

1 **Discovery of very high energy gamma rays from PKS 1424+240**
 2 **and multiwavelength constraints on its redshift**

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ABSTRACT

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We report the first detection of very-high-energy¹ (VHE) gamma-ray emission above 140 GeV from PKS 1424+240, a BL Lac object with an unknown redshift. The photon spectrum above 140 GeV measured by VERITAS is well described by a power law with a photon index of $3.8 \pm 0.5_{\text{stat}} \pm 0.3_{\text{syst}}$ and a flux normalization at 200 GeV of $(5.1 \pm 0.9_{\text{stat}} \pm 0.5_{\text{syst}}) \times 10^{-11} \text{ TeV}^{-1} \text{ cm}^{-2} \text{ s}^{-1}$, where stat and syst denote the statistical and systematical uncertainty, respectively. The VHE flux is steady over the observation period between MJD 54881 and 55003 (2009 February 19 to June 21). Flux variability is also not observed in contemporaneous high energy observations with the *Fermi* Large Area Telescope (LAT). Contemporaneous X-ray and optical data were also obtained from the *Swift* XRT and MDM observatory, respectively. The broadband spectral energy distribution (SED) is well described by a one-zone synchrotron self-Compton (SSC) model favoring a redshift of less than 0.1. Using the photon index measured with *Fermi* in combination with recent extragalactic background light (EBL) absorption models it can be concluded from the VERITAS data that the redshift of PKS 1424+240 is less than 0.66.

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Subject headings: BL Lacertae objects: individual (PKS 1424+240 =

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VER J1427+237); gamma rays: observations

¹ γ -ray emission above 100 GeV

1. Introduction

54

55 PKS 1424+240 was detected as a radio source by Condon et al. (1977). It was classified
56 as a blazar by Impey & Tapia (1988) from optical polarization studies. Fleming et al.
57 (1993) verified the polarization results and also reported non-thermal X-ray radiation,
58 further strengthening the classification.

59

Blazar emission is dominated by non-thermal radiation, which is thought to be related
60 to charged particle acceleration near a massive compact object in the center of the host
61 galaxy, or in outflowing relativistic jets. The SED is characterized by two peaks. The lower
62 peak is widely accepted to be synchrotron radiation from relativistic electrons and occurs
63 in the IR to X-ray bands. The higher energy peak is in the gamma-ray band, sometimes at
64 energies as high as a few TeV, and can be created via either inverse-Compton scattering by
65 relativistic electrons or hadronic interactions (for a review see Böttcher 2007, and references
66 therein). The position of the synchrotron peak of PKS 1424+240 has not been measured,
67 but it can be constrained from optical and X-ray data to be between 10^{15} Hz and 10^{17} Hz.
68 Depending on the definition used, PKS 1424+240 is either an intermediate-frequency-peaked
69 BL Lac (IBL) (Nieppola et al. 2006) or a high-frequency-peaked BL Lac (HBL) (Padovani
70 & Giommi 1996; Abdo et al. 2009a).

71

Gamma-ray emission from PKS 1424+240 was not detected by EGRET (Fichtel et al.
72 1994), but was recently observed with the *Fermi* LAT pair-conversion telescope (Abdo et
73 al. 2009b,c). The reported flux above 100 MeV of $(6.2 \pm 0.8) \times 10^{-8} \text{cm}^{-2} \text{s}^{-1}$ and hard
74 spectral index $\Gamma = 1.80 \pm 0.07$ ($dN/dE \propto E^{-\Gamma}$) triggered VERITAS observations.

75

The redshift of PKS 1424+240 is not known. Scarpa & Falomo (1995) have derived
76 a lower limit on the redshift of $z > 0.06$ and Sbarufatti et al. (2005) a limit of $z > 0.67$,
77 both assuming a minimum luminosity of the host galaxy. The latter authors also reported
78 evidence that the ratio of the nucleus to host luminosity is much larger than 100, which is

79 typical for BL Lac objects but complicates photometric determination of the redshift.

80 We report the detection of PKS 1424+240 in VHE gamma rays and contemporaneous
81 observations with *Fermi*, *Swift*, and the MDM observatory. Shortly after the VHE discovery
82 (Ong 2009), it was confirmed by the MAGIC collaboration (Teshima 2009). This discovery
83 marks the first *Fermi*-motivated VHE discovery.

