The TAO of the τ with CLEO II

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The TAO of the τ with CLEO-II

In a perfect world, this talk would be a complete review of present results and projections for:

- 1. Particle properties in the tauonic sector
 - (a) Measurements of the τ mass (1777.8±0.7±1.7 MeV-CLE092)
 - (b) Limits on the τ neutrino mass (CLEQ/ARGUS: <3! MeV)
 - (c) Measurements of the τ lifetime (-LE094): (2.91±0.03±0.07) (1 vs.1/IP: (2.94±0.07±0.12)10⁻¹³9 /vs 3 v/z: (2.91±0.04±0.07)
- 303.3 (2.ES=0.13=0.10) 2. Leptonic Decays (a) precision tests of Standard Model - $e\mu$ universality $\begin{cases} z \rightarrow e\nu\nu \\ \varphi \neq e\nu\nu \\ \varphi \neq$

 - (b) tests of Lorentz structure: V A and weak current parameters: •
 - i. Michel parameter ρ from dN/dx
 - ii. Low energy parameter η
 - iii. Polarization parameter ξ'
 - iv. Decay asymmetry parameters ξ and δ CLE094: USE でラボーそ い、てキラガキシ

> = hy: Find hy: -0.99 ± 0.06 ± 0.10 (Y helicity) (V-A expect 10)

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-14+G1C



Tau data 1322 events

$$\mathcal{L} = 1.1 \text{ fb}^{-1}$$

Tau HC 577 $\pi \cdot \pi$ 88%
58 $\pi \cdot \mu$ 10%
16 $\pi \cdot \rho$ 2%
Humung MC 2 pass, 99700 start. $\mathcal{E} = 2 \times 10^{-5}$
 $\sigma = 0.91 \text{ nb}$
 $\mathcal{E}\sigma\mathcal{L} = 20 \text{ events.}$
 $\delta\delta \rightarrow \pi^{\dagger}\pi^{-1}MC$ 9 pass, $\mathcal{I}^{\text{MC}} = 23 \text{ fb}^{-1}$
 $9 \times \frac{1.1}{23} = 0.4 \text{ events.}$
 $\delta\delta \rightarrow \mu^{\dagger}\mu^{-1}MC$ 0 pass, $\mathcal{I}^{\text{MC}} \simeq 170 \text{ pb}^{-1}$
Bottom line

$$\pi - \pi \sim 85\%$$

 $\pi - \mu \sim 10\%$
 $\pi - \rho \quad few\%$
non-tau $few\%$

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10. pl=1 for 2 // 12 consume (gagi)? 3" Z=3TT give sign of Lelicity



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The ν_{τ}

• the ν_{τ} has yet to be directly observed!

its existence is inferred from the energy spectrum of observed charged leptons or hadrons in decay products of τ

• it is distinct from ν_e , ν_{μ} :

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• ν_e , ν_μ beams are not observed to interact via charged current $\nu_l N \rightarrow \tau^- X$ in Fermilab E531

• $\Gamma(Z^0 \rightarrow \text{unobserved}) / \Gamma(Z^0 \rightarrow \nu \nu) = 3.00 \pm 0.05$

3. Semi-hadronic decays

- (a) spectral functions strange and non-strange (Eudence for 3h = (9.82 + 0.09 + 0.34)%; 3h (m) 4.25 + 0.09 + 0.26)% E-3p %; e.g.)
- (b) tests of QCD extraction of α_s ($\propto_s(M_2) = 0.14 \pm 0.003 \pm$
- (c) strange resonance dynamics production mechanism for KKX final states
 - i. Understanding K_1 states (connection to HQET)
 - ii. KK π from K^*K or $\rho\pi; \rho \to K\bar{K}$? Techer of al. $\tau \to t^*$.
- (d) Comparison with predictions of CVC \uparrow
- 4. Rare/forbidden decays
 - (a) lepton family number violation, neutrino mixing
 - (b) Second-class currents, etc., physics beyond the Standard Model! ーイングロン
- 5. Nostradamus comes to CLEO looking into the future, brightly
 - (a) Detector design of CLEO-III τ physics into the next millenium
 - (b) What are the limiting systematics of doing tau measurements?









