## A few comments to the ps workshop

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## Content

- BW, $\mathrm{S} / \mathrm{N}, \sigma_{\text {TTS }}$ of MCP-PMT
- Timing resolution limit
- Can G-APD arrays compete with MCP-PMTs for TOF ?
- SLAC test beam this year.
- Next MCP-PMT test at SLAC.
- Super-B detector in Italy.

MCP-PMT: BW, S/N and $\sigma_{T T S}$

## Hamamatsu MCP-PMT R3809U-50

## Hamamatsu data sheets



- $6 \mu \mathrm{~m}$ MCP hole diameter
- Useful photocathode dia.: 11 mm
- Rise time: ~150 ps => BW ~0.35/150ps ~2.3 GHz
- Single pixel device.
- MCP-to-anode capacitance: ~3pF
- $\sigma_{\mathrm{TTS}}=10-11 \mathrm{ps}$
- Light source jitter: ~ 5 ps (FWHM)


Hamamatsu C5594-44 amplifier 1.5 GHz BW, 63x gain


This is one of the fastest commercialy available photon detector

## Photek MCP-PMT

(J. Milnes and J. Howorth, Photek Ltd. information)

- 3.2 \& $6 \mu \mathrm{~m}$ MCP hole diameter
- Rise time: $\sim 66 \mathrm{ps}(3.2 \mu \mathrm{~m})=>\mathrm{BW} \sim 0.35 / 66 \mathrm{ps} \sim 5.3 \mathrm{GHz}$
- Rise time: ~95 ps (6 $\mu \mathrm{m})=>$ BW $\sim 0.35 / 95 \mathrm{ps} \sim 3.7 \mathrm{GHz}$
- Agilent 86100 C sampling scope $(18 \mathrm{GHz})$, average over 18 samples
- No amplifier used in this test, to my understanding
- 10 mm dia. single pixel anode
- Laser wavelength 650 nm


This is the fastest photon detector, to my knowledge

## Burle/Photonis MCP-PMT 85012

## J.V., log book \#3

## MCP-PMT 85012-501:



- $10 \mu \mathrm{~m}$ MCP hole diameter
- 64 pixel devices (use a single pad, ground the rest)
- $\mathrm{C}_{\text {Anode-to-ground }} \sim 5.5 \mathrm{pF}$
- A 1 GHz BW scope - limits the rise time
- MCP-PMT rise time: $\sqrt{ }\left\{500^{2}-350^{2}-230^{2}\right\} \sim 270 \mathrm{ps}$ $=>$ BW $_{\text {MCP-PMT }} \sim 0.35 / 270$ ps $\sim 1.3 \mathrm{GHz}$
- PiLas red laser diode ( 635 nm ):

$$
\sigma_{\mathrm{TTS}} \sim \sqrt{ }\left(32^{2}-13^{2}-11^{2}\right)=27 \mathrm{ps}
$$

Hamamatsu C5594-44 amplifier
1.5 GHz BW, 63x gain


## S/N MCP-PMT 85012

J.V., log book \#5

With a 1 GHz BW scope:
$1 \mathrm{mV} / \mathrm{div}, 1 \mathrm{~ns} /$ div


Run 386, Laser pulse, var. att. after MCP left in, $\sim 50 \mathrm{pe}^{-}$, no amplifier, Gain $\sim 1.4 \times 10^{5}$,
$10 \mathrm{mV} / \mathrm{div}, 1 \mathrm{~ns} / \mathrm{div}, 2.22 \mathrm{kV}, 4$ pads connected

$(\mathrm{S} / \mathrm{N})_{\mathrm{pp}} \sim 43 \mathrm{mV} / 0.5 \sim 80$

## Various timing schemes

## Timing strategies

## 1) High gain operation:

- Either no amplifier, or a small amplification only:
- One would expect much worse aging effects due to a high gain operation.
- Single pe- sensitivity (with an amplifier):
- In addition to the above comment, a very poor pulse recovery

2) Low gain, linear operation:

- Constant-fraction-discriminator (CFD).
- CFD + additional pulse height correction.

- A slight time-walk as number of photoelectrons corrected by the QTNT + ADC
- Waveform sampling (a'la Gary Varner's design from U. of Hawaii).
- The most powerful timing method.
- Double-threshold timing on both leading and trailing edges. $\qquad$
- Single threshold on both leading and trailing edges.
- The most simple.