84 2. Observations and Analysis of VERITAS Data

85 The VERITAS observatory, located in southern Arizona at 1.3 km a.s.l., is described in
86 detail in Weekes et al. (2002) and Holder et al. (2006).

87 PKS 1424+240 was observed with VERITAS between 2009 February 19 and June 21 at
88 zenith angles between 7° and 30° . The observations were performed in wobble mode (Fomin
89 et al. 1994) with a 0.5° offset, enabling simultaneous background estimation. About one
90 third of the data were taken during low levels of moonlight. About 65% of the observations
91 were conducted with only three telescopes due to the relocation of one telescope, which
92 began in May and was completed in August 2009. Of the 37.3 hours of data, 28.5 hours
93 survive standard data quality selection.

94 Events are reconstructed following the procedure in Acciari et al. (2008). The recorded
95 shower images are parameterized by their principal moments, giving an efficient suppression
96 of the far more abundant cosmic-ray background. Two separate sets of cuts are applied
97 to reject background events, hereafter called *soft* and *medium*. These cuts are applied
98 to the parameters *mean scaled width* (MSW), and *mean scaled length* (MSL), apparent
99 altitude of the maximum Cherenkov emission (shower maximum), and θ^2 , the squared
100 direction between the position of PKS 1424+240 and the reconstructed origin of the event.
101 Studies on independent data sets show that a shower-maximum cut significantly improves

102 the low energy sensitivity. *Soft* cuts have a higher sensitivity for sources with soft photon
 103 spectra because of a lower energy threshold resulting from a minimum *size* cut of 50
 104 photoelectrons. In the *medium* cuts a minimum *size* cut of 100 photoelectrons is applied.
 105 *Size* is a measure of the recorded photoelectrons from a shower and a good indicator of
 106 the energy of the primary. For the *soft*-cuts analysis the remaining cuts are $MSW < 1.06$,
 107 $MSL < 1.30$, shower maximum > 7 km, and $\theta^2 < (0.14^\circ)^2$, and $MSW < 1.04$, $MSL < 1.28$,
 108 shower maximum > 5 km, and $\theta^2 < (0.1^\circ)^2$ for the *medium* cuts. The cuts have been
 109 optimized *a priori* to yield the highest sensitivity for a source with 5% of the Crab Nebula
 110 gamma-ray flux. The results are independently reproduced with two different analysis
 111 packages explained in Cogan (2008) and Daniel (2008).

112 In the *soft*-cuts analysis, 1907 on-source events remain out of 1.25×10^7 triggered
 113 events. The background calculated with the reflected-region method (Berge et al. 2007) is
 114 1537 events, which leaves an excess of 370 events. Figure 1 shows the corresponding θ^2
 115 distribution. The statistical significance of the observed excess is 8.5 standard deviations,
 116 σ , calculated with Equation 17 of Li & Ma (1983), and including a trials factor of two
 117 for the two sets of cuts. In the *medium*-cuts analysis the post-trials significance is 4.8σ
 118 (329 on-source events with an estimated background of 244). The angular distribution of
 119 the excess events is consistent with a point source. The center of gravity of the excess is
 120 $14^{\text{h}} 27^{\text{m}}, 0^{\text{s}} \pm 7^{\text{s}}_{\text{stat}}, 23^\circ 47' 40'' \pm 2'_{\text{stat}}$ coinciding with the position of PKS 1424+240 in radio
 121 (Fey et al. 2004). The VERITAS source name is VER J1427+237.

122 Figure 2 shows the light curve of PKS 1424+240 in different energy bands for the time
 123 period overlapping the VERITAS observations. The flux measured by VERITAS above
 124 140 GeV is $\sim 5\%$ of the Crab Nebula flux. The VERITAS data from each dark period² are
 125 combined into a single bin to produce a light curve, which is consistent with a constant flux,

²The ~ 3 week observing period between full moons

126 $\chi^2=0.3$ for 3 degrees of freedom (d.o.f.). However, even a doubling in flux would have been
 127 difficult to detect. There is no evidence for strong flaring episodes on shorter timescales.