 $\mathbf{\hat{D}}_{3}^{\text{CLEO}} = 0.1455 \pm 0.0013 \pm 0.0059$

• Another important lepton family number violating decay: $\tau \rightarrow \ell \ell \ell$, where the final leptons are either e^{\pm} or μ^{\pm}

• providence •

- decays like $\tau^- \rightarrow \mu^- e^+ e^-$ are FCNC at one vertex; decays like $\tau^- \rightarrow \mu^+ e^- e^-$ are FCNC at two vertices!
- SINDRUM 1985: $\mathcal{B}(\mu \rightarrow eee) < 2.4 \times 10^{-12}$
- once again, mass-dependent couplings give interesting range for tau decays: $\mathcal{B}(\tau \rightarrow \ell \ell \ell) \lesssim 3 \times 10^{-6}$
 - ARGUS 1991: $\mathcal{B}(\tau \rightarrow \ell \ell \ell) < (15 20) \times 10^{-6}$ - CLEO II 1993: $\mathcal{B}(\tau \rightarrow \ell \ell \ell) < (3 - 8) \times 10^{-6}$ depending on mode, using 2M $\tau^+ \tau^-$ events

Exclusive semi-hadronic decays: $\tau \rightarrow \pi X$ and CVC

- $\tau^{\pm} \rightarrow \nu_{\tau} h^{\pm} \pi^{0}$ is the largest τ decay mode (most expts can't distinguish π^{\pm} from $K^{\pm} \Rightarrow h^{\pm}$) $\rightarrow \text{Normalize} \quad \tau \rightarrow \chi h^{\pm} n(\pi^{0}), nz2 \quad to \quad n=1$
- proceeds via the weak charged vector current: $W^{\pm} \rightarrow \pi^{\pm} \pi^{0}$
- EW unification: vector part of weak charged current is related to neutral EM current via isospin rotation: CVC
 - CVC relates $\Gamma(\tau^{\pm} \rightarrow \nu_{\tau} \pi^{\pm} \pi^{0})$ to $\sigma(e^{+}e^{-} \rightarrow \pi^{+}\pi^{-})$ radiative corrections have been calculated

• procedure:

- 1. extract isovector part of $\sigma(e^+e^- \rightarrow \pi^+\pi^-)$ up to $s = m_{\tau}^2$
- 2. fit to pion form-factor $F_{\pi}^{I=1}(q^2)$; compare with $F_{\pi}^{I=1}(q^2)$ extracted from $\tau \to \nu_{\tau} W$ charged current
- total $\mathcal{B}(\nu_{\tau}\pi^{\pm}\pi^{0}) \propto \int dq^{2} F_{\pi}^{I=1}(q^{2})$
- Standard Model predicts: $\mathcal{B}(\pi^{\pm}\pi^{0}\nu_{\tau})^{SM} = (24.58 \pm 0.93 \pm 0.50)$ (1st error: $e^{+}e^{-} \rightarrow \pi^{+}\pi^{-}$ data; 2nd: 2% uncertainty in rad-corr)
- Results after subtracting off $K^* \rightarrow K^{\pm}\pi^0$:
- Results after subtracting off $K^* \to K^{\pm}\pi^0$: (World Average) in agreement with CVC (4% precision)
- Aside: $T \rightarrow \overline{\pi}^{o} H \frac{1}{2} : \begin{bmatrix} 0.123 \pm 0.023 \pm 0.023 \end{bmatrix}$ • $m(\pi^{\pm}\pi^{0})$ spectrum from τ decays agrees well based
- with e^+e^- prediction; good evidence for $\rho'(1370)$



FIG. 1. The $M_{\gamma\gamma}$ distribution for the data (points) and the Monte Carlo (histogram), summed over all tags.