## Beam setup with two MCP-PMTs and a fiber splitter



## A laser-based result with two TOF counters

(SLAC-PUB-13073, Jan. 2008)
Two detector resolution (Npe $\sim 50$ pe ea.):


Each detector has Npe~50 pe:

| $\sigma_{\text {single detector }}$ | $\sim(1 / \sqrt{ } 2) \sigma_{\text {double detector }}$ |
| ---: | :--- |
|  | $\sim 7.2 \mathrm{ps}$ |



- Two Burle/Photonis MCP-PMTs with $10 \mu \mathrm{~m}$ MCP holes operating at $2.27 \& 1.88 \mathrm{kV}$.
- Ortec 9327Amp/CFD (two) with a -10 mV threshold and a walk threshold of +5 mV \& TAC566 \& 14 bit ADC114

Can one improve the CFD timing resolution with an additional pulse height correction?

## CFD timing pulse height correction with QTNT

J.V., log book 5


## QTNT circuit problems and advantages

J.V., log book 5

## Attenuator after the amplifier:

HPK amp. output, 200,V/div. 2ns/div



- Advantage of the QTNT circuit approach is that one does not integrate the pulse wiggles.
- Disadvantage of this approach is that its linearity depends on the trailing pulse shape.


## Pulse height calibration of the QTNT electronic



## Pulse height correction of the CFD timing

$\mathrm{N}_{\mathrm{pe}} \sim 50$ pe $:$
J.V., $\log$ book 5 , Laser test with a variable attenuator: $0-6 \mathrm{~dB}$

TOF1-TOF2 - uncorrected


## Correction:



TOF1-TOF2 - corrected with LeCroy ADC


- Observe only a slight improvement of the CFD timing resolution after a pulse height correction with the QTNT circuit
- Note: The above result is slightly worse than my best earlier laser test result ( $\sigma_{\text {single detector }} \sim 7.2 \mathrm{ps}$ ) because (a) a larger dynamic range now, (b) the attenuator is left in the circuit, and (c) may be the corrections are not perfect because I did not spend enough time on this.


## What is the resolution limit ?

## Time calibration of the electronics



- One of the best electronics performance, to my knowledge.


## What resolution do we expect to get ?

- A calculation indicates $\mathbf{N}_{\mathrm{pe}} \sim 50$ for $\mathbf{1 ~ c m}$-long Fused Silica radiator \& Burle/Photonis Bialkali photocathode:

- Expected resolution (for a single pad):
a) Beam (Radiator length $=10 \mathrm{~mm}+$ window):

$$
\begin{aligned}
& \sigma \sim \sqrt{ }\left[\sigma^{2}{ }_{\text {MCPPPMT }}+\sigma^{2}{ }_{\text {Radiator }}+\sigma_{\text {Pad broadenibng }}^{2}+\sigma_{\text {Electronics }}^{2}+\ldots\right]= \\
&=\sqrt{ }\left[\left(\sigma_{\text {TTS }} / \sqrt{ } N_{\mathrm{pe}}\right)^{2}+\left(\left(\left(12000 \mu \mathrm{~m} / \cos \Theta_{\mathrm{C}}\right) /(300 \mu \mathrm{~m} / \mathrm{ps}) / \mathrm{n}_{\text {group }}\right) / \sqrt{ }(12 \mathrm{Npe})\right)^{2}+\right. \\
&\left.+((6000 \mu \mathrm{~m} / 300 \mu \mathrm{~m} / \mathrm{ps}) / \sqrt{ }(12 \mathrm{Npe}))^{2}+(3.42 \mathrm{ps})^{2}\right] \sim \\
& \sim \sqrt{ }\left[3.8^{2}+3.3^{2}+0.75^{2}+3.42^{2}\right] \sim 6.1 \mathrm{ps}
\end{aligned}
$$

b) Laser $\left(\mathrm{N}_{\mathrm{pe}} \sim 50 \mathrm{pe}^{-}\right)$: $\qquad$

$$
\begin{aligned}
\sigma & \sim \sqrt{ }\left[\sigma^{2}{ }_{\text {MCP.PMT }}+\sigma_{\text {Laser }}^{2}+\sigma_{\text {Electronics }}^{2}+\ldots\right]= \\
& \left.=\sqrt{ }\left[\sigma_{\text {TTST }} / \sqrt{ } \mathrm{N}_{\mathrm{pe}}\right)^{2}+\sqrt{ }\left((\mathrm{FWHM} / 2.35) / \sqrt{ } \mathrm{N}_{\mathrm{pe}}\right)^{2}+(3.42 \mathrm{ps})^{2}\right] \sim \\
& \sim \sqrt{ }\left[3.8^{2}+1.8^{2}+3.42^{2}\right] \sim 5.4 \mathrm{ps}
\end{aligned}
$$

