

Phenomenology of Supersymmetric

Large Extra Dimensions :

Gravitino Effects in Selectron

Pair Production + Decay

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Supersymmetric Large Extra Dimensions

Original Model: Gravity in bulk

Arkani-Hamed,

Dimopoulos, Dvali

SM on brane

$M_0 \sim \text{TeV}$

Solve hierarchy problem w/o SUSY

Introduce Supersymmetry: Motivation

- Embed ADD in string theory
- SUSY stabilizes bulk radii
- Extra Dims provide new SUSY breaking mechanisms

Kaplan et al.
Chacko et al.
Di Clemente et al.
Arkani-Hamed et al.
⋮

SM supermultiplets on brane

Gravity supermultiplet in bulk

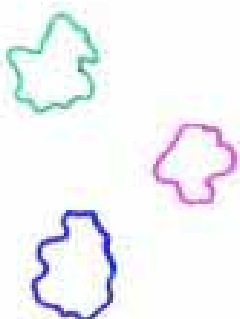
⇒ Graviton + Gravitino expand into KK Towers

Phenomenology: KK Gravitino Exchange

KK Gravitino Emission

see talk by
J. Song

The Model:



Near Brane

$D=10, N=1$

1 gravitino (Majorana-Weyl)

Bulk

$D=10, N=2$

Type II B

2 gravitinos (Majorana)

On Brane

$D=4, N=1$

4 Majorana spin- $3/2$: Assume

24 Majorana spin- $1/2$

3 have $\sim M_D$

1 is $m_0 \sim \Lambda_{string}^2 / m_{pl}$

4-D Effective Theory

$$E^{-1} \mathcal{L}_{\text{kin}} = \frac{i}{2} \bar{\Psi}_{\tilde{\lambda}} \Gamma^{\tilde{\lambda}\tilde{\nu}\tilde{\rho}} \nabla_{\tilde{\rho}} \Psi_{\tilde{\nu}}$$

$$\nabla_{\tilde{\rho}} = \partial_{\tilde{\rho}} + \frac{i}{4} \omega_{\tilde{\lambda}\tilde{m}\tilde{n}} \Sigma^{\tilde{m}\tilde{n}}$$

Expand: $\Psi_{\tilde{\lambda}}(x, y) = \sum_{\tilde{\lambda}} \frac{\tilde{\Psi}_{\tilde{\lambda}}^{\tilde{\lambda}}(x)}{\sqrt{V_{\tilde{\lambda}}}} e^{i\tilde{\lambda} \cdot \tilde{y} / R_c}$

Decompose $D=10$ Majorana-Weyl gravitino

Redefine fields to separate spin $1/2 + 1/2$ in \mathcal{L}_{kin}

Choose $\Gamma^m \chi_{\tilde{m}}^{\tilde{\lambda}} = 0$ to project out spin- $1/2$ comp.

$$\mathcal{L}_{\text{eff}}(x) = \int dy^{\tilde{\lambda}} \mathcal{L}(x, y)$$

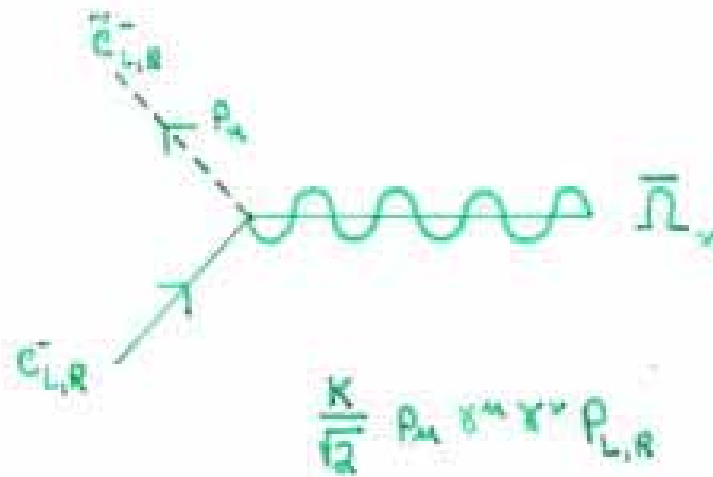
$$e^{-1} \mathcal{L}_{\text{eff}}(x) = -\frac{1}{2} \epsilon^{\mu\nu\rho\sigma} \bar{\omega}_{\tilde{m}}^{\tilde{\lambda}} (\delta_{\tilde{m}\tilde{n}} \partial_{\rho} \omega_{\sigma}^{\tilde{n}} - \frac{m_{\tilde{m}}^{\tilde{\lambda}}}{4} \bar{\omega}_{\tilde{m}}^{\tilde{\lambda}} [\gamma^{\mu}, \gamma^{\nu}] \omega_{\rho}^{\tilde{m}})$$

with $m_{\tilde{m}}^{\tilde{\lambda}} = \frac{\sqrt{\tilde{m} \cdot \tilde{m}}}{R_c}$