FIG. 2. The distributions of a variety of kinematical variables for the data (points) and the Monte Carlo (histogram), summed over all tags. (a) The scaled momentum of the charged particles, p_{π}/E_b (open circles and solid histogram), and the scaled energy of the π^0 , E_{π^0}/E_b (solid circles and dotted histogram), scaled by a factor 3 for clarity); (b) the scaled visible energy, E_{vis}/E_{cm} ; (c) the $\pi^{\pm}\pi^{0}$ invariant mass. The accepted region is to the right of the vertical lines in (a) and (b).

- $\underline{\tau \to K^0 \mathbf{X}}$
- 1. Precise $\mathcal{B}(\tau^- \to (\pi\pi)^- \nu_{\tau})$ and $\mathcal{B}(\tau^- \to (K\pi)^- \nu_{\tau})$ can be used as either
 - a test of the Das-Mathur-Okubo (DMO) sum rule, or
 - a check of $\left|\frac{V_{us}}{V_{ud}}\right|$ (expect a final $\approx 3\%$ error).



2. Study resonant structure of vector and axial-vector $S = \pm 1$ charged current

3. Use the ratio

$$\frac{\mathcal{B}(\tau^- \to K_1^-(1270)\nu_{\tau})}{\mathcal{B}(\tau^- \to K_1^-(1400)\nu_{\tau})}$$

to study the effects of K_1 mixing and $SU_f(3)$ symmetry breaking:





FIG. 6. (a) The $M(K_S^{\bullet}\pi^{-}\pi^{0})$ distribution for data (points) after sideband and background subtraction. The solid histogram is the result of a fit to the Monte Carlo predicted distributions for $K_1(1270)$ and $K_1(1400)$, and a fixed contribution from $K_S^{0}K^{-}\pi^{0}$ (according to the result in [8]). Their contributions are shown as hatched histograms in (b) $(K_1(1270))$, (c) $(K_1(1400))$ and (d) $(K_S^{0}K^{-}\pi^{0})$.

ν_{τ} Mass

- neutrinos are light!
- $m_{\nu_e} < 7.3 \text{ eV}/c^2$
- $m_{\nu_{\mu}} < 0.27 \text{ MeV}/c^2 (90\% \text{ CL})$
- Standard Model assumes massless neutrinos
- most extensions include neutrino mass; "seesaw model": e.g., $m_{\nu_{\tau}}/m_{\nu_e} = m_{\tau}^2/m_e^2$.
- neutrino mass is very important in astrophysics! - $m_{\nu_{\tau}} \sim 10 eV$ closes the universe

RECENT MEASUREMENTS

- Statistical error $\approx 10\%$ for $(K\pi)^-$ mode, > 25% for $(K\pi\pi)^-$ mode.
- Statistically limited.

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Mode	Expt	<i>B</i> (%)('94 Prelim.)
K*-	ARGUS	$0.97 \pm 0.15 \pm 0.12$
K^0h^-	ALEPH	$1.28 \pm 0.16 \pm 0.12$
$\overline{K}{}^{0}\pi^{-}$	ALEPH	$0.88 \pm 0.14 \pm 0.09$
<i>K</i> ⁰ <i>K</i> ⁻	ALEPH	$0.29 \pm 0.12 \pm 0.03$

$\overline{K}^{0}\pi^{-}\pi^{0}$	ALEPH	$0.33 \pm 0.14 \pm 0.07$
$K^{0}K^{-}\pi^{0}$	ALEPH	$0.05 \pm 0.05 \pm 0.01$

• TPC/2 γ sees more $K_1^-(1400)$ through $K^-\pi^+\pi^-$ channel:

au decay mode	B(%)('93 Prelim.)
$K_1^{-}(1400)$	$0.74_{-0.33}^{+0.40}$
$K_1^{-}(1270)$	0.43 ^{+0.40} -0.34

ANALYSIS

- K_S^0 detected by looking at secondary vertex.
- No particle ID on h^- or X^+ (= $e^+\nu_e, \mu^+\nu_\mu, \pi^+, K^+, \rho^+$). Assign π^- mass when calculating invariant mass.
- All π^0 's (signal or tag side) explicitly reconstructed.