SLAC test with Burle MCP-PMT, $10 \mu \mathrm{~m}: \quad \sigma_{\text {TTS }}(1 \mathrm{pad}) \sim 27 \mathrm{ps}$ (my data) SLAC test with Burle MCP-PMT, $10 \mu \mathrm{~m}: \quad \sigma_{\text {TTS }}$ (4 pads shorted) $\sim 40-50 \mathrm{ps}$ (my data) Nagoya test with HPK R3809U-50, $6 \mu \mathrm{~m}$ : $\sigma_{\text {TTS }} \sim \mathbf{1 0 - 1 1} \mathrm{ps}$ (Hamamatsu data) E-407 G-APD array: $\sigma_{\text {TTS }} \sim 100$ ps (Krizan's data for blue wavelength)


## Why G-APD does not compete with MCP-PMT at present?

## G-APD array $\sigma_{\text {TTS }}$

Measurements by Krizan's group

## Time resolution: blue vs red



|  |  | E 407 | S 137 | H100C | H050C |
| :---: | :---: | :---: | :---: | :---: | :---: |
| 635 nm: | H025C |  |  |  |  |
| 404 nm: $:$ | $\sigma_{\text {red }}(\mathrm{ps})$ | 127 | 182 | 145 | 212 |
| $\sigma_{\text {blue }}(\mathrm{ps})$ | $\mathbf{9 7}$ | $\mathbf{1 5 1}$ | $\mathbf{1 3 6}$ | $\mathbf{3 5 8}$ | $\mathbf{1 3 5}$ |

Hamamatsu
A possible TOF application:

- These G-APD arrays are not as good as the best single cell G-APDs.
- What is the reason? The passive quenching, or technology ?



## Single cell G-APD $\sigma_{\text {TTS }}$

Measurements by id Quantique


## TTS:



| Photocathode | Si |
| :---: | :---: |
| $\sigma_{\text {TTS }}$ | $\sim \mathbf{1 7} \mathrm{ps}$ |
| Low noise | $<20 \mathrm{~Hz}$ |
| Spectral range | $\mathbf{3 5 0 - 9 0 0} \mathrm{nm}$ |
| After pulsing probability | $<\mathbf{3 \%}$ |
| Dead time | $\sim \mathbf{5 0} \mathrm{ns}$ |
| Maximum rate | $\sim \mathbf{2 0 ~ M H z}$ |
| Active area | $\sim \mathbf{5 0} \boldsymbol{\mu m}$ |

- A a single G-APD cell id-100 is made by "id Quantique SA", Switzerland


## Single cell G-APD $\sigma_{\text {TTS }}$

Measurements by J.Va' vra, G-APD from Sopko, active quenching from Prochazka, CVUT Prague

G-APD:


Single photoelectron timing resolution:


- G-APD operating in a Geiger mode with active quenching and temperature control.
- With a PiLas red $(\lambda=635 \mathrm{~nm})$ laser diode operating in the single pe mode, I obtained:



## Next slide was added after the workshop

## Single pad vs. 4-pads shorted together

(this addendum is a response to a question Krzysztof asked me one day after the workshop)

## Single pad:

Run 213, single pad, 9327 mon, $2.33 \mathrm{kV}, \sim 50 \mathrm{pe}$, TOF1 Ortec 9327 amp out, $10 \mathrm{mV} / \mathrm{div}, 400 \mathrm{ps} /$ div


4-pads shorted together:


## 4-pads shorted together:

Run 261, four pads shorted, $9327 \mathrm{mon}, 2.33 \mathrm{kV}, \sim 50 \mathrm{pe}{ }^{-}$, TOF1


- A single pad capacitance to ground is $\sim 5.5 \mathrm{pF}$, 4 -pads shorted together represent $\sim 20 \mathrm{pF}$. Therefore, one would expect a degradation in rise time. Indeed that is what one observes.
- In addition, in the resolution plot, extrapolating towards Npe $\sim 1$, one observes $\sigma_{\text {TTS }}$ (4-pads) $\sim 40-50 \mathrm{ps}$, which is worse than $\sigma_{\text {TTS }}(1-$ pads $) \sim 27$ ps, shown on slide 6.
- The two conditions were the same except the pad connections


## SLAC test beam this year

## End Station A (ESA) beam line



- Configuration during the last FDIRC test.
- We use it as a secondary beam running electrons off the Be target.
- Use correctors XCOR32 \& YCOR33 to move the beam at our test end.


