

# Hadronic $B$ decays at *BABAR*

**Jong Yi**<sup>\*†</sup>

*Iowa State University, USA*

*E-mail:* jong@slac.stanford.edu

We present preliminary results on hadronic decays of  $B$  mesons, based on data recorded at the  $\Upsilon(4S)$  resonance with the *BABAR* detector at the PEP-II  $B$ -factory at SLAC. We measure branching fractions in  $CP$ -related analyses of  $B^- \rightarrow D(\rightarrow \pi^+\pi^-\pi^0)K^-$ ,  $B^0 \rightarrow \bar{D}^0(D^0)DK^+\pi^-$  and  $B^0 \rightarrow D_s^+ a_{0(2)}^-$  and in non- $CP$ -related analyses of  $B^+ \rightarrow D^{(*)+}K^0$ ,  $B^- \rightarrow D_s^{(*)-}\phi$  and  $B \rightarrow J/\psi\bar{D}$ . Because the results presented in this paper are preliminary, they are based on different amount of data samples.

*International Europhysics Conference on High Energy Physics*  
*July 21st - 27th 2005*  
*Lisboa, Portugal*

---

\*Speaker.

†on behalf of the *BABAR* Collaboration

## 1. The BABAR detector

The BABAR detector [1] at the PEP-II asymmetric-energy  $B$ -factory [2] at SLAC consists of a silicon vertex tracker (SVT) for precise decay vertex determination, a 40-layer drift chamber (DCH) for momentum and track angles measurement, a detector of internally reflected Cherenkov radiation (DIRC) for charged hadron identifications, and a CsI(Tl) electromagnetic calorimeter (EMC) for photon reconstruction and electron identification. A superconducting solenoid provides a magnetic field of 1.5T, and the iron of the flux return is instrumented with resistive plate chambers (IFR) to provide muon identification and neutral hadron reconstruction.

## 2. $CP$ related analyses

Decay modes of the type  $B \rightarrow DK$  is utilized in a theoretically clean method of measuring the angle  $\gamma = \arg(-V_{ud}V_{ub}^*/V_{cd}V_{cb}^*)$  in the CKM unitary triangle. Gronau and Wyler, for example, have proposed to constrain the relative phase,  $\gamma$ , of  $b \rightarrow u\bar{c}s$  to  $b \rightarrow c\bar{u}s$  processes.[3] However, extraction of  $\gamma$  in this method suffers from an eight-fold ambiguity due to *a priori* unknown strong phases. In addition, the  $b \rightarrow u\bar{c}s$  amplitude is suppressed by CKM element factor  $|V_{ub}V_{cs}/V_{cb}V_{us}| \approx 0.4$  and color suppression factor of  $0.2 - 0.5$ . We present results from two recent proposals for measuring  $\gamma$  in decays modes of  $B^- \rightarrow D(\rightarrow \pi^+\pi^-\pi^0)K^-$  and  $B^0 \rightarrow \bar{D}^0(D^0)DK^+\pi^-$ .

The value of  $\sin(2\beta + \gamma)$  can be extracted from the measurement of the time dependent  $CP$  asymmetry in  $B \rightarrow D^-X_{light}^+$  decays with, for instance,  $X_{light}^+ \rightarrow a_0^+, a_2^+$ . In this case the asymmetry is given by :  $\mathcal{A}_{CP}(\Delta t) = r \times \sin(2\beta + \gamma) \times \sin(\Delta m_d \Delta t)$  where  $r = \mathcal{B}(B^0 \rightarrow D^+X_{light}^-)/\mathcal{B}(B^0 \rightarrow D^-X_{light}^+)$ . The decay  $B^0 \rightarrow D^+X_{light}^-$  is doubly Cabibbo suppressed and difficult to measure directly. However, using  $SU(3)$  flavor symmetry, it is possible to infer the value of  $\mathcal{B}(B^0 \rightarrow D^+X_{light}^-)$  from the value of  $\mathcal{B}(B \rightarrow D_s^+X_{light})$ , the latter being less suppressed. We present the results of one analysis utilizing such a method, where  $X_{light}^+ = a_{0(2)}^+$ .

