EVIDENCE FOR PARITY VIOLATION IN THI: DECAYS
of the narrow states niar $1.87 \mathrm{GeV} / \mathrm{c}^{2}$ *
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## ABSTRACT

We have studied the Dalitz plot for the recently observed charged
 state is incompatible with a natural spin parity assignnent. This information, coupled with the carlier observation of the $K^{ \pm} \pi^{\mp}$ decay mode (a final state of natural spin parity) of the neutral state at $1865 \mathrm{McV} / \mathrm{c}^{2}$, suggests parity violation in the decays of these objects if they are members of the same isomultiplet as their proximity in mass suggests.

[^0]We have recently reported our obscrvation in $\mathrm{e}^{+} \mathrm{e}^{-}$annihilation of a narrow, charged state of mass $1876 \mathrm{MeV} / \mathrm{c}^{2}$ decaying into the exutic decay mode $K^{\mp} \pi^{ \pm} \pi^{ \pm}$. The proximity in mass of this state to the neut ral state decaying into $K \pi$ and $K 5 \pi$ at $1865 \mathrm{MeV} / \mathrm{c}^{2}$ suggests that they are members of the same isomultiplet. As such they are expected to have the same parity. Since the $K \pi$ final state is one of natural spin parity, a demonstration that the $K \pi \pi$ final state of the charged member of the isomultiplet is inconsistent with natural spin parity implies a parity violation in the decay. In this Letter we present evidence, based on a study of the $K^{\mp} \pi^{ \pm} \pi^{\prime \prime}$ Dalitz plot for such a parity violation, suggesting that the decay proceeds via the weak interaction as expected for the predicted $\left(0^{+}, D^{0}\right)$ isodoublet of charm. ${ }^{3}$

The present analysis is based on $K \pi \pi$ events observed among a sample of $\sim 44,000$ hadronic events taken from 3.9 GeV to 4.25 (iev center-of-mass energy. Thesc data were taken with the SLAC-LBL magnetic detector at SPEAR.

The $K \pi \pi$ combinations are selected with the aid of the time-of-flight system described in Ref. 2. In the present analysis we have used a modified form of the time-of-flight ( fO ) weighting technique described earlier. ${ }^{1,2}$ A given track in a multi-prong hadronic event is assigned a definite particle identity on the basis of the agrecnent between its observed TOF: over a $1.5-2.0$ meter flight path and that predicted for either a $n$ or a $k$ with a monentum as measured. Specifically we compute a $x^{2}$ value for both the $\pi$ and $K$ hypotheses ( $x_{\pi}^{2}$ and $x_{K}^{2}$ ) based on the observed and expected TOF and the 0.4 ns rms resolution of the rof system. Tracks satisfying the requirements $x_{K}^{2}<x_{\pi}^{2}, x_{K}^{2}<3$ are called kaons. Protons and anitprotons are separated from
kaons in a similar fashion. The remaining tracks are called pions. ${ }^{4}$ The above technique allows the direct study of scatter plots and in particular the balitz plot for the $K \pi \pi$ system.

In order to obtain a relatively clean sample of $k \pi \pi(1876)$ cuents we make use of the result that for the $E_{c m}$ region, $3.9<E_{\mathrm{cm}}<4.25 \mathrm{GeV}$; the recoil mass ( $M_{r e c}$ ) spectrum shows a sharp spike near 2 GeV .1 We thus used a data sample with the $F_{\text {cm }}$ region chosen as above coupledwith a cut 1.96 < Mrec < 2.04 $G e V / c^{2}$. Figure la and $l$ bhow the resulting exotic and nonexotic $k \pi m$ invariant mass distributions. A fit to the spectrum of lig. lb was appropriately scaled to serve as a background for Fig. la. Figure la shows a fit to a galussian peak over this background. Figure 2 a shows the (folded) Dalitz plot for $K^{\mp} \pi^{1 . \pm}$ events with the additional invariant mass (M) requirement $1.86<M<1.92 \mathrm{GeV} / \mathrm{c}^{2}$. We find a sample of 126 events in the Dalitz plot of Figure $2 a$ of which we estimate 58 arc background. In Figure 2 b we show a background Dalitz plot consisting of 112 nonexotic combinations $K^{\mp} \pi^{+} \pi^{-}$satisfying the same mass and missing mass cuts as the exotic combinations of Figure 2 a.

Both signal and background Dalitz plots are consistent with uniform population density. A uniformly populated Dalitz plot is incompatible with a $K \pi \pi$ final state of pure, natural spin parity. ${ }^{5}$ for the case of a natural spin-parity state decaying into three pseudoscalars one expects a depopulation (or zero) along the Dalitz plot boundary. This follows from the necessity of constructing the matrix clement from the vector product of the two independent center-of-mass momenta-a vector which vanishes on the Dalitz plot boundary where momenta are collinear. If, as in the case of $K^{+} \pi^{ \pm} \pi^{ \pm}$, two of the pseudoscalars are identical, one expects additional zeros. Since three
pseudoscalars cannot be in a $0^{+}$spin parity state, $1^{-}$and $2^{+}$exhaust natural spin parity combinations for spin less than 3 . For the case of $1^{-}$one expects an additional zero along the $y$-axis (symmetry axis), while in the case of $2^{+}$one expects a higher order zero at the top of the Dalitz plot. In order to quantitatively rule out the $k \pi \pi$ final states of $1^{-}$and $2^{+}$ we have utilized the phenomenological matrix elements of Zomach. ${ }^{5}$ These are the simplest matrix elements and are subject to multiplication by arbitrary form factors. Barring the presence of rapidly varying form factors, they can be expected to give a good approximation to the extent of the regions of depopulatiun, allowing a quantitative comparison with the experinental distribution.

