## INITIAL OPERATING EXPERIENCE AND PERFORMANCE OF THE SLAC ACCELERATOR"

R. B. Neal

Stanford Linear Accelerator Center, Stanford University, Stanford, California

#### Introduction

#### General Plan

At the date of this report somewhat more than four years have passed since the initial ground-breaking occurred at the SLAC site (July 9, 1962). From the viewpoint of construction, the overall results achieved are probably best observed in an air view of the site (Fig. 1) showing the principal buildings, shops, laboratories, the accelerator housings, and the research area facilities. The entire site comprises 480 acres located about two miles west of the main Stanford campus.

Buildings and structures having a total floor area of approximately 870,000 square feet have been constructed and are listed in Table I. Also given are the date of start of construction and the date of completion of each structure and its cost (to the nearest dollar) per square foot. The variation in unit cost results from the differing functional requirements of these structures and to a lesser extent to the state of the national economy at the time competitive bids were received in each instance. The costs given here include in each case the building and its standard facilities (heating, lighting, plumbing, ventilation, etc.) but not the cost of specialized equipment housed in the building. The total cost of the structures listed is \$22.9 million. Other miscellaneous small buildings add about \$2 million to this total. It is perhaps interesting to note that the average cost of all the listed space is \$26 per square foot.

In September 1961, Congress appropriated \$114 million for the construction of the two-mile accelerator. The integrated overall project construction costs and costs plus commitments versus time are shown in Fig. 2. These are in the form of the classical "lazy S" curves applicable to most construction projects. These curves represent the sum of SLAC and ABA (Aetron-Blume-Atkinson, the joint venture of three firms handling the architect-engineering-management work for SLAC) cost experiences. Consideringthe relatively small amount of construction remaining at this date, it now appears that the entire construction effort will be accomplished within the appropriated amount.

During the summer of this year, an electron beam was successfully accelerated through the two-mile machine and a maximum energy of 18.4 GeV was obtained. Extensive testing has shown that the accelerator meets specifications in almost all respects. The only notable exception to this statement is the failure to attain rated beam current due to the onset of "beam break-up" phenomena. A detailed account of these phenomena and of steps being taken to overcome them is being presented in a separate report<sup>1</sup> to this conference.

#### Description

The evolution of the design of the two-mile accel-<br> $\frac{2-6}{4}$  h high erator is given in several earlier reports.  $2^{-0}$  A brief review of the final design configuration of the accelerator and some of its principal components will be given here.

A schematic diagram depicting the entire twomile accelerator is shown in Fig. 3. The first item of interest starting at the west end of the accelerator is the main injector. A beam analyzing station (No. 1) is located about 40 feet downstream from the injector to permit easy set-up and analysis of the injected beam. An instrumentation section is located in a S-foot drift space at the end of each 333-foot sector to provide information on beam current, intensity, and profile. This region also contains the quadrupoles used for focusing the beam. A branch in the accelerator housing at the one-third point is provided for the future construction of a reduced energy research facility at this point. Just downstream from the beam take-off point is a facility to house a future injector. Next in line is the positron source, which has assumed increasing importance as the research program planning has evolved. Beam Analyzing Station No. 2, located in Sector 20, permits beam operation and analysis during periods when the beam cannot be run into the beam switchyard or research area. At the two-thirds point, another branch in the accelerator housing occurs to provide for a second future intermediate research facility. It is at this point that the proposed positronelectron storage ring<sup>7</sup> will be located if its construction is authorized by Congress. The beam switchyard at the east end of the accelerator contains the magnets, collimators, energy-defining slits, and other special beam handling and monitoring instrumentation required to determine the characteristics of the beam and to direct it into any of the three experimental areas. Two separate control rooms have been constructed. One of these, called the Central Control Room, contains all of the instrumentation and control equipment required for monitoring and controlling the components and systems of the two-mile accelerator. The other control room, known as the Data Assembly Building, contains similar equipment relating to monitoring and controlling the components and systems of the Beam Switchyard. Numerous control and data links extend between the two control rooms. Some data links also extend to the research area from the Data Assembly Building.

#### Accelerator Housing And Klystron Gallery

A cross section of the accelerator housing and the klystron gallery is shown in Fig. 4. Penetration tubes, 27 inches in diameter and spaced 20 feet apart, carry the waveguides, vacuum manifolds, water cooling pipes and instrumentation cables between the two housings. A personnel accessway is provided once per sector (333 feet).

