

τ Universality and Precision Measurements

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Universality and Precision Measurements

① 'Data'

i) "CC": τ_2 , Michel parameters, $W \rightarrow \tau\nu, \dots$

↓
important interplay... $\rightarrow x_{\text{eff}}^e$
'LEP vs SLD'

ii) "NC": $e^+e^- \rightarrow \gamma, Z \rightarrow \tau^+\tau^-$ vs $e^+\bar{e} / \mu^+\mu^-$

② 'Some Models' of Universality Violation

i) mix τ with exotics

ii) $Z-Z'$ mixing w/ Z' having gen. dep. couplings

iii) give τ a new intrinsic property: $x_{Z,\gamma}^\tau$

How are they doing? Sensitivity?

③ Conclusions

ZCF'94

TGR:220

τ_z is THE Classic test of τ Universality ...

→ input : m_τ , $B_{e,\mu}$ ⊕ SM RC ⊕ τ_μ

→ output : τ_z or G_τ/G_μ ∵ looks excellent !

remember...

$$\frac{G_\tau}{G_\mu} = \frac{1 - \Delta r}{1 - \Delta r_\mu} \xrightarrow{\text{SM}} 1 + O\left(\frac{\alpha}{\pi} \frac{m_\tau^2}{m_\mu^2}\right) = 1 + O(10^{-6}) !$$

• But ... BSM large non-universal's may exist which may be significant

$$\text{e.g., } \rightarrow 1 + O\left(\frac{\alpha}{\pi} \frac{m_\tau^2}{m_\mu^2} \tan^2 \rho\right) = 1 + O(10^{-3}) !!$$

↳ probed at τ cF !

Michel Parameters : is $\tau V_e W$ V-A ?

	"LE" †	"LEP" * (SWAIN)	"WA"
β	0.730 ± 0.026	0.798 ± 0.041 (A)	0.749 ± 0.032
δ	-	0.83 ± 0.17 (A, L3)	0.83 ± 0.17
ζ	0.90 ± 0.18	1.07 ± 0.17 (A, L3)	1.02 ± 0.13 *
η	$0.03 \pm 0.18 \pm 0.12$	-	$0.03 \pm 0.18 \pm 0.12$
h	1.25 ± 0.26	0.951 ± 0.058 (A)	0.965 ± 0.057 *

impressive, very nice, but so what ?

⇒ implications for precision
EW measurements

† All data 7/94 + before

* DPF '94 (conflict with Glasgow)
? (Paterson)

!!

..

$$\tilde{\tau}_e^{-1} = B_{ee}^{-1} \left[\frac{G_e^2 m_e^5}{192 \pi^3} \right] \cdot f(x_{e,e'}) \cdot \left\{ 1 + \frac{3}{5} \left(\frac{m_e^2}{M_W^2} \right) + O\left(\frac{m_e^4}{M_W^4}\right) \right\} \cdot \left[1 + \frac{\alpha(m_e)}{2\pi} \left(\frac{25}{4} - \pi^2 \right) + (3.3 \pm 1.8) \left(\frac{\alpha(m_e)}{\pi} \right)^2 + O\left(\frac{\alpha}{\pi} x_{e,e'}\right) \right]$$

↑
Data for
Comparison

↑
Phase
Space Correction
($x \equiv m_{e'}^2/m_e^2$)

↑
finite M_W
correction

$B(e \rightarrow e' \bar{\nu}_{e'} \bar{\nu}_{e'})$

$$O(10^{-8})$$

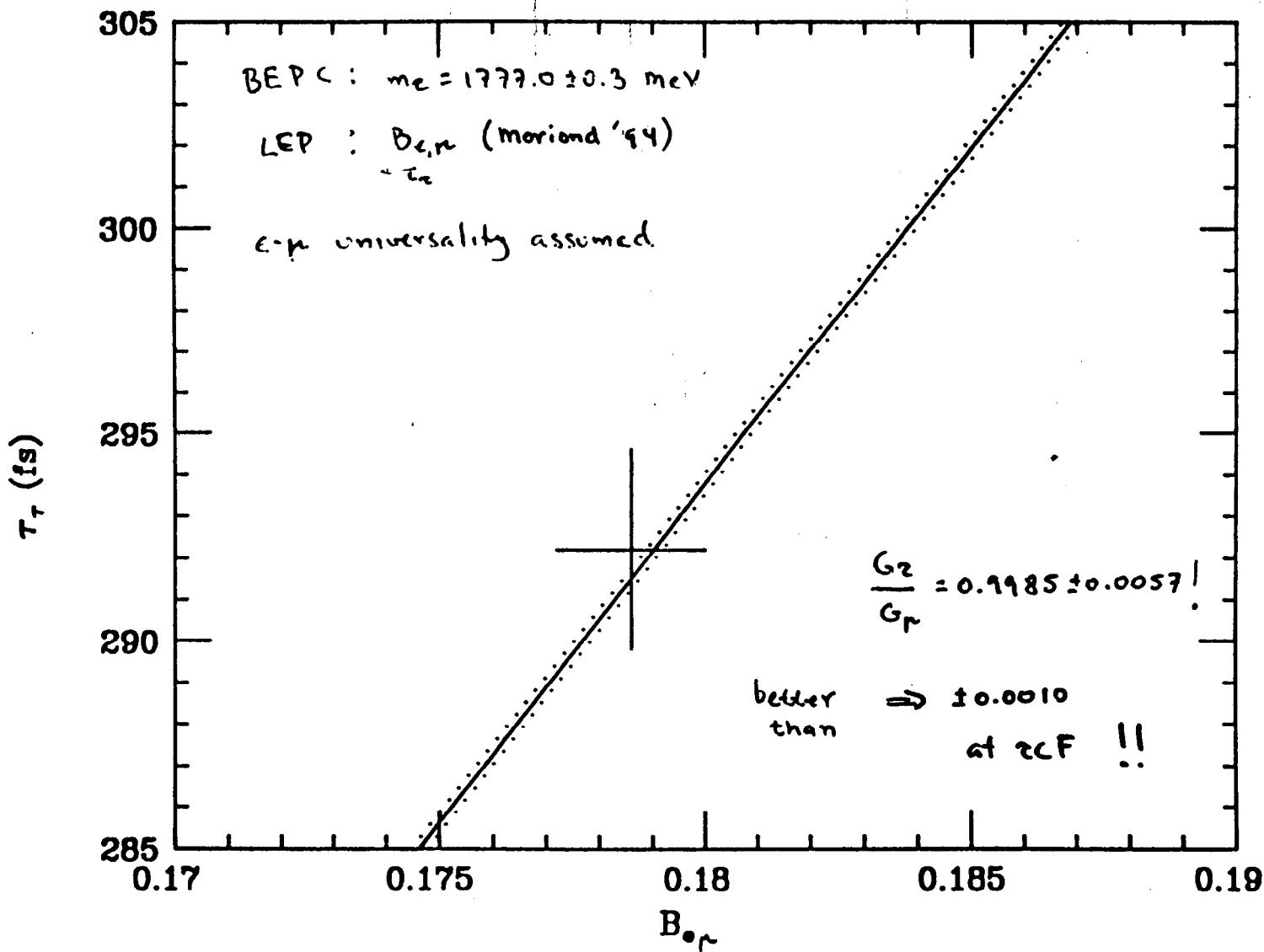
estimate $\xrightarrow{\text{known}}$

very well known QED
correction

This is the definition of $\underline{\underline{G_e}} ! !!$

BEPBC : DPF'94

$1776.96^{+0.18+0.20}_{-0.19+0.16}$



*

Patterson e Glasgow

$$\left. \begin{array}{l} m_\tau = 1777.00 \pm 0.26 \\ z_v = 290.8 \pm 1.5 \\ B_e = 17.66 \pm 0.11 \\ R_\mu = 17.18 \pm 0.12 \end{array} \right\}$$

$$\rightarrow 0.9956 \pm 0.0035$$

$\sim 1\sigma$ low

Patterson & Glasgow

v_τ HELICITY

- CONSIDER $\tau \rightarrow \pi v$

$\tau^- \rightarrow \nu_\tau \leftrightarrow \pi^-$ LEFT-HANDED ν_τ FAST π^-
 $\tau^- \rightarrow \pi^- \leftrightarrow \nu_\tau$ RIGHT-HANDED ν_τ SLOW π^-

