

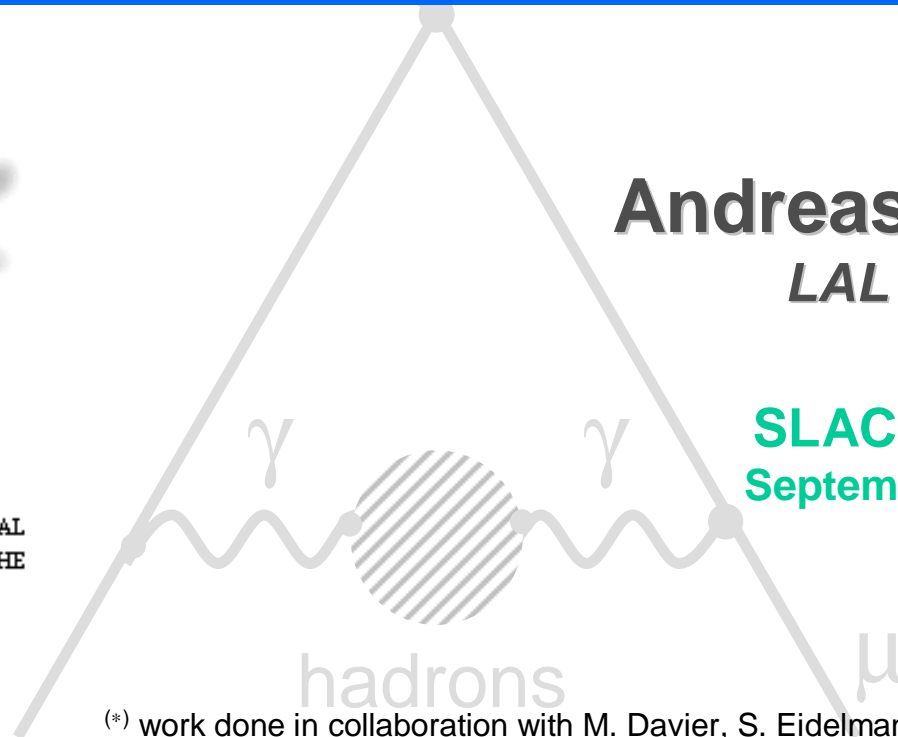
# Hadronic Vacuum Polarization

Confronting Spectral Functions from  $e^+e^-$  Annihilation and  $\tau$  Decays: Consequences for the Muon Magnetic Moment



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SLAC Seminar  
September 3, 2002



(\*) work done in collaboration with M. Davier, S. Eidelman and Z. Zhang [hep-ph/0208177]

# Outline

The Muon Magnetic Anomaly  
... and **Hadronic Vacuum Polarization**

$a_\mu$  [had]

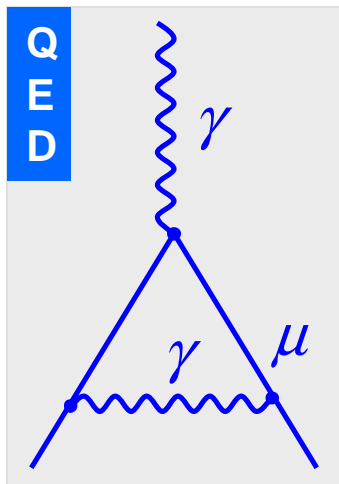
$\Delta\alpha_{\text{QED}}$  [had]

$e^+e^-$  versus  $\tau$  Data

Isospin Breaking  
and  
Radiative Corrections

Two New Standard Model Predictions  
for  $(g-2)_\mu$

# Magnetic Anomaly



QED Prediction:

