SLAC - PUB - 4613 April 1988 (T/E)

SEMILEPTONIC Λ_c^+ DECAYS WITH THE MARK II AT PEP*

SPENCER KLEIN[†] Stanford Linear Accelerator Center, Stanford, CA, 94309, USA

ABSTRACT

Semileptonic Λ_c^+ decays have been observed in 29 GeV e^+e^- annihilation. We observe the channels $\Lambda_c^+ \rightarrow e^+\Lambda X$ and $\Lambda_c^+ \rightarrow \mu^+\Lambda X$. The production rate for Λ_c^+ , times $\operatorname{Br}(\Lambda_c \rightarrow l^+\Lambda X)$ of 0.0031 \pm 0.0012 \pm 0.0010 per hadronic event for electrons, and 0.0024 \pm 0.0024 \pm 0.0007 per hadronic event for muons. When combined with the Lund model prediction for Λ_c^+ production, this gives semileptonic branching ratios of roughly 5%, about what is expected. When combined with the UCLA model prediction for Λ_c^+ production, this gives semileptonic branching ratios of 17% and 13% for electrons and muons, somewhat higher than expected.

Contributed to the Rencontre de Moriond: Current Issues in Hadron Physics, Les Arcs, France, March 13–19, 1988

*Work supported by the Department of Energy, contract DE-AC03-76SF00515. † Representing the Mark II/PEP5 Collaboration. The experimental results on semileptonic charmed baryon decays are very sparse; the only previous measurements are from Mark II at SPEAR.¹ Semileptonic Λ_c^+ decays are interesting because they may be easily compared with charmed meson semileptonic decays, providing a tests of baryonic form factors.

In high energy e^+e^- annihilation, Λ_c^+ production provides a strong test of fragmentation models; the predictions of the Lund² and UCLA³ models differ by a factor of three.

We have measured $\sigma(e^+e^- \to \Lambda_c X) \cdot \operatorname{Br}(\Lambda_c \to l\Lambda X)$, for both electrons and muons. Λ_c^+ are seen via their decay to a Λ plus a lepton, plus unobserved particles. The Λ are required to come from the origin and have with a momentum of at least 1.5 GeV/c. Electrons are required to have a momentum of at least 1.5 GeV/c and muons are required to have a momentum of at least 2.0 GeV/c. The leptons are required to be in the same hemisphere as the Λ . The invariant mass of the combinations are shown in Figure 1, separately for right sign $(l^+\Lambda, l^-\overline{\Lambda})$ and wrong sign $(l^+\overline{\Lambda}, l^-\Lambda)$. There are 17 right sign events and 5 wrong sign events in the region below the Λ_c^+ mass of 2.28 GeV/c². Above 2.28 GeV/c², there are slightly more wrong sign events.



Fig. 1. Right sign $(\Lambda l^+, \overline{\Lambda} l^-)$ and wrong sign $(\Lambda l^-, \overline{\Lambda} l^+)$ combinations.

The background for the Λ_c^+ signal comes from a variety of sources. In order of decreasing importance, they are:

- 1. K_{θ} and random track combinations misidentified as Λ .
- 2. Fake leptons: misidentified electrons, π/K decays, and hadron punchthrough.
- 3. Random background: A real Λ plus a real lepton, from incoherent sources.
- 4. A plus leptons, both from the same b decay. The most likely form for this is a B hadron decaying to a lepton plus a Λ_c^+ , with the Λ_c^+ subsequently decaying to a Λ . This produces a wrong sign combination. Less common are B mesons decaying to a lepton plus a baryon plus an $\overline{\Lambda}$, leading to a right sign event. The b decay background is peaked at higher mass than Λ_c^+ decays, and the wrong sign excess seen at high masses may be from b decays. Because b decays are rarer than charm decays, and because of the higher average b mass, this background source is small.

Since most of the background populates the right and wrong sign plots equally, the wrong sign combinations provide a measurement of the background.

The efficiency to detect a Λ_c^+ depends on the exact decay modes. Some modes that one could consider are: $\nu\Lambda$, $1 \nu\Sigma^0$, $1 \nu\Sigma^{*0}$, $1 \nu\Lambda\pi^0$, and $\nu\Lambda(\pi\pi)^0$. Other modes are Cabbibo suppressed, or otherwise limited. We can eliminate all of these modes except $\nu\Lambda$ on various theoretical grounds. The Λ_c^+ is isospin 0. The modes containing a single Σ or $\Lambda\pi^0$ are isospin 1, and should be heavily suppressed. We can also compare with D meson semileptonic decays. In D meson semileptonic decays, the hadronic part of the final state is overwhelmingly a single particle.⁴

In Λ_c^+ semileptonic decays, the charmed quark emits a W, changing to a strange quark. In the baryon, the strange quark will have a large perpendicular momentum. Any extra quark production should occur between the strange quark and the nonstrange quarks. In that case, the final state would not include a Λ .

