# Measurement of $|V_{ub}|$

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Inclusive approaches (endpoint,  $M_X$ ,  $q^2$ ) Exclusive approaches ( $B \rightarrow \pi 1 \nu$ ,  $B \rightarrow \rho 1 \nu$ ) Conclusion



# The V<sub>ub</sub> element of the CKM matrix $\begin{pmatrix} Vud & Vus & Vub \\ Vcd & Vcs & Vcb \\ Vtd & Vts & Vtb \end{pmatrix} = \begin{pmatrix} 1-\lambda^2/2 & \lambda & A\lambda^3(\rho-i\eta) \\ -\lambda & 1-\lambda^2/2 & A\lambda^2 \\ A\lambda^3(1-\rho-i\eta) & -A\lambda^2 & 1 \end{pmatrix}$

 $|V_{ub}|$  determines a circle of radius<sup>2</sup> =  $\rho^2 + \eta^2$  for the apex of the Bjorken triangle. *Very important for indirect constraints on the CKM triangle and for detecting New Physics.* 



# How are $|V_{ub}|$ measurements different from $|V_{cb}|$ measurements?

Cannot observe the whole spectrum of  $b \rightarrow u \ v$  unlike  $b \rightarrow c \ v$ . Backgrounds are too large hence restricted to small portions of phase space.

Heavy  $\rightarrow$ light FF not Heavy  $\rightarrow$ Heavv FF: very little simplification from HQET. W or  $q^2$ range is much larger (dependence on FF is much greater).

# $|V_{ub}|$ from the Lepton Endpoint



A large theoretical extrapolation is required to obtain |Vub|

#### CLEO 2001: $|V_{ub}|$ from leptons beyond the b $\rightarrow$ c endpoint



#### BABAR 2002: $|V_{ub}|$ from leptons beyond the endpoint



Agrees well with CLEO  $(0.143\pm0.010\pm0.014) \times 10^{-3}$ 

### Model Dependence in $|V_{ub}|$ from inclusive decay.

#### CLEO 1993: Model dependence

TABLE IV. Partial branching fractions and corresponding values of  $|V_{ub}/V_{cb}|^2$  and  $|V_{ub}/V_{cb}|$  for the strict-cut analysis in the momentum interval 2.3 to 2.6 GeV/c.

Model	$10^6 \Delta B_{ub}(p)$	$10^2  V_{ub}/V_{cb} ^2$	$\left V_{ub}/V_{cb} ight $
ISGW	$121 \pm 17 \pm 15$	$1.02 \pm 0.20$	$0.101 \pm 0.010$
KS	$115 \pm 16 \pm 15$	$0.31\pm0.06$	$0.056 \pm 0.006$
WSB	$122 \pm 17 \pm 16$	$0.53\pm0.11$	$0.073 \pm 0.007$
ACCMM	$154 \pm 22 \pm 20$	$0.57 \pm 0.11$	$0.076 \pm 0.008$

CLEO 2001, BABAR 2002: No longer use a model for extrapolation. Instead rely on the  $b \rightarrow s \gamma$  shape function..

(Discussed in lectures by Ligeti)  $B \rightarrow \text{light quark shape function, SAME (to lowest order in <math>\Lambda_{OCD}/m_b$ ) for  $b \rightarrow s \gamma$  ( $B \rightarrow X_s \gamma$ ) and  $b \rightarrow u \ln (B \rightarrow X_u \ln)$ .





#### Convolute with light cone shape function.



Fraction of  $b \rightarrow u\ell v$ spectrum above 2.2 is  $0.13 \pm 0.03$ 



Extrapolation from endpoint region using  $b \rightarrow s \gamma$  *Idea: Use the shape function in*  $b \rightarrow s \gamma$  *to determine the*  $b \rightarrow u l v$  *shape function. Then get the fraction of events in the*  $b \rightarrow u l v$ *endpoint region.* 



# $V_{ub}$ Extrapolation from endpoint region using b $\rightarrow$ s $\gamma$

TABLE II: Results for five momentum intervals. Uncertainties on yields,  $f_u$ , and branching fractions are statistical and systematic. The first uncertainty on the total branching fraction is from the measurement of  $\Delta \mathcal{B}_u(p)$  and the second is from  $f_u$ . The first two uncertainties on  $|V_{ub}|$  are from the branching fraction and the third and fourth are from theory.