Interactions obtained via general Noether procedure

$$\mathcal{L}_I = -\frac{\kappa}{\sqrt{2}} |\epsilon| \left[(\partial_\mu \Phi_L) \bar{\Psi}_\nu \gamma^\mu \gamma^\nu \chi_L + (\partial_\mu \Phi_R) \bar{\Psi}_\nu \gamma^\mu \gamma^\nu \chi_R \right] + h.c.$$

$$\kappa = \sqrt{8\pi G_N}$$



Contributions to $e^+e^- \rightarrow \tilde{e}_{L,R}^+ \tilde{e}_{L,R}^-$

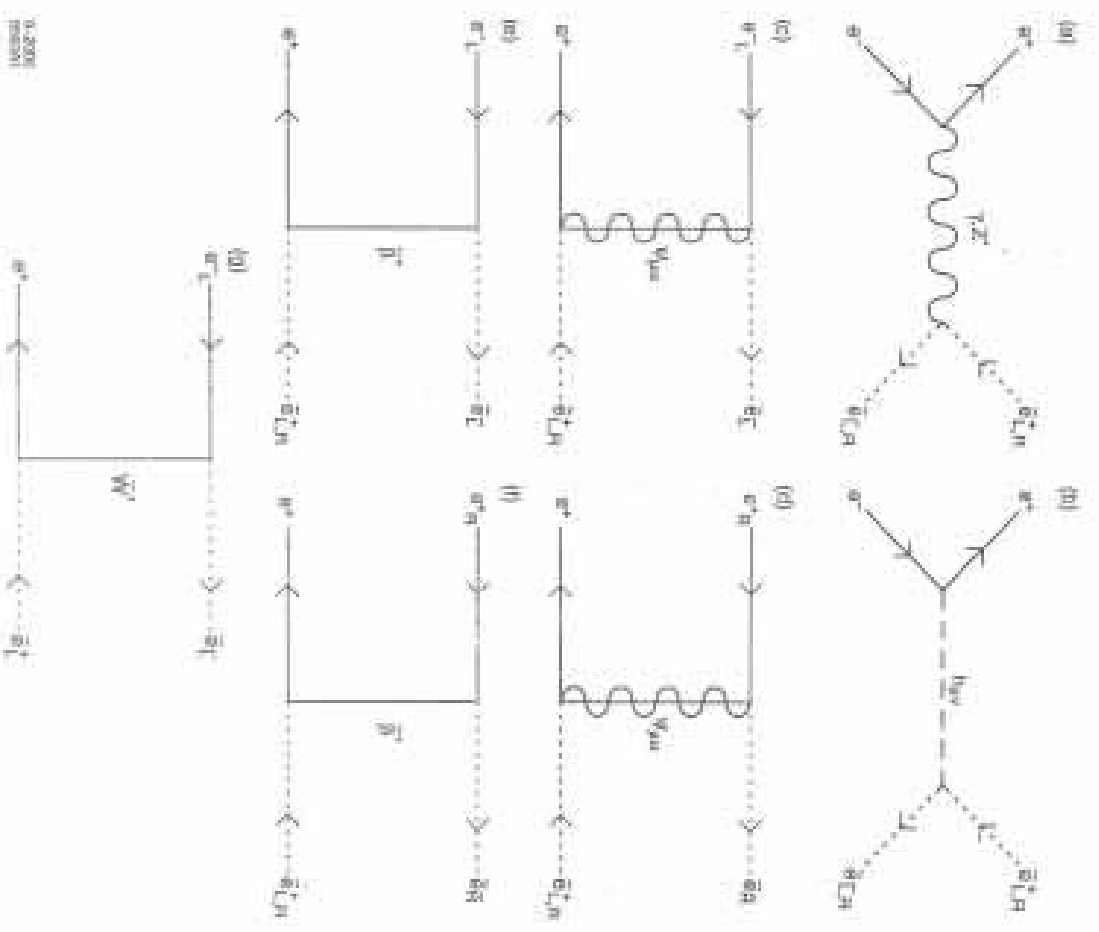


Figure 2: Various processes contributing to selectron pair production.

$$e^+ e^- \rightarrow e^{+n} e^{-n}$$

Polarized Scattering:

| | $\vec{e}_L^- \vec{e}_L^+$ | $\vec{e}_R^- \vec{e}_L^+$ | $\vec{e}_L^- \vec{e}_R^+$ | $\vec{e}_R^- \vec{e}_R^+$ |
|-------------------|---|------------------------------------|------------------------------------|--|
| $\vec{e}_L^- e^+$ | S-channel $\epsilon_1, \epsilon_2, \epsilon_n$ t-channel $\tilde{W}, \tilde{B}, \tilde{G}_n$ | | t-channel \tilde{B}, \tilde{G}_n | S-channel $\gamma, \tilde{Z}, \tilde{G}_n$ |
| $e^- e^+$ | S-channel $\epsilon_1, \epsilon_2, \epsilon_n$ | t-channel \tilde{B}, \tilde{G}_n | | S-channel $\gamma, \tilde{Z}, \tilde{G}_n$ t-channel \tilde{B}, \tilde{G}_n |

$$M_{G_n} = \frac{2}{N_c} (t-u) (K_1 - K_2)_{\mu\nu} \bar{e} \gamma^\mu e \quad \text{Dim} = 8$$

$$M_{G_n}^c = \frac{\kappa^2}{2} \sum_n \frac{t - m_n^2}{t - m_n^2} \bar{e} \gamma_\mu \gamma_\nu P_{\tilde{A}, \mu^T} \chi_\nu \gamma_\rho e$$

