# CHARMLESS HADRONIC $B$ DECAYS AT BELLE and BABAR 

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#### Abstract

I report on recent measurements from the Belle and BABAR collaborations on the decay of the $B$ meson to hadronic final states without a charm quark.


## 1 Introduction.

The study of the branching fractions and angular distributions of $B$ meson decays to hadronic final states without a charm quark probes the dynamics of both the weak and strong interactions, and plays an important role in understanding $C P$ Violation (CPV) in the quark sector. $C P$ Violation at the $B$ factories is described graphically by a triangle with sides formed from the CKM matrix elements $V_{q d} V_{q b}^{*}(q=u, c, t)$ and internal angles $\alpha, \beta, \gamma\left(\right.$ or $\left.\phi_{2}, \phi_{1}, \phi_{3}\right)$. Discrepancies in the measured values of the sides and angles could be an indication of New Physics beyond the Standard Model (SM) due to enhanced branching fractions or modified $C P$ asymmetries. The experimental measurements of branching fractions, $C P$ asymmetries, polarization and phases (both weak and strong) can be compared to theoretical models based on, for example, QCD factorization, $\mathrm{SU}(3)$ symmetry and Lattice QCD.

The results presented below assume charge-conjugate states and all branching fraction upper limits (UL) are at the $90 \%$ confidence level (C.L.). The time-integrated $C P$ asymmetry is defined as $A_{C P}=\left(N_{b}-N_{\bar{b}}\right) /\left(N_{b}+N_{\bar{b}}\right)$ where $N_{b}\left(N_{\bar{b}}\right)$ is the number of $B$ mesons containing a $b(\bar{b})$ quark. The latest results are based on a total dataset of $467 \times 10^{6} B \bar{B}$ pairs for $B A B A R$ and $657 \times 10^{6} B \bar{B}$ pairs for Belle, unless indicated.

## 2 Decays involving two-body final states.

The last few years have seen considerable advancement in the prediction of the branching fractions and polarizations of $B$ meson decays to Vector-Vector ( $V V$ ), Vector-Scalar ( $V S$ ) and Vector-Tensor $(V T)$ final states. In general, there has been good agreement between theory and experiment on branching fractions (with some notable exceptions) but the polarization measurements have presented a challenge. The $V V$ states are expected to be almost fully longitudinally polarized $\left(f_{L} \sim 1\right)$ and this should remain true in the presence of penguin loop decays. However, penguin-dominated decays seem to have a smaller $f_{L}\left(\right.$ e.g. $f_{L} \sim 0.5$ for $\left.B \rightarrow \phi K^{*}\right) 1$.

Belle has recently measured the decay $B^{-} \rightarrow K^{* 0} K^{-}$which is dominated by $b \rightarrow d s \bar{s}$ gluonic penguin diagrams. They measure a yield of $47.7 \pm 11.1$ events, corresponding to a branching fraction $\mathcal{B}\left(B^{-} \rightarrow K^{* 0} K^{-}\right)=(0.68 \pm 0.16 \pm 0.10) \times 10^{-6}$ with a $4.4 \sigma$ significance 2 . The event yield for $B^{-} \rightarrow K_{2}^{* 0}(1430) K^{-}$is measured to be $23.4 \pm 12.1$ with an upper limit on the branching fraction of $\mathcal{B}\left(B^{-} \rightarrow K_{2}^{* 0}(1430) K^{-}\right)<1.1 \times 10^{-6}$. A similar analysis has
been done for $B^{0}$ decays to the $V V$ final states $\rho^{0} K^{* 0}$ and $f_{0} K^{* 0} 3$. Unlike an earlier BABAR analysis ${ }^{4}$, Belle sees no evidence for $\rho^{0} K^{* 0}$ and $f_{0} K^{* 0}$ (and, consequently, do not measure $f_{L}$ ) but observes $B^{0} \rightarrow \rho^{0} K^{+} \pi^{-}$and sees first evidence for $B^{0} \rightarrow f_{0} K^{+} \pi^{-}$and $B^{0} \rightarrow \pi^{+} \pi^{-} K^{* 0}$, with branching fractions (significance) of $(2.8 \pm 0.5 \pm 0.5) \times 10^{-6}(5.0 \sigma),\left(1.4 \pm 0.4_{-0.4}^{+0.3}\right) \times 10^{-6}$ $(3.5 \sigma)$, and $\left(4.5_{-1.0-1.6}^{+1.1+0.9}\right) \times 10^{-6}(4.5 \sigma)$, respectively. BABAR has measured $B$ meson decay to an $\omega$ accompanied by a $K^{*}, \rho$ or $f_{0}$. Five measurements have a significance above $5 \sigma$, with another two above $3 \sigma$. This has allowed $B A B A R$ to measure both $f_{L}$ and $A_{C P}$. The $V V$ branching fractions agree with theory predictions and the asymmetries are consistent with zero, as expected, while $f_{L} \sim 0.5$ except for $\omega \rho^{+} \sim 0.9$. The results ${ }^{[5]}$ are summarized in Table 1 .

