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## ABSTRACT

Data from $\mathrm{e}^{+} \mathrm{e}^{-}$annihilation into hadrons at 4.8 GeV center-of-mass energy were used to search for charmed mesons in the mass range 1.5 to $4.0 \mathrm{GeV} / \mathrm{c}^{2}$. We looked for narrow peaks in the invariant mass distributions for $\mathrm{K}^{\mp} \pi^{+}, \mathrm{K}_{\mathrm{S}}^{0} \pi^{+} \pi^{-}, \pi^{+} \pi^{-}$, $\mathrm{K}^{+} \mathrm{K}^{-}, \mathrm{K}^{\mp} \pi^{ \pm} \pi^{ \pm}, \mathrm{K}_{\mathrm{S}}^{0} \pi^{ \pm}, \mathrm{K}_{\mathrm{S}}^{0} \mathrm{~K}^{\ddagger}$, and $\pi^{+} \pi^{-} \pi^{ \pm}$. We Sresent upper limits for the inclusive production cross section times the branching ratio for charmed mesons having these decay modes.
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[^0]Since the discoveries of the $\psi(3095)^{1}$, the $\psi(3684)^{2}$ and the broad enhancement ${ }^{3}$ at 4.15 GeV center-of-mass energy ( $E_{\mathrm{cm}}$ ) in the total cross section for $e^{+} e^{-}$annihilation into hadrons, there has been renewed interest in determining whether the increase in $R=\sigma\left(e^{+} e^{-} \rightarrow\right.$ hadrons $) /$ $\sigma\left(e^{+} e^{-} \rightarrow \mu^{+} \mu^{-}\right)$is associated with the onset of charmed particie ${ }^{4-7}$ production.

In order to search for inclusive production of charmed mesons, we looked for narrow peaks in the invariant mass distributions of the following combinations of particles: $\bar{K}^{-} \pi^{ \pm}, K_{S}^{0} \pi^{+} \pi^{-}, \pi^{+} \pi^{-}, K^{+} K^{-}$, $K^{\mp} \pi^{ \pm} \pi^{ \pm}, K_{S}^{0} \pi^{ \pm}, K_{S}^{0} K^{ \pm}$, and $\pi^{+} \pi^{-} \pi^{ \pm}$. These are expected to be dominant decay modes of the conjectured charmed mesons $D^{+}, D^{0}$, and $F^{+}$. 4, 7-11

We have used the data from hadron production by $e^{+} e^{-}$annihilation at $\mathrm{E}_{\mathrm{cm}}=4.8 \mathrm{GeV}$ taken with the SIAC-LBL solenoidal magnetic detector at the SIAC electron-positron colliding beam facility SPEAR. The detector, trigger requirements, and the selection of events from $e^{+} e^{-}$ annihilation into hadrons have been described previously. ${ }^{3,12}$ The data sample consisted of 9914 hadronic events taken with an integrated luminosity of $787.2 \mathrm{nb}^{-1}$.

All particles associated with an event vertex were used in forming the combinations. For the $\pi^{+} \pi^{-}$and $\pi^{+} \pi^{-} \pi^{ \pm}$combinations and the $K^{+} K^{-}$ combinations each particle was assumed to have the mass of a charged pion or a charged kaon. For $\mathrm{K}^{\dagger} \pi^{ \pm}$the combination was entered twice in the mass histogram -- once for each of the two possible particle mass assignments. For $\mathrm{K}^{\mp} \pi^{ \pm} \pi^{ \pm}$combinations the particle with charge opposite that of the other two particles in a three-particle, charge $\pm$ I combination was assigned a kaon mass. No attempt was made to identify charged kaons
by time of flight since such identification was not reliable for particles with momentum greater than $600 \mathrm{MeV} / \mathrm{c}$. Combinations containing a $K_{S}^{0}$ were formed by finding pairs of oppositely-charged particles having invariant mass between 470 and $520 \mathrm{MeV} / \mathrm{c}^{2}$ assuming pion masses and then including with this pair one or two other particles which were assumed to have pion or kaon masses. A distinct ${ }^{\prime} K_{S}^{O}$ signal was observed with a signal-to-background ratio of $0.25 \pm 0.03$. The observed invariant mass distributions are shown in Figs." 1 and 2.