Summation of KK Gravitino States

$$P_{\vec{n}, \mu\nu}^{\vec{n}} = i(K + m_{\vec{n}}) \left(\frac{K^{\mu\nu}}{m_{\vec{n}}^2} - \eta^{\mu\nu} \right) - \frac{i}{3} \left(\gamma^{\mu} + \frac{K^{\mu}}{m_{\vec{n}}} \right) (K - m_{\vec{n}}) \left(\gamma^{\nu} + \frac{K^{\nu}}{m_{\vec{n}}} \right)$$

Integrate over density of states

$$\int_{m_0}^{\Lambda_c^2} dm_{\vec{n}}^2 \rho(m_{\vec{n}}^2) \frac{P_{\vec{n}, \mu\nu}^{\vec{n}}}{t - m_{\vec{n}}^2}$$

Evaluated via Appell's hypergeometric function

Leading order term:

$$\int \sim \frac{M_{Pl}^2}{\Lambda_c^2}$$

\Rightarrow independent of m_0

$$\frac{\chi^2}{a} \lesssim \frac{P_{\vec{n}, \mu\nu}^{\vec{n}}}{t - m_{\vec{n}}^2} \rightarrow \frac{-i 8\pi}{5 \Lambda_c^2} \left(\eta^{\mu\nu} - \frac{1}{3} \gamma^{\mu} \gamma^{\nu} \right)$$

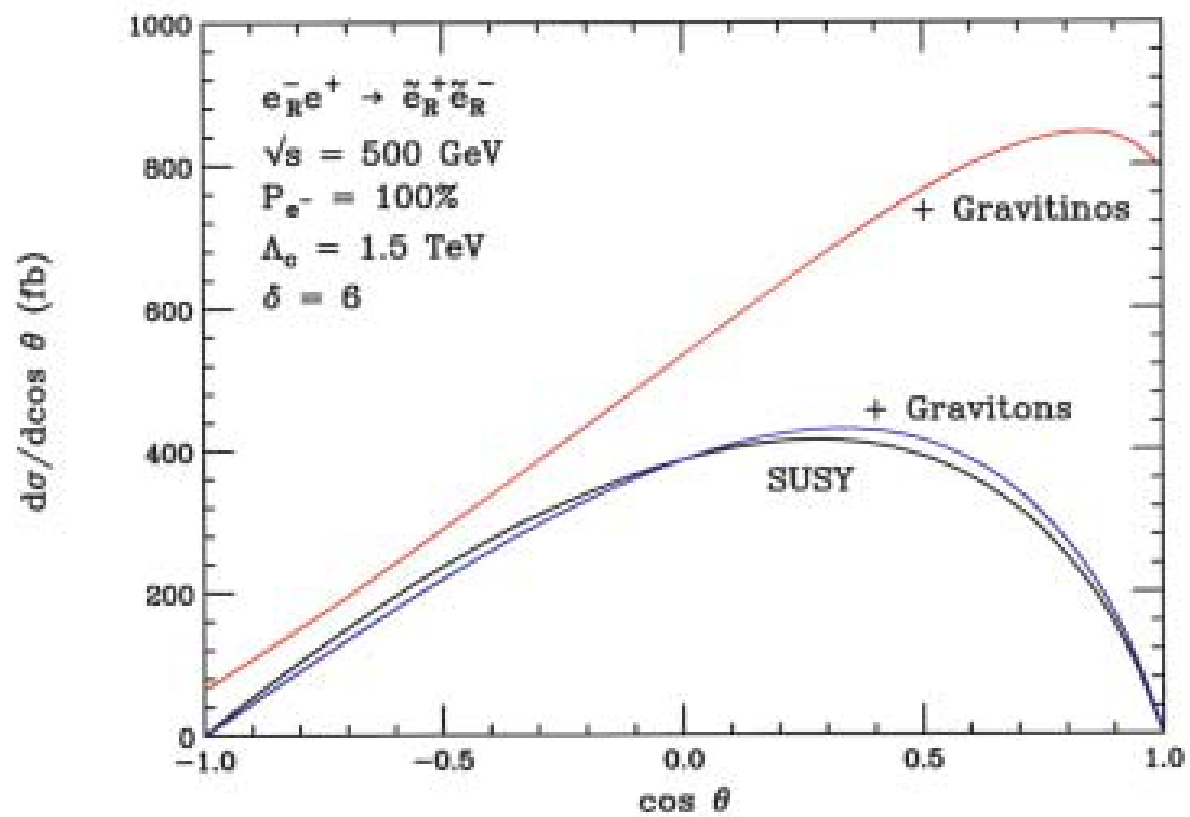
is Dimension - 6 !!!