Table 1: Branching fraction central value $(\mathcal{B})$ and upper limit (UL) in units of $10^{-6}$, significance S in standard deviations, longitudinal polarization $\left(f_{L}\right)$ and $C P$ asymmetry $\mathrm{A}_{C P}$ for the Vector-Vector ( $V V$ ), Vector-Scalar $(V S)$ and Vector-Tensor ( $V T$ ) decays of $B \rightarrow \omega K^{*}, \omega f_{0}$ and $\omega \rho$.

| Mode | Decay | $\mathrm{S}(\sigma)$ | $\mathcal{B}$ | UL | $f_{L}$ | $\mathrm{~A}_{C P}$ |
| :---: | :--- | :---: | :---: | :---: | :---: | :---: |
| $V V$ | $\omega K^{* 0}$ | 4.1 | $2.2 \pm 0.6 \pm 0.2$ | - | $0.72 \pm 0.14 \pm 0.02$ | $+0.45 \pm 0.25 \pm 0.02$ |
| $V V$ | $\omega K^{*+}$ | 2.5 | $2.4 \pm 1.0 \pm 0.2$ | 7.4 | $0.41 \pm 0.18 \pm 0.05$ | $+0.29 \pm 0.35 \pm 0.02$ |
| $V S$ | $\omega(K \pi)_{0}^{* 0}$ | 9.8 | $18.4 \pm 1.8 \pm 1.7$ | - | - | $-0.07 \pm 0.09 \pm 0.02$ |
| $V S$ | $\omega(K \pi)_{0}^{*+}$ | 9.2 | $27.5 \pm 3.0 \pm 2.6$ | - | - | $-0.10 \pm 0.09 \pm 0.02$ |
| $V T$ | $\omega K_{2}^{*}(1430)^{0}$ | 5.0 | $10.1 \pm 2.0 \pm 1.1$ | - | $0.45 \pm 0.12 \pm 0.02$ | $-0.37 \pm 0.17 \pm 0.02$ |
| $V T$ | $\omega K_{2}^{*}(1430)^{+}$ | 6.1 | $21.5 \pm 3.6 \pm 2.4$ | - | $0.56 \pm 0.10 \pm 0.04$ | $+0.14 \pm 0.15 \pm 0.02$ |
| $V V$ | $\omega \rho^{0}$ | 1.9 | $0.8 \pm 0.5 \pm 0.2$ | 1.6 | - | - |
| $V V$ | $\omega f_{0}$ | 4.5 | $1.0 \pm 0.3 \pm 0.1$ | 1.5 | - | - |
| $V V$ | $\omega \rho^{+}$ | 9.8 | $15.9 \pm 1.6 \pm 1.4$ | - | $0.90 \pm 0.05 \pm 0.03$ | $-0.20 \pm 0.09 \pm 0.02$ |

