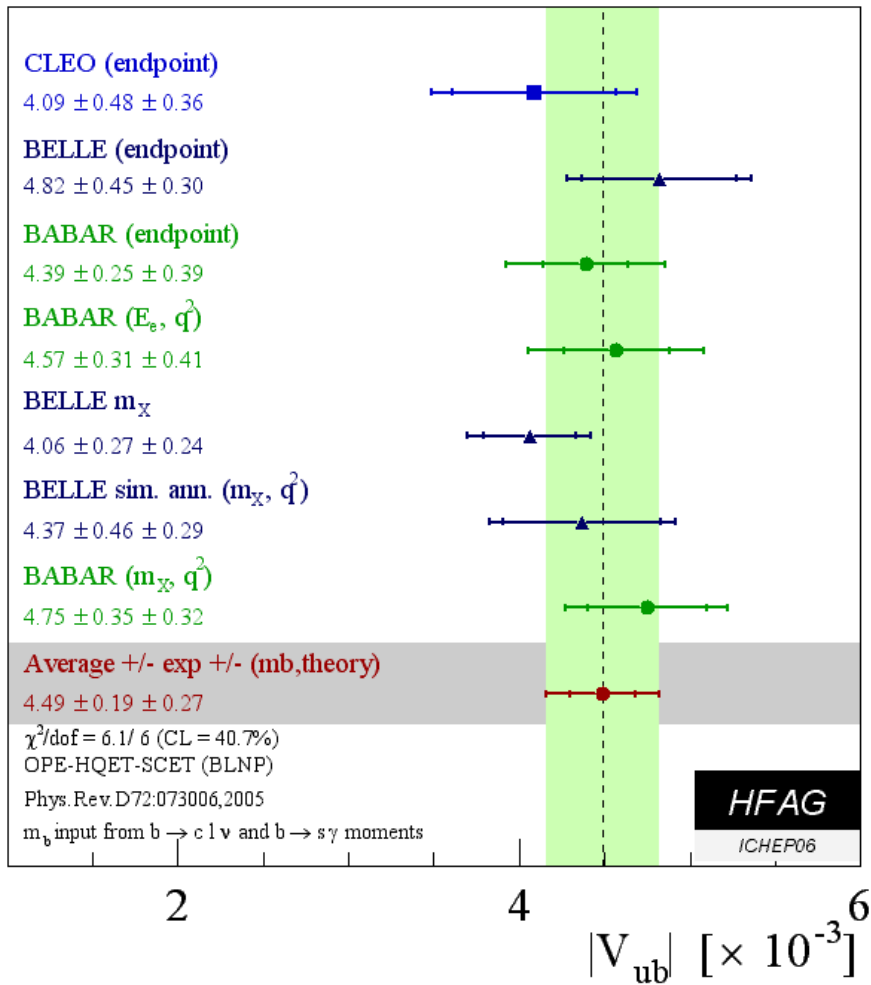
Bayesian Statistics at Work:  
the Troublesome  
Extraction of the CKM Angle  $\alpha$

(J. Charles, A. Höcker, H. Lacker  
F.R. Le Diberder and S. T'Jampens)





BACKUP SLIDES



## FLAVOR STRUCTURE

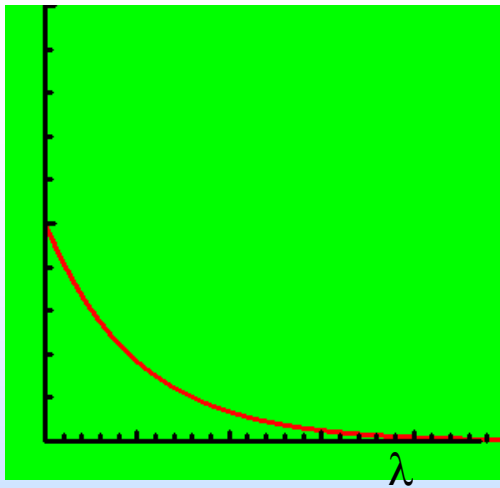
ELECTROWEAK  
STRUCTURE

	$b \rightarrow s$	$b \rightarrow d$	$s \rightarrow d$
$\Delta F=2$ box	$\Delta M_{B_s}$ $A_{CP}(B_s \rightarrow \psi \phi)$	$\Delta M_{B_d}$ $A_{CP}(B_d \rightarrow \psi K)$	$\Delta M_K, \epsilon_K$
$\Delta F=1$ 4-quark box	$B_d \rightarrow \phi K, B_d \rightarrow K\pi, \dots$	$B_d \rightarrow \pi\pi, B_d \rightarrow \rho\pi, \dots$	$\epsilon'/\epsilon, K \rightarrow 3\pi, \dots$
gluon penguin	$B_d \rightarrow X_s \gamma, B_d \rightarrow \phi K,$ $B_d \rightarrow K\pi, \dots$	$B_d \rightarrow X_d \gamma, B_d \rightarrow \pi\pi, \dots$	$\epsilon'/\epsilon, K_L \rightarrow \pi^0 l^+ l^-, \dots$
$\gamma$ penguin	$B_d \rightarrow X_s l^+ l^-, B_d \rightarrow X_s \gamma$ $B_d \rightarrow \phi K, B_d \rightarrow K\pi, \dots$	$B_d \rightarrow X_d l^+ l^-, B_d \rightarrow X_d \gamma$ $B_d \rightarrow \pi\pi, \dots$	$\epsilon'/\epsilon, K_L \rightarrow \pi^0 l^+ l^-, \dots$
$Z^0$ penguin	$B_d \rightarrow X_s l^+ l^-, B_s \rightarrow \mu\mu$ $B_d \rightarrow \phi K, B_d \rightarrow K\pi, \dots$	$B_d \rightarrow X_d l^+ l^-, B_d \rightarrow \mu\mu$ $B_d \rightarrow \pi\pi, \dots$	$\epsilon'/\epsilon, K_L \rightarrow \pi^0 l^+ l^-,$ $K \rightarrow \pi \nu \nu, K \rightarrow \mu\mu, \dots$
$H^0$ penguin	$B_s \rightarrow \mu\mu$	$B_d \rightarrow \mu\mu$	$K_{L,S} \rightarrow \mu\mu$