- TAU HELICITY CORRELATION

$\tau \leftrightarrow \tau$ FROM SPIN 1 γ^*, Z

MOMENTUM CORRELATIONS IN $\pi^- \nu_\tau \pi^+$ EVENTS
 CAN YIELD ν_τ HELICITY

ALEPH

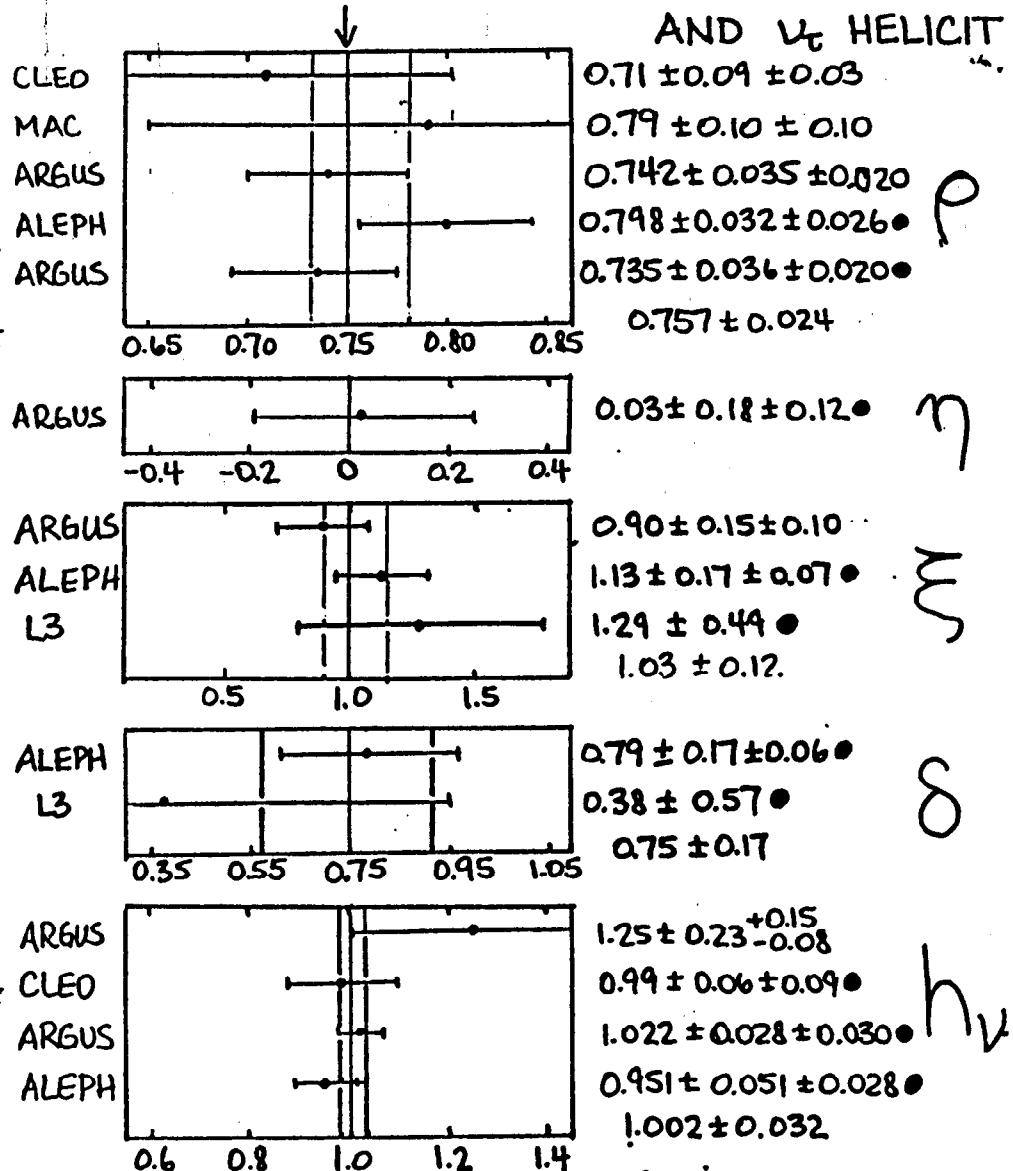
- ISOLATE $\tau\tau$ SAMPLE: $e \nu \mu, e \nu \pi \dots$
- SIMULTANEOUS FIT FOR $p, \eta, \epsilon, \delta, h_{\nu\tau}, p_T$

ARGUS

- USE $\tau \rightarrow \nu \nu$ VS. $\tau \rightarrow \nu \nu$ ENERGY AND ANGULAR CORRELATIONS TO EXTRACT $h_{\nu\tau}$

CLEO

- USE $\tau \rightarrow \pi \nu$ VS $\tau \rightarrow \pi \nu$ TO EXTRACT $h_{\nu\tau}$ [2]



• = NEW

LEP P_T data is a major source of the conflict with SLD

over $\sin^2 \theta_{\text{eff}}$: $P_T^{\text{LEP}} (\text{Glasgow}) \rightarrow x_{\text{eff}} = 0.2325 \pm 0.0009$
 $\text{ALR (SLD)} \rightarrow " = 0.2294 \pm 0.0010$

However: LEP actually measures $\begin{cases} \bar{s} \cdot P_T & \text{for } \tau \rightarrow e \\ h \cdot P_T & \text{for } \tau \rightarrow \pi, \eta, \eta' \end{cases}$

+ Assumes V-A for $\tau \nu_L W$; i.e., $\bar{s} = h = 1$.

{ ... which has prompted them to measure $\bar{s}, h, \delta, g, \dots$!! }
 (double correlations)

If f is the 'weight' of the $\tau \rightarrow e$ data, then

$$\text{"} P_T \text{"} = \frac{2(1-4x_0)}{1+(1-4x_0)^2} = k \cdot P_T^{\text{true}} = k \cdot \frac{2(1-4x)}{1+(1-4x)^2}$$

↑
 'apparent'
 averaged result ↓
 V-A Hypothesis
 $x_{\text{eff}}^{\text{true}} = 0.2123 \pm 0.0009$

↑ actual x_{eff}

where: $k \equiv f \bar{s} + (1-f)h$, $f_{\text{measured}} = 0.2868$

(no new result
but probably
 ~ 0.25 now)

$$\therefore k < 1 \rightarrow x_{\text{eff}} < x_0$$

'model-independent' approach \rightarrow take \bar{s}, h from data +
allow 1σ variation.

$$\rightarrow x_{\text{eff}} = 0.2311 \pm 0.0009 \quad [A_{FB}^L \rightarrow 0.23107 \pm 0.00090]$$

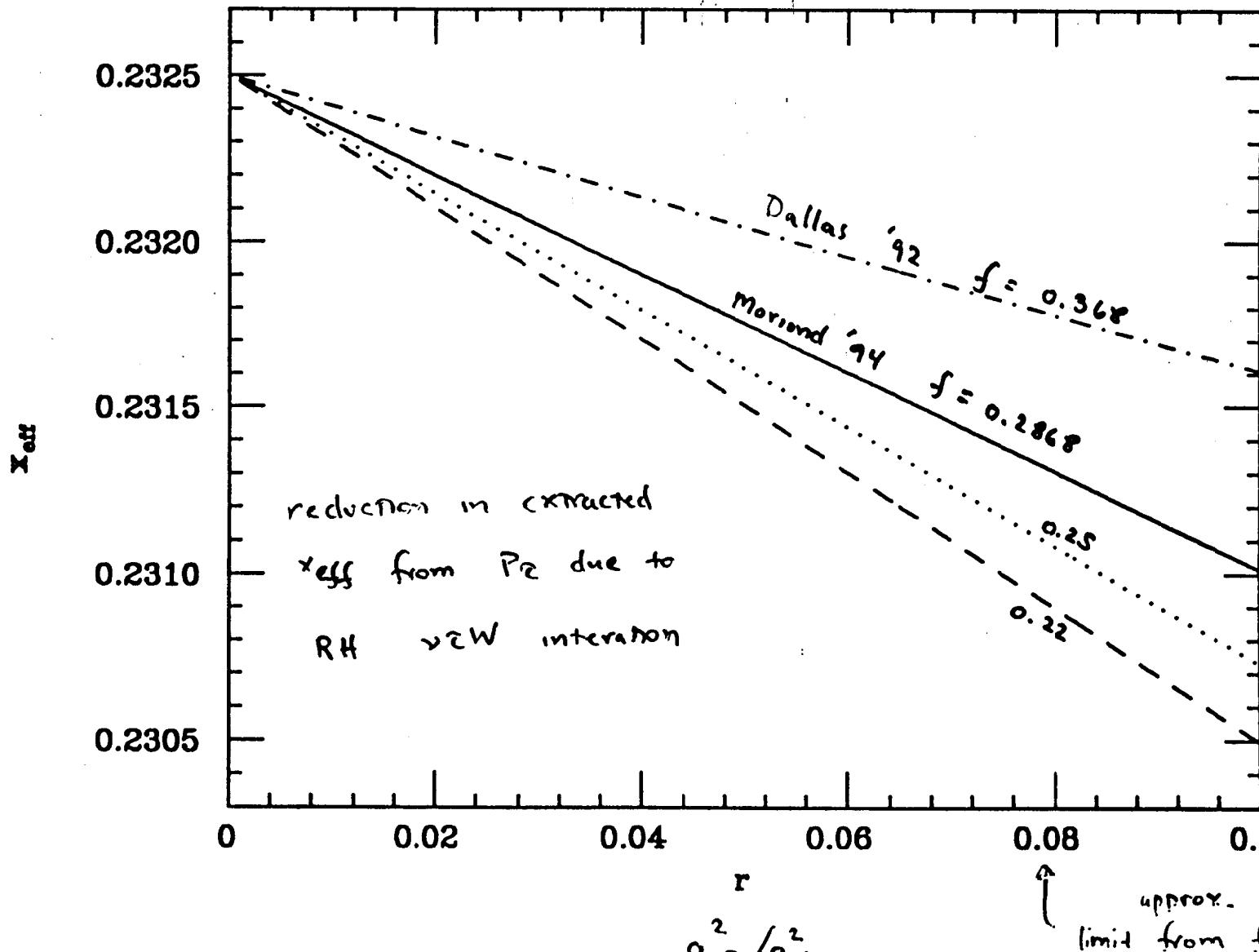
'Toy model': i) $e \rightarrow \mu$ universality + purely LH CC's
 ii) τ may have a RH piece
 iii) $e \nu \tau$ univ. in Ξ couplings...