Specify SUSY Model (D=4)

Choose GMSB: Naturally light zero-mode gravitino

2 sample sets of parameters:

| | $m_{\tilde{\chi}_1^0}$ | $m_{\tilde{\chi}_2^0}$ | $\tilde{\chi}_1^0$ |
|--------|------------------------|------------------------|--------------------|
| Set I | 217 GeV | 108 GeV | Bino-like |
| Set II | 210.5 GeV | 104.5 GeV | Higgsino-like |



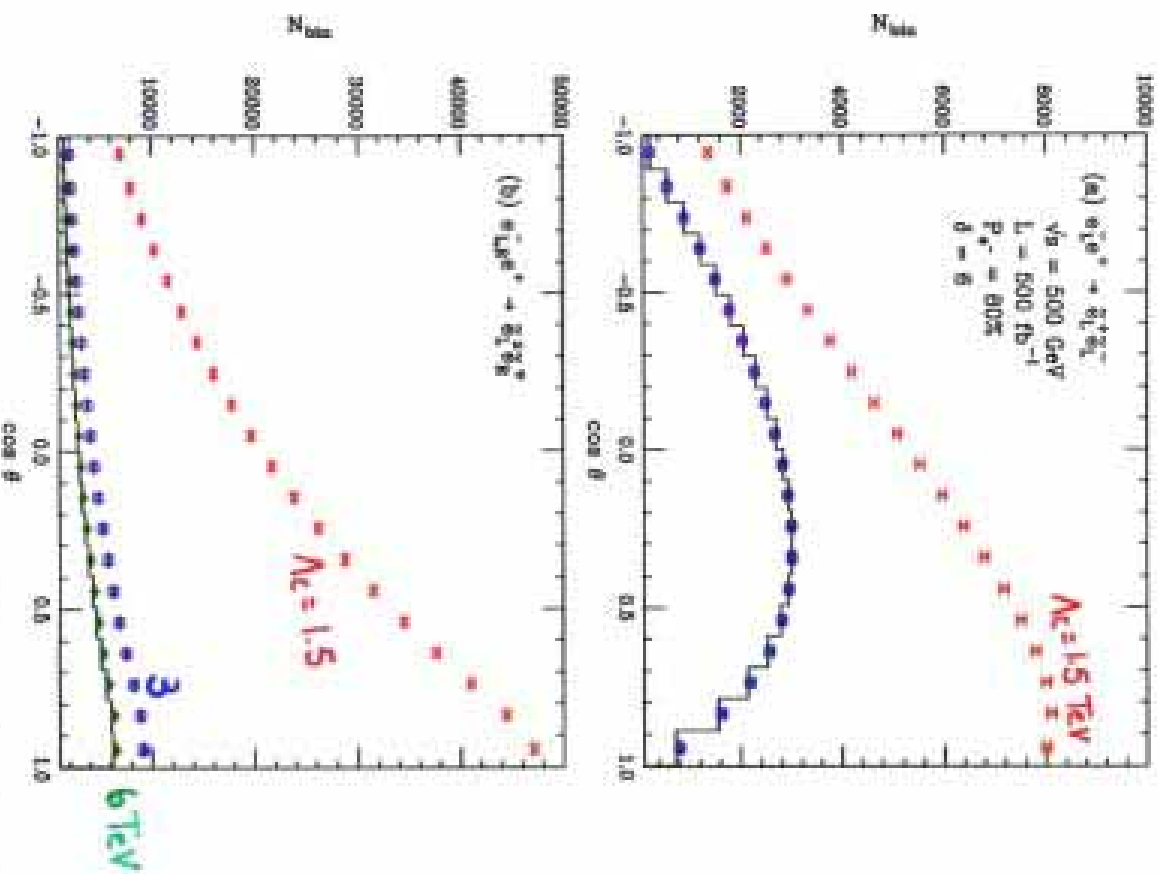
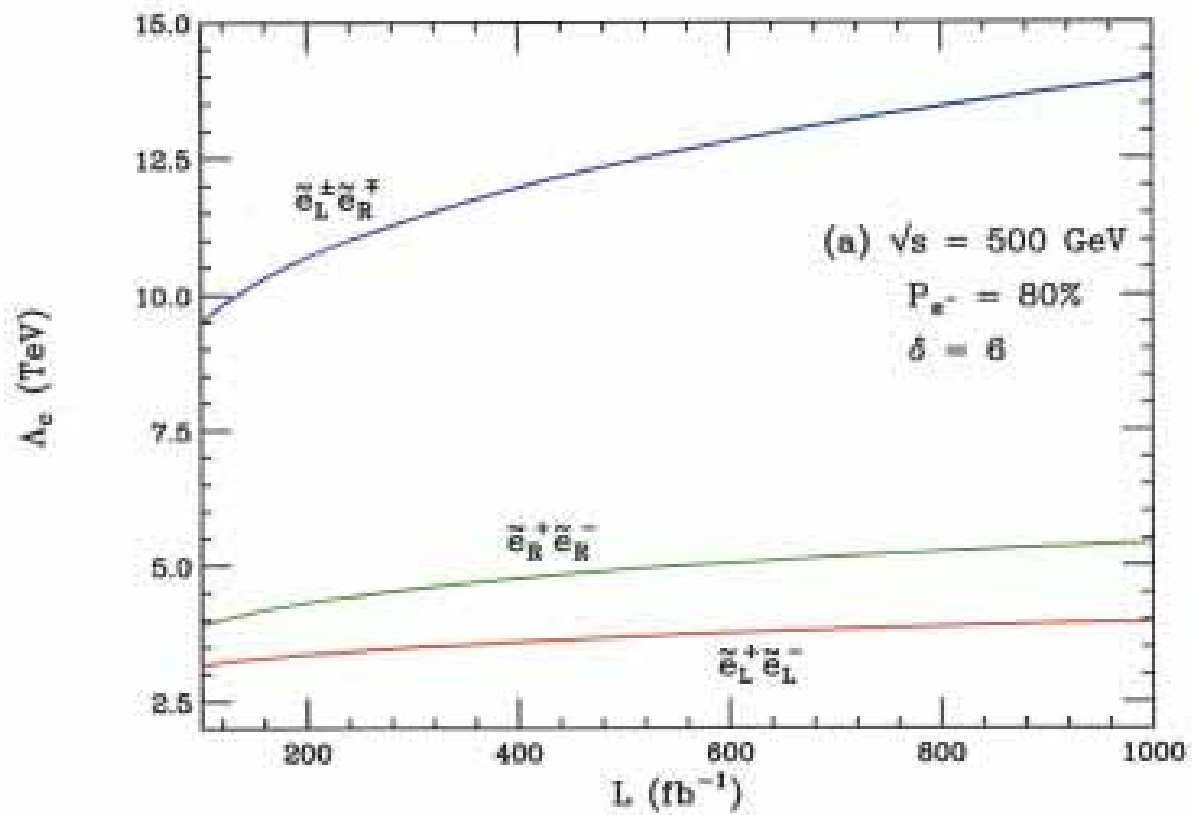
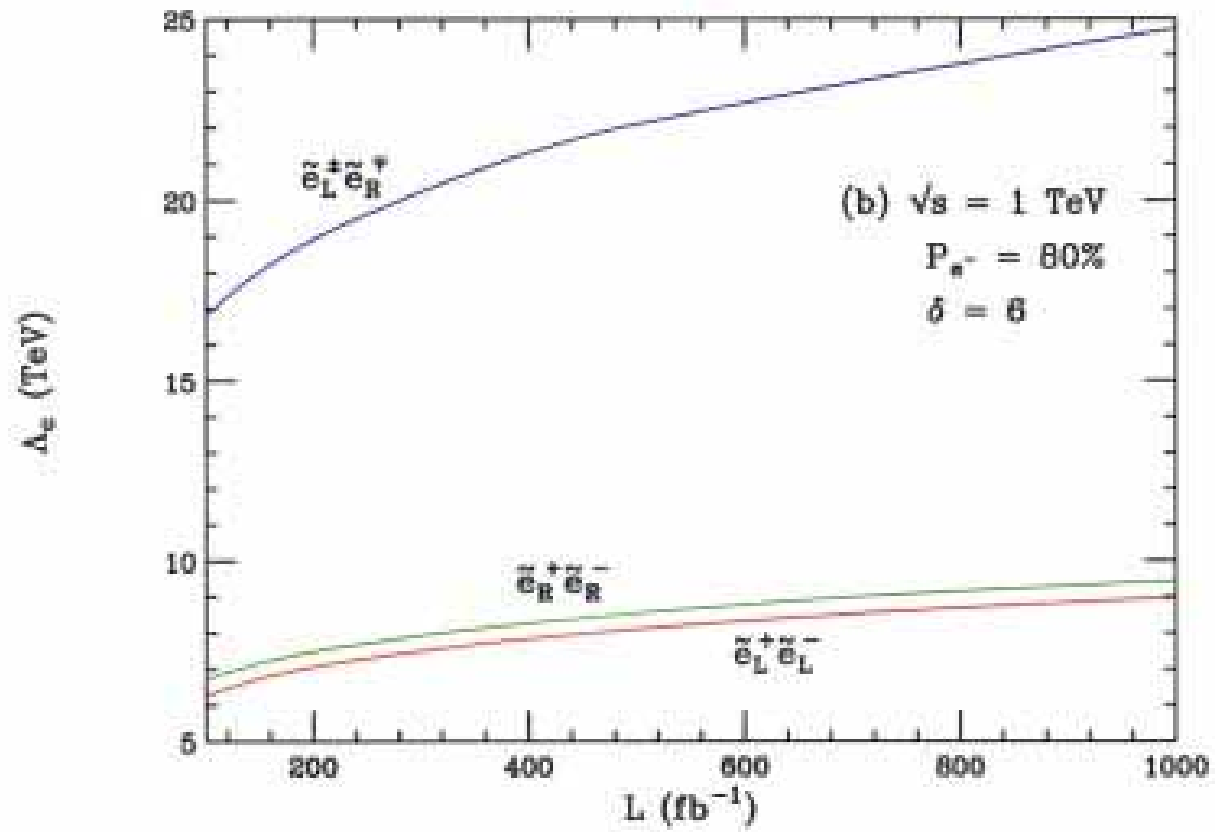


