

Radiative Penguin decays at the B Factories

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CANADA*

Abstract. Recent results from the B -Factories on radiative decays such as $b \rightarrow s(d)\gamma$, $b \rightarrow s\ell\ell$ and leptonic decay $B^0 \rightarrow \tau^+\tau^-$ are reviewed.

Keywords: Rare decays, Leptonic decays

PACS: 12.15.Hh, 13.20.-v, 13.20.He, 13.30.Ce, 13.35.Dx

INTRODUCTION

The flavor changing neutral current (FCNC) rare B decays such as $b \rightarrow s(d)\gamma$ and $b \rightarrow s\ell\ell$ play a crucial role in testing the Standard Model (SM) as well as probing physics beyond the SM. Such processes proceed via one loop (penguin or box) diagrams. The effective hamiltonian for $b \rightarrow s$ transition is given by $H_{eff} = -4\frac{G_F}{\sqrt{2}}V_{ts}^*V_{tb}\sum_i C_i(\mu)O_i(\mu)$ [1], where G_F and $V_{ts,tb}$ are the Fermi constant and the Cabibbo-Kobayashi-Maskawa (CKM) matrix elements, respectively. The FCNC operators, O_i , can be considered as the vertices of the effective theory and the Wilson coefficients, C_i , play the role of coupling constants. Contributions in the loops of these decays from the Higgs boson or supersymmetric particles can lead to an enhancement in their rates which range between 10^{-4} to 10^{-6} in the SM. A comparison between the measurement and the SM prediction provides a stringent constraint on such new particles. At lowest order $b \rightarrow s(d)\gamma$ is described by the coefficient C_7 (photon penguin operator) and is computed up to next-to-leading order (NLO). The decay $b \rightarrow s\ell\ell$ is defined by a Z -penguin and a W -box diagram in addition to the radiative penguin. Its amplitude depends on the coefficients C_7, C_9 and C_{10} (vector and axial vector operators) which are computed next-to-next-to-leading order (NNLO).

The Babar and Belle detectors are located at the PEP-II and KEKB asymmetric energy e^+e^- colliders, respectively. $B\bar{B}$ pairs are produced coherently from the $\Upsilon(4S)$ decays. Up to now, approximately 400 fb^{-1} and 600 fb^{-1} of data have been recorded by the Babar and Belle experiments, respectively.

$$b \rightarrow s\gamma$$

The signature of a $b \rightarrow s\gamma$ decay is an energetic photon of $E_\gamma^* = (2 - 2.7) \text{ GeV}$. Due to the confinement of the b -quark inside the B meson, the energy spectrum is smeared instead of producing a monochromatic energy of $E \approx m_b/2$. A fit to E_γ^* , or to the first and second moments of the photon energy spectrum, can be used to extract information on the Heavy Quark Expansion (HQE) parameters, m_b (b -quark mass) and μ_π^2 (b -quark

kinetic energy inside the B meson). The E_γ spectrum is also input to semileptonic B decay analyses to extract $|V_{cb}|$ and $|V_{ub}|$. The inclusive decay $B \rightarrow X_s \gamma$ is studied with two methods. In the fully inclusive method, the photon energy spectrum is measured without reconstructing the X_s system. In the semi exclusive method, the X_s system is reconstructed from a kaon (K or K_s) and up to four pions (one or two π^0 allowed). The reconstructed X_s system is then combined with a photon to form a B meson.

Babar and Belle have measured the branching fraction and the moments from the photon energy spectra using 82 fb^{-1} and 140 fb^{-1} , respectively [2, 3]. The branching fraction results are summarized in Fig.1(left) which also includes the earlier measurement by CLEO at 9.1 fb^{-1} [4]. The world average of $B(b \rightarrow s\gamma) = (3.55 \pm 0.26) \times 10^{-4}$ is published by the Heavy Flavor Averaging Group (HFAG) [5]. It is in good agreement with the SM prediction of $B(b \rightarrow s\gamma) = (3.61^{+0.37}_{-0.49}) \times 10^{-4}$ [6]. A fit to all available $b \rightarrow s\gamma$ and $B \rightarrow X_c \ell v$ data results in the HQE parameters of $m_b = 4.590 \pm 0.039 \text{ GeV}$ and $\mu_\pi^2 = 0.401 \pm 0.040 \text{ GeV}^2$ (in the kinetic scheme) [7]. As shown in Fig.1(right), the moments obtained from the E_γ spectra are consistent among the different measurements as well as with the theory predictions that rely on data from $B \rightarrow X_c \ell v$ moments. The $b \rightarrow s\gamma$ decay mode has also been measured in many exclusive final