$$J_L = \frac{4G}{\sqrt{2}} J_2^+(L) [g_{LL} J_z(L) + g_{LR} J_z(R)] + \text{h.c.}$$

$$\rightarrow r = g_{LR}^2 / g_{LL}^2 \quad \underline{\text{small}}$$

		<u>WA</u>
michel's change ...	$\zeta = \frac{3}{4} (1+r)^{-1} < 3/4$	✓
	$\delta = \frac{3}{4} (1+3r)^{-1} < 3/4$	✗ (OK, with 4S)
	$\xi = (1+3r)/(1+r) > 1$	✓ ←
	$h = (1-r)/(1+r) < 1$	✓ ←
	$\eta = 0$	✓

- Use r to parameterize $h, \xi \neq 1$ (Plot)
- RHC can lower x_{eff} extracted from P_T !!!
perhaps to 0.2311 ± 0.0009
- LEP will try to get better h, ξ (ζ, δ, η) info
- BUT : CCF will remove this problem completely
- $\Rightarrow \Delta \xi \rightarrow \pm 0.025 \quad \Delta h \rightarrow \pm 0.004 !$
- Great in itself, but will clarify precision EW measurement
• help with LEP/SLD data comparisons...
may be even more important as precision improves
to $\Delta \sin^2 \theta_W = \pm 0.0003$ level



↑
approx.
limit from fit
to existing data.
(Michel's)

$\{SLD\}$
 $\{LEP\}$ data \oplus $e-\mu$ universality imply

$$\Rightarrow \left\{ \begin{array}{l} \frac{P_\tau}{P_{\nu\mu}} = 1.0048 \pm 0.0045 \quad (!) \\ \frac{P_\tau}{P_e} = 0.9591 \pm 0.0744 \end{array} \right. \quad \begin{array}{l} \frac{A_{FB}^e}{A_{FB}^{\nu\mu}} = \frac{A^e}{A^{\nu\mu}} = 1.367 \pm 0.169 \\ \uparrow \qquad [1.0253 \pm 0.0681] \\ A_{FB}^e = \frac{3}{4} A_e A_\tau \quad \underline{\text{assumed}} \end{array}$$

(not neces. $= A_\tau$!)

- For discussion I will assume $\tau e \tau \nu$ is LH.
- LEP collaborations also assume $Z \bar{e} e$ is V,A ... (plot)

$\Rightarrow \frac{P_\tau}{P_{\nu\mu}}$ generally provides the best constraints on τ Universality violating models

Some Possibilities

- i) τ mixes with some other, 'exotic' $Q=-1$ fermion (E) thru angle ϕ $\{ \phi_{L,R} \leftrightarrow s_{L,R}, c_{L,R} \}$ - 'old standby'

Some Typical Cases	T_{3L}	T_{3R}	
	0	0	vector-like singlet
	0	$-\frac{1}{2}$	mirror fermion
	$-\frac{1}{2}$	$-\frac{1}{2}$	vector-like doublet
	-1	0	isotriplet

- $s_L \neq 0$ REDUCES G_F/G_F \rightarrow $s_L \lesssim 0.10$ $\text{at } 95\% \text{ CL}$

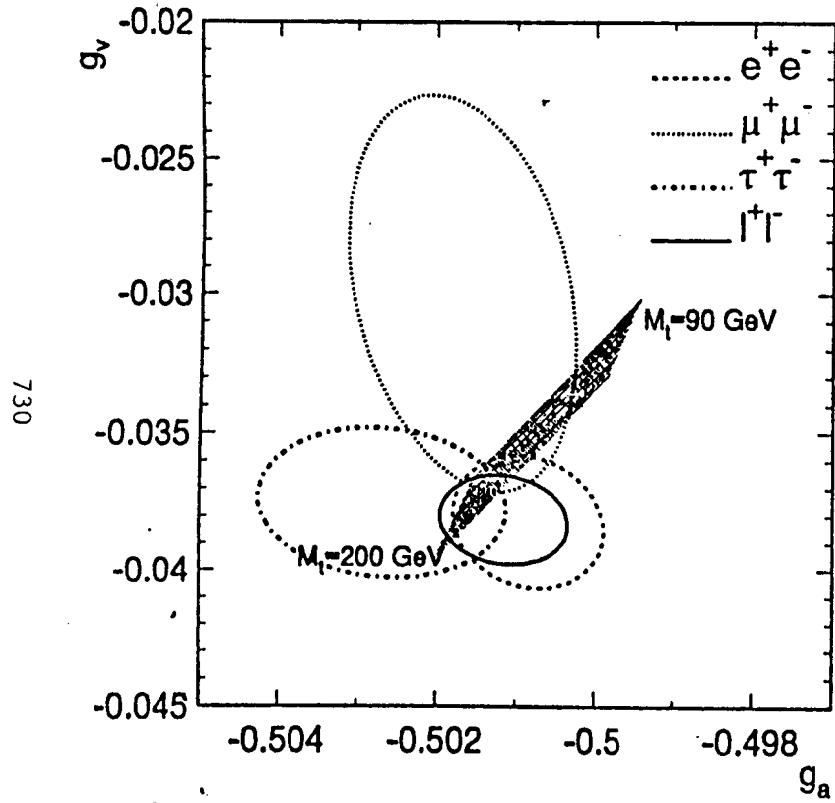
s_R has no direct CC effect .. assume $s_L = s_R = s \dots$

Schaile e Glasgow '94

Combined fits

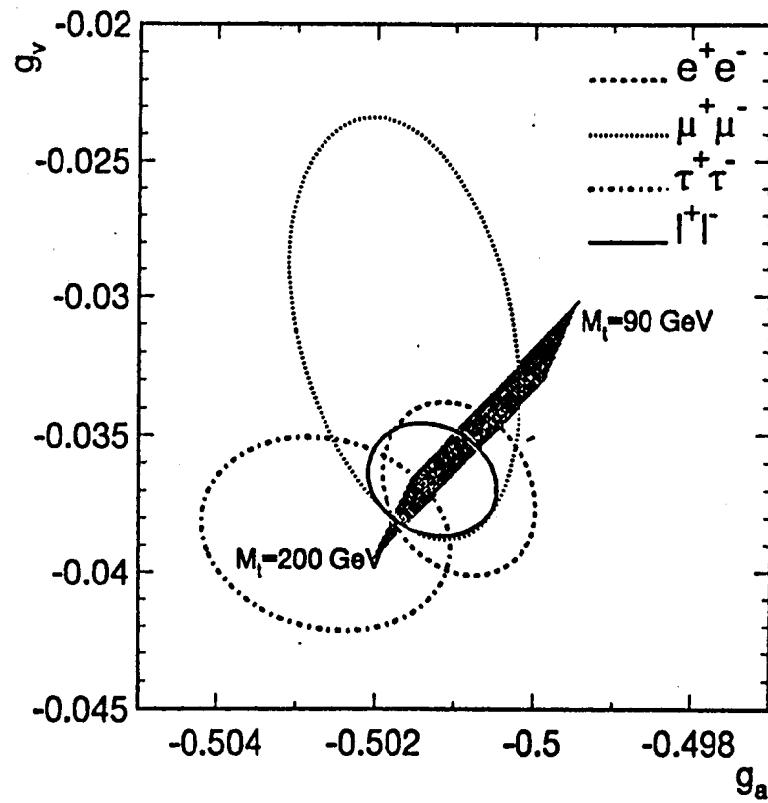
LEP + SLC

68% C.L.