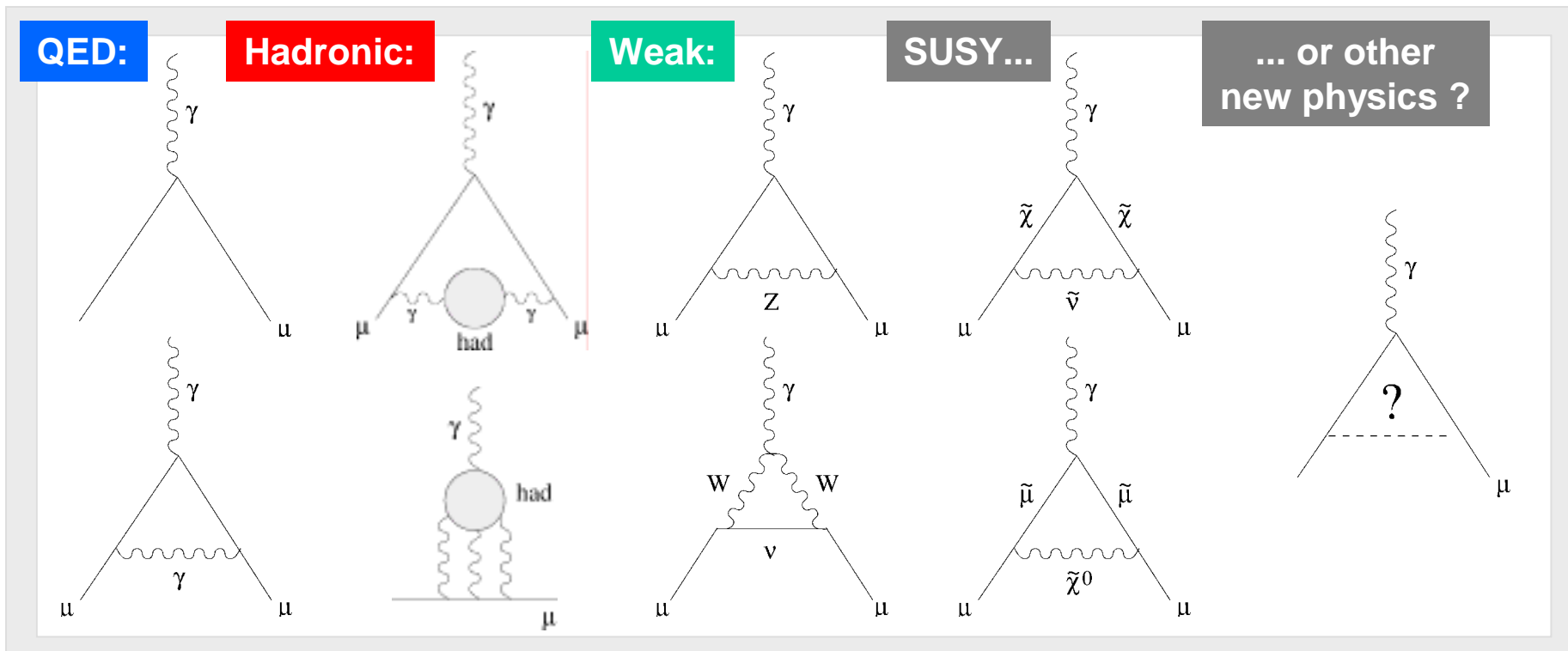
Computed up to 4<sup>th</sup> order  
[Kinoshita *et al.*]  
(5<sup>th</sup> order estimated  
[Mohr, Taylor])

$$\Gamma_{\mu} = e\gamma_{\mu} + a_{\ell} \frac{ie}{2m} \sigma_{\mu\nu} q_{\nu}$$

$$a_{\ell} = \frac{\alpha}{2\pi} = 0.001161\dots$$

**Schwinger** 1947  
(Nobel price 1965)

$$a_{\mu}^{\text{QED}} = \sum_{n=1} \left( \frac{\alpha}{\pi} \right)^n \approx \left( \begin{array}{l} 11614098.1 + 41321.8 \\ +3014.2 + 36.7 + 0.6 \end{array} \right) \times 10^{-10}$$



# The Muonic $(g-2)_\mu$

Contributions to the Standard Model (SM) Prediction:

$$a_\mu \equiv \left( \frac{g-2}{2} \right)_\mu = a_\mu^{\text{QED}} + a_\mu^{\text{had}} + a_\mu^{\text{weak}}$$