FIND: 0 (1, don.) (2 T - → Fo h- TO V = = (0.519±0.025±0.062) 70 (100 0' 1270 (1) T-→, T·K Z: (0.123 ± 0.023 ± 0.023) % VCVC τ-> FOK TOVE: (0.129±0.050±0.032)% - MKKTT CONSISTENT W/ K*K (3) Z-> K° K°TT 2 = (C.CE3 + C.CI7 = C.OI7) & Show equal K°KTI i <u>All</u> K*K 700



FIG. 8. The $K, K^-\pi^{\bullet}$ invariant mass distribution for $\tau^- \to K^{\bullet}K^-\pi^{\bullet}\nu_{\tau}$ candidates. The data points are obtained by taking 50 MeV "slices" in invariant mass and fitting the background subtracted σ_K distribution for those events. The solid histogram shows the invariant mass spectrum from a full Monte Carlo simulation of $\tau^- \to e_1\nu_{\tau}$, $e_1 \to K^{\bullet}K$ decays. The dashed histogram shows the mass distribution for a simulation of $\tau^- \to e_7\nu_{\tau}$, $\rho \to K^{\bullet}K^-$ decays.



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FIG. 5. The K_*K^- invariant mass of $\tau^- \to K^*K^-\nu_\tau$ events. The filled circles, with error bars, are for the data. The solid histogram is the result of a fit to the sum of $\tau^- \to \rho^-\nu_\tau$, $\rho \to K_*K^-$ (dashed) and $\tau^- \to K^*K^-\nu_\tau$ phase space (dot-dashed) Monte Carlo models.

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TABLES

Decay	€ (%)	\mathcal{B} (in units of 10^{-6})			
channel		Previous	This analysis		
$\overline{\tau^- \rightarrow e^- e^+ e^-}$	20.4	13	3.3		
$\tau^- \rightarrow \mu^- e^+ e^-$	19.6	14	3.4		
$\tau^- \rightarrow \mu^+ e^- e^-$	19.9	14	3.4		
$\tau^- \rightarrow e^- \mu^+ \mu^-$	18.8	19	3.6		
$\tau^- \rightarrow e^+ \mu^- \mu^-$	19.4	16	3.5		
$\tau^- \rightarrow \mu^- \mu^+ \mu^-$	15.9	17	4.3		
$\tau^- \rightarrow e^- \pi^+ \pi^-$	15.5	27	4.4		
$\tau^- \rightarrow e^- \tau^- K^+ \qquad 14.6$		58	4.6		
$\tau^- \rightarrow e^- \pi^+ K^-$	14.9	29	7.7		
$\tau^- \rightarrow e^+ \tau^- \tau^-$	15.5	17	4.4		
$\tau^- \rightarrow e^+ \tau^- K^-$	15.1	20	4.5		
$\tau^- \rightarrow \mu^- \tau^+ \tau^-$	9.1	36	7.4		
$\tau^- \rightarrow \mu^- \tau^- K^+$	7.4	77	15		
$\tau^- \rightarrow \mu^- \overline{\tau}^+ K^-$	7.8	77	8.7		
$\tau^- \rightarrow u^+ \pi^- \pi^-$	9.8	39	6.9		
$\tau^- \rightarrow \mu^+ \pi^- K^-$	7.7	40	20		
$\tau^- \rightarrow e^- a^0$	16.2	19	4.2		
$\tau^- \rightarrow e^- K^{*0}$	10.7	38	6.3		
$\tau^- \rightarrow e^- \bar{k}^{*0}$	10.5		11		
$\tau^- \rightarrow \mu^- \rho^0$	11.9	29	5.7		
$\tau^- \rightarrow \mu^- F^{*=0}$	7.2	45	9.4		
$r \rightarrow \mu h$	7.8		8.7		

TABLE I. The detection efficiencies, ϵ , and upper limits for the branching fractions at 90% CL. B, together with the most restrictive limits from previous experiments [1,4].