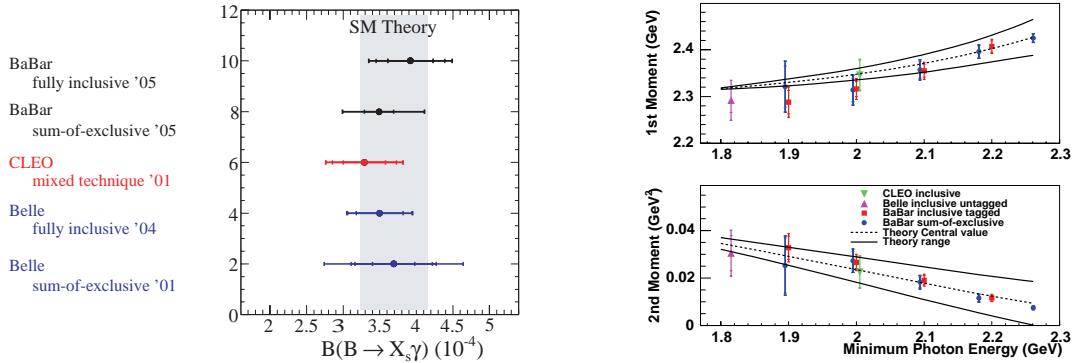


FIGURE 1. (left) Results on branching fraction of $b \rightarrow s\gamma$. (right) First and second moments from photon spectra of $b \rightarrow s\gamma$.

states. For example, the most recent Babar result on $B \rightarrow K\eta\gamma$ confirms the earlier Belle measurement. Both charged and neutral decays are observed at better than 5σ significance. The combined average limits are $B(B^+ \rightarrow K^+\eta\gamma) = (9.3 \pm 1.1) \times 10^{-6}$, and $B(B^0 \rightarrow K^0\eta\gamma) = (10.3^{+2.3}_{-2.1}) \times 10^{-6}$ [5]. Babar has also reported an upper limit $B(B^{+(0)} \rightarrow K^{+(0)}\eta'\gamma) < 4.2(6.6) \times 10^{-6}$ at the 90% C.L. [8].

$$b \rightarrow d\gamma$$

The decay $b \rightarrow d\gamma$ has similar sensitivity to the new physics (NP) as $b \rightarrow s\gamma$, but its branching fraction is suppressed by the ratio of $|V_{td}/V_{ts}|^2$. The exclusive decay $B \rightarrow (\rho, \omega)\gamma$ were first observed by Belle using $386 \times 10^6 B\bar{B}$. Assuming isospin symmetry, a simultaneous fit to three decay modes ($B \rightarrow (\rho^+, \rho^0, \omega)\gamma$) results in $B(B \rightarrow (\rho, \omega)\gamma) = (1.32^{+0.34}_{-0.31}(\text{stat})^{+0.10}_{-0.09}(\text{sys})) \times 10^{-6}$ [9]. Based on $211 \times 10^6 B\bar{B}$ events, Babar observed no evidence for these processes and placed combined upper limits of $B(B \rightarrow (\rho, \omega)\gamma) <$

1.2×10^{-6} at the 90% C.L. [10]. The SM predictions for these branching fractions are in the range of $(0.8 - 1.8) \times 10^{-6}$ [11]. The UTfit collaboration has performed a fit to the $b \rightarrow d\gamma$ data to extract the $|V_{td}/V_{ts}|$ ratio. The fit results in the value of $|V_{td}/V_{ts}| = 0.16 \pm 0.02$ [12]. The most recent value of $|V_{td}/V_{ts}| = 0.208^{+0.008}_{-0.007}$ is reported by CDF Collaboration [13] from the $\Delta m_s/\Delta m_d$ measurement.

$$b \rightarrow s\ell\ell$$

In addition to the decay rate of $b \rightarrow s\ell\ell$, other observables such as the forward-backward asymmetry (A_{FB}) can be used to identify the sign of the Wilson coefficient C_7 and the magnitudes of C_9 and C_{10} . The CP asymmetry, A_{CP} , and the ratio of the muon to electron rate, $R_{K^{(*)}} = B(B \rightarrow K^{(*)}\mu\mu)/B(B \rightarrow K^{(*)}ee)$, are sensitive to NP.