128 Figure 3 shows the differential photon spectra derived with the *soft*-cuts and
 129 *medium*-cuts analyses, with one overlapping flux point at 260 GeV. The fraction of events
 130 that are used both in the last bin in the *soft*-cuts analysis and in the second bin in the
 131 *medium*-cuts analysis is about 2%, small enough to allow a combined fit of the flux points
 132 from the two analyses, with the more significant *soft*-cuts result at 260 GeV used in the
 133 fit. The combined spectrum is well parameterized ($\chi^2=2.2$ for 4 d.o.f.) by a power law
 134 $dN/dE = F_0 \cdot (E/E_0)^{-\Gamma}$, where the photon index Γ is $3.8 \pm 0.5_{\text{stat}} \pm 0.3_{\text{syst}}$ and F_0 is
 135 $(5.1 \pm 0.9_{\text{stat}} \pm 0.5_{\text{syst}}) \times 10^{-11} \text{ TeV}^{-1} \text{ cm}^{-2} \text{ s}^{-1}$ for $E_0 = 200 \text{ GeV}$. The combined spectrum is
 136 consistent with the fit of the *soft*-cuts points alone, albeit with half the uncertainty on the
 137 photon index.

138 3. Multiwavelength Observations

139 Gamma-ray observations with *Fermi*-LAT (100 MeV to 300 GeV), X-ray and
 140 optical observations with *Swift* XRT (0.2–10 keV) and UVOT (170–650 nm), and optical
 141 observations in the R, V and I bands at the MDM observatory were obtained simultaneously
 142 or quasi-simultaneously with the VERITAS observations.

143 The LAT pair-conversion telescope on board the *Fermi* Gamma-ray Space Telescope
 144 continuously monitors the entire sky between 100 MeV and several hundred GeV (Atwood
 145 et al. 2009). The LAT data overlapping with the VERITAS observations were analyzed by
 146 selecting “diffuse” class events that have the highest probability of being photons. Further
 147 event selection was done by only accepting events that come within a 15° radius from
 148 PKS 1424+240 and have energies between 0.1 and 300 GeV. Events with zenith angles above

149 105° were excluded to limit contamination by gamma rays coming from the Earth’s albedo.

150 The analysis of the photon spectrum and light curve were performed with the standard
 151 likelihood analysis tools available from HEASARC `ScienceTools v9r15p2`. Accidental
 152 coincidences with charged cosmic rays in the detector were accounted for using instrument
 153 response functions `P6_V3_DIFFUSE`. The background model used to extract the gamma-ray
 154 signal includes a Galactic diffuse emission component and an isotropic component³. The
 155 isotropic component includes contributions from the extragalactic diffuse emission as well as
 156 from residual charged particle backgrounds. The spectral shape of the isotropic component
 157 was derived from residual high latitude events after the Galactic contribution had been
 158 modeled. The background model also takes into account unresolved gamma-ray sources in
 159 the region of interest, thus avoiding a bias in the spectral reconstruction. To further reduce
 160 systematic uncertainties in the analysis, the normalization and spectral parameters in the
 161 background model were allowed to vary freely during the spectral point fitting.

162 The *Fermi*-LAT flux measurements are shown in the broadband SED in Figure 4. The
 163 flux values are unfolded by assuming an underlying power-law, giving an integrated flux
 164 over the 0.1–300 GeV band $(7.04 \pm 0.96_{\text{stat}} \pm 0.38_{\text{sys}}) \times 10^{-8} \text{cm}^{-2} \text{s}^{-1}$, and a differential
 165 photon spectral index $\Gamma_{\text{LAT}} = 1.73 \pm 0.07_{\text{stat}} \pm 0.05_{\text{sys}}$. The light curve of the integral flux
 166 above 100 MeV is plotted with 10-day bins in Figure 2. A fit with a constant yields a
 167 $\chi^2 = 11.5$ for 11 d.o.f., suggesting no variability.