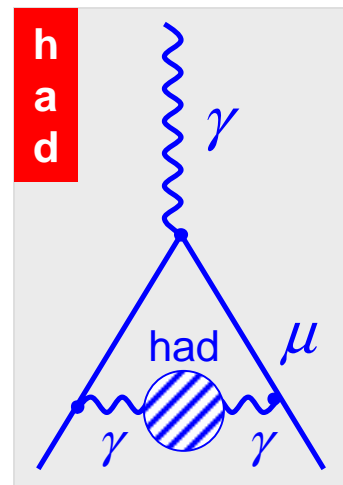
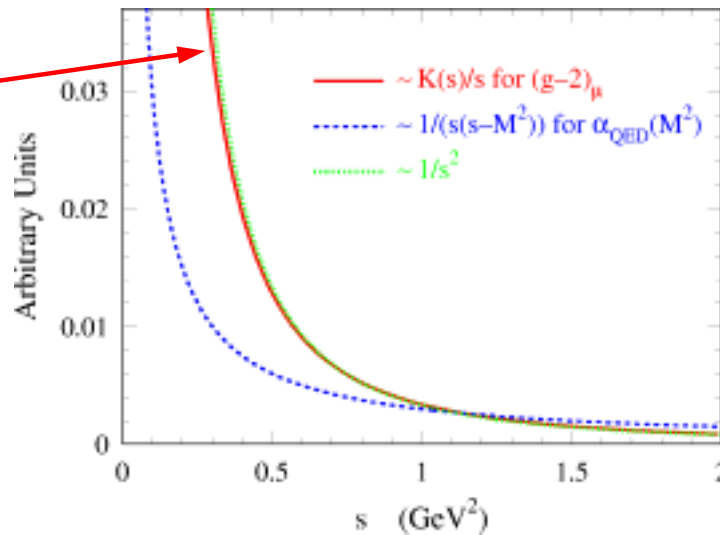
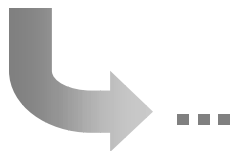
The Situation 1995

Source	$\alpha(a_\mu)$	Reference
QED	$\sim 0.3 \times 10^{-10}$	[Schwinger '48 & others]
Quarks and Hadrons	$\sim (15 \oplus 4) \times 10^{-10}$	[Eidelman-Jegerlegner '95 & others]
Z, W exchange	$\sim 0.4 \times 10^{-10}$	[Czarnecki <i>et al.</i> '95 & others]

Dominant uncertainty from lowest order hadronic piece. Cannot be calculated from QCD ("first principles") – but: we can use experiment (!)

$$a_\mu^{\text{had}} = \frac{\alpha^2}{3\pi^2} \int_{4m_\pi^2}^{\infty} ds \frac{K(s)}{s} R(s)$$

"Dispersion relation"



# Hadronic Vacuum Polarization

Define: photon vacuum polarization function  $\Pi_\gamma(q^2)$

$$i \int d^4x e^{iqx} \langle 0 | T J_{\text{em}}^\mu(x) (J_{\text{em}}^\nu(x))^\dagger | 0 \rangle = - (g^{\mu\nu} q^2 - q^\mu q^\nu) \Pi_\gamma(q^2)$$

Ward identities: only vacuum polarization modifies electron charge

$$\alpha(s) = \frac{\alpha(0)}{1 - \Delta\alpha(s)} \quad \text{with:} \quad \Delta\alpha(s) = -4\pi\alpha \text{Re} [\Pi_\gamma(s) - \Pi_\gamma(0)]$$

Leptonic  $\Delta\alpha_{\text{lep}}(s)$  calculable in QED. However, quark loops are modified by long-distance hadronic physics, cannot (yet) be calculated within QCD (!)

Way out: Optical Theorem (unitarity) ...

... and the subtracted dispersion relation of  $\Pi_\gamma(q^2)$  (analyticity)

Born:  $\sigma^{(0)}(s) = \sigma(s) (\alpha / \alpha(s))^2$

$$12\pi \text{Im} \Pi_\gamma(s) = \frac{\sigma^{(0)}[e^+e^- \rightarrow \text{hadrons}]}{\sigma^{(0)}[e^+e^- \rightarrow \mu^+\mu^-]} \equiv R(s)$$

$$\text{Im} [ \text{diagram with shaded circle} ] \propto | \text{diagram with hadron loop} |^2$$

$$\Pi_\gamma(s) - \Pi_\gamma(0) = \frac{s}{\pi} \int_0^\infty ds' \frac{\text{Im} \Pi_\gamma(s')}{s'(s' - s) - i\epsilon} \quad \Rightarrow \quad \Delta\alpha_{\text{had}}(s) = -\frac{\alpha s}{3\pi} \text{Re} \int_0^\infty ds' \frac{R(s')}{s'(s' - s) - i\epsilon}$$