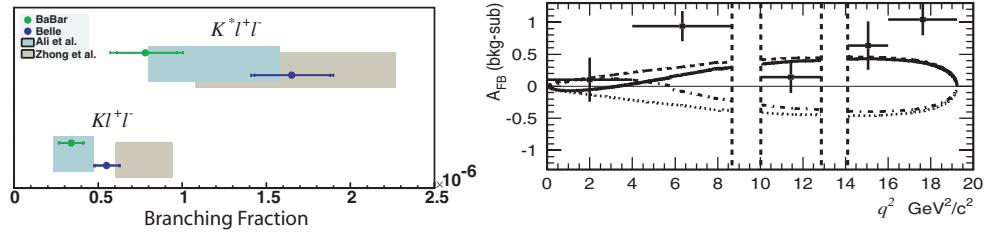


FIGURE 2. (left) Comparison of the $B \rightarrow K^{(*)}ll$ branching fraction measurements and the two SM predictions. (right) A_{FB} in $B \rightarrow K^{*}ll$ measured by Belle Collaboration. The solid line is for negative A_7 , dashed line is for positive A_7 , dot-dashed is for positive A_{10} and dotted is for positive A_7 and A_{10} .

Semi exclusive decay $b \rightarrow s\ell\ell$ is reconstructed from two leptons ($\ell = e, \mu$), a kaon and up to four pions (one or two π^0 allowed). Babar and Belle have measured the branching fraction summing up 10 and 18 hadronic states [14, 15], respectively. In order to avoid the resonance regions, the branching fraction has been reported in two q^2 (dilepton mass squared) regions: $q^2 = [1 - 6] \text{ GeV}^2$ and $q^2 > 14 \text{ GeV}^2$. For $q^2 = [1 - 6] \text{ GeV}^2$, branching fraction measured by Babar [Belle] $B(B \rightarrow s\ell\ell) = (1.8 \pm 0.9)[1.5 \pm 0.6] \times 10^{-6}$ is consistent with the SM prediction of $(1.6 \pm 0.2) \times 10^{-6}$ [16]. For $q^2 > 16 \text{ GeV}^2$, the branching fraction results from Babar and Belle are $(5.6 \pm 2.0) \times 10^{-6}$ and $(4.1 \pm 1.1) \times 10^{-6}$, respectively. These result agree well with the SM expectation of $(4.4 \pm 0.7) \times 10^{-6}$ [16].

The exclusive decay $B \rightarrow K^{(*)}ll$ has been studied extensively by Babar [17] and Belle [18] using 208 fb^{-1} and 253 fb^{-1} of data. The measured branching fractions are summarized on Fig.2(left). The CP asymmetries, $A_{CP}^K = (-0.07 \pm 0.22)$ and $A_{CP}^{K^*} = (0.03 \pm 0.23)$ measured by Babar are very small, consistent with the SM expectation. For the ratio, Babar has reported $R_K = (1.06 \pm 0.49)$ and $R_{K^*} = (0.91 \pm 0.45)$. For these ratios, Belle obtains $R_K = (1.38 \pm 0.41)$ and $R_{K^*} = (0.98 \pm 0.32)$. The results of R_K and R_{K^*} agree well with the SM expectation of unity [19] and 0.75 [20], respectively.

The $A_{FB}(q^2)$ measurement from Belle [21] is shown in Fig.2(right). The plot shows NP scenarios as well as the SM one. The Belle results exclude the NP scenarios where the sign of C_9C_{10} is flipped with respect to the SM. The models with C_7 and $-C_7$ agree well with data, hence no conclusion can be made on the sign of C_7 . Babar has measured A_{FB} in two regions. At large q^2 , the C_9C_{10} is found to be consistent with the SM. At

low q^2 , a lower limit $A_{FB} > 0.19$ at 95% C.L. has been set. The fraction of longitudinal polarization, F_L , in the K^* mode has the same importance as A_{FB} in measuring the sign of C_7 . Babar has measured F_L , however data does not favor either of the C_7 sign.

LEPTONIC DECAYS $B^0 \rightarrow \tau^+ \tau^-$

The recent result from Babar on $B^0 \rightarrow \tau^+ \tau^-$ is also presented. This decay proceeds via W -box and penguin diagrams. Due to the large τ mass, the branching fraction is less suppressed with respect to electron and muon modes ($e : \mu : \tau = 10^{-7} : 10^{-3} : 1$). However, the multiple neutrinos in the final states make it harder to measure. The SM prediction is $B(B^0 \rightarrow \tau^+ \tau^-) = 1.2 \times 10^{-7} [f_B/200 \text{ MeV}]^2 [|V_{td}|/0.007]^2$ where f_B is the B meson decay constant [1]. Using 210 fb^{-1} , no evidence was found for this decay. At the 90% C.L., the upper limit of $B^0 \rightarrow \tau^+ \tau^- < 4.1 \times 10^{-3}$ has been obtained [22].

CONCLUSION

Such rare decays, pioneered by CLEO, are brought to a more mature state by the B -Factories. The results shown here are based on one-fourth to one-half of the total integrated luminosity. Both theoretical and experimental updates will further constraint physics beyond the SM.

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