#### Accelerator Structure and Support Girder

A complete 10-foot accelerator section<sup>8</sup> is shown in Fig. 5. This structure is fabricated by a special brazing technique from separately machined copper disks and rings. The dimensions of the separate parts vary over the 10-foot length to result in a constant

 $*$ Work supported by the US Atomic Energy Commission.

gradient design, i.e. , the axial electric field which accelerates the electrons is constant over the entire section. Four 10-foot sections are mounted on a 40foot aluminum girder which serves not only as a support for the sections but also as a portion of a continuous "light pipe" through which a laser beam from a source located at the east end of the accelerator is transmitted for alignment purposes $^{\circ}$  through a distance of two miles to a detector located at the west end. A retractable Fresnel target located in one end of each 40-foot girder, as shown in Fig. 6, serves to focus the laser beam on the detector. Deviations of the accelerator from correct alignment are indicated by the displacement of the laser image at the detector. Realignment is accomplished by adjustment of the jacks supporting the girders.

#### High Power Klystrons

A total of 245 klystrons is required to fill all the sockets along the two-mile accelerator. These tubes are rated at 21 megawatts peak and 21 kilowatts average power output. They have a design capability somewhat higher than these ratings. In view of the importance and criticalness of these tubes to the entire accelerator program and the large number required to fill the accelerator sockets and the repair cycle "pipe-line," several parallel procurements have been undertaken. A basic tube design with permanent magnet focusing was first developed at SLAC. Later four outside companies fabricated tubes meeting the same basic specifications but not identical to the SLAC tube in mechanical and electrical design details. However, all of these tubes are electrically and mechanically interchangeable in the SLAC sockets. These developments led to three outside procurements, as noted in Table II. These procurements did not all start at the same time. Klystron models manufactured by SLAC and the four commercial companies are shown in Fig. 7. These tubes are all designed to be repairable. After repair they are put through the same acceptance tests as new tubes and must meet original specifications. As discussed in more detail later, operating results and failure rates have been sufficiently favorable that we have been able to negotiate an extended warranty agreement with one of the outside suppliers wherein, after initial purchase of a tube, SLAC pays a fixed hourly rate for the first 1500 hours of operation of the tube and a decreased rate thereafter. Under this arrangement, the manufacturer agrees to replace a failed tube with a tube meeting original specifications at no additional cost to SLAC.

Installation of klystrons in the gallery is shown in Fig. 8. Also shown in the same figure are the modulators and the ion-getter pumps and manifolds which evacuate the waveguides and accelerator.

The power output from each klystron is divided into four equal parts and feeds four successive lofoot accelerator sections. The general arrangement is illustrated in Fig. 9. Waveguide feeds to adjacent sections are from opposite sides to compensate for residual coupler asymmetries. The provision of a waveguide valve just above each klystron allows the klystron to be replaced without affecting the accelerator vacuum or interfering with beam operation.

#### Other Components and Systems of the Accelerator

Time limitations permit only a brief review of the few more accelerator components and systems. The main injector system is illustrated in Fig. 10. The 80-kV, 2-ampere gun is followed by a pre-bunching cavity, a disk-loaded  $v_n = 0.75$  c buncher, and a 10 foot accelerator section. For phase coherence, all of these are fed by a single high power klystron. Electron bunching within 5 electrical degrees (0.15 cm) is achieved. Other devices are provided to focus, steer and collimate the beam, and to measure its intensity, transverse position, phase spread and profile. The injector is capable of providing three independent timeinterlaced beams, each adjustable in intensity and pulse length.

At the end of each 333-foot sector there is an instrument section, as shown in Fig. 11, containing devices for measuring the beam characteristics (intensity, position, and profile) and for steering and focusing the beam. The position monitors are of the microwave type utilizing two rectangular cavities which are excited in the TM120 mode by an off-axis beam. Two types of intensity monitors are used. One is a cylindrical microwave cavity operating in the  $TM_{010}$  mode which provides, with the aid of appropriate electronic circuitry, an output proportional to the logarithm of the charge contained in the beam pulse. The other intensity monitor, utilizing a toroid with a ferrite core, provides an output directly proportional to the charge in the beam pulse. Both of these devices, as well as the position monitors, can sense independent pulses and therefore are capable of measuring the characteristics of several time-interlaced beams. Focusing is provided by means of three quadrupoles whose coils are connected in series and whose polarities are such that they constitute a quadrupole triplet. In order to strengthen the overall focusing along the accelerator for the purpose of increasing the beam break-up threshold,  $\frac{1}{1}$  the quadrupole triplets will shortly be rearranged in the form of doublets with the shorter quadrupoles located in the early sectors and the longer quadrupoles in the later sectors. Larger power supplies will be provided for the downstream doublets.