CLE094: Searches for neutrinoless T-decay => Beyond S.M. physics

Also:	مر د ت بر ⇒ و	r≤4 ×<5	.2 × 10 5× 10 ⁻¹	- 6		
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TABLES

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From CLEO-II and a half to CLEO-III

BaBar/CLEO-III detector design driven by **B-physics**, not tau-physics.

Must have:

• Particle ID: 4σ separation π/K at 2.8 GeV/c for $B \to \pi\pi/B \to K\pi \ (D \to \pi\ell\nu/D \to K\ell\nu)$

- CLEO-II, present: 1.8 σ in Rel. Rise

- Lepton ID down to 400 MeV/c for electrons, 800 MeV/c for muons
 - 3-dim vertex resolution ≤ 50µ (improve S:N for charm, measure D* widths?)
 - Momentum resolution: $(\frac{dp}{p})^2 \le (0.0015 \ p)^2 + (0.005)^2$

- Important to know momentum resolution for measurement of $m_{
u_{ au}}$.

• Photon resolution and segmentation at least as good as CLEO-II ($D^{*+} \rightarrow D^+ \pi^0$, charm decays w/ neutrals)

- Ability to handle high rates > 10² at L2 with small dead times
- Silicon vertexing (~ 50μ in two-track vertex in $r-\phi$)
- Hermiticity: Important for rejection of $q\bar{q} \rightarrow \pi^{0}$'s



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5/fb equivalent



Determination of Systematic Errors

"I have seen the future of τ physics and it is *precision* physics" - D. Marsh, 1975

At CLEO, ϵ_{MC} vs. ϵ_{data} increasingly THE dominant error

Need to measure ϵ as function of:

• Momentum $(h^{\dagger},\pi^{\circ})$

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- Particle type (π⁺, k⁺, p⁺, e⁺, μ⁺)
- Dip angle wrt beam axis (h^+, π°)
- Event parameters: thrust and mulitplicity (h⁴,π[•])

1. Absolute charged particle reconstruction efficiency

- (a) Track-<u>finding</u> efficiency (topological BR's, e.g.)

MinIPeak in calorimeter tags charged particle, count how often one is observed

- -> with 2 independent tracking systems (Si/DR or UD/DR)
- High-energy showers in CC from Bdecay 90% from e^{+-} . count how often one is track matched.
- Multiplicity correlations data vs. MC
 - Use ratio of evts with net charge ± 1 relative to net charge = 0 in data vs. MC $N_{\bullet} \approx N_{\bullet}(1-\epsilon)$

 $- N_{\ell vs 2}/N_{\ell vs 3}$

 $-\gamma\gamma \rightarrow h^+h^-$; count $N_{h^+}/N_{h^+h^-}$.

For τ decay, quoted systematic error varies from 0.4% ($\tau \rightarrow \rho^- \nu_\tau$) to 1.25% ($\tau \rightarrow 5$ -prong)

· Have observed E loss due to evt. multiplicity/confusion

(b) Product of track-finding and track-fitting efficiency (for finding the yield underneath $\eta \rightarrow \pi^+\pi^-\pi^0$ peak, e.g.) Depends on modeling hit resolution function, e.g. $D^{*+} \rightarrow D^0\pi^+; D^0 \rightarrow K^-\pi^0(\pi^+)/D^0 \rightarrow K^-\pi^0\pi^+$ using "satellite peak" (hoth $\varepsilon_{\pi^o}, \varepsilon_{\pi^\pm})$ $K^-\pi^0\pi^+$ using "satellite peak" (hoth $\varepsilon_{\pi^o}, \varepsilon_{\pi^\pm})$ G: How offen are even in Kin^{*}(n) plak fully reconstructed w/ π^o ? $\pi^0 \rightarrow e^+e^-\gamma$; use shower due to e, etrace back helical path to origin and produce a π^0 peak. How often is track found within some window?