The invariant mass resolutions and detection efficiencies for particles having the decay modes considered were calculated using a Monte Carlo program which incorporated the geometric acceptance, the trigger efficiency, and all other known inefficiencies of the detector. Also included was the momentum resolution, which was measured to be $\sigma(\mathrm{p}) / \mathrm{p}=$ $1.6 \mathrm{p} \%+0.6 \%$ added in quadrature, where p is the momentum in $\mathrm{GeV} / \mathrm{c}$. The charmed mesons were assumed to be produced with angular and momentum distributions as given by Lorentz invariant phase space and to decay isotropically. The detection efficiency was taken to be the number of particles detected with mass within the fwhm measurement resolution of the true mass, with background subtracted, divided by the number of particles produced. For $K_{S}{ }_{S}$ s the efficiency included the branching ratio for $K_{S}{ }_{\mathrm{S}} \rightarrow$ $\pi^{+} \pi^{-}$. The efficiencies and mass resolutions are indicated in Figs. I and 2. The efficiencies used were the lower limits obtained after considering the effects of binning and the effects of different assumed modes of charmed meson production (e.g., $D^{+} D^{-}$and $D^{+} D^{-}+n \pi$ ).

We found no convincing narrow peaks in the invariant mass distributions. The particles for which we are searching were assumed to have
decay widths much less than our mass resolutions. Two four standard deviation peaks were observed -- at $2.05 \mathrm{GeV} / \mathrm{c}^{2}$ in $\mathrm{K}_{\mathrm{S}}^{0} \mathrm{~K}^{+}$and at 2.40 $\mathrm{GeV} / \mathrm{c}^{2}$ in $\mathrm{K}^{\mp}{ }_{\pi^{-}}^{ \pm} \pi^{ \pm}$. There was no evidence for these peaks in the approximately 2500 hadronic events at $\mathrm{E}_{\mathrm{cm}}=5.0 \mathrm{GeV}$.

Upper limits on the inclusive cross section for the production of a charmed meson times the branching ratio to the decay mode considered were calculated. The background was estimated by fitting a smooth curve to each mass distribution. The fitted curves are represented by the solid lines in Figs. 1 and 2. We calculated $90 \%$ confidence level upper limits for the number of combinations above background in each group of adjacent bins as determined by the fwhm mass resolution. Each upper limit for a number of combinations was converted to an upper limit for cross section times branching ratio by dividing by the efficiency and the integrated luminosity. In Table $I$ we present the largest $90 \%$ confidence upper limit for each of the decay modes for each of three mass regions: 1.5 to $1.85 \mathrm{GeV} / \mathrm{c}^{2}, 1.85$ to $2.4 \mathrm{GeV} / \mathrm{c}^{2}$, and 2.4 to $4.0 \mathrm{GeV} / \mathrm{c}^{2}$. The mass region of particular interest for the production of charmed mesons ranges from $1.85 \mathrm{GeV} / \mathrm{c}^{2}$ (half the mass of the $\psi(3684)$ ) to $2.4 \mathrm{GeV} / \mathrm{c}^{2}$ (half of $E_{c m}$ ) since charmed mesons are expected to be produced in charmedanticharmed pairs. We have also calculated combined largest $90 \%$ confidence upper limits for pairs of mass distributions in which the combinations have the same charge and assumed strangeness. In this case we have used efficiencies for combinations including $K^{0}$ or $\bar{K}^{0}$ which are one half the efficiencies for the $K_{S}^{O}$ combinations. The combined upper limit can be less than either of the individual upper limits if dips in one mass distribution occur at the same masses as peaks in the other distribution.

The combined upper limits are also listed in Table 1.
If one assumes that all of the increase in $R$ is due to charmed meson production, these limits are inconsistent with models in which the lowest mass charmed mesons decay nonleptonically according to a conventional current-current weak interaction. ${ }^{4-11}$ The dominant decays of the lowest mass charmed mesons are expected to be nonleptonic decays into two and three hadrons. $7,8,13$ For charmed mesons of $2 \mathrm{GeV} / \mathrm{c}^{2}$ mass the branching ratios for these decays are estimated ${ }^{14}$ to add up to approximately $80 \%$ after allowing for a $10 \%$ branching ratio into leptonic and semileptonic modes. ${ }^{7}$ The excess hadron cross section "above the three-color quark model prediction of $R=2$ is $10.7 \pm 1.5 \mathrm{nb}$ at $\mathrm{E}_{\mathrm{cm}}=4.8 \mathrm{GeV} .^{3}$ If this cross section is due to production of pairs of charmed mesons resulting in equal cross sections for $e^{+} e^{-} \rightarrow D^{+}$(or $D^{-}$) + anything, $e^{+} e^{-} \rightarrow$ $\mathrm{F}^{+}$(or $\mathrm{F}^{-}$) + anything, and $\mathrm{e}^{+} \mathrm{e}^{-} \rightarrow \mathrm{D}^{0}$ (or $\bar{D}^{0}$ ) + anything, the cross section for inclusive production of $D^{+}$or $D^{-}$, for example, would be 7.I nb. Of the two-body decay modes of $D^{+}$, the dominant mode is expected to be $\bar{K}^{0} \pi^{+}, 7,9$ although this mode may be suppressed as it is in the limit of exact $\operatorname{SU}(3) .^{7,8}$ of the three-body decays of $D^{+}$, the dominant modes are expected to be $7,8 \mathrm{~K}^{-} \pi^{+} \pi^{+}, \overline{\mathrm{K}}^{0} \overline{\mathrm{~K}}^{0} \mathrm{~K}^{+}, \overline{\mathrm{K}}^{0} \pi^{+} \eta$, and $\bar{K}^{0} \pi^{+} \pi^{0}$ which would occur in the ratio ${ }^{15} 4: 4: 3: 1$ (without correcting for phase space). Our combined upper limit for $\bar{K}^{0} \pi^{ \pm}$, $K^{0} \pi^{ \pm}$, and $K^{+} \pi^{ \pm} \pi^{+}$is 0.51 nb which gives an upper limit on the branching ratio:

$$
\frac{\Gamma\left(D^{+} \rightarrow \bar{K}^{0} \pi^{+} \text {or } K^{-} \pi^{+} \pi^{+}\right)}{\Gamma\left(D^{+} \rightarrow \mathrm{a} 11\right)}<\frac{0.51 \mathrm{nb}}{7.1 \mathrm{nb}}=7.2 \% .
$$

This upper limit on the $D^{+}$branching ratio into $\overline{\mathrm{K}}^{0} \pi^{+}$and $K^{-} \pi^{+} \pi^{+}$is at least a factor of three lower than estimates based on the conventional
phenomenological model. ${ }^{1}$
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14. See Table IV of Reference 7 .
15. H. Harari, private commanication.

TABLE I. Largest Upper Limits at the $90 \%$ Confidence Level for Inclusive Production Cross Section Times Branching Ratio (nb)

| Decay Mode | Mas <br> 1.50 to 1.85 | Region ( $\mathrm{GeV} / \mathrm{c}$ 1.85 to 2.40 | $2.40 \text { to } 4.00$ |
| :---: | :---: | :---: | :---: |
| $\mathrm{K}^{-} \pi^{+}$and $\mathrm{K}^{+} \pi^{-}$ | 0.25 | 0.18 | 0.08 |
| $\mathrm{K}_{\mathrm{S}}^{0} \pi^{+} \pi^{-}$ | 0.57 | 0.40 | 0.29 |
| $\pi^{+} \pi^{-}$ | 0.13 | 0.13 | 0.09 |
| $\mathrm{K}^{+} \mathrm{K}^{-}$ | 0.23 | 0.12 | 0.10 |
| $\mathrm{K}^{-} \pi^{+} \pi^{+}$and $\mathrm{K}^{+} \pi^{-} \pi^{-}$ | 0.51 | 0.49 | 0.19 |
| $\mathrm{K}_{S} \mathrm{~J}^{+}$and $\mathrm{K}_{S}^{\mathrm{O}} \pi^{-}$ | 0.26 | 0.27 | 0.09 |
| $\mathrm{K}_{\mathrm{S}}^{\mathrm{O}} \mathrm{K}^{+}$and $\mathrm{K}_{\mathrm{S}}^{\mathrm{O}} \mathrm{K}^{-}$ | 0.54 | 0.33 | 0.09 |
| $\pi^{+} \pi^{-} \pi^{+}$and $\pi^{+} \pi^{-} \pi^{-}$ | 0.48 | 0.38 | 0.18 |
| $\mathrm{K}^{\overline{+}} \pi^{ \pm}, \overline{\mathrm{K}}^{0} \pi^{+} \pi^{-}$and $\mathrm{K}^{0} \pi^{+} \pi^{-}$ | 1.16 | 0.90 | 0.58 |
| $\mathrm{K}^{+} \mathrm{K}^{-}$and $\pi^{+} \pi^{-}$ | 0.23 | 0.16 | 0.15 |
| $\mathrm{K}^{\overline{+}} \pi^{ \pm} \pi^{ \pm}, \overline{\mathrm{K}}^{0} \pi^{ \pm}$and $\mathrm{K}^{0} \pi^{ \pm}$ | 0.64 | 0.51 | 0.30 |
| $\bar{K}^{0} \mathrm{~K}^{ \pm}, \mathrm{K}^{0} \mathrm{~K}^{ \pm}$and $\pi^{+} \pi^{-} \pi^{ \pm}$ | 1.10 | 0.76 | 0.29 |

## Figure Captions

1. Observed invariant mass distributions in $25 \mathrm{MeV} / \mathrm{c}^{2}$ bins for charge 0 combinations. The solid lines represent smooth curves fitted to the data.
2. Observed invariant mass distributions in $25 \mathrm{MeV} / \mathrm{c}^{2}$ bins for charge $\pm$ I combinations. The solid lines represent smooth curves fitted to the data.


Fig. 1


Fig. 2


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