

TESTING THE STANDARD MODEL WITH RADIATIVE PENGUIN DECAYS AT BABAR*

J. M. TUGGLE[†]

*University of Maryland
College Park, MD 20742
E-mail: jtuggle@slac.stanford.edu*

I discuss two recent results in $b \rightarrow s\gamma$ decays from BABAR. The first is a measurement of the branching fraction and photon energy spectrum in the B meson frame of the decay $B \rightarrow X_s\gamma$. The second result probes the photon polarization via time-dependent CP violation in neutral B decays to $K^{*0}\gamma$.

1. Introduction

Processes that are suppressed in the standard model (SM) are natural places to search for new physics. In $b \rightarrow s\gamma$ decays, deviations from the SM can appear in the overall branching fraction, the CP and isospin asymmetries, and the photon polarization. Furthermore, the photon energy spectrum contains information on the kinematics of the b quark inside the B meson. I discuss two recent BABAR results that measure all of these observables.

2. Standard Model Predictions

Because the $b \rightarrow s\gamma$ process only occurs through a loop diagram at leading order in the SM, it is sensitive to the possible existence of non-SM particles contributing to the loop. The most recent calculation of the SM branching fraction for $B \rightarrow X_s\gamma$ is at the next-to-next-to-leading order in QCD corrections¹: $\mathcal{B}(B \rightarrow X_s\gamma) = (3.15 \pm 0.23) \times 10^{-4}$ for photon energy $E_\gamma > 1.6$ GeV in the B rest frame. The CP asymmetry is expected to be 0.6% for $b \rightarrow s\gamma$ decays². No isospin asymmetry is expected. The shape of the photon energy spectrum can be used to extract the mass of the b quark and its fermi momentum inside the B meson using the kinetic scheme³.

*Report numbers: BaBar-PROC-08/003, SLAC-PUB-13213.

[†]Representing the BaBar Collaboration.

Due to the chirality of the weak interaction, the photon produced in the decay is expected to be left-handed for \bar{B}^0 decays and right-handed for B^0 decays. A low rate of “wrongly-polarized” photons allowed by nonzero quark masses is expected, permitting interference between decay and $B^0 - \bar{B}^0$ mixing in $b \rightarrow s\gamma$ final states that are accessible to both B flavors. The amount of this interference, denoted S , is expected to be about -0.02 in the SM^{4,5}, although possible hadronic corrections could allow values as large as -0.1 ^{6,7}. Any significant measured deviation from the predictions in this section would be an indication of new physics.

3. $B \rightarrow X_s\gamma$ Recoil Technique

Complementary to other inclusive $b \rightarrow s\gamma$ analyses, *BABAR* has developed an analysis technique that uses fully-reconstructed hadronic decays of the non-signal B . This analysis uses 210 fb^{-1} of data⁸. Hadronic decays are defined by 1114 exclusive decay channels of the form $B \rightarrow D^{(*)}Y^\pm$, where Y^\pm is a combination of up to nine pions and kaons⁹. Although this allows determination of the $b \rightarrow s\gamma$ photon in the B frame, the efficiency to reconstruct the recoil system is 0.3%.

The recoil system is checked for consistency with a B decay using two variables, $m_{\text{ES}} = \sqrt{s/4 - p_B^2}$ and $\Delta E = E_B - \sqrt{s}/2$, where, in the e^+e^- frame, s is the total energy-squared and E_B and p_B are the energy and momentum of the recoil system. We require $|\Delta E| < 60 \text{ MeV}$ and $m_{\text{ES}} > 5.2 \text{ GeV}/c^2$. The remaining particles in the event must contain an isolated photon of energy greater than 1.3 GeV in the B frame. The invariant mass of the photon combined with another photon in the event must not be consistent with a π^0 or η meson. Background from $e^+e^- \rightarrow q\bar{q}$ decays ($q = \{u, d, s, c\}$) is suppressed using a Fisher discriminant of event-shape variables. The surviving events are divided into 100-MeV-wide bins in photon energy from 1.3–2.7 GeV. The peaking yield in each bin is determined by a fit to m_{ES} . Because the region below 1.9 GeV is dominated by $B\bar{B}$ backgrounds, the Monte Carlo (MC) sample used for this background is scaled to the data yield in this range, then subtracted from the entire range up to 2.7 GeV. A total of 119 ± 22 signal $B \rightarrow X_s\gamma$ events are observed over a background of $145 \pm 9 B\bar{B}$ events.