$p \; (\text{GeV}/c)$	Yield	$\Delta \mathcal{B}_u(p)(10^{-4})$	$f_u$	$\mathcal{B}(B \to X_u \ell \nu) \ (10^{-3})$	$ V_{ub} (10^{-3})$
2.0-2.6	${\bf 3538 \pm 279 \pm 1470}$	$4.22 \pm 0.33 \pm 1.78$	$0.266 \pm 0.041 \pm 0.024$	$1.59 \pm 0.68 \pm 0.28$	$3.87 \pm 0.83 \pm 0.35 \pm 0.15 \pm 0.12$
2.1 - 2.6	$2751\pm191\pm584$	$3.28 \pm 0.23 \pm 0.73$	$0.198 \pm 0.035 \pm 0.020$	$1.66 \pm 0.39 \pm 0.34$	$3.95 \pm 0.46 \pm 0.40 \pm 0.16 \pm 0.16$
2.2 - 2.6	$1901\pm122\pm256$	$2.30 \pm 0.15 \pm 0.35$	$0.130 \pm 0.024 \pm 0.015$	$1.77 \pm 0.29 \pm 0.38$	$4.08\pm0.34\pm0.44\pm0.16\pm0.24$
2.3 - 2.6	$1152\pm80\pm61$	$1.43 \pm 0.10 \pm 0.13$	$0.074 \pm 0.014 \pm 0.009$	$1.94 \pm 0.22 \pm 0.43$	$4.27 \pm 0.24 \pm 0.47 \pm 0.17 \pm 0.34$
2.4 - 2.6	$499\pm57\pm14$	$0.64 \pm 0.07 \pm 0.05$	$0.037 \pm 0.007 \pm 0.003$	$1.74\pm0.24\pm0.38$	$4.05\pm0.28\pm0.45\pm0.16\pm0.45$

Optimal interval is 
$$2.2 < p_L < 2.6 \text{ GeV}$$

 $BF(2.2-2.6 \text{ GeV}) = (2.30\pm0.15\pm0.35) \times 10^{-4}$ 

 $f_u (2.2-2.6 \text{ GeV}) = 0.130 \pm 0.024 \pm 0.05$ 

$$|V_{ub}| = (3.07 \pm 0.12) \times 10^{-3} \times \left[\frac{\mathcal{B}(B \to X_u e\nu)}{0.001} \frac{1.6 \text{ ps}}{\tau_B}\right]^{\frac{1}{2}}$$

#### CLEO: $|V_{ub}|$ from Lepton Endpoint (using $b \rightarrow s\gamma$ )

 $|Vub| = (4.08 \pm 0.34 \pm 0.44 \pm 0.16 \pm 0.24) \times 10^{-3}$ 

stat  $b \rightarrow s \gamma$   $\Gamma$  theory  $\Lambda/M_B$  theory



Subleading
 corrections large
 C. Bauer, M. Luke, T. Mannel
 A. Leibovich, Z. Ligeti, M. Wise

 Method for partial inclusion of subleading corrections: Neubert

Quark-hadron duality ?

How can we improve our knowledge of  $|V_{ub}|$  from inclusive decays ?

Use all of phase space (LEP)

Use favorable regions of  $M_X$  or  $q^2$ (DELPHI, CLEO) and new theoretical strategies

Fraction with  $E_L > 2.2 \text{ GeV} (\sim 15\%)$ ; fraction with  $M_X < M_D (\sim 70\%)$ , fraction with  $q^2 > (M_B - M_D)^2 (\sim 20\%)$ 

## ALEPH $|V_{ub}|$ Measurement





20 kinematic variables in a NN

#### OPAL $|V_{ub}|$ Measurement



Huge background suppressed with 7 variable Neural Net Small signal extraction depends on  $b \rightarrow c\ell\nu$  model! S/B = 0.05

## Published Inclusive |Vub| Determinations

			., –		
Exp.	Method	S/B	$ V_{ub}  \ [10^{-3}]$	$\sigma_{b \to c} \\ ( V_{ub} )$	$\sigma_{th} \ ( V_{ub} )$
ALEPH	Neural Net	0.07	$4.12 \pm .67 \pm .62 \pm 0.35$	15%	9%
OPAL	Neural Net	0.05	$4.00 \pm .71 \pm .59 \pm 0.40$	15%	10%
DELPHI	$M_X$	0.10	$4.07 \pm .65 \pm .47 \pm 0.39$	12%	10%
L3	$\pi - \ell \operatorname{Cut}$	0.22	$5.7 \pm 1.0 \pm 1.3 \pm 0.5$	22%	10%
LEP	Average		$4.09 \pm 0.37 \pm 0.44 \pm 0.34$		9–15%
CLEO	$E_\ell$ endpoint	0.39	$4.12 \pm 0.34 \pm 0.44 \pm 0.33$	7%	10–15%