LEP only

68% C.L.



Γ_e , A_{FB}^L , P_τ , A_{LR}

Assumes $V-A$ $\tau u w$ coupl.
and $Z e^+ e^-$ is v, a only
(factorization)

$$\left. \begin{aligned} v_\tau &= \left[-\frac{1}{2} + 2x \right] + s_L^2 (T_{3L} + \frac{1}{2}) + s_R^2 T_{3R} \\ a_\tau &= \left[-\frac{1}{2} \right] + s_L^2 (T_{3L} + \frac{1}{2}) - s_R^2 T_{3R} \end{aligned} \right\} \begin{array}{l} \text{mixing-modified} \\ \text{Z}\tau\tau \text{ couplings} \\ (\text{Figs}) \end{array}$$

$$\rightarrow \left. \begin{array}{ll} T: \lesssim 0.08 & VS: \lesssim 0.05 \\ m: \lesssim 0.03 & VO: \lesssim 0.05 \end{array} \right\} \text{better than CC...}$$

ii) τ has an 'anomalous-magnetic moment'-type coupling to Z ($\sim \frac{g}{2c_W} \frac{x_\tau^2}{2m_\tau} \bar{\tau} \sigma_{\mu\nu} \gamma^\nu \tau Z^\mu$)

- old idea for $Z \rightarrow \gamma$ $\left. \begin{array}{l} Z \rightarrow \tau\tau\gamma: |k_\tau| < 0.11 \\ \text{TRISTAN better!} \end{array} \right\}$ trifolts + minder (TDR)
- 'positive' indication for $Z \rightarrow b\bar{b}$
(another talk) (no 'CC' influence)

- fit to Γ , $P_\tau + A_{FB}^{\tau\tau}$ w/ $e\text{-}\mu$ universality

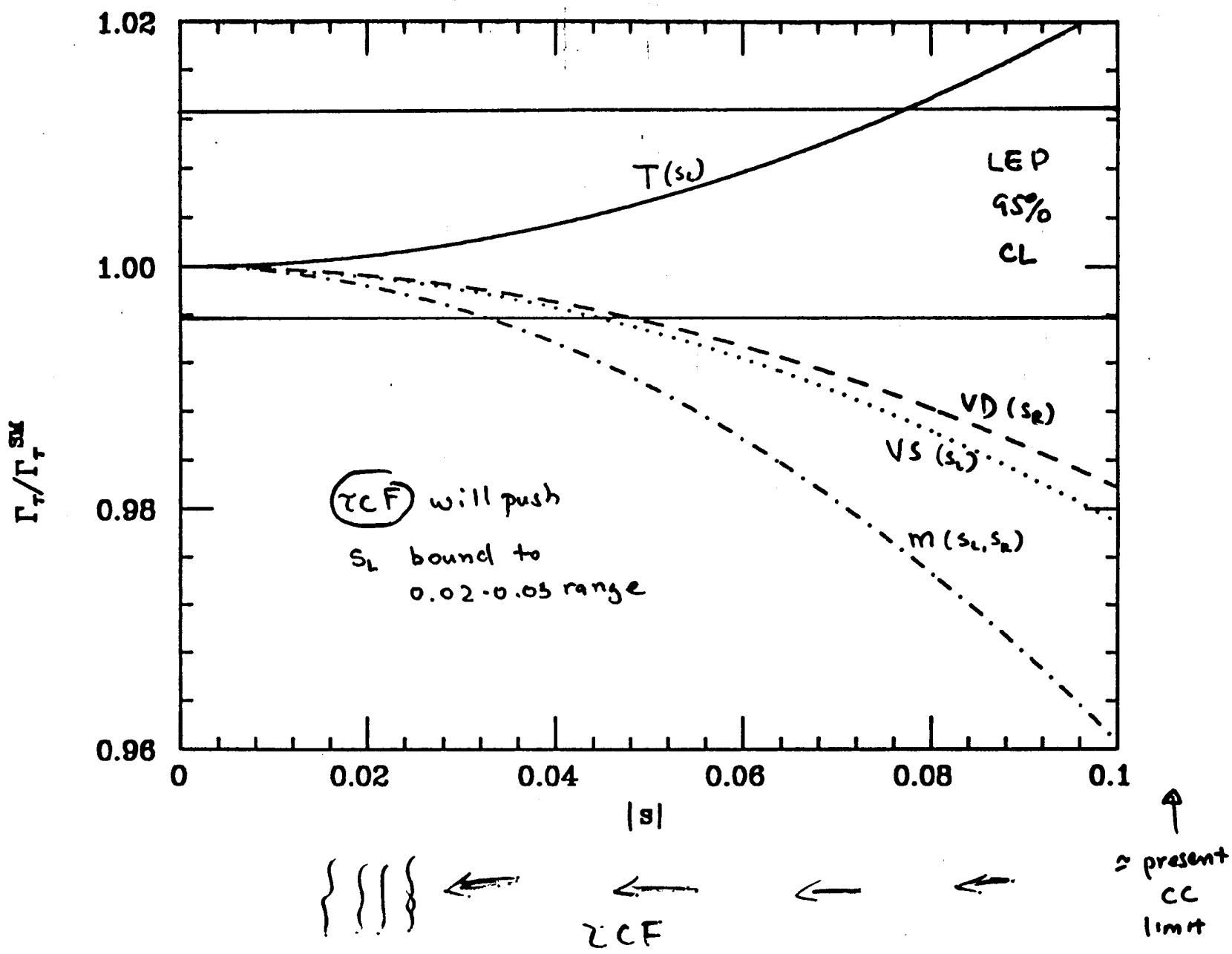
$$\rightarrow x_\tau^2 = (-1.3^{+4.1}_{-1.5}) \cdot 10^{-3} \quad \in 95\% \text{ CL}$$

$$[0(\alpha/2\pi)]$$

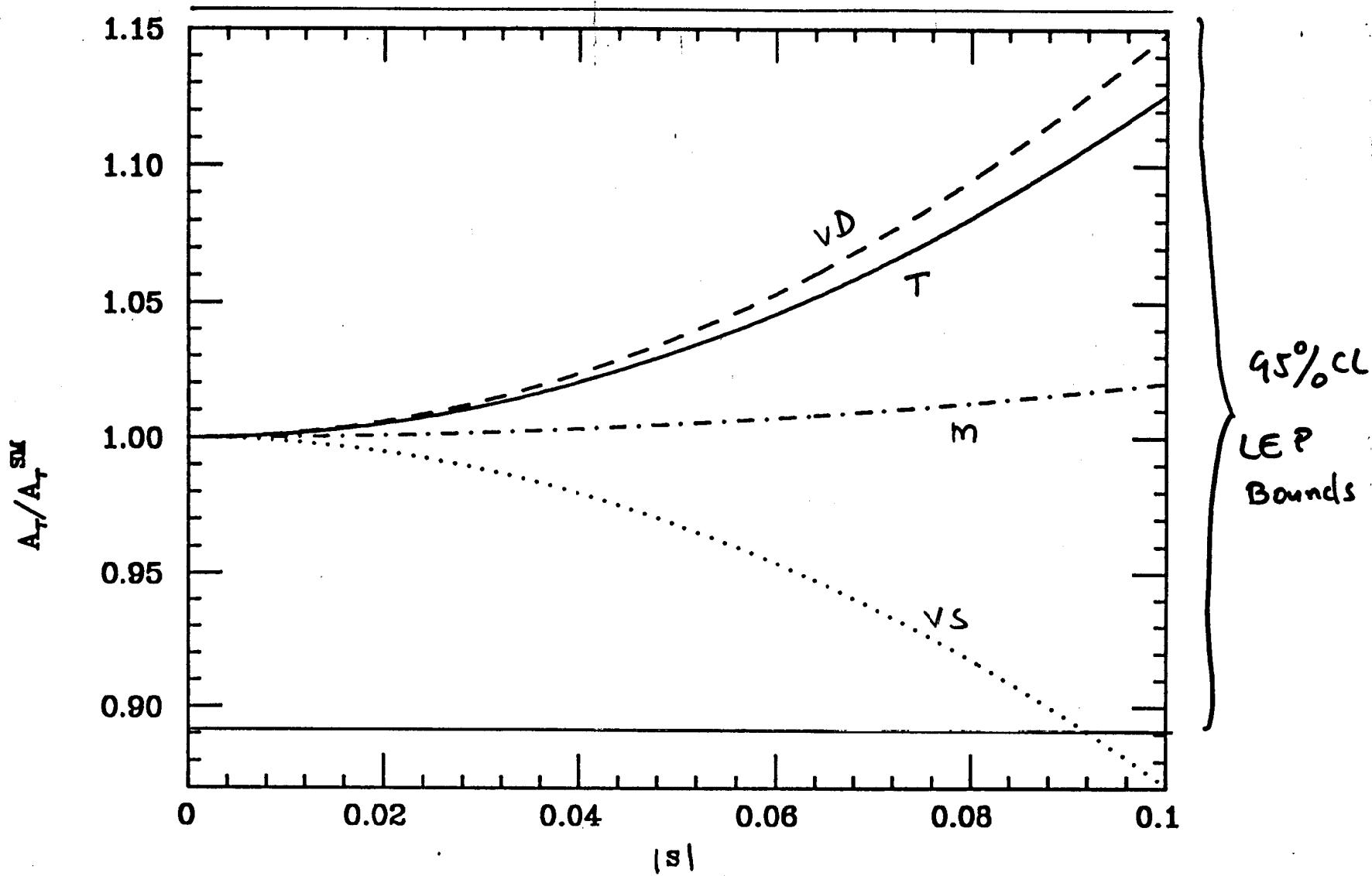
- As before ... limit ^{mostly} dominated by $\Gamma_{\tau, e, \mu}$ data