The positron source located at the one-third point along the accelerator includes equipment located on three separate girders having a total length of 40 feet. This equipment is shownin Fig. 12. The positron sources proper are located on the first girder. There are two sources each consisting of one layer of gold, two of silver, and eight of copper, with cooling water passages between. Each layer is approximately l/8 inch thick. The total material in each source constitutes about 3.5radiation lengths. One of the sources (the "wand" radiator) is used to give an intermittent source of positrons. It consists of a small target about 0.5 inch wide which is driven across the beam path on command in a time equivalent to about nine machine pulses (at 360 pps). The center pulse of this group results in a positron pulse; the other eight are eliminated by gating of the main injector. All the remaining pulses (up to 351) may be electrons if desired. If

a continuous low intensity positron beam is desired, the wand can be caused to remain stationary in the incident beam, The "wand" source and its associated hardware (still incomplete) is shown in Fig. 13. The second source, shown in Fig. 14, is in the form of a water-cooled wheel which is caused to undergo a nutating motion by a crank mechanism and an external motor drive system. With an incident electron beam of 100 kW on the wheel radiator, it is predicted that a positron average current of approximately  $1.5 \mu A$ can be accelerated in an energy band of about 1% and a transverse phase space of approximately 0.3 mc-cm. A magnetic lens system is used to improve the match between the source emittance and the accelerator phase space acceptance. The radiator is located in a 20 kilogauss axial magnetic field which decreases rapidly to 2.4 kilogauss about 2 feet downstream from the radiator and then remains constant for the next 25feet. Acceleration begins 2.5 feet downstream from the radiator and the positron energy at 25 feet is about 75 MeV, at which point the solenoidal focusing is replaced by a series of 13 quadrupole triplets whose spacing increases with energy until this focusing system merges with the regular machine triplet system located at the end of each sector. A pulsed rf deflector located downstream from the converter is used to produce an angular deflection of the positron and electron beams. Since the beams are 180 degrees apart in phase, they are both deflected by the same angle; a pulsed steering magnet is then used to restore the direction of whichever beam is desired to the axis while the other unwanted beam isdeflected even farther from the axis. The second and third of the positron source girders are shown in Fig. 15 as they appeared in the assembly area just prior to installation in the accelerator housing.

#### The Beam Switchyard

A diagram of the Beam Switchyard indicating the location of the various pulsed and dc magnets, steering devices, quadrupoles , collimators, slits, and other beam handling equipment is shown in Fig. 16. Also shown is the cross section of the switchyard housing at various locations along its approximately lOOO-foot length. The beam path is in the lower level of the two-level structure and is separated from the upper level by means of removable concrete shielding blocks. The upper level contains utility and instrumentation runs, tracks for maintenance cars, and three overhead cranes.

The problems of operation and maintenance in the switchyard are severe because of the high-level radiation environment. Provisions for remote handling of components in case repair or replacement is necessary have been essential. These factors have imposed design requirements which make the achievement of a vacuum as low as that of the accelerator  $\alpha$ (  $< 10^{-7}$  torr) very unlikely. Fortunately, a much poorer vacuum (10<sup>-9</sup> to 10<sup>-+</sup> torr) is adequate in this area. At these pressures and with the large volume to be pumped, diffusion pumps are the most feasible and economic solution. The entire siwtchyard is pumped by eight stations located at ground level about 40 feet above the top of the switchyard. Each pumping station utilizes a 6-inch oil diffusion pump and a refrigerated trap. Connection of each pumping station to the switchyard is accomplished by means of a

20-inch-diameter stainless steel manifold. The beam switchyard vacuum system is effectively isolated from the accelerator vacuum system by means of a differential pumping system, which includes a special 10-foot  $\arcc{e}$  accelerator section refrigerated to -30 $^{\circ}$ C to trap migrating oil molecules.