> Quoted systematic error varies from $2\% (p > p_{curl}, barrel)$ to $5\% (p < p_{curl})$ Note that efficiency outside of barrel region well-reproduced in MC (isotropy of charged tracks from B's - will this be done at τcF with D° or D^{\dagger} decays at -threshold

Note: Potentially BIG TCF advantage: Nultiplicity low, particles well i.d.'d => calibrate efficiencies VERY well using min techniques as f(particle type, cons) e.g. U=K*K*T(TT)







2. Absolute γ finding efficiency

- Use $D^{*0} \rightarrow D^0 \pi^0$; $\pi^0 \rightarrow \gamma(\gamma)$ relative to case of fully reconstructed π^0
- Use flatness of dN/dE_{γ} from $\pi^{0} \rightarrow \gamma\gamma$ to extrapolate efficiency in region 30 MeV to 2 GeV. (shape of \mathcal{E}_{r} curve)
- Use flatness of $dN/dcos\theta$ spectrum to extrapolate endcap (noise??) ξ
- Quoted error of 2% per photon, includes uncertainty on photon line shape (shower containment uncertainties, e.g.)
- 3. Absolute π^0 finding efficiency
 - $D^{*+} \rightarrow D^0 \pi^+; D^0 \rightarrow K^- \pi^+ (\pi^0) / D^0 \rightarrow K^- \pi^0 \pi^+$ using "satellite peak"
- mass recoiling against $\pi^+\pi^-\pi^+\pi^-$ in $\Upsilon \to \pi^+\pi^-\pi^+\pi^-\pi^0$ events vs. number of times π^0 fully reconstructed.

• $\eta \to \pi^0 \pi^0 \pi^0 / \eta \to \gamma \gamma$, efficiency-corrected

BR compared with PDG value

Quoted error of 5% per π^0

- 4. Ratio of $\frac{\pi^+}{\pi^0}$ efficiencies:
 - Assume equality of cross-sections of \equiv 's: $\equiv^- \rightarrow \Lambda \pi^- / \equiv^0 \rightarrow \Lambda \pi^0$
 - R(Ispin ratio) in *D**' decay:

$$R = \frac{\Gamma(D^{*+} \to D^0 \pi^+)}{\Gamma(D^{*+} \to D^+ \pi^0)} = 2.17 \pm 0.08$$

• $\Upsilon(2S) \rightarrow \Upsilon(1S)\pi^+\pi^-: \Upsilon(2S) \rightarrow \Upsilon(1S)\pi^0\pi^0=2.0$

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$$\eta \rightarrow \pi^+ \pi^- \pi^0 / \eta \rightarrow \pi^0 \pi^0 \pi^0$$
, cf. PDG

Quoted error of 6% on ratio (8% for V_{cb} analysis)

The World According to Dave - 10 GeV vs. τcF

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Physics	τcF	10 GeV/Systematic
Tau mass	\$\$\$	\$
leptonic decays	\$\$	\$\$/fakes,QED
lifetime		\$\$
Lorenz structure	\$\$	\$\$
spin structure	\$\$	\$\$
hadronic decays	\$\$	\$ (qā, γγ)
'nclusive/exclusive BR's	\$\$	\$ (q \bar{q} , $\gamma\gamma$)
spectral functions/ α_s	\$\$	\$ (q \bar{q} , $\gamma\gamma$)
spectral functions/ K_1 's	\$\$	$(q\bar{q}, \gamma\gamma)$
Rare decays/L-viol.	\$\$	\$\$
$ u_{ au}$ mass	\$\$	\$\$

Conclusions

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 τ physics very active at CLEO - World- class precision for almost all measure-ments

BUT: progress requires high statistics, good control of systematics

- The era of very high data rates means that control of systematics will become increasingly important, and will require knowing, at the 0.1% level, systematics from:
 - 1. Absolute tracking efficiencies
 - Knowledge of fake rates for leptons, e.g.

3. Knowledge of backgrounds

Monte Carlo tuning will become a fulltime preoccupation!