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The positron source for the NLC project utilizes a 6.2 GeV electron beam interacting in a high-Z positron production target. The electromagnetic shower in the target results in large energy deposition which can cause damage to the target. Optimization of the collection system is required to insure long-term operation of the target with needed high positron yield into the 6-dimensional acceptance of the subsequent pre-damping ring. Positron tracking through the accelerating system indicates a dilution of the initial positron phase space density. Results of the studies on the development of the collector systems design are presented.

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The positron source for the NLC project utilizes a 6.2 GeV electron beam interacting in a high-Z positron production target. The electromagnetic shower in the target results in large energy deposition which can cause damage to the target. Optimization of the collection system is required to insure long-term operation of the target with needed high positron yield into the 6-dimensional acceptance of the subsequent pre-damping ring. Positron tracking through the accelerating system indicates a dilution of the initial positron phase space density. Results of the studies on the development of the collector systems design are presented.

## 1 INTRODUCTION

The positron injector includes a production target followed by a short solenoid with a strong magnetic field (flux concentrator), a 250 MeV linac with 0.5 Tesla focusing solenoids, a 1.73 GeV linac with quadrupole focusing, and an energy compressor system before injection into the positron pre-damping ring (PPDR). The preliminary design of injector is presented in Ref. [1], [2], [3]. The new set of parameters is listed in Tables 1, 2. The main changes in the requirements to the injector are imposed by a new value of PPDR normalized transverse acceptance of 0.045 $\pi$ m rad instead of the previous value of 0.09 $\pi$ m rad. The ultimate goal of the collector system is to provide the highest number of positrons within the 6-dimensional acceptance of the pre-damping ring.

## 2 POSITRON YIELD

As positron capture is restricted by the acceptance of the pre-dumping ring, it is convenient to select an energy-invariant 6-dimensional phase space volume and observe the dilution of positron phase space density inside this volume. The 6-dimensional volume is defined in canonical conjugate variables $(x, P_x)$, $(y, P_y)$, $(z - z_s, p - p_s)$, where $P_x$ and $P_y$ are the canonical momentum

$$P_x = p_x - \frac{e B_z}{2} y, \qquad P_y = p_y + \frac{e B_z}{2} x , \qquad (1)$$

$p_x$, $p_y$ are the mechanical momentum, $B_z$ is a longitudinal magnetic field, and $z_s$, $p_s$ define the dynamics of synchronous particles. To insure proper capture of positrons into PPDR, the edge normalized emittance of the positron beam is selected to be $\varepsilon_x$, $\varepsilon_y \le 0.03$ $\pi$ m rad, bunch length $\sigma_z$ = 6.97 mm and energy spread $\Delta E/E$ = 2% at the PPDR injection energy of E = 1.98 GeV. The value of energy spread corresponds to momentum spread of the beam $\Delta(\beta\gamma) = 80$. The positron yield is defined as a ratio of positrons in a volume of phase space per incident electron at the target, $Y_{e^+} = N_{e^+}/N_{e^-}$.

Table 1. Parameters of the electron beam drive for the positron target.

| Parameter | Value |
|---|---|
| Energy | 6.2 GeV |
| Bunch spacing | 1.4 / 2.8 ns |
| Bunch energy variation | 1% FW |
| Single bunch energy spread | 1 % FW |
| Normalized emittance | $10^{-4}$ $\pi$ m rad |
| Transverse size, $\sigma_x$ | 1.6 mm |
| Bunch length, $\sigma_z$ | 5 mm |
| Particles/Bunch | $1.2 / 2.4 \cdot 10^{10}$ |
| Train population uniformity | 1 % FW |
| Bunch-to-Bunch Pop. Unif. | 2 % rms |
| Number of Bunches | 190 / 95 |
| Repetition Rate | 120 Hz |
| Beam Power | 271 kW |

Table 2. Positron beam parameters.

| Parameter | Value |
|---|---|
| Energy | 1.98 GeV |
| Bunch spacing | 1.4/2.8 ns |
| Bunch energy variation | 1% FW |
| Single bunch energy spread | 2 % FW |
| Normalized emittance | 0.03 $\pi$ m rad |
| Bunch length, $\sigma_z$ | 10 mm |
| Particles/Bunch | $0.9/1.8 \cdot 10^{10}$ |
| Train population uniformity | 1 % FW |
| Bunch-to-Bunch Pop. Unif. | 2 % rms |
| Number of Bunches | 190/95 |
| Repetition Rate | 120 Hz |
| Beam Power | 58 kW |

## 3 BEAM DYNAMICS IN INJECTOR

The initial positron distribution was obtained via Monte-Carlo simulation of electron-positron pair production utilizing the code EGS [4]. The positron beam created by an electromagnetic shower in solid target