An artist's conception of the exterior of the beam switchyard and its surroundings is shown in Fig. 17.

The high average power ( $\approx$  1 MW) carried by the electron beam requires the use of large and complicated beam handling devices. Typical of such devices is the adjustable energy-defining slit<sup>11</sup> shown in Fig. 18. This unit is 16 feet long. Two such devices, with one rotated 90 degrees about the beam axis with respect to the other, are utilized as an adjustable collimator to define the size and position of the beam as it enters the beam switchyard.

#### Beam Switchyard Instrumentation

Instrumentation is provided in the beam switchyard to measure the beam characteristics (intensity, pulse shape, energy, energy spectrum, transverse position, and profile). Other instruments are used to protect various sensitive devices from being damaged by high temperatures or radiation effects due to excessive beam interception. A typical instrumentation grouping is shown in Fig. 19. All of the instruments illustrated can be installed, aligned, or removed by personnel located in the upper level using remote operation tools.

A general impression of some of the individual instruments in the beam switchyard can be gained from Figs. 20-26.

Figure 20 shows a 3-inch-aperture beam current monitor which incorporates two ferrite toroids. One of the toroids is used in the current monitoring system and the other in a precision integrating system. These devices have a sensitivity of 2 mV/mA and an effective rise time of 30 nsec.

The microwave-type beam position monitor shown in Fig. 21 is similar to that used along the accelerator except that the aperture has been increased to 2 inches. To prevent unwanted coupling between the x and y cavities, it was found necessary to separate them by about 4 inches by means of a 2-inch-aperture pipe. The sensitivity of these monitors is  $0.17 \text{ mW/(mA mm)}$ .

Three different techniques for viewing the beam profile are employed in the devices shown in Figs. 22, 23, and 24. A synchrotron light pick-up station is shown in Fig. 22. Light radiated by the electron beam as it is deflected in the first dc magnet of the beam transport system is transmitted by means of metal mirrors to a television camera located in an enclosure on top of the beam switchyard. Using a vidicon camera, this monitor can view beams having a current of  $1 \mu A$ or higher. The lower limit can be extended to about  $10^{-2}$   $\mu$ A by using a more sensitive orthocon camera. An advantage of this type of monitor is that it does not intercept the beam. However, it is somewhat less sensitive than the monitors to be described next and requires a suitable magnetic field in the near vicinity of the position where the profile is viewed. A "carrousel" type of zinc sulphide screen profile monitor is shown in Fig. 23. Since the individual screens lose luminescence after exposure to an integrated beam current of about 10  $\mu$ A hr/cm<sup>2</sup>, the carrousel device containing 48 independent screens has been designed. The carrousel can be rotated and the individual screens

can be raised or lowered by remote control. Other types of zinc sulphide screen monitors in the forms of individual flaps and film rolls have also been developed. The third type of profile monitor, utilizing Cerenkov light, is shown in Fig. 24. A cell filled with argon or helium at a pressure slightly above atmospheric can be moved into the beam by a remotely controlled pivoting action. The gas cell has 3-milthick aluminum entrance and exit windows. The Cerenkov light cone is reflected vertically by a l-milthick aluminum mirror located inside the cell at an angle of 45 degrees to the beam. The beam image is observed by means of a television camera looking at one part of the radiated cone. The Cerenkov monito has a sensitivity of about  $10^{-3}$   $\mu$ A/cm<sup>2</sup> (with helium or  $3 \times 10^{-4} \mu\text{A/cm}^2$  (with argon).

Secondary emission monitors (SEM) are used in the beam switchyard for spectrum analysis, monitoring of beam interception by the jaws of the slits and collimators, to provide beam centering information at the high power beam dumps, and to provide signals for the equipment protection system. Aluminum foils having a secondary emission coefficient of about 4% are used in all of these applications. The spectrum analyzer used just upstream from the main energydefining slits is shown in Fig. 25. Two foil holders each carrying 6 foils made of l-mil aluminum can be moved into or out of the beam by a remotely controlled pivoting action. The foils vary in width from 6 mm (0.1% resolution) for those nearest to the beam axis to 48 mm $(0.8\%$  resolution) for those farthest from the axis.