For $J^{P}=1^{-}$the matrix element is constructed from an axial vector symetric under the exchange of the two pions. The essential form of such a quantity is $\left(T_{\pi_{1}}-T_{\pi_{2}}\right) \vec{\pi}_{1} \times \vec{\pi}_{2}$, where $\vec{\pi}$ represents a pion momentum in the rest frame of the $K \pi \pi(1876)$, and $T_{\pi}$ represents its hinetic energy. For the case of unpolarized porduction one then expects an intensity $\mathrm{I}_{1}$ - $\mathfrak{q}$ iven by

$$
I_{1}-\propto\left|T_{\pi_{1}}-T_{\pi_{2}}\right|^{2}\left|\vec{\pi}_{1} \times \vec{\pi}_{2}\right|^{2}
$$

To compare the distribution of $I_{1}$ - with the data, we have divided the Dalitz plot into two discrimination regions divided by a contour of constant $I_{1}-$. The particular contour was chosen so that an equal number of events would be found in each region for a phase space decay of the state $k \pi \pi(1876), 5$ as detcrmined by a Monte-Carlo calculation. Owing to the approximitely miform
 Figures 3 a and 3 b show the $K^{\mp} n^{+} \pi^{+}$invariant mas spectra for cents with

Dalitz variables lying inside the two $1^{-}$discrimination regions as indicated by the shaded area in the respective inserts.

A fit to a Gaussian signal over the scaled background of fig. Ib reveals $34 \pm 8$ signal events in the peripheral region compared to $38+9$ signal cvent:: in the central region. Such a division is consistent with equal population with a $x^{2}$ of 0.1 for one degree of freedom (DF) or a confidence level $C L=75^{\circ}$. On the other hand, a Monte-Carlo simulation of $K \pi \pi$ decays using the intensity distribution $I_{1-}$ gives an expected population division of $1: 8.2$ for peripheral to central region. This is effectivcly ruled out with a $x^{2}$ of 18.1 ( $\mathrm{CL}=2 \times 10^{-5}$ ).

For $2^{+}$we construct a symmetric, tracelcss, second-rank tensor which is also symmetric under the exchange of the two pions. We use $A^{i j}=\Delta \pi^{i} q^{j}+\Delta \pi^{j}{ }_{4}$ Wherc $\Delta \pi$ is the difference of the pion momenta and $q$ is their cross product. For unpolarized production one expects an intensity given by:

$$
I_{2}+\propto \sum_{i} \sum_{j} A^{i j} A_{j i}=\left|\vec{\pi}_{1}-\vec{\pi}_{2}\right|^{2}\left|\vec{\pi}_{1} \times \vec{\pi}_{2}\right|^{2}
$$

Here we again divide the Dalitz plot into two regions, using a contur of constant $I_{2}+$ chosen to give equal population for phase space decay. $I_{2}+$ depopulates the peripheral region relative to the central region by $1: 5.6$. Figure 3 b and 3 c show the $\mathrm{K}^{\mp} \pi^{ \pm} \pi^{+}$invariant mass spectra for events with Dalitz variables in the shaded $2^{+}$discrimination regions. Our fits give $31 \pm 9$ events in the peripheral regions and $35 \pm 10$ events in the contral region. This result is again consistent with equal population with a $x^{2}$ of 0.1 for one $D F\left(C L=75^{\circ}\right)$, and inconsistent with $\mathrm{I}_{2}+$ with a $x^{2}$ of 9.4 for one $D F(C L=0.002)$. The ubserved sample population of the $2^{+}$peripheral discrimination region indicates the absence of a general boundary zero. The
absence of such a zero argues against natural spin parity final states of spin 3 and greater as well.

In summary the distribution in the lalitz plot is incompatible with the zeros expected for spin parity $1^{-}$or $2^{+}$for the $k \pi \pi(1876)$. Parity violation then follows from the observation that the presuned isomultiplet state at 186.5 $\mathrm{MeV} / \mathrm{c}^{2}$ decays into $K \pi$, a natural spin parity statc.

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6. A phase space decay is possible for the $K \pi \pi$ spin parity assignment of $0^{-}$.

FIGURE CAPTIONS
Fig. 1. The $K \pi \pi$ mass distributions with the cuts designed to enhance the signal-to-background ratio: $E_{c m}=3.90-4.25 \mathrm{GeV}$ and $M_{\mathrm{rec}}=1.90$ - $2.04 \mathrm{GeV} / \mathrm{c}^{2}$. (a) Exotic combination $\mathrm{K}^{\mp} \pi^{ \pm} \pi^{ \pm}$; (b) non-exotic combination $\mathrm{K}^{ \pm} \pi^{+} \pi^{-}$.

Fig. 2. Dalitz plots, folded around $y$-axis, for the $K \pi \pi$ system with the mass cuts $M=1.86-1.92 \mathrm{GeV} / \mathrm{c}^{2}$ and the cuts given for Fig. 1. (a) Exotic combination $\mathrm{K}^{\dagger} \pi^{ \pm} \pi^{ \pm}$; (b) non-exotic combination $\mathrm{K}^{ \pm} \pi^{+} \pi^{-}$. llere $Q=T_{k}+T_{\pi_{1}}+T_{\pi_{2}}$.
Fig. 3. $M\left(K^{\dagger} \pi^{ \pm} \pi^{ \pm}\right)$distributions for the same data sample as in fig. 2. (a) "peripheral" and (b) "central" regions (on the folded plot) for a contour of a $1^{-}$matrix element as indicated by the shaded regions of the inserts, (c) "peripheral" and (d) "central" regions for a contour of a $2^{+}$matrix element. The solid curves are fits to a Gaussian signal over the scaled backgrounds of Fig. lb.


Fig. 1


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Fig. 2


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Fig. 3


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