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possesses large momentum spread. The typical value of transverse positron emittance after the interaction of the incident 6.2 GeV electron beam of $\sigma_{x,y}$ = 1.6 mm with a W-Re target of 4 radiation lengths (RL) is $\varepsilon_o$ = 0.135 $\pi$ m rad. The production target is in the field of a tapered solenoid with a maximum field of $B_t$ = 1.2 Tesla. The presence of a magnetic field at the target does not significantly increase the normalized beam emittance:

$$\varepsilon = \sqrt{\varepsilon_0^2 + (\frac{e\, B_t\, R_t^2}{2\, m\, c})^2} \approx \varepsilon_o \,. \qquad (2)$$

The created beam enters the strong longitudinal magnetic field of a flux concentrator. The flux concentrator is a solenoid with sharp increase of magnetic field up to peak value of 5.8 Tesla at a distance of 5 mm from the target and an adiabatic decrease of the field over a distance of 15 cm [5]. The strong magnetic field at the target is required to confine the emitted positron beam with its large momentum spread.

The flux concentrator is followed by a 250 MeV linac with an effective accelerating gradient of 25 MeV/m and transverse focusing by a 0.5 Tesla solenoid field. The transverse acceptance of the linac with solenoid focusing is energy independent:

$$\alpha = \frac{eB_z}{2\, mc}\, a^2 = 0.06\ \pi \text{ m rad} \,. \qquad (3)$$

After acceleration to 250 MeV positrons are accelerated in a linac with quadrupole focusing and an accelerating gradient of 15 MeV/m.

Tables 3 and Fig. 1 show the results of the positron beam dynamics simulation in the injector. The positron yield within the 6D phase space volume drops from the value of 4.7 at the target to 0.76 at the end of the linac. One of the reasons for the dilution of the beam phase space density in the linac is the beam energy spread, which results in chromaticity-induced beam emittance growth. Calculations were done for two cases, using longitudinal focusing: for a linac with solenoidal focusing with B = 0.5 Tesla, and a linac with focusing by periodic permanent magnets (PPM). The field of the PPM was approximated by function $B_Z = B_o \sin(2\pi z/L)$, where $B_o$ = 0.5 Tesla is the maximum field and L = 50 cm was chosen to be the period of the PPM structure. Particle tracking indicates that in the case of longitudinal focusing, the positron yield at the end of the linac is increased to $Y_{e^+}$ = 1.15.

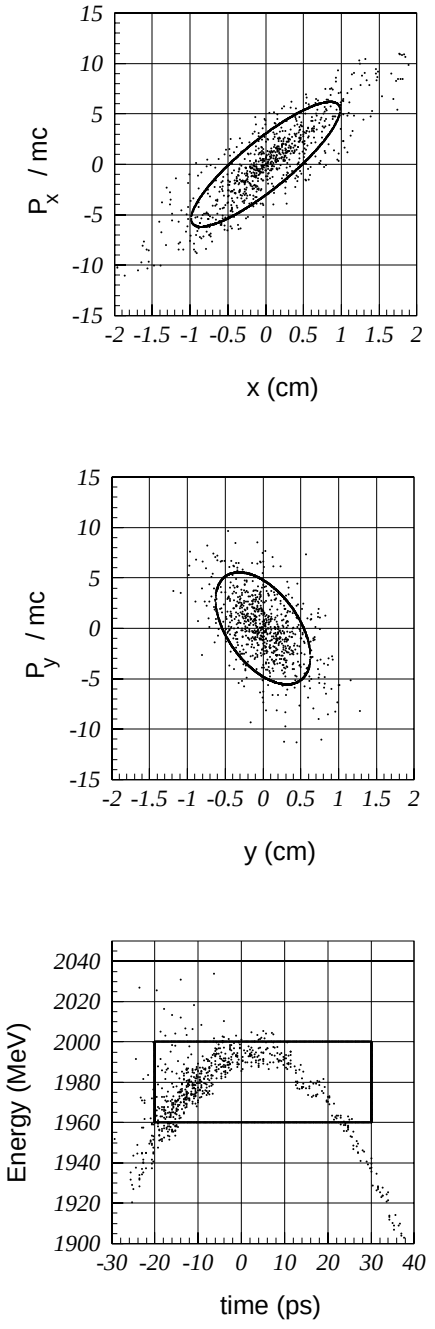


Fig. 1. Positron distribution at 1.98 GeV.