Figure 7: (a-b) Polarized binned angular distributions for each helicity configuration, taking an 80% polarization of the initial electron beam. The solid histogram represents the $D = 4$ bin-to-bin model, while the “data” points correspond to the effects of a supersymmetric bulk with $A_s = 1.5, 3.0, 6.0$ TeV from top to bottom. The $A_s = 6.0$ TeV case is only discernable in Figure (b).

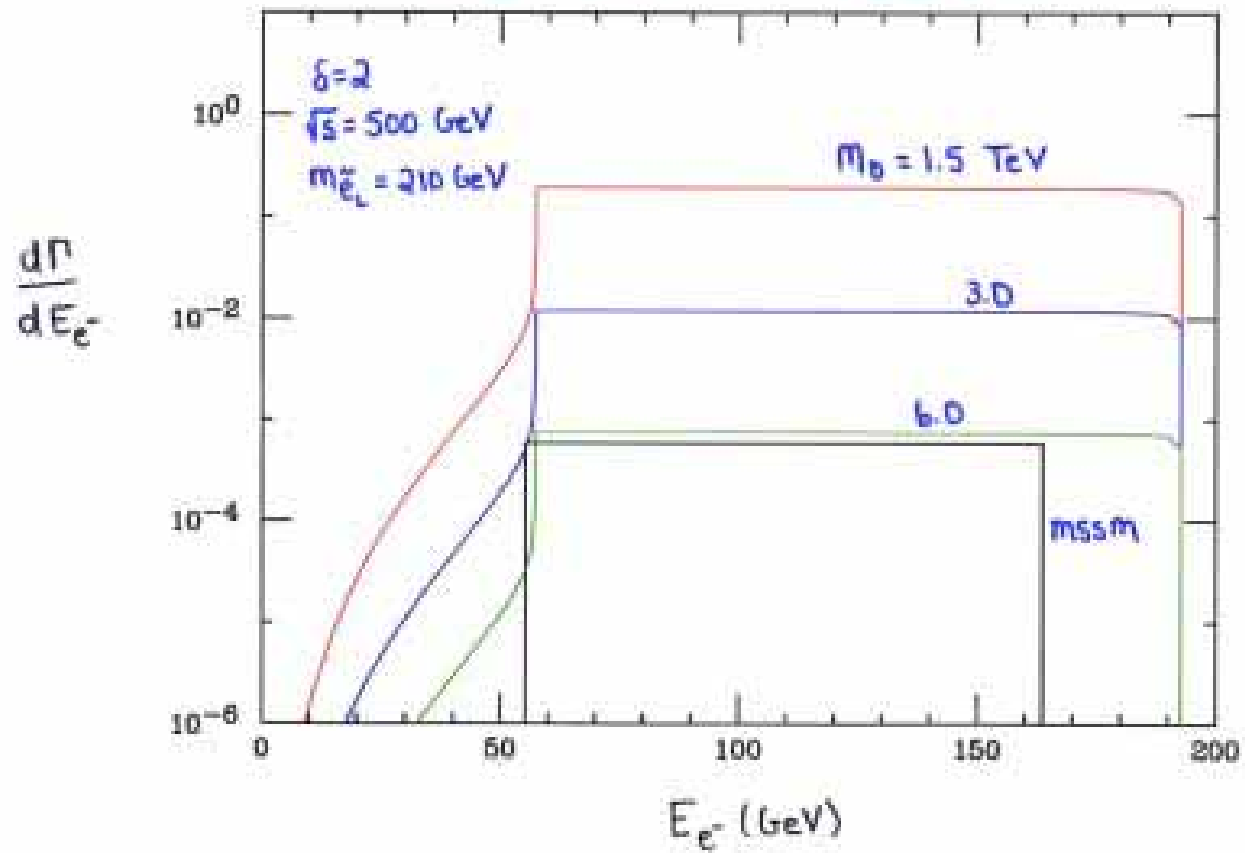
Search Reach for Λ_c



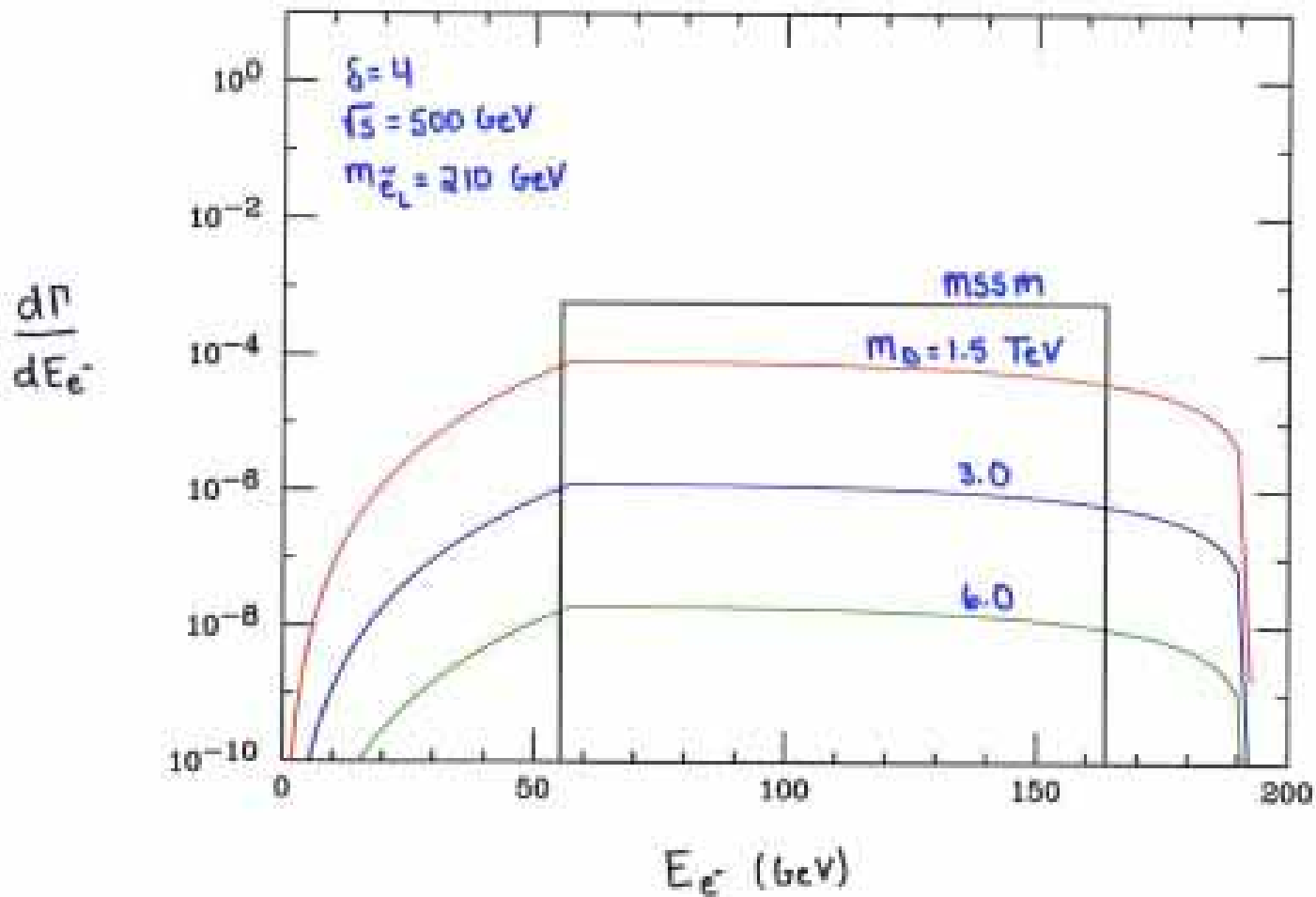
Search Reach for M_H/Λ_e



$$e^+e^- \rightarrow \tilde{e}_L^+ \tilde{e}_L^- + \tilde{e}_L \rightarrow e + \tilde{\nu}^{(e)}$$

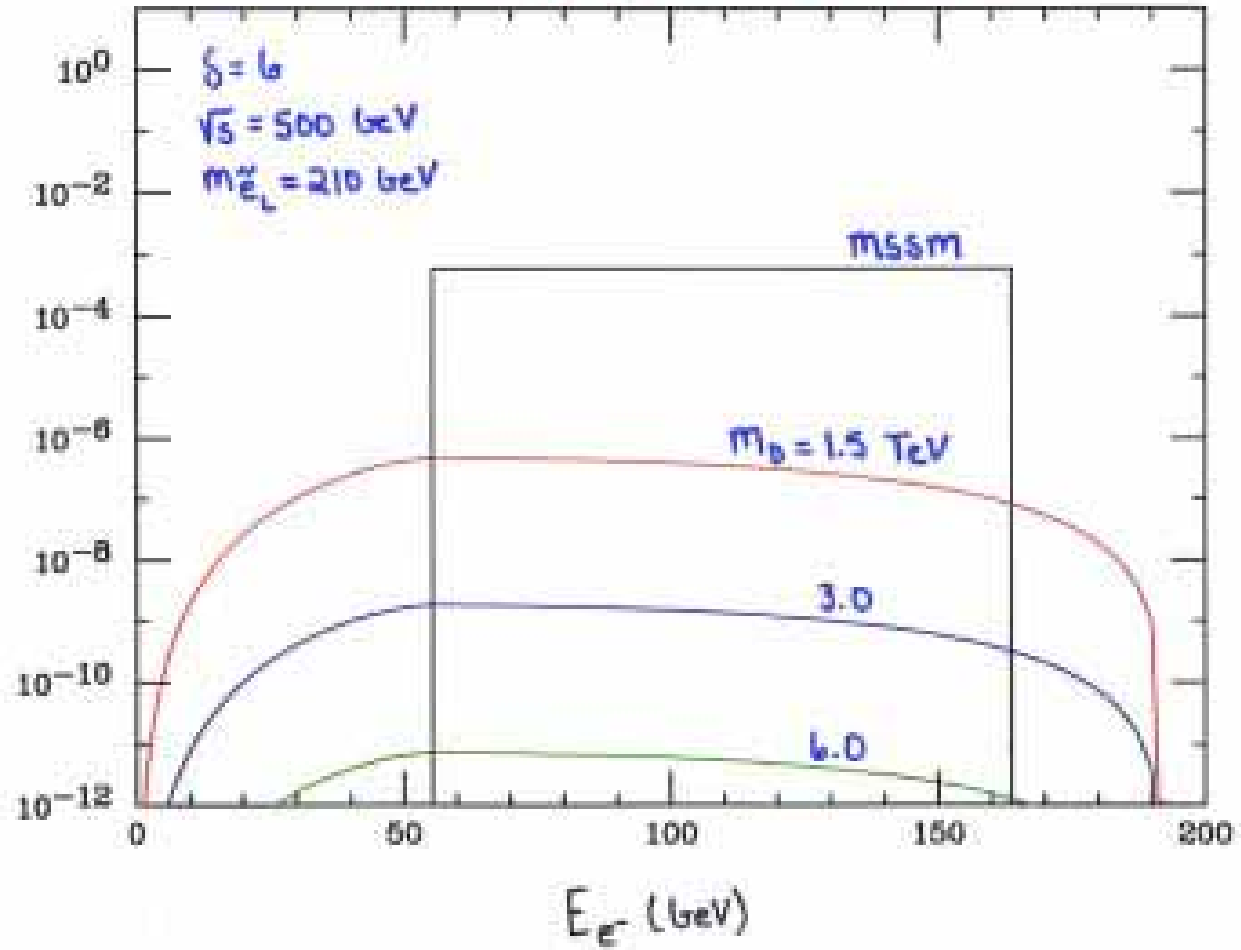


$$e^+e^- \rightarrow \tilde{e}_L^+ \tilde{e}_L^- + \tilde{e}_L \rightarrow e + \tilde{G}^{(0)}$$