## 3 Decays involving three-body final states.

An interesting use of the decay to final states with three particles is the search by Belle for the exotic state $\mathrm{X}(1812)$ in the decay $B^{+} \rightarrow K^{+} X(1812), X(1812) \rightarrow \omega \phi$. This is similar to the observation by Belle of the $\mathrm{Y}(3940)$ resonance in $B^{+} \rightarrow K^{+} \omega \psi$. ${ }^{6}$. Belle observe $N_{K^{+} \omega \phi}=22.1_{-7.2}^{+8.3}$ events leading to a branching fraction for the Dalitz plot of $\mathcal{B}\left(B^{+} \rightarrow K^{+} \omega \phi\right)=$ $\left(1.15_{-0.38-0.13}^{+0.43+0.14}\right) \times 10^{-6}(2.8 \sigma)$ and an upper limit $<1.9 \times 10^{-6}$. Assuming the $\mathrm{X}(1812)$ masses and width from BES ${ }^{7}$, Belle searches for a near-threshold enhancement in the $M_{\pi^{+} \pi^{-} \pi^{0} K^{+} K^{-}}$ mass spectrum. No significant yield is seen and an upper limit of $3.2 \times 10^{-7}$ is placed on the product branching fraction $\mathcal{B}\left(B^{+} \rightarrow K^{+} X(1812), X(1812) \rightarrow \omega \phi\right)^{8}$.

BABAR has also looked at rare processes in Dalitz plots. Previous measurements have shown that almost $50 \%$ of the events in $B^{0} \rightarrow K^{+} K^{-} \pi^{+}$can be assigned to an ill-defined resonance, called $f_{X}(1500)$ by BABAR. If this is an even-spin resonance, the rate for $f_{X}(1500) \rightarrow K_{S}^{0} K_{S}^{0}$ would be expected to be half the rate for $f_{X}(1500) \rightarrow K^{+} K^{-}$. They see $15 \pm 15$ events in the whole Dalitz plot placing an upper limit on the total branching fraction of $\mathcal{B}\left(B^{+} \rightarrow K_{S}^{0} K_{S}^{0} \pi^{+}\right)<$ $5.1 \times 10^{-7}$. This makes the even-spin hypothesis unlikely but interpretation is difficult as the exact quantum numbers of the $f_{X}(1500)$ are unknown 9 .

Some MSSM models could enhance the branching fractions of SM-suppressed decays from the SM values of $\sim 10^{-16}$ to $\sim 10^{-6}$. BABAR has searched for $B^{-} \rightarrow K^{+} \pi^{-} \pi^{-}$and $B^{-} \rightarrow K^{-} K^{-} \pi^{+}$ and placed upper limits of $9.5 \times 10^{-7}$ and $1.6 \times 10^{-7}$, respectively, on the branching fractions $\frac{10}{}$.

The decay $B^{+} \rightarrow \pi^{+} \pi^{+} \pi^{-}$can in principle be used to extract the CKM angle $\gamma$ by measuring the interference between $\pi^{+} \pi^{-}$resonances and the $\chi_{c 0}$ resonance which has no $C P$ violating phase. It can also be helpful in understanding broad resonances and nonresonant backgrounds
that are present in $B^{0} \rightarrow \pi^{+} \pi^{-} \pi^{0}$ and so improve our measurement of the CKM angle $\alpha$. BABAR's results ${ }^{11}$ for $B^{+} \rightarrow \pi^{+} \pi^{+} \pi^{-}$are summarized in Table 2. No significant direct $C P$ asymmetry is measured and, although some resonances are significant, no evidence is found for $\chi_{c 0}$ and $\chi_{c 2}$, leading to branching fraction upper limits for $B^{+} \rightarrow \chi_{c 0} \pi^{+}<1.5 \times 10^{-5}$ and $B^{+} \rightarrow \chi_{c 2} \pi^{+}<2.0 \times 10^{-5}$, making the measurement of $\gamma$ in this mode unlikely at Belle or BABAR.