Two final hints come from the data. A search for the decay model $\Lambda \pi^+\pi^-$ yielded no candidates. From this, I conclude that the two-pion decay mode is negligible. Finally, the data shown in Figure 1 is in reasonable agreement with the Lund prediction for $l\nu\Lambda$, indicating that the decay is not dominated by other final states.

In what follows, we will take $\Lambda_c^+ \rightarrow l\nu\Lambda$ to be the dominant mode, and allow for other decay modes as systematic errors.

Figure 2 shows the (wrong sign) background subtracted momentum spectrum of the Λ + lepton pairs. The normalized Lund model prediction is superimposed.



Fig. 2. Momentum of the Λ + lepton combination. The solid curve is the Lund model prediction, normalized to the data in the 4 to 8 GeV/c region.

We also measure the total production rate. With radiative corrections, we find for the electron subsample (11 signal, 2 background):

$$\sigma(e^+e^- \rightarrow \Lambda_c X) \cdot \operatorname{Br}(\Lambda_c \rightarrow e\Lambda X) = 1.2 \pm 0.5 \pm 0.4 \text{ pb},$$

or $0.0031 \pm 0.0012 \pm 0.0010$ per hadronic event.

1. 31.

- For the muons (6 signal, 3 background),

$$\sigma(e^+e^- \to \Lambda_c X) \cdot \operatorname{Br}(\Lambda_c \to \mu \Lambda X) = 1.0 \pm 1.0 \pm 0.3 \text{ pb},$$

or $0.0024 \pm 0.0024 \pm 0.0007$ per hadronic event.

To put these numbers in perspective, it is useful to consider some Monte Carlo predictions. The Lund model,² for example, predicts that at 29 GeV, 0.06 Λ_c^+ should be produced per-hadronic event. With this assumption, we find $Br(\Lambda_c \to e\Lambda X) = 5.1 \pm 2.0 \pm 1.7\%$ and $Br(\Lambda_c \to \mu\Lambda X) = 4.0 \pm 4.0 \pm 1.2\%$.

We can also consider the predictions of the UCLA model.^{3]} Since the UCLA model bases hadron production rates on their mass, and since Λ_c^+ are heavy, the predicted rate is much lower, 0.018 Λ_c^+ per hadronic event. This yields $Br(\Lambda_c \to e\Lambda X) = 17 \pm 7 \pm 6\%$ and $Br(\Lambda_c \to \mu\Lambda X) = 13 \pm 13 \pm 4\%$.

These rates can be compared with the Mark II/SPEAR measurements of $Br(\Lambda_c \to eX) = 4.5 \pm 1.7\%$ and $Br(\Lambda_c \to e\Lambda X) = 1.1 \pm 0.8\%$.

These results were found by measuring the increased rate of proton and Λ production in e^+e^- annihilation as the e^+e^- energy was increased above the Λ_c^+ production threshold.

Particle	Br(e X)	Lifetime $(10^{-13}s)$
D+	18.2 ± 1.7	10.5 ± 0.3
\mathbf{D}^{0}	7.0 ± 1.1	4.3 ± 0.1
· \Lambda_c^+	?	1.9 ± 0.2

Table 1. Charmed particle semileptonic branching ratios and lifetimes.

These numbers may also be compared with the comparable figures for D meson semileptonic decays shown in Table 1. Neglecting the hadronic form factors, the charmed particle lifetimes should be proportional to the semileptonic branching ratios. Hadronic effects should be of order two or three, so, $Br(\Lambda_c \to e\Lambda X)$ should be small, probably less than 10%. So, it appears that the UCLA model prediction for Λ_c^+ production is low, while the Lund model prediction is reasonable.

One interesting application of this Λ_c^+ signal is to search for $\Sigma_c \to \Lambda_c \pi$, for both Σ_c^{++} and Σ_c^0 , using the same Δm technique used to find D^* . Because of the missing neutrino, the Δm band is much wider than for D^* . However, because of the purity of the Λ_c^+ signal, the technique works. Unfortunately, with the limited statistics present, the resulting upper limits are uninteresting.

In conclusion, we have observed semileptonic Λ_c^+ decays in 29 GeV e^+e^- annihilation.

REFERENCES

- 1. E. N. Vella et al., Phys. Rev. Lett. 48, 1515 (1982).
- 2. B. Andersson et al., Nucl. Phys. B197, 45 (1981) and references therein.
- 3. C. D. Buchanan and S. B. Chun, Phys. Rev. Lett. 59, 1997 (1987).
- 4. R. Morrison, these Proceedings.