168 Target of opportunity observations of nearly 16 ksec, distributed over ten observing
 169 periods, were obtained with *Swift* (Gehrels et al. 2004) following the detection of VHE
 170 emission from PKS 1424+240. The data reduction and calibration of the XRT data
 171 were completed with the HEASoft v6.6.3 standard tools. The XRT data were taken in

³<http://fermi.gsfc.nasa.gov/ssc/data/access/lat/BackgroundModels.html>

172 photon-counting mode and contained modest pile up for nine of the observations, which
 173 was taken into account by masking a region with 3-6 pixel radius around the source. The
 174 outer radius chosen for the signal region was 20 pixels and a background region of similar
 175 size was chosen about 5 arcminutes off source.

176 X-ray energy spectra could be extracted from all observing periods and are well
 177 described by an absorbed power law using the fixed Galactic column density of neutral
 178 hydrogen from Dickey & Lockman (1990) ($N_{\text{H}} = 0.264 \times 10^{21} \text{ cm}^{-2}$). The fit spectral index
 179 varies between 2.1 and 2.9 (photon index between 3.1 and 3.9) with a typical statistical
 180 uncertainty of 0.1, while the normalization changes between 1.40×10^{-2} and 0.74×10^{-2}
 181 photons $\text{keV}^{-1} \text{cm}^{-2} \text{s}^{-1}$ at 1 keV with a typical uncertainty of $0.07 \times 10^{-2} \text{ keV}^{-1} \text{cm}^{-2} \text{ s}^{-1}$.
 182 For the modeling of the SED we use the average spectrum shown in Figure 4. The light
 183 curve shows that the X-ray flux is variable over the ten days of observation. A fit to a
 184 constant flux yields a χ^2 of 60 for 9 d.o.f.. UVOT observations were taken in the six V,
 185 B U, W1, M2 and W2 bands and were calibrated using standard techniques (Poole et al.
 186 2008). The reddening has been accounted for by interpolating the absorption values from
 187 Schlegel et al. (1998) with a galactic spectral extinction model (Fitzpatrick 1999) obtaining
 188 0.663, 0.968, 0.922 mag for the three UV bands W1, M2, and W2 and an assumed redshift
 189 of $z=0$. The corresponding light curves are shown in Figure 2.

190 Data in the optical bands were also obtained with the 1.3m telescope and 4K
 191 imager of the MDM observatory located on the west ridge of Kitt Peak near Tucson,
 192 Arizona. The CCD was operated in unbinned mode, which produces an image scale of
 193 0.315 arcseconds/pixel. 2-4 images were obtained in the V, R, and I bands during each
 194 observation. Physical magnitudes were computed from differences in the instrumental
 195 magnitudes from the three standard stars in Fiorucci & Tosti (1996), assuming that the
 196 magnitudes quoted in that paper are exact. The magnitudes were then corrected for

197 Galactic extinction using extinction coefficients calculated following Schlegel et al. (1998),
 198 taken from NED⁴, and were then converted into νF_ν fluxes. During the 14-day span of the
 199 optical photometry, the visual brightness increased by 14% and colors became slightly bluer.

200 4. Redshift Upper Limit

201 The observed gamma-ray spectrum above 100 GeV is affected by the absorption of
 202 gamma rays via pair conversion with EBL photons (Nikishov A. I. 1962; Gould & Schröder
 203 1967). Depending on the redshift, this effect can result in a significant softening of the
 204 spectrum. We estimate an upper limit of the redshift of PKS 1424+240 by assuming an
 205 intrinsic VHE spectrum and making use of the recent advances in extragalactic background
 206 light (EBL) modeling.

207 We assume that the intrinsic spectrum above 140 GeV can be described by a power law.
 208 The hardest photon index that we consider is 1.7, which is the value from the simultaneous
 209 *Fermi* observations. The use of Fermi observations allows a model independent estimate
 210 of the hardest possible intrinsic spectrum (see also Abdo et al. 2009d). The power law
 211 with an index of 1.7 is absorbed using recent EBL models from Franceschini et al. (2008),
 212 Gilmore et al. (2009), and Finke et al. (2009). After absorption the shape of the spectrum
 213 is fit to the VERITAS spectrum with the normalization as a free parameter, and the best
 214 estimate of the redshift is determined by minimizing χ^2 . For an intrinsic index of 1.7 this
 215 best fit redshift is $z = 0.5 \pm 0.1_{\text{stat}} \pm 0.1_{\text{syst}}$ with a $\chi^2 = 4$ and 5 d.o.f. . The systematic
 216 uncertainty is estimated from the differences in the EBL models.