# Bayes at work

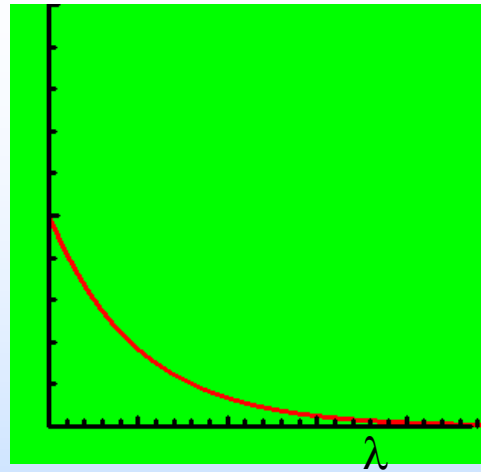
Zero events seen

$$P(n; \lambda) = e^{-\lambda} \lambda^n / n!$$



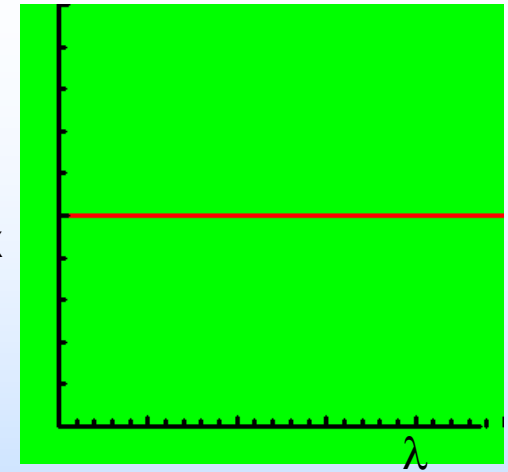
Posterior  $P(\lambda)$

=



$P(0 \text{ events} | \lambda)$   
(Likelihood)

x



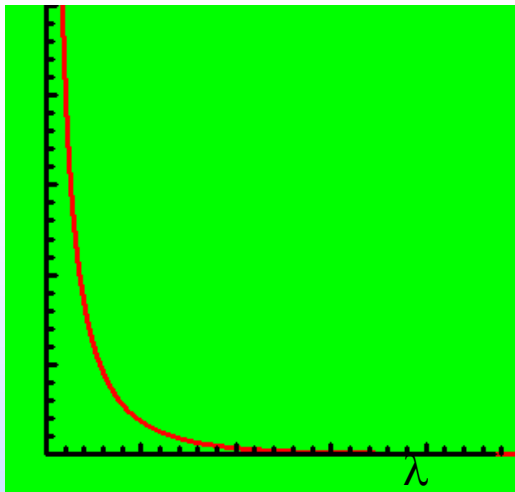
Prior: uniform

$$\int_0^3 P(\lambda) d\lambda = 0.95$$

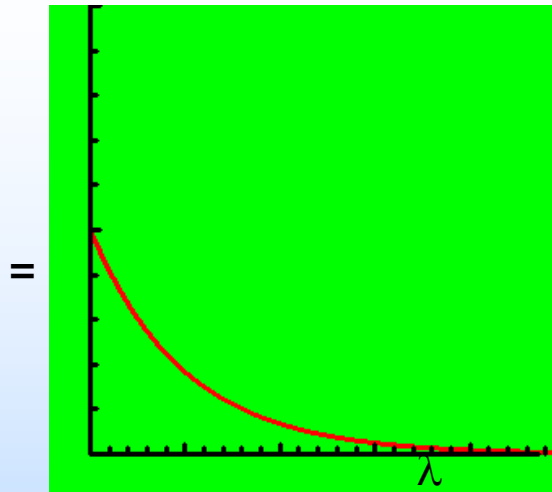
Same as Frequentist limit -  
Happy coincidence

# Bayes at work again

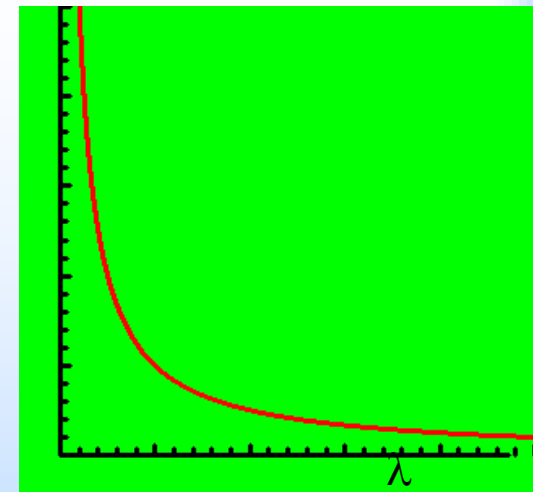
Is that uniform prior really credible?



Posterior  $P(\lambda)$



$P(0 \text{ events}|\lambda)$



Prior: uniform in  $\ln \lambda$

Upper limit totally different!

$$\int_0^3 P(\lambda) d\lambda \gg 0.95$$

# Bayes: the bad news

- The prior affects the posterior. It is your choice. That makes the measurement subjective. This is BAD. (We're physicists, dammit!)
- A Uniform Prior does not get you out of this.