The systematic uncertainties arise mainly from the $B\bar{B}$ background shape. We apply E_γ -dependent corrections to account for shape differences resulting from vetoing π^0 and η mesons. The remaining backgrounds have a roughly linear slope in E_γ , which is varied by $\pm 30\%$. Relatively

smaller uncertainties due to the m_{ES} fits and the assumed E_γ shape are also included.

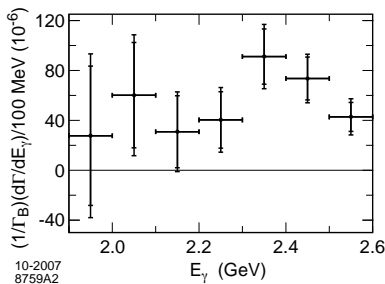


Figure 1. The partial branching fractions $(1/\Gamma_B)(d\Gamma/dE_\gamma)$ with statistical (inner) and total (outer) uncertainties.

The photon spectrum is shown in Fig. 1. After lowering the branching fraction by 4% to account for $b \rightarrow d\gamma$ decays, we find

$$\mathcal{B}(B \rightarrow X_s \gamma) = 3.66 \pm 0.85_{\text{stat}} \pm 0.60_{\text{syst}} \times 10^{-4}.$$

The CP and isospin asymmetries (A_{CP} and Δ_{0-} respectively) in the signal region, $E_\gamma > 2.2$ GeV, are not corrected for $b \rightarrow d\gamma$ contamination, and are found to be

$$A_{CP} = \frac{\mathcal{B}(B \rightarrow X_{s,d}\gamma) - \mathcal{B}(\bar{B} \rightarrow X_{s,d}\gamma)}{\mathcal{B}(B \rightarrow X_{s,d}\gamma) + \mathcal{B}(\bar{B} \rightarrow X_{s,d}\gamma)} \frac{1}{1 - 2\omega} = 0.10 \pm 0.18_{\text{stat}} \pm 0.05_{\text{syst}},$$

where ω is the flavor mistag probability, and

$$\Delta_{0-} = \frac{\Gamma(\bar{B}^0 \rightarrow X_{s,d}\gamma) - \Gamma(B^- \rightarrow X_{s,d}\gamma)}{\Gamma(\bar{B}^0 \rightarrow X_{s,d}\gamma) + \Gamma(B^- \rightarrow X_{s,d}\gamma)} = -0.06 \pm 0.15_{\text{stat}} \pm 0.07_{\text{syst}}.$$

Using the kinetic scheme, we determine the heavy quark parameters $m_b = 4.46_{-0.23}^{+0.21}$ GeV/ c^2 and $\mu_\pi^2 = 0.64_{-0.38}^{+0.39}$ GeV 2 .

4. Photon Polarization in $B^0 \rightarrow K^{*0}\gamma$

The $b \rightarrow s\gamma$ photon polarization can be inferred through time-dependent CP violation in decays to final states accessible by both B^0 and \bar{B}^0 . This is a preliminary result for time-dependent CP violation in $B^0 \rightarrow K^{*0}(\rightarrow K_s^0\pi^0)\gamma$ decays using 391 fb^{-1} of data¹⁰.

The proper-time difference between the decay of the signal B and the tagging B , $\Delta t \equiv t_{\text{sig}} - t_{\text{tag}}$, follows the probability distribution

$$\mathcal{P}_{\pm}(\Delta t) = \frac{e^{-|\Delta t|/\tau_B}}{4\tau_B} \times [1 \pm S \sin(\Delta m_d \Delta t) \mp C \cos(\Delta m_d \Delta t)],$$

where the upper and lower signs correspond to the tag-side B having flavor B^0 and \bar{B}^0 respectively, τ_B is the B^0 lifetime, Δm_d is the $B^0 - \bar{B}^0$ mixing frequency, and $C = -A_{CP}$. Any major deviation from $S = 0$ would be a signal of new physics.