217 Instead of assuming no break in the photon spectrum, a more likely scenario is that the
 218 intrinsic spectrum softens with increasing energy. In this case an index of 1.7 is an upper

⁴ <http://nedwww.ipac.caltech.edu>

219 limit of the true photon index and the corresponding upper limit on the redshift is $z < 0.66$
 220 with a 95% confidence level.

221 5. Spectral Modeling

222 The spectral energy distribution, comprising data from all of the observations, is shown
 223 in Figure 4. We model the SED using an improved version of the leptonic one-zone jet
 224 model of Böttcher & Chiang (2002). These calculations include time-dependent particle
 225 injection and evolution, and they allow for quasi-equilibrium solutions in which a slowly
 226 varying broken power-law electron distribution arises from a single power-law injection
 227 function, $dn_{\text{inj}}/d\gamma \propto \gamma^{-q}$ with a low- and high-energy cutoff γ_1 and γ_2 , respectively. All
 228 model fits presented here are in the fast-cooling regime, with the cooling break at γ_1 . We
 229 define the magnetic-field equipartition ϵ_B as $\epsilon_B \equiv L_B/L_e$ with L_B the Poynting flux derived
 230 from the magnetic energy density and L_e the energy flux of the electrons propagating
 231 along the jet. The corresponding partition fraction for an electron-proton plasma assuming
 232 $L_p = 10 \times L_e$ of cold protons would be one order of magnitude lower. For an in-depth
 233 description of this quasi-equilibrium jet model, see Acciari et al. (2009).

234 There are few observational constraints on the model parameters for PKS 1424+240
 235 and the redshift is unknown. No superluminal motion has been resolved in this object, and
 236 it has not been monitored well enough to firmly establish a minimum variability timescale
 237 to constrain the size of the emitting region. The different sizes of the emission region R_B
 238 assumed here are compatible with the X-ray variability timescale of about a day. We
 239 therefore consider a range of plausible redshifts and adopt model parameters which were
 240 typically adequate for modeling other VHE blazars. The redshifts investigated range from
 241 $z = 0.05$, similar to the redshift of the nearby HBLs Mrk 421 and Mrk 501, to $z = 0.7$. This
 242 covers the redshift range determined in the previous section and is just above the lower

243 limit set by Sbarufatti et al. (2005), $z > 0.67$.

244 The shape of the high-energy part of the electron spectrum is well constrained by the
 245 rather steep slope of the X-ray spectrum, which has an average photon index $\Gamma_{\text{X-ray}} \sim 3.7$.
 246 In all fits, the relativistic electrons are injected into the emission region with a fixed $q = 5.1$.
 247 Lacking direct constraints on the viewing angle θ_{obs} , it was chosen such that the Doppler
 248 factor $D = (\Gamma[1 - \beta_{\Gamma} \cos \theta_{\text{obs}}])^{-1} = \Gamma$, where Γ is the bulk Lorentz factor of the emitting
 249 material, and $\beta_{\Gamma}c$ is the velocity. The model parameters that were varied are shown in Table
 250 1. Figure 4 shows the fits, after EBL absorption using the model of Gilmore et al. (2009).

251 The SED modeling shows that a reasonable fit can in principle be achieved for any
 252 redshift in the considered range. However, the inset in Figure 4 illustrates that above
 253 $z \sim 0.2$, the model VHE gamma-ray spectrum becomes increasingly too steep compared
 254 with the observed VERITAS spectrum. Furthermore, for redshifts $z > 0.4$ the models
 255 require unreasonably large Doppler factors of $D > 50$. We note that in particular for the
 256 lowest redshift considered, $z = 0.05$, a good fit can be achieved with almost equipartition
 257 between magnetic-field and electron energy densities.

258 An attempt to improve the fit in the gamma-ray bands, by including an external
 259 Compton component, results in a steeper VHE gamma-ray spectrum. This is in conflict
 260 with the VERITAS spectrum and a worse representation of the *Fermi* spectrum. We
 261 therefore conclude that a leptonic fit to the SED of PKS 1424+240 during the VERITAS
 262 observation is possible with a pure SSC model very close to equipartition, in particular if
 263 the redshift of the source is $z < 0.1$.