### 2.1 Study of $B^- \rightarrow D(\rightarrow \pi^+\pi^-\pi^0)K^-$ decays

The decays  $B \rightarrow D^{(*)0}K^{(*)}$  can be used to measure the angle  $\gamma$  taking advantage of the interference between  $b \rightarrow u\bar{c}s$  and  $b \rightarrow c\bar{u}s$  decay amplitudes. Different approaches have been developed, and among which,  $\gamma$  measurements involving  $D$  decays to multi-body, using a Dalitz plot analysis technique as described in reference [4]. In this analysis, we study the decay mode  $B^- \rightarrow DK^-$  with the  $D$ -decay:  $D \rightarrow \pi^+\pi^-\pi^0$  which is Cabibbo suppressed. This yields a much smaller event sample compared to Cabibbo allowed decay but its interfering  $D^0$  and  $\bar{D}^0$  amplitudes have similar magnitudes. Due to these interferences, the production rate might be different from the product  $\mathcal{B}_{prod} \equiv \mathcal{B}(B^- \rightarrow D^0K^-) \times \mathcal{B}(D^0 \rightarrow \pi^+\pi^-\pi^0) = (4.1 \pm 1.6) \times 10^{-6}$ [5]. From a sample of 229 million of BB pairs, we found  $133 \pm 23$  signal events which correspond to a branching ratio of  $\mathcal{B}(B^- \rightarrow D_{\pi^+\pi^-\pi^0}K^-) = (5.5 \pm 1.0 \pm 0.7) \times 10^{-6}$ . We determine the raw asymmetry and do not find any significant deviation from zero:  $\mathcal{A}_{CP}^{raw} = 0.02 \pm 0.16 \pm 0.03$ . The  $\gamma$  extraction needs a full Dalitz analysis of the  $D$ -decay.

## 2.2 Study of three-body $B^0 \rightarrow \bar{D}^0(D^0)DK^+\pi^-$ decays

In the three-body decays of  $B \rightarrow DK\pi$ , the CKM suppressed  $b \rightarrow u\bar{c}s$  processes contain color-allowed diagrams, which could result in larger rates and more significant  $CP$  violation effect. The 8-fold ambiguity in strong phase could be reduced 2-fold by using the Dalitz plot. To assess the feasibility of this method, which was proposed by R. Aleksan, *et al*[6], we study the decays  $B^0 \rightarrow \bar{D}^0(D^0)DK^+\pi^-$ , where  $D^0 \rightarrow K^-\pi^+, K^-\pi^+\pi^0, K^-\pi^+\pi^-\pi^+$  using 205  $fb^{-1}$  of on-resonance data. In this measurement, the contribution from  $B \rightarrow D^*K$  is excluded. While the branching fraction for the CKM favored mode is determined to be  $\mathcal{B}(B^0 \rightarrow \bar{D}^0K^+\pi^-) = (8.6 \pm 1.5 \pm 1.0) \times 10^{-5}$ , no significant signal events are found for the CKM suppressed mode. We, therefore, set the upper limit for the suppressed mode with 90% C.L.:  $\mathcal{B}(B^0 \rightarrow D^0K^+\pi^-) < 1.9 \times 10^{-5}$ . As we do not observe a significant signal for the  $b \rightarrow u$  mode, measuring the CKM angle  $\gamma$  with this mode is determined to be very difficult.

## 2.3 Search for $B \rightarrow D_s^+ X_{light}$ with $X_{light} \equiv a_0^-, a_2^-$

In extracting  $\sin(2\beta + \gamma)$  from time dependent  $CP$  asymmetry in  $B \rightarrow D^- X_{light}^+$  decays, we analyze the case where  $X_{light}^+ = a_{0(2)}^+$ . In this case,  $r = \mathcal{B}(B^0 \rightarrow D^+ X_{light}^-) / \mathcal{B}(B^0 \rightarrow D^- X_{light}^+)$  may be quite large. This is due to the small coupling constant of the  $W$  to the  $a_0$  scalar meson ( $a_2$  meson of spin 2) which decreases the production rate of the Cabibbo allowed decay  $B^0 \rightarrow D^- a_0^+(a_2^+)$ . The factorization hypothesis predicts a similar rate for Cabibbo allowed and Cabibbo suppressed decays, which results in  $r \sim 1$ [7]. These decays are not yet in the reach of the experiment (branching ratios around  $10^{-6}$ ); nevertheless, the theoretical predictions can be tested with the measurement of the branching ratio of the decay  $B^0 \rightarrow D_s^+ a_0^-(a_2^-)$  expected at larger values:  $\mathcal{B}(B^0 \rightarrow D_s^+ a_0^-(a_2^-)) \approx 7.5(1.5) \times 10^{-5}$  [7]. From a sample of 230 million of  $B\bar{B}$  pairs, we measure these two branching ratios. The  $a_0^-(a_2^-)$  is reconstructed in  $a_{0,2}^- \rightarrow \eta(\rightarrow \gamma\gamma)\pi^+$  which has a branching ratio of the order of 100% (only 15% for the  $a_2^-$  which lowers the experimental sensitivity). We do not find any significant signal and quote the upper limits of  $\mathcal{B}(B^0 \rightarrow D_s^+ a_{0(2)}^-) < 4.0(25) \times 10^{-5}$  which shows a discrepancy of at least a factor two with the theoretical predictions.