Ŷ)

To prevent thermal or radiation damage to components accidentally struck by the high power electron beam, an equipment protection system consisting of more than 20 protection collimators along with associated ionization chambers and thermometers is used in the beam switchyard. An out-of-limit signal from any of these sensing devices will either shut off the beam or give a warning signal to the operator. A typical water-cooled protection collimator and its two associated ion chambers are shown in Fig. 26. The size and shape of the aperture of eacn protection collimator depends upon the corresponding beam configuration at the point where the protective action takes place.

#### The Research Area

A general view of the research area is shown in Fig. 27. The beam line to End Station A has been completed and the beam line to End Station B is scheduled to be completed by November 1, 1966. A central (C) beam line has been designed and will be completed by June, 1967. End Station A will be used primarily for elastic and inelastic scattering of electrons and positrons and for photoproduction experiments. Three large spectrometers capable of analyzing scattered particles with maximum energies of 1.8, 8, and 20 GeV, respectively, are nearing completion. A two meter streamer chamber and associated magnet will be located just downstream from End Station A. This device is on rails and can be moved into or out of the beam as desired. After passing through the A experimental area, the electron beam will be transported in an evacuated pipe covered by concrete shielding to a high power dump ("Beam Dump East"). End Station B

is the center of experiments involving secondary particles. Initially, three secondary beams of  $\mu$  mesons, monochromatic photons, and K mesons, respectively, are being constructed. A one-meter bubble chamber and a large spark chamber, both under construction at SLAC, will be used in conjunction with the secondary particle experiments on the B side. An 82-inch bubble chamber is now undergoing modification at Berkeley and will be moved to SLAC around July, 1967. It will be used in conjunction with the C-Beam experimental program.

#### Operating Results

Starting in January, 1965, when a 1.5-GeV beam was accelerated through the first two sectors (666 feet), the two-mile machine has been activated in a series of steps. Some of the more significant events and dates are shown in Table III. The most recent advance noted was the transmission on September 20, 1966, of a beam through the accelerator, the beam switchyard, and End Station A to the farthest point the beam can reach, Beam Dump East.

The performance of the accelerator in its operation to date is summarized in Table IV. Comparison of the experimental results with design parameters shows that, with the exception of peak beam current which has been limited by beam break-up phenomena to about 40% of the design level, all design parameters have been met or bettered.

Typical energy spectra are shown in Fig. 28. The spectrum of higher energy was obtained for light beam loading (2.3~mA peak beam current). A higher beam current (15-mA) results in the lower energy spectrum having 0.44 GeV less energy than the first. Note that the lower energy spectrum contains some electrons having energies as high as the lightly loaded spectrum. These are the electrons which pass through the accelerator during the earliest part of the pulse before a significant amount of stored rf energy has been absorbed by the beam. The spectrum width at half maximum is 1.33%, of which about 0.9% is attributable to the resolution of the measurement devices.

Compensation of the beam loading exhibited in Fig. 28 can be achieved by delaying the trigger to one or more of the accelerator sectors. In this case, the first electrons acceleratedduring the beam pulse pass through some sectors which are only partly filled with rf energy and therefore they gain less energy than they would if the sectors were completely filled. The compensation achievable with this technique is demonstrated in Fig. 29, where the dashed line shows the uncompensated spectrum and the solid line shows the compensated spectrum resulting from trigger delay to one sector by approximately  $0.5 \mu$ sec.

The energy spectra shown in Figs. 28 and 29 represent the output of one foil of the spectrum analyzers described earlier as the deflecting magnet current is swept through a range corresponding to the energy spread of the electrons in the beam. The simultaneous output of all of the foils gives an instantaneous representation of the energy spectrum as shown in Fig. 30.

Beam transmission through the accelerator has been found to be quite good. About 90% of the current measured close to the input end of the machine, say at the beam analyzing station No. 1 located 40 feet from

the injector, is preserved through the entire two-mile length. This favorable result arises from the effective performance of the beam position and intensity monitors, the steering and focusing systems, and the long ion chamber. The microwave position monitors have been able to indicate the transverse position of the beam with respect to the accelerator axis within  $\pm$  0.5 mm. The long ion chamber, which is an argon plus  $CO<sub>2</sub>$ -filled 1-5/8-inch coaxial line installed alongside the accelerator, enables the operator to detect beam losses and from the times of arrival of the ionization signals to resolve their location within one or two hundred feet. The oscilloscope presentation of the long ion chamber signal is shown in Fig. 31, where the ordinate is the ion chamber reading proportional to beam interception and the abscissa is the sector number where the beam loss occurs. The solid-line curve shows the original ion chamber reading and the other curves show the ion chamber reading as the steering and focusing is improved in successive steps by the operator.