6. Summary

264

265 We report the detection of PKS 1424+240 in VHE gamma-rays. The observation with
 266 VERITAS was motivated by the release of the first *Fermi* source lists (Abdo et al. 2009b,c)
 267 and this is the first time that *Fermi* observations have led to the discovery of a new source
 268 in the adjacent VHE band.

269

The VHE spectrum of PKS 1424+240 has a photon index of $3.8 \pm 0.5_{\text{stat}} \pm 0.3_{\text{sys}}$, whereas
 270 the spectrum in the *Fermi* energy range has a photon index of $1.73 \pm 0.07_{\text{stat}} \pm 0.05_{\text{sys}}$,
 271 indicating a break in the spectrum at several tens of GeV. The break can be explained
 272 by a one-zone SSC model assuming a wide range of redshifts or could result from EBL
 273 absorption if the redshift is about 0.5 and the intrinsic photon index is 1.7, from which a
 274 redshift upper limit of 0.66 is inferred. The modeling favors a lower redshift but cannot
 275 exclude that PKS 1424+240 is among the most distant sources detected in the VHE regime.

276

PKS 1424+240 is the third extragalactic source detected in the VHE regime with
 277 an unknown or uncertain redshift. It is evident that increased efforts are needed to
 278 determine the redshifts of VHE detected blazars. A redshift measurement will allow a
 279 better understanding of the source-intrinsic mechanisms and the absorption effects which
 280 go along with the gamma-ray propagation.

281

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287

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296 *Facilities:* VERITAS, Swift, Fermi.

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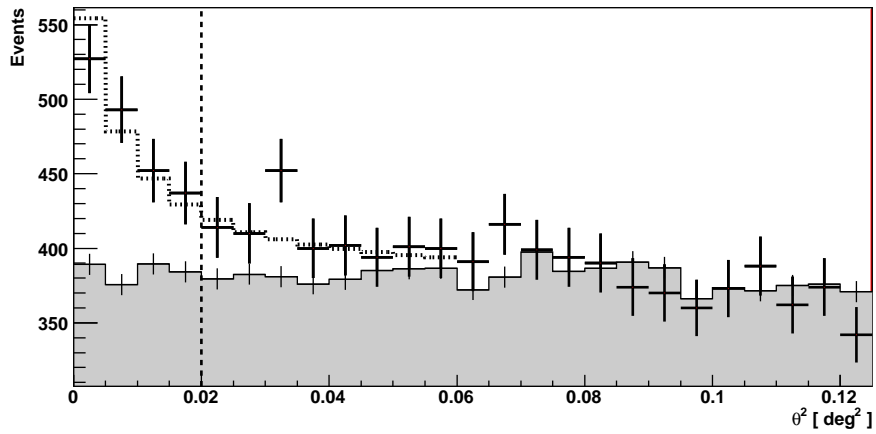


Fig. 1.— Distribution of θ^2 for VERITAS events selected with *soft* cuts. The points with error bars denote the on-source events. The background is shown by the shaded histogram. The dashed vertical line shows the applied θ^2 -cut. The expected distribution for a point source is given by the dotted line.