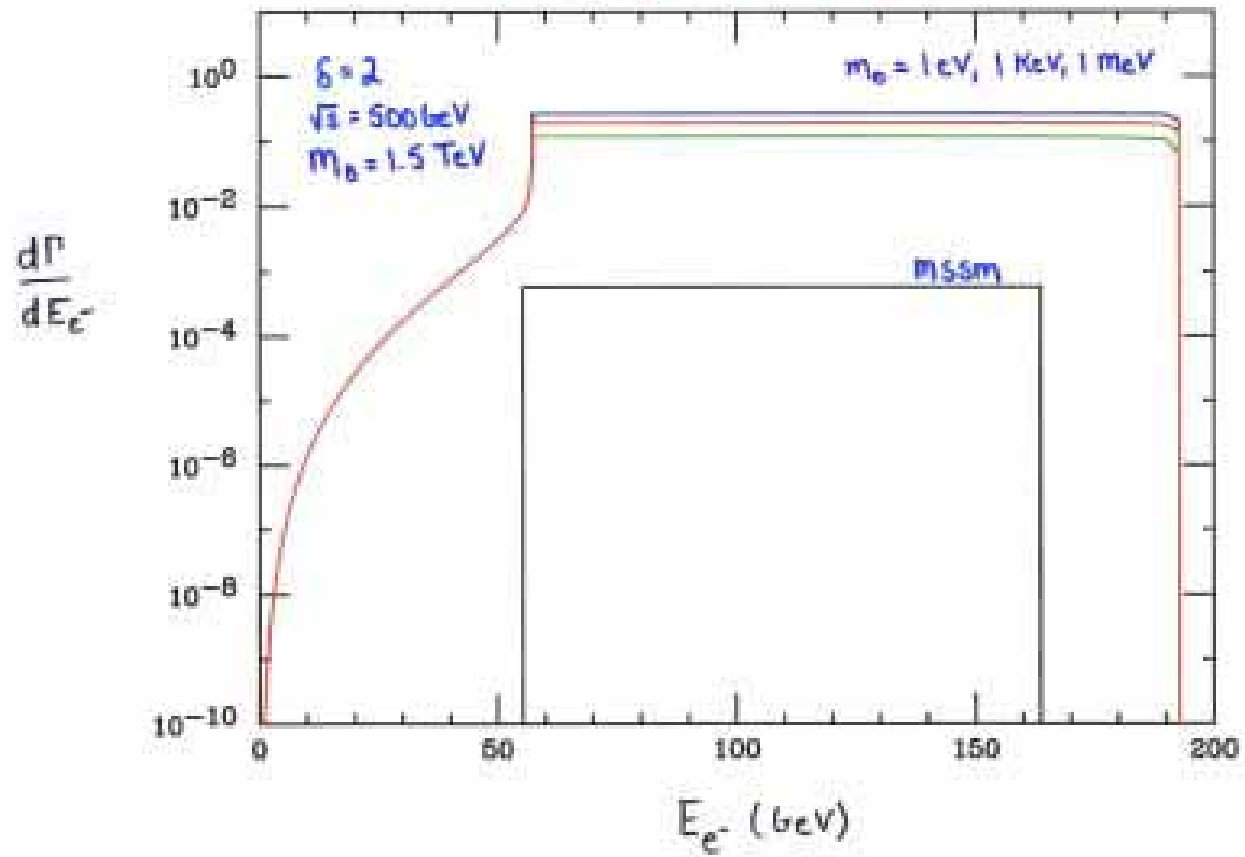


$$e^+e^- \rightarrow \tilde{e}_L^+ \tilde{e}_L^- + \tilde{e}_L \rightarrow e + \tilde{G}^{(0)}$$

$$\frac{d\Gamma}{dE_e}$$



$$e^+e^- \rightarrow \tilde{\nu}_L^* \tilde{\nu}_L^- + \tilde{\nu}_L \rightarrow e + \tilde{G}^{(H)}$$



Conclusions

- Have studied pheno of supersymmetric bulk Gravity bulk supermultiplet \rightarrow KK tower of gravitinos
- Derived effective 4-d theory
- Studied \tilde{G}_{KK} exchange in $e^+e^- \rightarrow \tilde{e}^- \tilde{e}^+$
Leading term is Dimension -6!
- Strong search reach : $\Lambda_c \gtrsim 20 \text{ TeV}$!
- Distorts \tilde{e} decay distributions : $\tilde{e} \rightarrow e + \tilde{G}^{(n)}$
- Famous 'edges' which provide precise sparticle mass measurements may not be so clean....