~ one order of magnitude BETTER than NEW k_τ^{γ}
bound from TRISTAN data.. (hi-precision $\tau\text{-}F$ may surpass TRISTAN)

Note: leptonic asymmetries for both scenarios still not quite competitive with Γ 's $\{(\pm P_\tau) \text{ for } x_\tau^2\}$
more important for x_τ^2

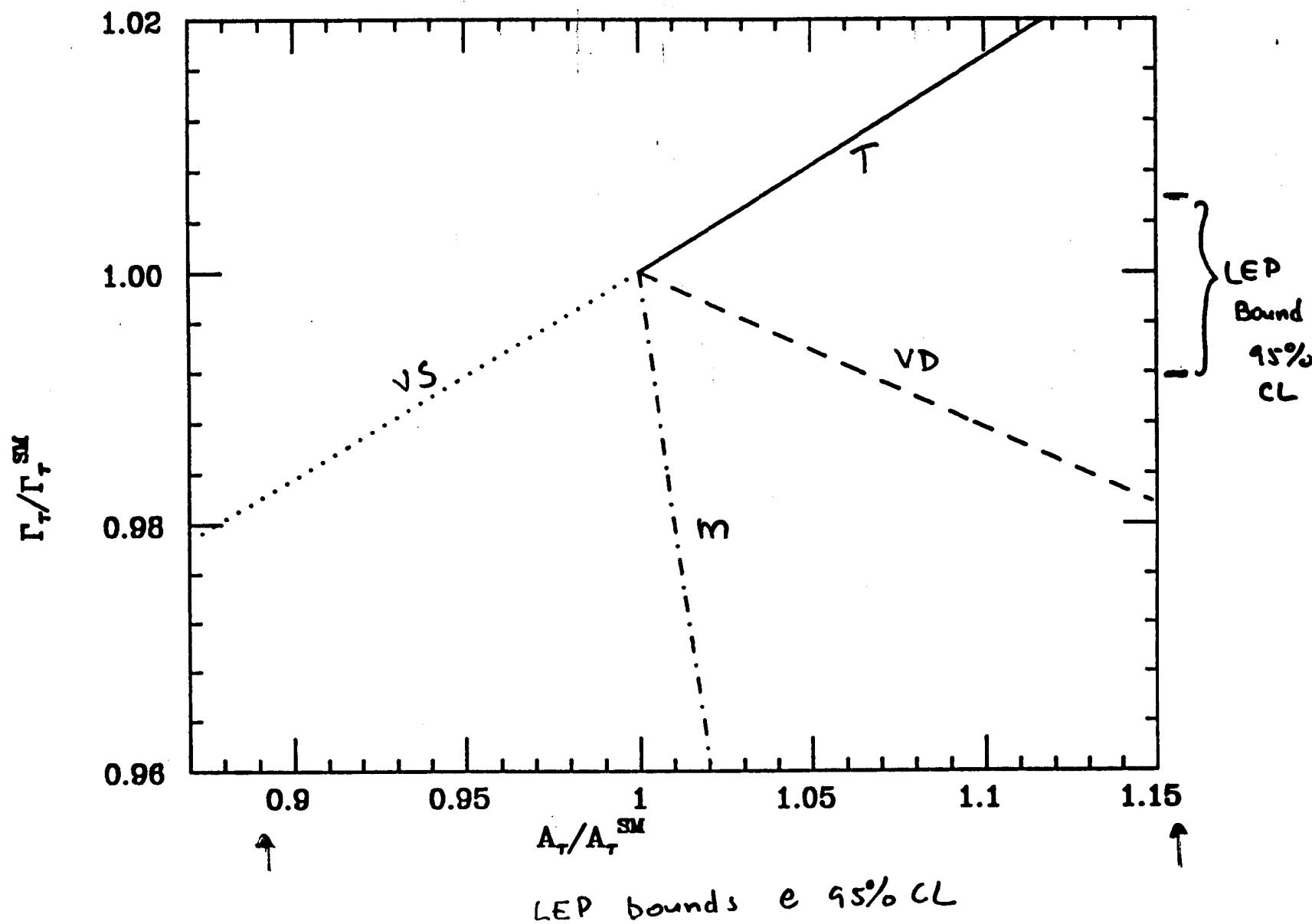


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no bounds yet....

734



Note that in the quark sector the limits on the mixings of the b are surprisingly stringent. This is because both s_L^b and s_R^b tend to decrease the theoretical value of the partial decay width $Z \rightarrow b\bar{b}$ ($\Gamma_b \simeq [\Gamma_b]_{SM}(1 - 2.3s_L^b)^2 - 0.4s_R^b)^2$) thus worsening the already existing $\sim 3\sigma$ discrepancy between the SM expectation and the experimental value.

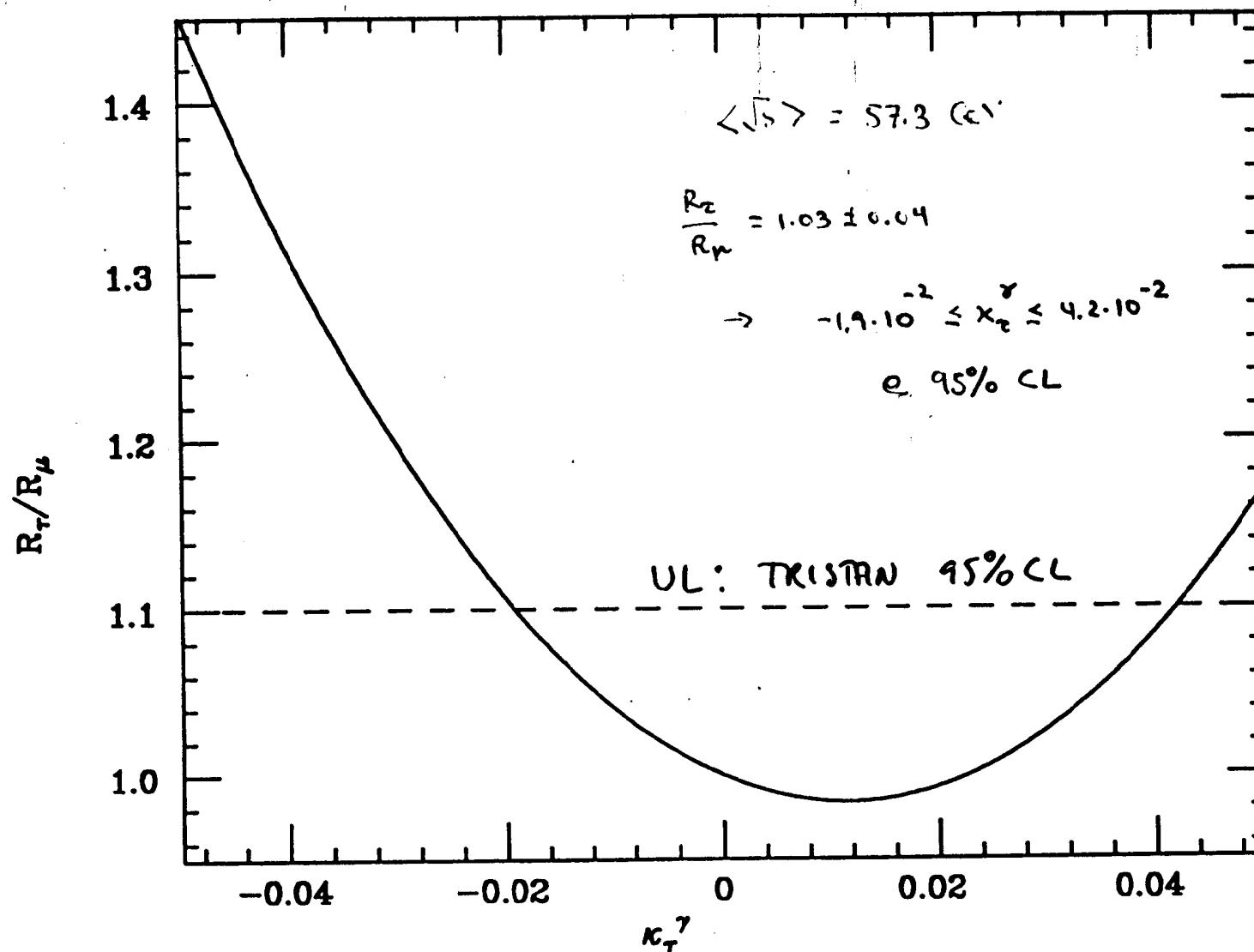
Table 3: 90% C.L. upper limits on the ordinary-exotic flavor diagonal mixing angles for individual fits (one angle at a time is allowed to vary) and joint fits (all angles allowed to vary simultaneously). Only the results for $\Lambda = 2$, corresponding to neutrinos mixed with heavy singlets are shown. Some of the results are preliminary. A complete list of final results will be given in Ref. [9]. In most cases the observables more effective in setting the constraints are the LEP measurements of partial widths and asymmetries²⁰. The limits have been derived assuming the values $m_t = 170$ GeV and $M_H = 200$ GeV for the top and the Higgs mass in the theoretical expressions.