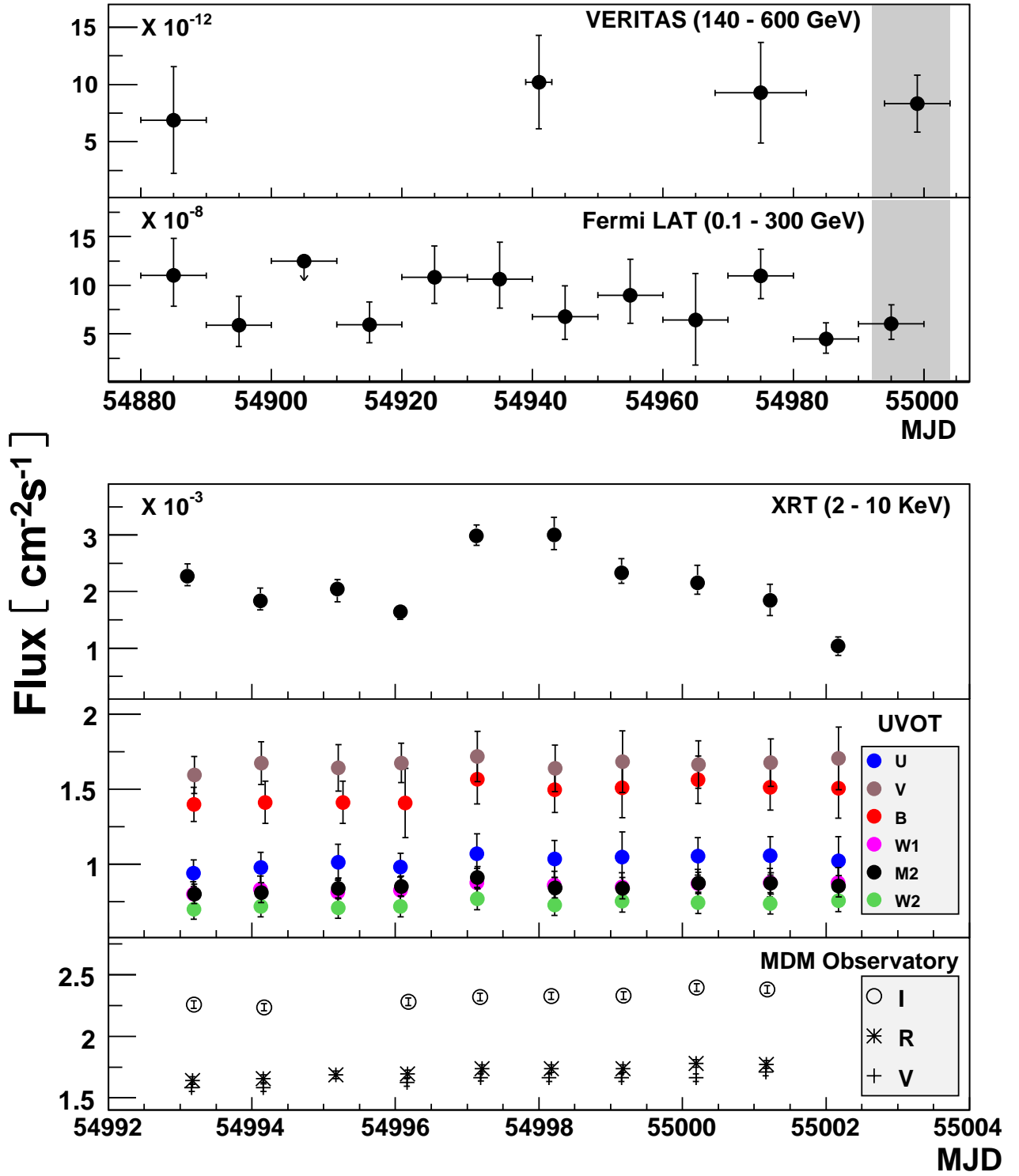


Fig. 2.— Light curves of PKS 1424+240 in VHE gamma rays (VERITAS), HE gamma rays (*Fermi*-LAT), X-rays (*Swift* XRT), UV (*Swift* UVOT) and optical (*Swift* UVOT, MDM). The X-ray, UV and optical light curves cover the time period indicated in the upper two light curves by the shaded region. The horizontal bars in the VHE and HE light curves