## 3. Non $CP$ -related analyses

### 3.1 Search for rare quark-annihilation decays $B^- \rightarrow D_s^{(*)-} \phi$

In the Standard Model (SM), the decay  $B^- \rightarrow D_s^{(*)-} \phi$  occurs through annihilation of the two quarks in the  $B$  meson into a virtual  $W$ . Such a process is highly suppressed, and calculations of the  $B^- \rightarrow D_s^- \phi$  branching fraction give predictions of  $3 \times 10^{-7}$  using a perturbative QCD approach [8],  $1.9 \times 10^{-6}$  using factorization [9], and  $7 \times 10^{-7}$  using QCD-improved factorization [9]. Since the current experimental limits are about three orders of magnitude higher than the SM expectations, searches for  $B^- \rightarrow D_s^{(*)-} \phi$  could be sensitive to contributions from new physics. Enhancement of such kind has been calculated to be as high as  $8 \times 10^{-6}$  in a two-Higgs doublet model and  $3 \times 10^{-4}$  in minimal supersymmetric model with  $R$ -parity violation (RPV-MSSM), depending on the details of the parameters of the new physics.[9] Based on a sample of 234 million  $B\bar{B}$  pairs, we reconstruct  $B^- \rightarrow D_s^{(*)-} \phi$ , where  $D_s^{*-} \rightarrow D_s^- \gamma, D_s^- \rightarrow \phi \pi^-, K_s^0 K^-, K^{*0} K^-$ . We observe no

significant signal and set upper limits with 90% C.L. of  $\mathcal{B}(B^- \rightarrow D_s^- \phi) < 1.8 \times 10^{-6}$  and  $\mathcal{B}(B^- \rightarrow D_s^{*-} \phi) < 1.1 \times 10^{-5}$ . These limits are lower than the prediction based on RPV-MSSM.

### 3.2 Search for the rare decays $B^+ \rightarrow D^{(*)+}K^0$

This decay is expected to occur via a pure annihilation diagram. Such processes provide interesting insights into the internal dynamics of B mesons. This kind of diagram cannot be calculated in QCD factorization since both the quarks play a role. These amplitudes are expected to be suppressed, with respect to the amplitudes of spectator quark trees, by a factor  $f_B/m_B \approx 0.04$  and have never been observed. Some studies indicate, though, that processes with a spectator quark can contribute to annihilation-mediated decays by rescattering.[10] The branching ratio is expected to be either at the level of the current sensitivity ( $10^{-5}$ ) if large rescattering occurs, or three orders of magnitude below if not.[10] We reconstruct the two decay modes  $B^+ \rightarrow D^{*+}K_s^0$  and  $B^+ \rightarrow D^+K_s^0$  with a sample of 226 million of  $B\bar{B}$  pairs. We do not see any significant excess of signal and therefore set the upper limits with 90% C.L.:  $\mathcal{B}(B^+ \rightarrow D^+K_s^0) < 0.5 \times 10^{-5}$  and  $\mathcal{B}(B^+ \rightarrow D^{*+}K_s^0) < 0.9 \times 10^{-5}$ .

### 3.3 Search for $B \rightarrow J/\psi D$ Decays

The spectra of the momentum of inclusive  $J/\psi D$  mesons in the  $Y(4S)$  rest frame observed by CLEO and by BABAR, compared with calculations using non-relativistic QCD (NRQCD), show an excess at low momentum, corresponding to a branching fraction of approximately  $6 \times 10^{-4}$ . Many hypothesis have been proposed to explain this result but no experimental evidence has been found to support them. The presence of bucc components (intrinsic charm) in the  $B$ -meson wave function has also been suggested to enhance the branching ratio of decays such as  $b \rightarrow J/\psi \bar{D}(\pi)$  to the order of  $10^{-4}$  while perturbative QCD predicts a branching ratio for  $B \rightarrow J/\psi \bar{D}$  to the order of  $10^{-8}$  to  $10^{-9}$ . We test the decay channels  $B \rightarrow J/\psi D$  with a sample of 124 million of  $B\bar{B}$  pairs. We do not find any evidence of signal and obtain upper limits of  $1.3 \times 10^{-5}$  for  $B^0 \rightarrow J/\psi \bar{D}^0$  and  $1.2 \times 10^{-4}$  for  $B^+ \rightarrow J/\psi D^+$  at 90% C.L. Therefore, we conclude that intrinsic charm is not supported as the explanation of low momentum  $J/\psi$  excess in  $B$  decays.

## References

- [1] BABAR Collaboration, B. Aubert *et al.*, *Nucl. Instr. and Methods A* **479**, 1-116 (2002).
- [2] *PEP-II - An Asymmetric B Factory*, Conceptual Design Report, SLAC-R-418, LBL-5379 (1993).
- [3] M. Gronau, D. Wyler, *Phys. Lett. B* **265**, 172 (1991).
- [4] A. Giri, Y. Grossman, A. Soffer and J. Zupan, *Phys. Rev. D* **68**, 054018 (2003).
- [5] S. Eidelman *et al.* [Particle Data Group], *Phys. Lett. B* **592**, 1 (2004).
- [6] R. Aleksan, T.C. Petersen, A. Soffer, *Phys. Rev. D* **67**, 096002 (2003).
- [7] M. Diehl, G. Hiller, hep-ph/0105194.
- [8] C.D. Lu, *Eur. Phys. J. C* **24**, 121 (2002).
- [9] R. Mohanta, *Phys. Lett. B* **540**, 241 (2002).
- [10] M. Blok, M. Gronau and J.L. Rosner, *Phys. Rev. Lett.* **78**, 3999 (1997).