	Individual	Joint		Individual	Joint
$(s_L^e)^2$	0.0016	0.0054	$(s_L^u)^2$	0.0022	0.012
$(s_R^e)^2$	0.0020	0.0018	$(s_R^u)^2$	0.010	0.023
$(s_L^\mu)^2$	0.0013	0.0049	$(s_L^d)^2$	0.0026	0.016
$(s_R^\mu)^2$	0.0019	0.0040	$(s_R^d)^2$	0.0066	0.019
$\rightarrow (s_L^t)^2$	0.0011	0.0037	$(s_L^s)^2$	0.0036	0.019
$\rightarrow (s_R^t)^2$	0.0018	0.0034	$(s_R^s)^2$	0.021	0.59 ?
$(s_L^{v_e})^2$	0.0053	0.0053	$(s_L^c)^2$	0.0044	0.024
$(s_L^{v_\mu})^2$	0.0020	0.0052	$(s_L^b)^2$	0.0017	0.0031
$\rightarrow (s_L^{v_\tau})^2$	0.0055	0.017	$(s_R^b)^2$	0.0091	0.015

3.3. Limits on a Z_1 from E_6 from ordinary-ordinary fermion mixing

E_6 GUTs are well known examples of theories where additional fermions and new neutral gauge bosons are simultaneously present. For a general breaking of E_6 (rank 6) to the SM (rank 4) it is possible to define a whole class of Z_1 bosons corresponding to a linear combination of the two additional Cartan generators. We parametrize this combination in terms of an angle β . Fermions are assigned to the fundamental 27 representation of the group which contains 12 additional states for each generation, among which we have a vector doublet of new leptons $(N \ E^-)_L^T$, $(E^+ \ N^c)_L^T$. Non-diagonal mass terms with the standard $(\nu \ e^-)_L^T$ and e_L^c leptons will give rise respectively to ordinary-ordinary and ordinary-exotic mixings, and in particular will induce LFV L and R chiral couplings between the first and second generation, allowing for LFV processes as $\mu \rightarrow eee^{12}$ and $\mu-e$ conversion in nuclei¹³. Very stringent experimental limits exist for both these processes^{18,24}. Due to the expected suppression of the ordinary-exotic mixings, the LFV couplings in the R sector can be neglected. We stress that the assumption that the only source

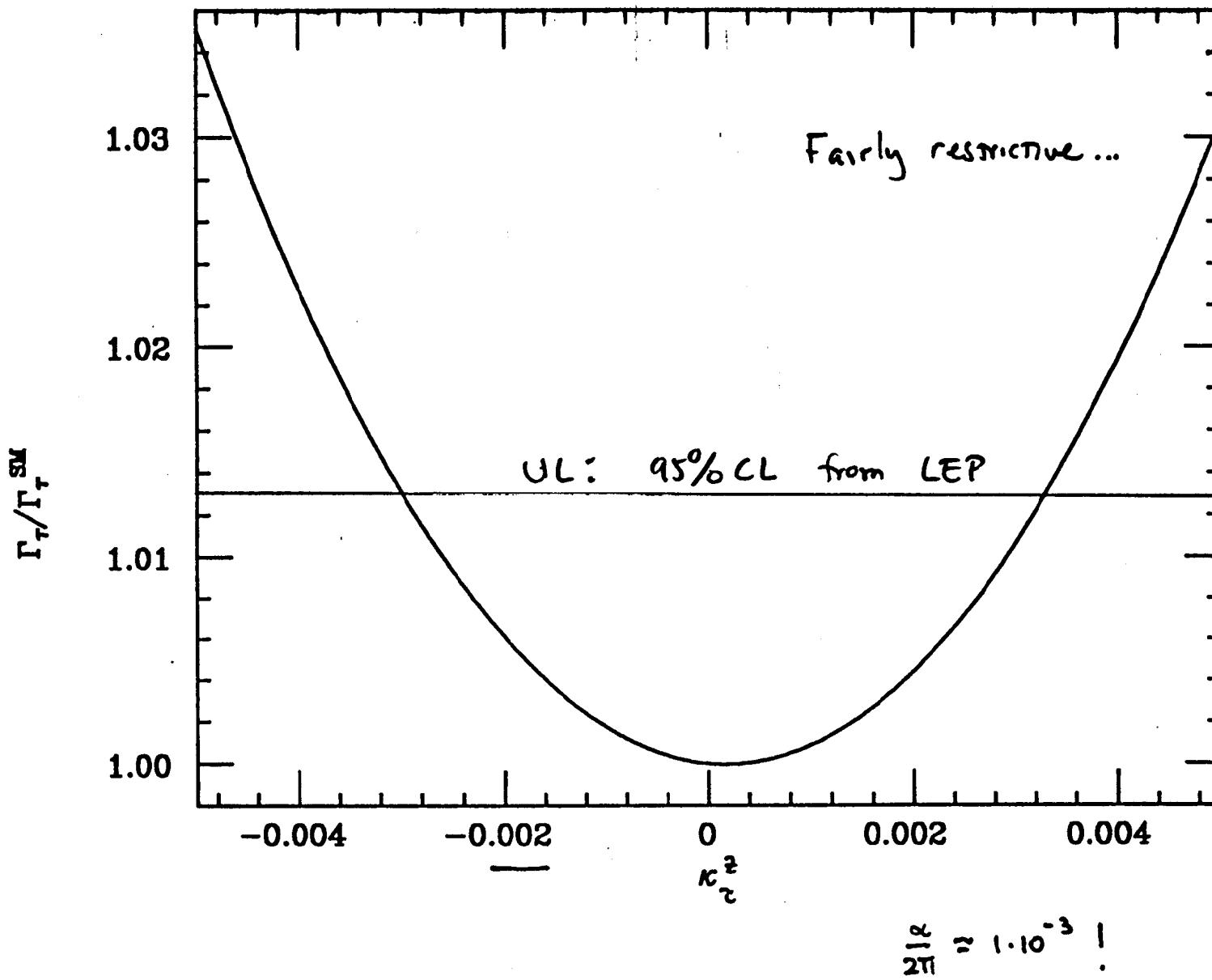
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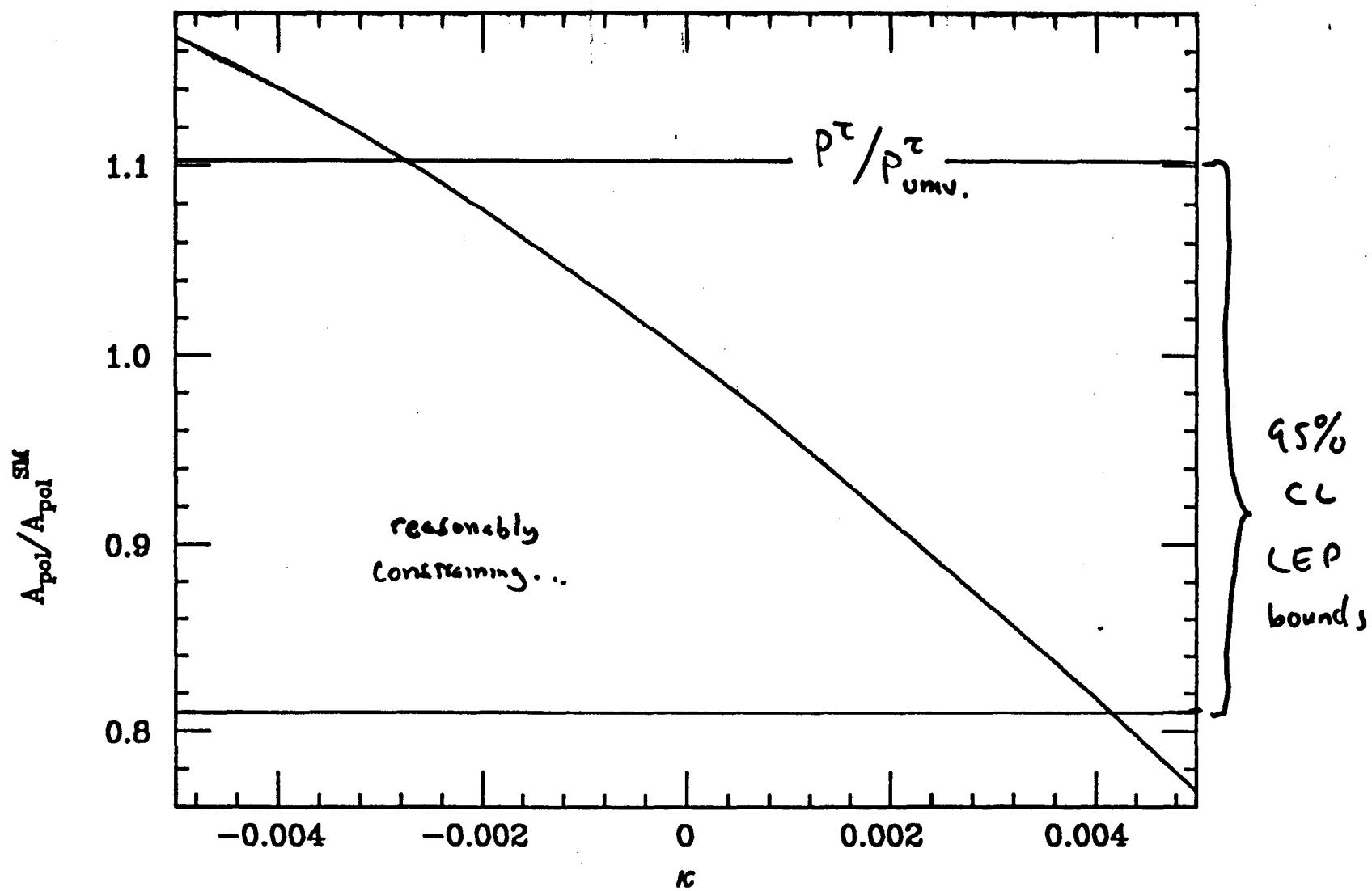