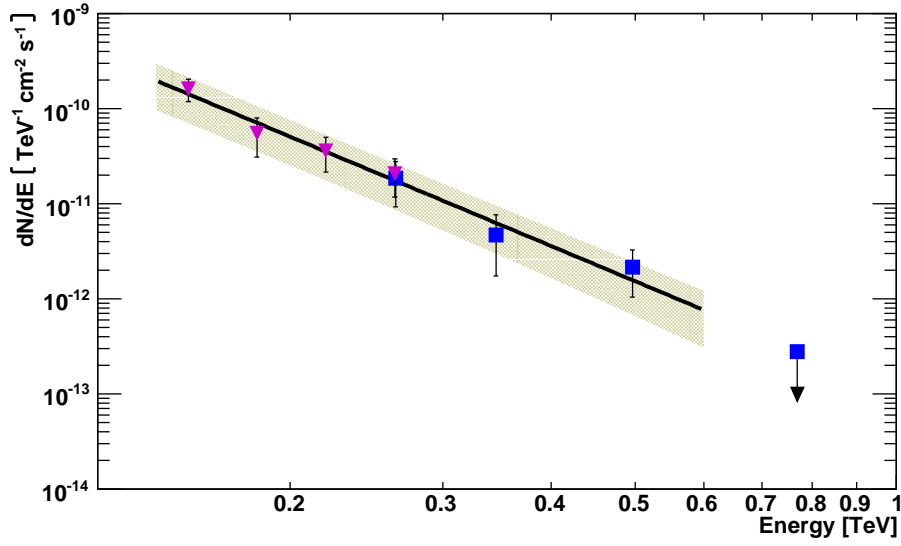


Fig. 3.— The time averaged differential photon spectrum of PKS 1424+240 measured by VERITAS between February 19 and June 21, 2009. The triangles are from the *soft*-cuts analysis and the squares from the *medium*-cuts analysis. The flux point at 260 GeV is reconstructed in both analysis. The solid lines shows the fit with a power law. The shaded area shows the systematic uncertainty of the fit, which is dominated by a 20% uncertainty on the energy scale.

Table 1: SSC fit parameters for PKS 1424+240 as a function of assumed redshift.

Parameter	$z = 0.05$	$z = 0.10$	$z = 0.2$	$z = 0.3$	$z = 0.4$	$z = 0.5$	$z = 0.7$
L_e [10^{43} erg s $^{-1}$]	1.60	4.12	10.7	18.9	29.2	47.1	88.8
L_B [10^{43} erg s $^{-1}$]	1.66	5.47	16.9	31.1	45.9	49.8	66.2
γ_1 [10^4]	3.7	3.7	3.6	3.4	3.2	3.6	3.7
γ_2 [10^5]	4.0	4.0	4.0	4.0	4.5	4.0	4.0
D	15	18	25	30	35	45	60
B [G]	0.37	0.31	0.25	0.24	0.25	0.18	0.14
ϵ_B	1.04	1.33	1.59	1.65	1.57	1.06	0.75
R_B [10^{16} cm]	1.2	2.2	3.4	4.0	4.0	4.5	5.0

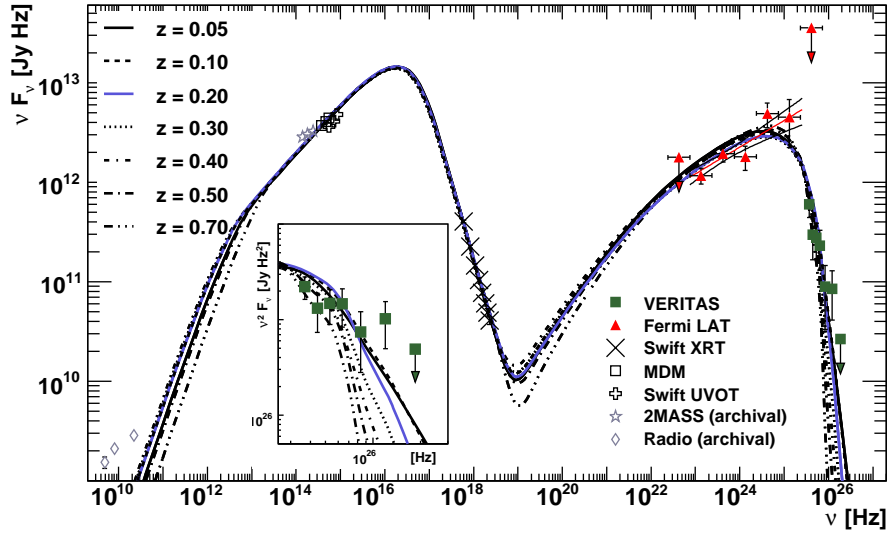


Fig. 4.— SED of PKS 1424+240. The lines show SSC-model fits assuming different redshifts. The inset shows a zoom of the SED on the VERITAS data in a $\nu^2 F_\nu$ representation. The *Fermi* data are presented together with their corresponding power-law fit and one standard deviation uncertainty. The upper limits correspond to 95% confidence level.