#### Battaglia and Gibbons

LEP



Hadronic Invariant Mass Spectrum for b→u Decay

- singularity is smeared out by b quark light-cone distribution function  $f(k_+)$
- rate is sensitive to details of  $f(k_+)$  unless  $m_X^2 >> \Lambda_{QCD} m_b$  (bad for  $m_X < m_D!$ ) introduces model dependence unless we know  $f(k_+)$

#### DELPHI 2000: Analysis with $M_X < 1.6 \text{ GeV}$



 $|V_{ub}| = (4.07 \pm 0.65(exp) \pm 0.47(b \rightarrow c) \pm 0.39(theo)) \times 10^{-3}$ 

#### Lepton Invariant Mass Spectrum for b u Decay



Lepton q<sup>2</sup> spectrum is insensitive to Fermi motion (and has less model dependence.)

	Representative cuts:			
	(a) $q^2 > 6 \text{ GeV}^2$ , $m_X <$	$m_D$	46% of rate	
M.Luke:	(b) $q^2 > 8 \text{ GeV}^2$ , $m_X <$	1.7 GeV	33% of rate	
	(c) $q^2 > 11 \text{ GeV}^2, m_X <$	1.5 GeV	18% of rate	
Uncertaint	y Size (in $V_{ub}$ )	Improven	nent?	
$\Delta m_b$	±80 MeV: 7%, 8%, 10%	RG improv spectra, lat	ved $\Upsilon$ sum rules, moments of <i>B</i> d ttice	lecay
	3%, 3%, 4%			
$lpha_{_S}$	2%, 3%, 7%	full two-lo	oop calculation	
$1/m_b^3$ (weak annihi	3%, 4%, 8% lation)	compare <i>B</i> compare S	$B^{\pm}$ , $B^0$ S.L. width of $D^0$ , $D_S$ , lattice	

#### *CLEO:* $B \rightarrow X I_V$ with Neutrino Reconstruction



 $|V_{ub}| = (4.05 \pm 0.18 \pm 0.58 \pm 0.25 \pm 0.21 \pm 0.56) \ 10^{-3}$ stat sys b > c b > u theory

 $|V_{ub}|$  from M<sub>X</sub> or q<sup>2</sup> with fully reconstructed B tags (MC)



FIG. 1: (a) Hadronic mass  $(m_{had})$  distribution for 1000  $fb^{-1}$  data found with CLEO III fast MC. The solid histogram is the  $m_{had}$  distribution of  $b \to u\ell\nu$ , and the dashed histogram is the  $m_{had}$  distribution of  $b \to c\ell\nu$ . (b)  $q^2$  distribution for 1000  $fb^{-1}$  data found with CLEO III fast MC. The solid histogram is the  $q^2$  distribution of  $b \to u\ell\nu$ , and the dashed histogram is the  $q^2$  distribution of  $b \to c\ell\nu$ .

		$m_{ m had}$				$q^2$					
year	$\mathcal{L}_{int}(fb^{-1})$	S	B	$\delta \mathrm{V}^{\mathrm{e}}_{\mathrm{u}}$	xpt. .b (	(%)	S	B	$\delta \mathrm{V}^{\mathrm{e}}_{\mathrm{u}}$	хрt. .b (	(%)
				stat.	sys.	$\mathbf{tot.}$			stat.	sys.	tot.
2002	100	335	127	3.2	2.2	3.9	127	7	4.6	3.0	5.5
2005	500	1675	635	1.5	1.5	2.1	635	36	2.0	1.2	2.3
2010	2000	6700	2540	0.7	1.5	1.7	2538	144	1.0	1.2	1.6

Shipsey andLee

# Exclusive Approaches to $|V_{ub}|$

*Measure*  $BF(B \rightarrow \pi l v)$ ,  $BF(B \rightarrow \rho l v)$  or  $BF(B \rightarrow \omega l v)$ .