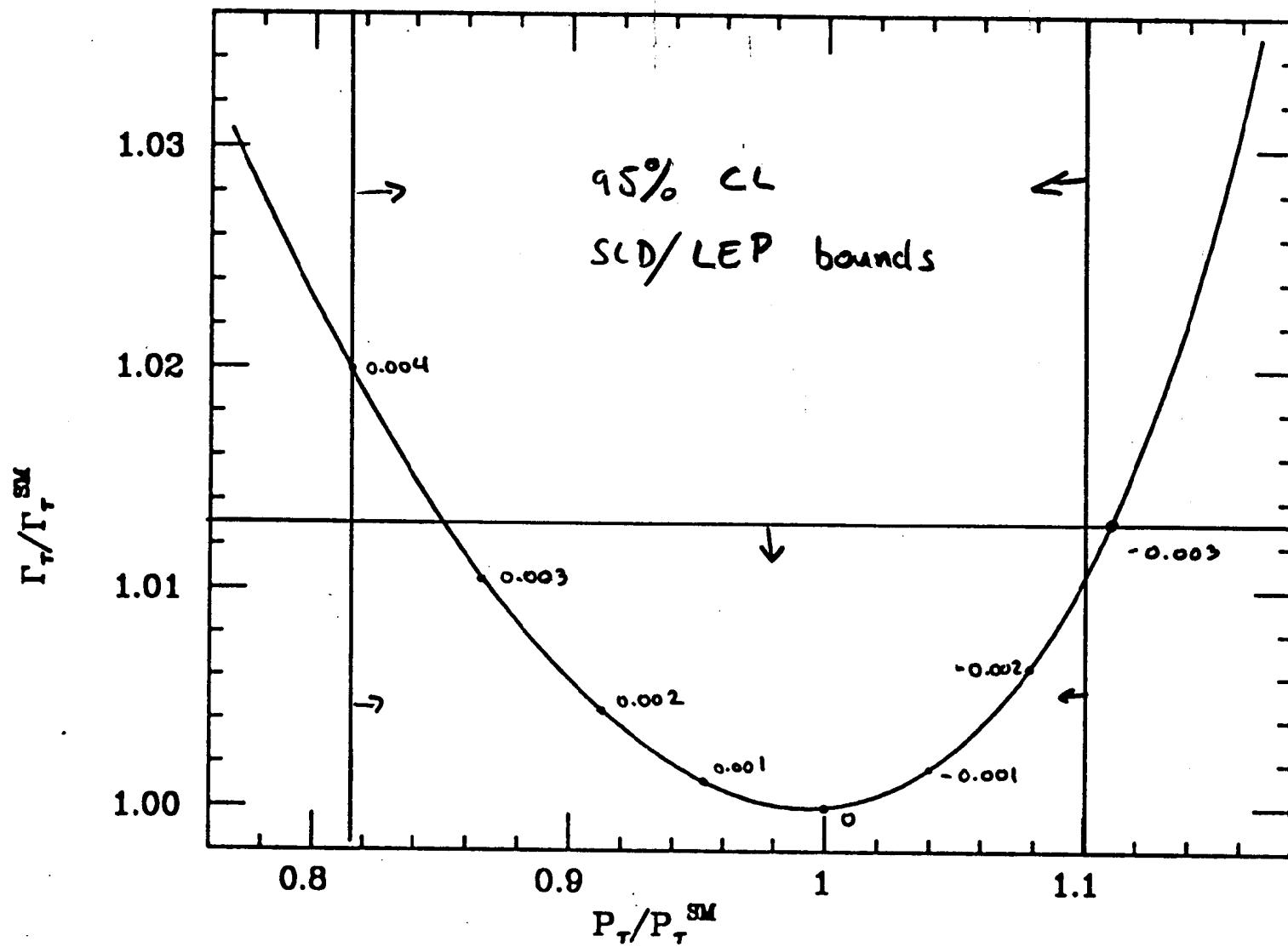
hi-star TCF may be
competitive

$\alpha_{\pi\pi}^2$
 $|x_\pi^7| < 0.11$ e 95% CL

737







- (iii) $Z-Z'$ mixing w/ Z' having generation dependent couplings
 - τ and b couplings changed relative to first two generations {Li + Ma}

$$\Gamma_i \sim N_i^2 \left\{ 1 + (1 - 4x_i(1/3))^2 \right\} \text{ with } \begin{cases} x_{12} = \frac{x}{N_{12}} \\ x_3 = \frac{x}{N_3} \end{cases}$$

- $N_{12}^2 \sim 1 + 0.0077 \cdot \left(\frac{\epsilon}{0.008}\right) (1-p) \cdot p \geq 1 \Rightarrow \frac{G_\tau}{G_\mu} = 1 + \epsilon$
- $N_3^2 \sim 1 - 0.0077 \cdot \left(\frac{\epsilon}{0.008}\right) (1-p)^2 \leq 1 \Rightarrow \frac{G_\tau}{G_\mu} = 1 + \epsilon$

$$\therefore x_3 \geq x \geq x_{1,2} \quad \& \quad \Gamma_3 \leq \Gamma \leq \Gamma_{1,2} \therefore \begin{array}{l} \Gamma_b / \Gamma_b^{\text{SM}} < 1 \\ \Gamma_\tau / \Gamma_\tau^{\text{SM}} < 1 \text{ hmm} \\ \text{Not likely} \end{array}$$