Tests to date have demonstrated the feasibility of accelerating at least three beams of different energies in a time-interlaced manner. The pulse length and intensity of each of these beams can be independently controlled. Spectra of three beams having energies of 1.17, 1.79, and 2.42 GeV are shown in Fig. 32. These relatively low energy beams were measured at beam analyzing station No. 2 located at Sector 20. The irregular base line is due to scattered low energy electrons at the analyzing station. Figure 33 shows the presentation observed by the operator on the "Linear Q" oscilloscope when two interlaced beams having energies of 11 and 5.65 GeV, respectively, are accelerated. The heights of the dots above the base line are proportional to the beam current measured at the end of each sector. These measurements are made by means of the toroid-type beam current monitors described earlier.

The beam break-up phenomena which presently limit the maximum current achievable with the twomile accelerator are discussed in detail in a separate  $report<sup>1</sup>$  being presented to this conference and will not be treated comprehensively here. Various means of curing or mitigating the beam break-up problem are under investigation. These include microwave suppression and feed-back schemes, the addition of sextupole or octupole devices, and improved quadrupole focusing along the machine. The latter measure is the only one which so far has been found to be definitely efficacious. As discussed earlier steps are now underway to increase the effective quadrupole focusing strength along the two-mile accelerator and it is expected that these measures will lead to a significant increase in the beam breakup current threshold.

During the past several months an endurance run involving 14 sectors of the accelerator was conducted. In order to gain information as to the performance and life of accelerator components (especially klystrons) as a function of the operating level, seven different levels were maintained as indicated in Table V. The variables were the pulse repetitition rate which ranged from 60 to 360 pulses per second and the klystron beam voltage ranging from 200 to 245 kV. The corresponding peak power outputs and the equivalent electron beam energies (assuming 230 klystrons in operation) are also shown in the table.

Experimental results are the total failures under each condition in approximately 1200 hours of operation and the average number of klystron station faults per minute (extrapolated to 230 stations). The number of failures for the various conditions ranged from 0 to 6 and the fault rate from 0.2 to 2.0 per minute. The larger number in each instance occurred for the operating condition resulting in maximum peak and average power output from the klystrons. Another indication arising from the experimental results is the fraction of the time the klystrons will, on the average, be available under each operating condition. These latter results are shown in Fig. 34. These results, while Preliminary, would seem to suggest that an attempt to achieve energies higher than about 20 GeV (at the maximum repetition rate) by increasing the klystron power output will meet with diminishing returns due to decreasing klystron availability.

Tests of the accelerator and beam switchyard are continuing on a 10-shift-per-week basis. Installation in the B side of the beam switchyard and in the research area takes place regularly during the day shift so that beam operation into these areas can now occur only during the two night shifts. However, accelerator beam tests up to Beam Analyzing Station No. 2 in Sector 20 are regularly conducted during the day shift without interference with other work. During the last two weeks, runs have started for preliminary checkout of equipment in End Station A. The actual research program in both End Stations A and B will start around December, 1966.