To determine Δt , the decay vertex of the signal side is calculated by constraining the K_s^0 momentum to intersect with the beam spot, within errors. The decay vertex of the tagged B is reconstructed¹¹ based on the tracks in the event that are not used for the signal B . We constrain the sum of the decay times of both B mesons in the event to be twice the B lifetime with an uncertainty of $\sqrt{2}\tau_{B^0}$.

We select events with a photon of energy $1.5 < E_\gamma < 3.5$ GeV in the e^+e^- frame, using the same photon quality requirements as in Sec. 3. We impose requirements on K_s^0 and π^0 masses, the K_s^0 flight length divided by its uncertainty, and the event shape. We require $|\Delta t| < 20$ ps and the Δt uncertainty $\sigma_{\Delta t} < 2.5$ ps. Four variables are combined into a maximum likelihood fit to separate signal from continuum and $B\bar{B}$ backgrounds: m_{ES} , ΔE , K^{*0} mass, and L_2/L_0^a . The S and C parameters are extracted based on the signal Δt distribution.

To evaluate the systematic uncertainty associated with assuming $S = C = 0$ for the $B\bar{B}$ background, we vary its value of S by ± 0.4 and C by ± 0.3 . We also include uncertainties in the fit model due to differences between MC and data, and MC statistics. Finally, data-MC differences in the Δt resolution function are evaluated using $B^0 \rightarrow J/\psi K_s^0$ events.

The results of the fit are shown in Fig. 2. We find 316 ± 22 signal events, and measure

$$S = -0.08 \pm 0.31_{\text{stat}} \pm 0.05_{\text{syst}} \quad \text{and} \quad C = -0.15 \pm 0.17_{\text{stat}} \pm 0.03_{\text{syst}}.$$

5. Summary

The $B \rightarrow X_s \gamma$ recoil technique has much less background than other inclusive analyses, allows the determination of the photon energy spectrum

^a $L_i = \sum_j |p_j^*| \cos \theta_j^*$, where p_j^* is the momentum in the e^+e^- frame of each particle j not used to reconstruct the signal B candidate, and θ_j^* is the angle between p_j^* and the thrust axis of the signal B candidate in the same frame.

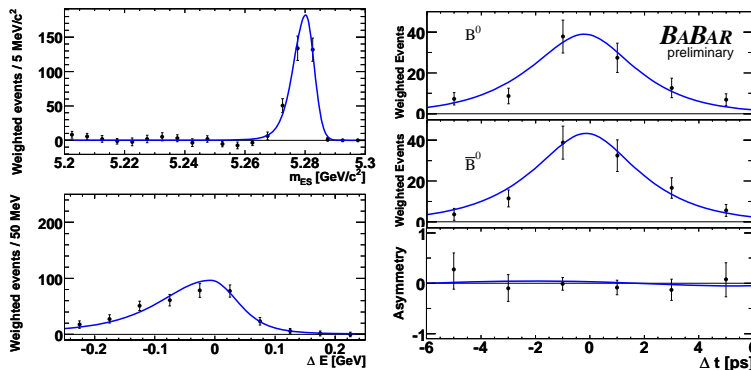


Figure 2. On the left are the signal distributions for m_{ES} (top) and ΔE (bottom), using the event weighting technique described in the preprint. On the right are the signal distributions for Δt , with B_{tag} tagged as B^0 (top) or \bar{B}^0 (center), and the asymmetry (bottom).

in the B frame, and provides information on the b quark kinematics inside the B meson. Time-dependent CP violation in $b \rightarrow s\gamma$ decays is the only known tool to measure the photon polarization at the B -Factories, offering another way to look for new physics. These measurements test the standard model in all aspects of the $b \rightarrow s\gamma$ interaction: the rate, CP asymmetry, isospin asymmetry, photon spectrum, and photon polarization. The measured values are consistent with the standard model.

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