Table 3. Positron yield along the injector.

| | After target | After flux concentrator | At 250 MeV | At 1.9 GeV |
|---|---|---|---|---|
| Positron yield | 12 | 4.7 | 2.2 | 1.9 |
| Fraction of positrons within 6-D phase space | 0.39 | 0.37 | 0.74 | 0.40 |
| Positron yield within 6-D phase space | 4.70 | 1.76 | 1.60 | 0.76 |

## 4 DYNAMICS OF POLARIZED POSITRONS

One option for the positron injector being considered is the collection and acceleration of polarized positrons. Polarized positrons may be produced by targeting helically polarized gammas, which could be made by Compton backscattering [6] or with a helical undulator [7], on a thin target. To study the effects of positron depolarization, the particle tracking code was modified to include Thomas-BMT equation [8], describing the precession of the spin vector $\vec{S}$ :

$$\frac{d\vec{S}}{dt} = \frac{e\,\vec{S}}{m\gamma} \times \left[(1+G\gamma)\vec{B}_{\perp}+(1+G)\vec{B}_{II}+(G\gamma + \frac{\gamma}{1+\gamma})\frac{\vec{E}\times\vec{\beta}}{c}\right], \quad (4)$$

where G is the anomalous magnetic moment of the positron, $\vec{E}$ is the electrical field, and $\vec{B}_{\perp}$ and $\vec{B}_{II}$ are components of the magnetic field perpendicular and parallel to the particle velocity. The spin advance at a small distance $\delta z$ is described as a matrix [9] :

$$\begin{vmatrix} S_x \\ S_y \\ S_z \end{vmatrix} = \begin{vmatrix} 1- a(B^2+C^2) & ABa + Cb & ACa - Bb \\ ABa -Cb & 1- a(A^2+C^2) & BCa +Ab \\ ACa + Bb & BCa - Ab & 1- a(A^2+B^2) \end{vmatrix} \begin{vmatrix} S_{x,o} \\ S_{y,o} \\ S_{z,o} \end{vmatrix}, \quad (5)$$

$$A = \frac{D_x}{D_o},\ B = \frac{D_y}{D_o},\ C = \frac{D_z}{D_o},\ D_o = \sqrt{D_x^2 + D_y^2 + D_z^2}\ , \quad (6)$$

$$a = 1 - \cos(D_o\,\delta z)\ ,\ \ b = \sin(D_o\,\delta z)\ , \quad (7)$$

where components $D_x$, $D_y$, $D_z$ are defined by the equations:

$$D_x = \frac{e}{m\,\gamma\,v}\left[(1+G\gamma)(B_x - x'B_z) + (1+G)x'B_z + \frac{v}{c^2}\left(\frac{\gamma}{1+\gamma} + G\gamma\right)(E_y - y'E_z)\right], \quad (8)$$

$$D_y = \frac{e}{m\gamma v}\left[(1+G\gamma)(B_y - y'B_z) + (1+G)y'B_z + \frac{v}{c^2}\left(\frac{\gamma}{1+\gamma} + G\gamma\right)(x'E_z - E_x)\right], \quad (9)$$

$$D_z = \frac{e}{m\,\gamma\,v}\left[(1+G\gamma)(-x'B_x - y'B_y) + (1+G)\cdot (x'B_x+B_z+y'B_y) + \frac{v}{c^2}\left(\frac{\gamma}{1+\gamma} + G\gamma\right)(y'E_x - E_y x')\right]. \quad (10)$$

Fig. 2 illustrates the polarization of positrons emerging from 0.2 RL W-Re target as a function of their energy, calculated by the EGS code extended to study polarization propagation in the target [10]. Polarization is defined as the probability to find the spin of the positron along the direction of the positron momentum.

During beam transport and acceleration, the spin vector precesses, resulting in the depolarization of the beam. We define the longitudinal polarization as an average of the product of the longitudinal component $S_z$ and the value of polarization, P, over all positrons:

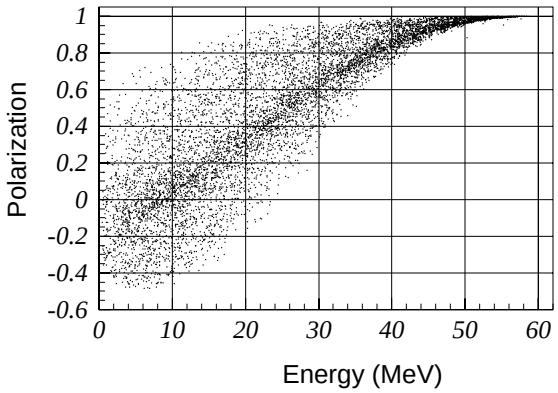


Fig. 2. Initial positron polarization from 60 MeV γ-rays.

Table 4. Polarization of positrons at 250 MeV captured within the 6-D phase space.

| Energy range, MeV | Positron capture within 6D phase space | $<P_z>$ |
|---|---|---|
| 260-280 | 0.028 | 0.92 |
| 255-280 | 0.038 | 0.86 |
| 245-280 | 0.084 | 0.68 |
| 242-280 | 0.11 | 0.60 |
| 240-280 | 0.14 | 0.53 |

$$<P_z> = \frac{1}{N}\sum_{i=1}^{N} S_z^{(i)}\,P^{(i)}\ . \quad (11)$$

Table 4 illustrates the distribution of polarized positrons at 250 MeV. Simulation shows that 11% of the produced positrons have an average value of polarization of 60 %. Additional study is required to increase the number of captured positrons keeping the value of polarization high, and determine the needed γ - flux.

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