#### References

- 1. E. V. Farinholt, 2. D. Farkas, W. B. Hermannsfeldt, H. A. Hogg, R. F. Koontz, C. J. Kruse, G. A. Loew, and R. H. Miller, "Beam Break-Up Experiments at SLAC," 1966 Linac Conference (Los Alamos).
- 2. R. B. Neal and W.K.H. Panofsky, Proceedings of the International Conference on High Energy Accelerators, (CERN, Geneva, 1956), Vol. I, p. 530.
- 3. R. B. Neal, Proceedings of the International Conference on High Energy Accelerators, (CERN, Geneva, 1959), p. 349.
- 4. K. L. Brown, A. L. Eldredge, R. H. Helm, J. H. Jasberg, J. V. Lebacqz, G. A. LOeW, R. F. Mozley, R. B. Neal, W.K.H. Panofsky, T. F. Turner, Proceedings of the International Conference on High Energy Accelerators, (Brookhaven, 1961), p. 79.
- 5. W.K.H. Panofsky, Proceedings of the International Conference on High Energy Accelerators, (Dubna, 1963), p. 407.
- 6. J. BaIlam, G. A. Loew, R. B. Neal, Proceedings of the Fifth International Conference on High Energy Accelerators (to be published), (Frascati, 1965).
- 7. S. Berman et al., "Proposal for a High-Energy Electron-Positron Colliding-Beam Storage Ring at the Stanford Linear Accelerator Center" (June, 1965).
- 8. R. P. Borghi, A. L. Eldredge, G. A. Loew, and R. B. Neal, "Design and Fabrication of the Accelerating Structure for the Stanford Two-mile Accelerator," to be published in Advances in Microwaves, Vol. I, (Academic Press, New York, N.Y.)
- 9. W. B. Herrmannsfeldt, The First National Particle Accelerator Conference, (Washington, D. C. , 1965), IEEE Transactions on Nuclear Science, Vol. NS-12, No. 3, p. 929. Also: W.B. Hermannsfeldt, "SLAC Alignment System," 1966 Linac Conference, (Los Alamos).
- 10. E. L. Hoyt and E. L. Garwin, The First National Particle Accelerator Conference, (Washington, D.C., 1965), IEEE Transactions on Nuclear Science, Vol. NS-12, No. 3, p. 740.
- 11. D. R. Walz, J. Jurow, and E. L. Garwin, The First National Particle Accelerator Conference, (Washington, D. C. , 1965), IEEE Transactions on Nuclear Science, Vol. NS-12, No. 3, p. 867.



# TABLE I

I

 $\boldsymbol{t}$ 

 $\mathcal{L}$ 

## TABLE II



# KLYSTRON PROCUREMENT CONTRACTS

245 Klystrons are required to fill all accelerator sockets.

ù.

# TABLE III

 $\mathbf{I}$ 

# Key Dates



# TABLE IV

## COMPARISON OF SIGNIFICANT SLAC DESIGN PARAMETERS AND PERFORMANCE (September 1966)



a This number had been adjusted to take into account non-standard accelerator girders.

b Based on 90% availability of klystrons.

ę

c Improvements are still required to make the system more reliable.



 $\gamma_{\rm eff} \chi_{\rm eff} \sigma_{\rm gas} = 0.5$  .

## ENDURANCE RUN

العادرج

 $\mathfrak{t}$ 



 $\small \begin{array}{cl} \texttt{P} & = & \texttt{EXPECTED} \ \texttt{A VERAGE} & \texttt{POWER} \ \texttt{OUTPUT} \ \texttt{ATP} & \texttt{KLYSTRON} \ \texttt{BEAM} & \texttt{VOLTAGE} & \texttt{V}_k \end{array}$ 

 $V_{\sim}$  = COMPUTED NO-LOAD ELECTRON BEAM UP ENERGY FOR 230 STATIONS OPERATING

### List of Figures

- 1. Air view of SLAC site showing the two-mile accelerator, the research facilities, and the principal laboratories and shops.
- 2. Integrated total costs and commitments for overall SLAC construction project.
- 3. Overall layout of the two-mile accelerator.
- 4. Cross-section of Accelerator Housing and Klystron Gallery.
- 5. View of completed lo-ft. accelerator section.
- 6. End of 40-ft. girder which carries four lO-ft. accelerator sections, Retractable Fresnel alignment target is shown inside the girder.
- 7. Klystron models manufactured by SLAC and four commercial companies.
- 8. Klystron Gallery interior.
- 9. General arrangement of klystrons, connecting waveguides and four 10-ft. accelerator sections supplied by each klystron.
- 10. Profile view of injector system.
- 11. Standard instrument section located at end of each 333-ft. sector.
- 12. Positron source and its associated focusing and instrumentation equipment.
- 13. "Wand" type positron source.
- 14. "Wheel" type positron source.
- 15. Special girders down-stream from positron source supporting focusing and instrumentation devices.
- 16. Plan view of Beam Switchyard showing principal components and cross section of housing at different points.
- 17. Exterior view of Beam Switchyard.
- 1s. High power energy-defining slit.
- 19. Typical instrumentation grouping in Beam Switchyard.
- 20. S-inch aperture beam current monitor.
- 21. Microwave-type beam position monitor.
- 22. Synchrotron light type beam profile monitor.
- 23. Carrousel type zinc sulphide screen profile monitor.
- 24. Cerenkov light profile monitor.
- 25. Spectrum analyzer.
- 26. Protection collimator with ion chambers.
- 27. Layout of Research Area,
- 28. Typical energy spectra showing effect of beam loading.
- 29. Compensation of spectrum broadening due to beam loading by trigger delay to 1 sector.
- 30. Energy spectrum as observed by means of foil-type spectrum analyzer.
- 31. Improvement of beam steering with aid of PLIC.
- 32. Energy spectra of 3 time-interlaced beams at beam analyzing station No. 2 (Sector 20).
- 33. Beam intensity at end of each sector for two inter-laced beams.
- 34. Relative availability of klystrons under different operating conditions.