$x_2 > x_{e,\mu} \checkmark ? \quad \{0 \leq p \leq 1\}$

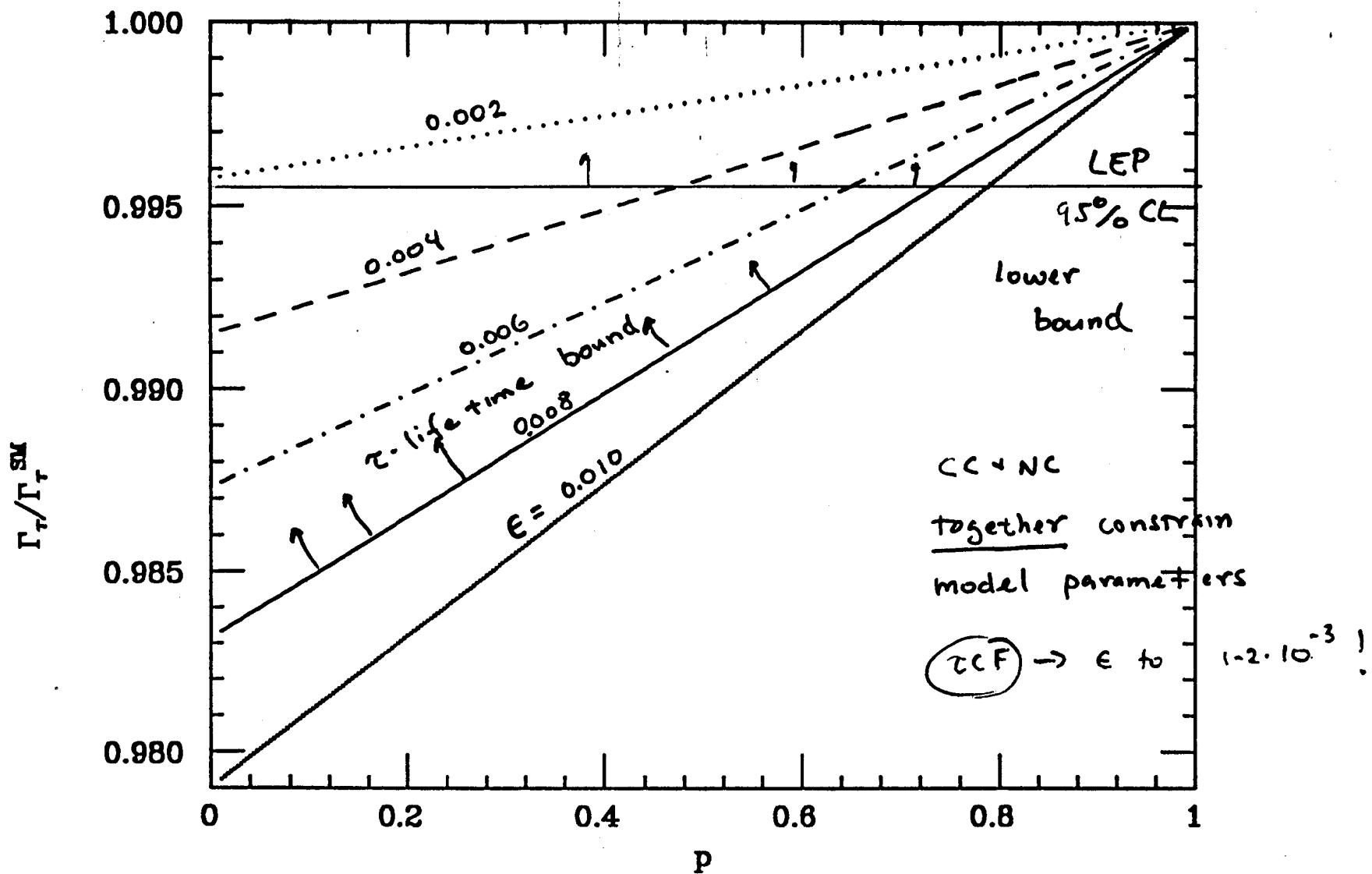
→ Pushes Γ_b (Γ_τ) 'right' (wrong) way ∴

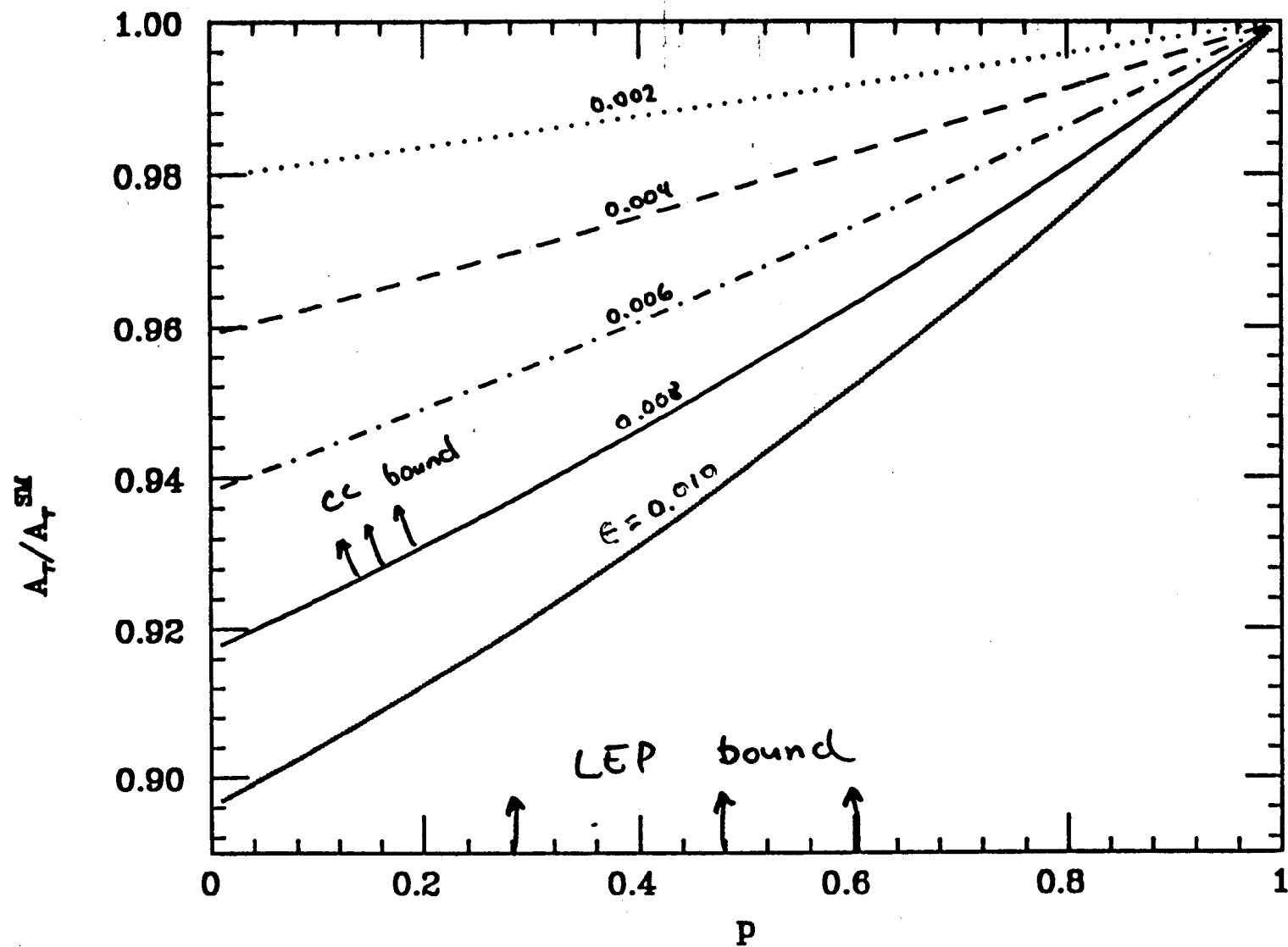
Γ_τ highly restrictive ; not asymmetries

- Strong interplay between precision $\overline{\text{CC}} + \text{NC}$ constraints
 \hookrightarrow TCF will force $\epsilon \rightarrow \sim (1-2) \cdot 10^{-3}$ level...
- Remember too (but beyond scope of THIS talk...)

Many non-universality models also predict FC τ decays, e.g., $\tau \rightarrow \mu e \bar{\nu}$, $\mu \nu \bar{\nu}$, $3e, 3\mu$, etc., at some level!

L





Conclusions

- i) Both L.E.E. + LEP/SLD are closing-in on the possibility of τ non-universality by probing the $\tau v_\tau w + 2\tau\bar{\tau}$ couplings - complementary! • CC vs NC •
- ii) However some models 'predict' such small effects that only TCF can see them... still lots of room
- iii) τ -Universality + assumptions about τ 's have influence elsewhere - test of EW RC in SM \leftrightarrow 'SLD vs LEP' perhaps only resolvable with TCF...
- iv) Non-universality + FC τ decays are linked in many models ∴ hi-lumi required beyond LEP to be visible
- v) only TCF can probe BSM loop effects ... necessary to truly establish universality