Table 2: Branching fraction $(\mathcal{B}), C P$ asymmetry $\mathrm{A}_{C P}$, and Fit Fraction for the decay $B^{+} \rightarrow \pi^{+} \pi^{+} \pi^{-}$with the resonance decaying to $\pi^{+} \pi^{-}$. The errors are statistical, systematic and model-dependent, respectively.

| Decay | Fit Fraction (\%) | $\mathcal{B}\left(\times 10^{-6}\right)$ | $\mathrm{A}_{C P}(\%)$ |
| :--- | :---: | :---: | :---: |
| $\pi^{+} \pi^{+} \pi^{-}$Total | - | $15.2 \pm 0.6 \pm 1.2_{-0.3}^{+0.4}$ | $3.2 \pm 4.4 \pm 3.1_{-2.0}^{+2.5}$ |
| $\pi^{+} \pi^{+} \pi^{-}$nonresonant | $34.9 \pm 4.2 \pm 2.9_{-3.4}^{+7.5}$ | $5.3 \pm 0.7 \pm 0.6_{-0.5}^{+1.1}$ | $-14 \pm 14 \pm 7_{-3}^{+17}$ |
| $\rho^{0}(770) \pi^{ \pm} ; \rho^{0} \rightarrow \pi^{+} \pi^{-}$ | $53.2 \pm 3.7 \pm 2.5_{-7.4}^{+1.5}$ | $8.1 \pm 0.7 \pm 1.2_{-1.4}^{+0.4}$ | $18 \pm 7 \pm 5_{-14}^{+2}$ |
| $\rho^{0}(1450) \pi^{ \pm} ; \rho^{0} \rightarrow \pi^{+} \pi^{-}$ | $9.1 \pm 2.3 \pm 2.4_{-4.5}^{+1.9}$ | $1.4 \pm 0.4 \pm 0.4_{-0.7}^{+0.3}$ | $-6 \pm 28 \pm 20_{-35}^{+12}$ |
| $f_{2}(1270) \pi^{ \pm} ; f_{2} \rightarrow \pi^{+} \pi^{-}$ | $5.9 \pm 1.6 \pm 0 . ._{-0.0}^{+2.0}$ | $0.9 \pm 0.2 \pm 0.1_{-0.1}^{+0.3}$ | $41 \pm 25 \pm 13_{-8}^{+12}$ |
| $f_{0}(1370) \pi^{ \pm} ; f_{0} \rightarrow \pi^{+} \pi^{-}$ | $18.0 \pm 3.3 \pm 2.6_{-3.5}^{+4.3}$ | $2.9 \pm 0.5 \pm 0.5_{-0.5}^{+0.7}(<4.0)$ | $72 \pm 15 \pm 14_{-8}^{+7}$ |
| $f_{0}(980) \pi^{ \pm} ; f_{0} \rightarrow \pi^{+} \pi^{-}$ | - | $<1.5$ | - |
| $\chi_{c 0} \pi^{ \pm} ; \chi_{c 0} \rightarrow \pi^{+} \pi^{-}$ | - | $<0.1$ | - |
| $\chi_{c 2} \pi^{ \pm} ; \chi_{c 2} \rightarrow \pi^{+} \pi^{-}$ | - | $<0.1$ | - |

## $4 C P$ Violation and the CKM angle $\alpha\left(\phi_{2}\right)$.

The precision of the measurement of the CKM angle $\alpha\left(\phi_{2}\right)$ continues to improve. In the absence of penguin loops in the decays, the angle $\alpha$ can be measured in the time-dependent decay of $B^{0} \rightarrow \rho \rho$ and $B^{0} \rightarrow \pi \pi$. However the penguin contribution, particularly in $\pi^{0} \pi^{0}$, is not small and so the measured $\alpha_{\text {eff }}$ differs from the true $\alpha$ by $\Delta \alpha=\alpha-\alpha_{e f f}$. $\Delta \alpha$ can be constrained by performing an Isospin analysis on the decays $B^{0} \rightarrow \rho^{0} \rho^{0}, B^{ \pm} \rightarrow \rho^{ \pm} \rho^{0}$ and $B^{0} \rightarrow \rho^{+} \rho^{-}$. Table 3 summarizes the measurements from $B A B A R{ }^{12}$, where the $C P$ parameters are quoted for the longitudinally polarized ( $C P$-even) component of the $V V$ decays. When combined, $\Delta \alpha$ is constrained to be between $-1.8^{\circ}$ and $6.7^{\circ}\left(68 \%\right.$ C.L.). The angle $\alpha$ is measured to be $\left(92.4_{-6.5}^{+6.0}\right)^{o}$ and can be compared to the recent result from Belle ${ }^{13}$ of $\alpha=(91.7 \pm 14.9)^{\circ}$. A similar analysis using $B \rightarrow \pi \pi$ decays produces a looser constraint $|\Delta \alpha|<43^{\circ}$, which results in an exclusion range for $\alpha$ between $23^{\circ}$ and $43^{\circ}$ at the $90 \%$ C.L. The result of combining these measurements using the CKMfitter programme ${ }^{14}$ with earlier measurements of $B \rightarrow \pi \rho$ are shown in Fig. [1.