 $\sim$ 

 $\overline{1}$ 

 $\mathbf{y} = \mathbf{y} \in \mathbb{R}^{d \times d}$  , where

in the contract of



 $\mathbf{1}$ 

 $\sim$ 

 $\sim 10^{-1}$ 

 $\label{eq:2.1} \mathcal{F}(\mathcal{S}_{\mathbf{a}}^{\mathbf{a}}(\mathbf{a}^{\mathbf{a}})) = \mathcal{F}(\mathbf{a}^{\mathbf{a}}(\mathbf{a}^{\mathbf{a}}), \mathbf{a}^{\mathbf{a}}(\mathbf{a}^{\mathbf{a}})) = \mathcal{F}(\mathbf{a}^{\mathbf{a}}(\mathbf{a}^{\mathbf{a}}), \mathbf{a}^{\mathbf{a}}(\mathbf{a}^{\mathbf{a}})) = \mathcal{F}(\mathbf{a}^{\mathbf{a}}(\mathbf{a}^{\mathbf{a}}), \mathbf{a}^{\mathbf{a}}(\mathbf$ 

 $\mathcal{R}^{(1)}$  and  $\mathcal{R}^{(2)}$  and  $\mathcal{R}^{(3)}$ 



 $\alpha$  .

 $\sim$ 

그 고양 12

ti tufi.







)<br>N

EIMAC SPERRY LITTON SLAC RCA 353 -5-A



 $\alpha$ 



þ



 $\label{eq:2.1} \mathcal{L}_{\mathcal{A}}(\mathcal{A}) = \mathcal{L}_{\mathcal{A}}(\mathcal{A}) = \mathcal{L}_{\mathcal{A}}(\mathcal{A})$ 

 $\mathcal{A}(\mathbf{r})$  and  $\mathcal{A}(\mathbf{r})$  and  $\mathcal{A}(\mathbf{r})$  and

t frisc

 $\sim 100$ 

 $\mathbf{H}^{\mathrm{eff}}$ 

 $\mathcal{O}(\log n)$  , where  $\mathcal{O}(\log n)$ 

 $\mathcal{L}^{\text{max}}_{\text{max}}$ 

 $\gamma_{\mathcal{P}}$  spaces.



92 P.M

Grina i

 $\mathbf{1}$ 

 $\sim$ 



 $\sim 100$ 

 $\sim$ 

 $2.35 + 1.7$ 

r Teoreta (1970)

 $\mathcal{A}(\mathcal{A})$  and  $\mathcal{A}(\mathcal{A})$  and  $\mathcal{A}(\mathcal{A})$ 



39°



 $\bar{1}$ 

 $\bar{1}$ 





 $\sim 1000$ 

 $\mathcal{L}_{\text{max}}$  and  $\mathcal{L}_{\text{max}}$  and  $\mathcal{L}_{\text{max}}$  and  $\mathcal{L}_{\text{max}}$ 

 $\mathbf{T}$ 

 $\alpha$  .

2.2597

 $\sim 100$ 

 $\sim$ 



1994. PR

 $\mathcal{L}^{\mathcal{L}}$  and  $\mathcal{L}^{\mathcal{L}}$  are  $\mathcal{L}^{\mathcal{L}}$  . The contribution

 $\mathcal{A}^{\mathcal{A}}$ 



















j.



ENERGY, GeV

 $538 \cdot 1 \cdot B$ 





 $538 - 8 - A$ 



 $\sigma$ 

 $\rightarrow$   $\ensuremath{\mathrm{d}}\rho$ 

 $\bar{1}$ 

 $\rightarrow$ 

 $\gamma_{\rm c}$ 



J

ţ

 $\frac{1}{2}$ 



age of the court

 $\overline{1}$ 

 $593 - 4 - A$