With more statistics, can then measure  $d\Gamma/dq^2$  ( $B \rightarrow \pi l v$ ) or even form factors for  $BF(B \rightarrow \rho l v)$ .

A key experimental ingredient is the use of detector hermiticity to deduce the v momentum

Variables for v reconstruction of exclusive semileptonic decays (used for  $B \rightarrow \pi(\rho) l v$ )

$$p_{miss} = -\sum p_{i}$$

$$E_{miss} = 2 E_{beam} - \sum E_{i}$$

$$M^{2}_{miss} = E^{2}_{miss} - P^{2}_{miss}$$

$$p_{v} = (p_{miss}, |p_{miss}|)$$

$$\Delta E \equiv E_{beam} - (E_{\pi} + E_l + E_{\nu})$$
$$M_B \equiv \sqrt{E^2_{beam} - |p_{\pi} + p_l + p_{\nu}|^2}$$

where Ebeam = 5.29 GeV

#### Detector hermiticity requirements (cont'd)

e.g. Only one lepton and  $|Q_{tot}| = \pm 2$ . Also require  $|M_{miss}|^2 < 3.0 \text{ GeV}^2$ 

Fiducial cut on  $p_{miss}$  is important at the B factories. E.g.  $17^0 < \theta_{miss} < 150^0$ 



# Belle: $B^0 \rightarrow \pi^+ l^- \nu$ signal



 $\Delta E$ 

Signal is extracted from a 2D fit to  $\Delta E$  and  $p_L$ 

# Determination of $|V_{ub}|$ from BF(B<sup>0</sup> $\rightarrow \pi 1 \nu$ )

Models or lattice calculations are needed to determine detection efficiency as well as convert the BF into a value of  $|V_{ub}|$ 

		(Khodjamirian <i>et al</i> .)
model	UKQCD	LCSR
Reference	PLB 486, 111 (2000)	PRD 62, 114002 (2000)
good for	large $q^2$	small $q^2$
$\gamma_{\pi}$	$9^{+3+2}_{-2-2}$	$7.3 \pm 2.5$
effi. (%)	2.9	3.1
$\mathcal{B}(B^0 \to \pi^- \ell^+ \nu_\ell)$	$(1.35\pm0.11\pm0.21)\times10^{-4}$	$(1.31\pm0.11\pm0.20)\times10^{-4}$
$ V_{ub} $	$(3.11 \pm 0.13 \pm 0.24 \pm 0.56) \times 10^{-3}$	$(3.58 \pm 0.15 \pm 0.28 \pm 0.63) \times 10^{-3}$





# CLEO:Determination of $|V_{ub}|$ from BF(B<sup>0</sup> $\rightarrow \pi l \nu$ )



 $q^2$  distribution of  $B \rightarrow \pi l V$ 



#### BABAR B $\rightarrow \rho^{0(+)}$ l v signal

LOLEP



Figure 3: Continuum-subtracted projections of the ISGW2 fit result for the  $B^+ \rightarrow \rho^0 e^+$ channels in the LOLEP and HILEP electron energy regions; the contributions are the direc and crossfeed components of the signal (unhatched region, above and below the dashed line respectively); the background from  $b \to ue\nu$  other than  $B \to \rho e\nu$  and  $B \to \omega e\nu$  modes (double hatched region); the background from  $b \rightarrow ce\nu$  and other backgrounds (single-hatched region)

#### BABAR B $\rightarrow \rho^{0(+)}$ l v signal



HILEP

2.3<p<sub>1</sub> <2.7 GeV

#### LOLEP 2.0<p<sub>L</sub><2.3 GeV

Table 2: Summary of data yields for the  $B^0 \to \rho^- e^+ \nu$  and  $B^+ \to \rho^0 e^+ \nu$  modes with electron energies between 2.3 and 2.7 GeV (HILEP), and between 2.0 and 2.3 GeV (LOLEP). The yields presented in this table were obtained using the ISGW2 form-factor. The downfeed background includes all  $B \to X_u e\nu$  modes except for  $\rho$ ,  $\omega$ , and  $\pi$ . The crossfeed signal contribution corresponds to events from the other signal modes with  $\rho^0$ ,  $\omega$ , or  $\pi$  and is constrained to the signal in the fit. All errors are statistical only.