Table 3: Branching fraction $(\mathcal{B})$, longitudinal polarization $\left(f_{L}\right)$, direct $C P$ asymmetry $\left(\mathrm{C}_{L}\right), C P$ asymmetry in the interference between mixing and decay $\left(\mathrm{S}_{L}\right)$ and $C P$ asymmetry $\mathrm{A}_{C P}$ for the decays $B^{0} \rightarrow \rho^{+} \rho^{-}, B^{0} \rightarrow \rho^{0} \rho^{0}$ and $B^{+} \rightarrow \rho^{+} \rho^{0}$ measured by BABAR.

|  | $B^{0} \rightarrow \rho^{+} \rho^{-}$ | $B^{0} \rightarrow \rho^{0} \rho^{0}$ | $B^{+} \rightarrow \rho^{+} \rho^{0}$ |
| :--- | :---: | :---: | :---: |
| $\mathcal{B}\left(\times 10^{-6}\right)$ | $25.5 \pm 2.1_{-3.9}^{+3.6}$ | $0.92 \pm 0.32 \pm 0.14$ | $23.7 \pm 1.4 \pm 1.4$ |
| $f_{L}$ | $0.992 \pm 0.024_{-0.013}^{+0.026}$ | $0.75 \pm 0.14 \pm 0.04$ | $0.950 \pm 0.015 \pm 0.006$ |
| $\mathrm{C}_{L}$ | $0.01 \pm 0.15 \pm 0.06$ | $0.2 \pm 0.8 \pm 0.3$ | - |
| $\mathrm{S}_{L}$ | $-0.17 \pm 0.20_{-0.06}^{+0.05}$ | $0.3 \pm 0.7 \pm 0.2$ | - |
| $\mathrm{A}_{C P}$ | - | - | $-0.054 \pm 0.055 \pm 0.01$ |



Figure 1: Constraints on $\alpha\left(\phi_{2}\right)$ from $B \rightarrow \pi \pi, B \rightarrow \rho \pi$ and $B \rightarrow \rho \rho B A B A R$ and Belle measurements compared to the prediction from the global CKM fit from CKMfitter. Similar results are available from the UTfit group 14 .

Belle has seen direct $C P$ in $B^{0} \rightarrow \pi^{+} \pi^{-}$but BABAR does not, reporting only that $C_{\pi^{+} \pi^{-}}=$ $-0.25 \pm 0.08 \pm 0.02$ with a significance of just $2.2 \sigma$. However, both experiments see significant direct $C P$ in $B^{0} \rightarrow K^{+} \pi^{-}$with $B A B A R$ reporting $A_{C P}=-0.107 \pm 0.016_{-0.004}^{+0.006}$ with $6.1 \sigma$ significance, to be compared to $-0.094 \pm 0.018 \pm 0.008$ from Belle. Both experiments also measure $A_{C P}$ for $B^{ \pm} \rightarrow K^{ \pm} \pi^{0}$ to be slightly positive but consistent with zero. $A_{C P}$ should be similar for both $K \pi$ modes but Belle reports a $4.4 \sigma$ difference and BABAR sees a similar discrepancy 15 .

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