## Running conditions at present

- The ESA secondary electron test beam momentum was set to $10 \mathrm{GeV} / \mathrm{c}$, with LCLS beam energy of $14.5 \mathrm{GeV} / \mathrm{c}$. Previously we were always running $28 \mathrm{GeV} / \mathrm{c}$ primary electron beam. Until this run it was not clear (a) if it is even possible to run parasiting with the LCLS operation, and (b) if the particle yield is sufficient at such a low LINAC energy. We proved that it is possible, and that we get a good rate of $0.2-0.3 \mathrm{ppp}$ with a momentum byte of +$\mathbf{0 . 2 \%}$, that we get a good beam spot of $\sigma=\mathbf{1 - 2} \mathbf{~ m m}$ at far end of ESA, and good cleanliness judging from the lead glass spectrum.
- Monitoring of the primary LINAC beam by MCC operators:
- Monitoring the primary beam - pick up electrodes on Be target $\longrightarrow$
- Plus a usual LINAC monitoring
- Monitoring of the secondary beam by MCC operators:
- Scalars indicating a particle flux going through our test

- Our own monitoring histograms, such beam spots in the hodoscopes, lead glass spectra, rates, etc.
- Particle yields:
- The following graph shows the yield for positive polarity $\longrightarrow$
- For negative beam one gets mostly electrons
- Generally MCC people encourage negative beam polarity i.e., one needs some political umpf to re-cable magnets into the positive polarity and spend a week to tune the beam; the last test to do this was the Glast experiment.


## Summary of beam parameters

a) Height from the floor:
b) Total left-right clearance:
c) Beam spot size at the bar entrance:
d) Beam position knowledge:
e) Beam divergence:
f) Particle type:
g) Polarity:
g) Particle flux during the test:
h) Rate:
i) Secondary beam momentum:
j) Primary beam momentum:
k) Target:

1) Production angle:
m) Timing signals:

7 feet $\pm 2$ inches
$>10$ meters
$\sim 1 \mathrm{~mm}$
$\pm 1 \mathrm{~mm}$ (after tuning correctors)
$\pm \sim 0.6 \mathrm{mrad}$ (based on the hodoscopes)
Electrons (mostly)
Negative (typically)
~ 0.2 ppp
$10-30 \mathrm{~Hz}$
$10 \mathrm{GeV} / \mathrm{c}$
$14.5 \mathrm{GeV} / \mathrm{c}$ (when LCLS is controlling)
~1 ft-long 0.3 r.l. Beryllium
$0.5^{0}$
AB01-8-09, AB01-8-10 (programmable)

## Latest 2007 Beam Test Setup

- Instrumentation available:
- 2 x-y scint. fiber hodoscopes
- START \#1 counter to monitor flux
- Time start from the LINAC RF signal, but correctable with a local START \#1 counter
- Lead glass to monitor beam multiplicity (very important in SLAC's beam)


Hodoscopes \#1\&2 (scint. fibers)


## Next MCP-PMT tube to test at SLAC in the beam

## Existing MCP-PMTs with improvements - May test



- Pulse height-corrected CFD timing (with the QTNT scheme)


## Burle/Photonis MCP-PMT 85014 - for July SLAC tests



- $\quad \sim 6.9 \mathrm{~mm}$ thick quartz radiator
- Cathode-to-MCP gap: ~864 $\mu \mathrm{m}$
- MCP-to-Anode gap: $\sim 2.7 \mathrm{~mm}$
- Charge spread on ~16 pads => no more suitable for a few channel electronics test.
- Gary's electronics (need 16 channels)


## Super-B detector in Italy

## PID systems in Super-B



- Converging on two PID systems:

Barrel Focusing DIRC \& Forward TOF

## Timing at a level of $\sigma \sim 15 \mathrm{ps}$ can start competing with the RICH techniques

Example of various Super-B factory PID designs:

Calculation done for Flight Path
Length $=2 \mathrm{~m}$


- The PID performance of a forward TOF system with $\sigma \sim \mathbf{1 5}-20 \mathrm{ps}$ is equivalent to the PID performance of the BaBar DIRC.
- Adding a TOF system would improve the hermeticity of the PID coverage.