	B <sup>0</sup> -	$\rightarrow \rho^- e^+ \nu$	$B^+ \rightarrow \rho^0 e^+ \nu$		
	HILEP	LOLEP	HILEP	LOLEP	
On-resonance yield	2302	39349	2213	40155	
Direct signal	$510 \pm 63$	$718 \pm 89$	$324 \pm 40$	$440 \pm 55$	
Crossfeed signal	$262 \pm 32$	$538\pm73$	$363 \pm 42$	$725\pm86$	
$\mathbf{Downfeed}$	$203 \pm 55$	$2278 \pm 403$	$226 \pm 92$	$2435\pm430$	
b  ightarrow ce  u	$414 \pm 5$	$33859 \pm 438$	$367\pm5$	$34366\pm458$	
$e^+e^-  ightarrow q ar q$	$917 \pm 73$	$1928\pm106$	$912\pm73$	$2063 \pm 110$	
Fake electrons	$12 \pm 3$	$80 \pm 9$	$18 \pm 4$	$76 \pm 9$	

# Model dependence in $|V_{ub}|$ BABAR B $\rightarrow \rho l v$ signal



Does it make sense to take the average of models ?

How can we improve our knowledge of  $V_{ub}$  from exclusive decays ?

A considerable amount of model dependence is due to FF uncertainties.



Lattice Calculations of  $B \rightarrow \pi l \nu$  Form Factor



(from A. Kronfeld, hep-ph/0010074)

Need to measure  $d\Gamma/dq^2$  for  $B \rightarrow \pi l v$  at high  $q^2/low p_{\pi}$ 

#### Future Improvements in $|V_{ub}|$ from the Lattice

$$|V_{ub}|^{2} = \frac{12\pi^{2}}{G_{F}^{2}m_{B}} \frac{1}{T_{B}(p_{\min}, p_{\max})} \int_{p_{\min}}^{p_{\max}} dp \frac{d\Gamma_{B \to \pi\ell\nu}}{dp}$$



 $\Delta V_{ub} \approx 15-18\%$  + quenching error

(A. El-Khadra et. al., PRD64, 014502)

To understand these errors: see Lectures by P. Lepage

#### MC simulation of the $B \rightarrow \rho l \nu$ Dalitz plot



#### Model dependence in $B \rightarrow \rho l \nu$ form factors



A tight  $p_L$  cut makes it difficult to distinguish models.

FF model	$\tilde{\Gamma}_{\rm thy}~({\rm ps}^{-1})$	$\Gamma(E_\ell>2.3~{ m GeV})/\Gamma~(\%)$	$\Gamma(2.0 < E_{\ell} < 2.3 \text{ GeV})/\Gamma$ (%)
ISGW2	14.2	35	33
LCSR	16.9	24	28
UKQCD	16.5	27	30
Wise/Ligeti+E791	19.4	31	34
Beyer/Melikhov	16.0	27	30

#### Model dependence of $B \rightarrow \rho l \nu$ form factors



FIG. 4.  $d\Gamma/dq^2$  distributions for each of the three terms in Eq. 9: (a) the terms proportional to  $|H_-|^2$  and  $|H_+|^2$  and (b) the  $|H_0|^2$  term.

#### BELLE: $B^0 \rightarrow \omega e^+ \nu$ signal selection

Lattice calcs of  $B \rightarrow \rho$  FFs cannot handle the finite width of the  $\rho$ 

Less exp problem with  $B \rightarrow \pi \pi 1 \nu$ 



Belle-CONF-0242

#### BELLE: $B^0 \rightarrow \omega e^+ \nu$ signal

N(events) in the signal region with  $0.76 < m(3\pi) < 0.81$ 

$$222\pm15 \text{ (total)}$$

$$48\pm10 (b \rightarrow c)$$

$$2\pm2 \text{ (fake)}$$

$$47\pm21 \text{ (cont.)}$$

$$MC \text{ est.}$$

•Excess in  $m(3\pi)$ after side-band subtraction = 59 ± 15 events



# Summary of recent $|V_{ub}|$ determinations



Review by Battaglia

# Conclusions

An improved method for  $|V_{ub}|$  determination using leptons in the endpoint region has been introduced. The uncertainty in the extrapolation is reduced by using the shape function measured in  $b \rightarrow s \gamma$ .

An inclusive method using optimized cuts on  $q^2$  and  $M_X$  appears promising.

Prospects for improved  $|V_{ub}|$  in  $B \rightarrow \pi l v$  using high statistics measurements of  $d\Gamma/dq^2$  and FFs determined from the lattice appear good.