... and equivalently for  $a_\mu[\text{had}]$

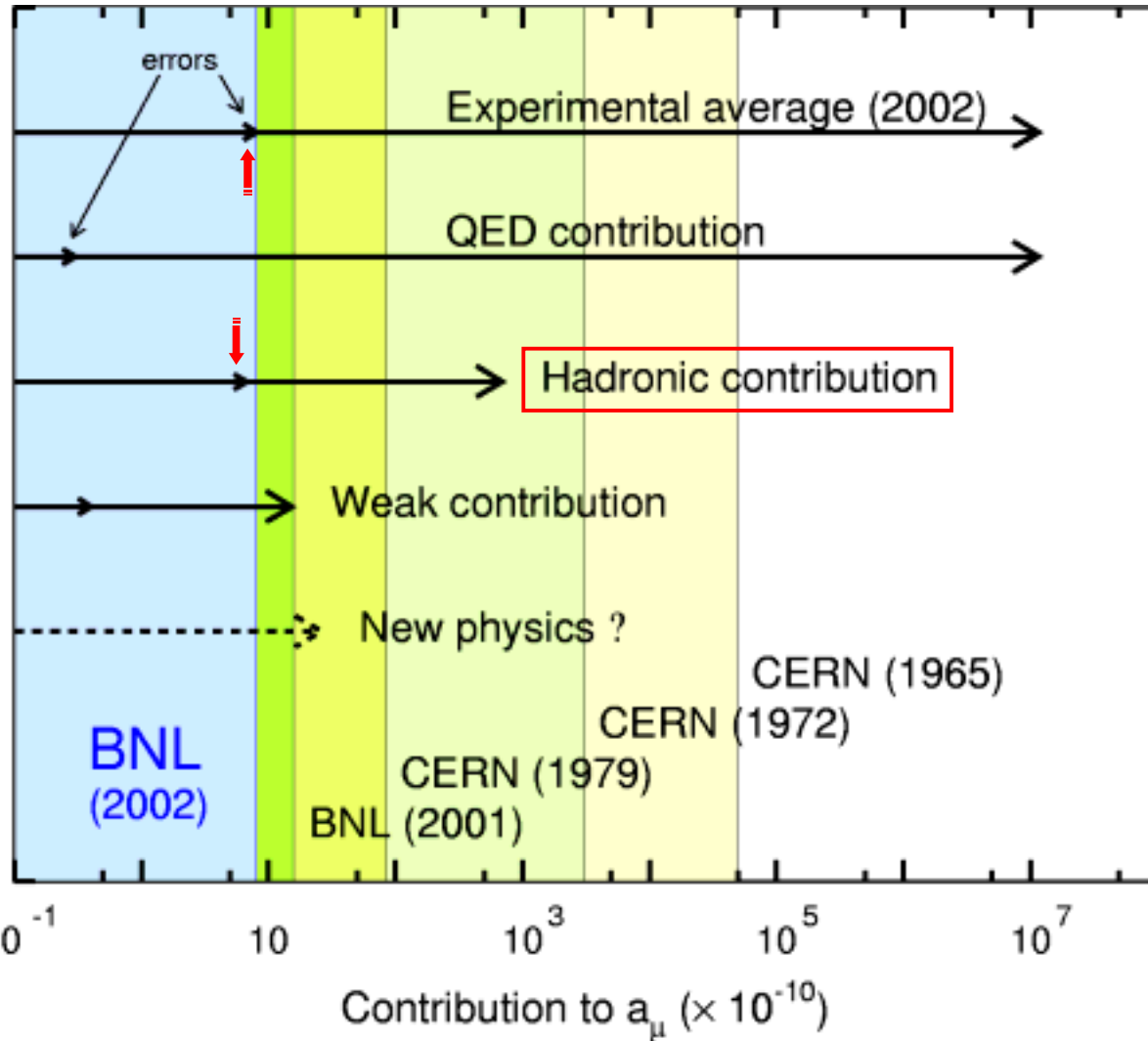
# Why Do We Need to Know it so Precisely?

## Experimental Progress on Precision of $(g-2)_\mu$



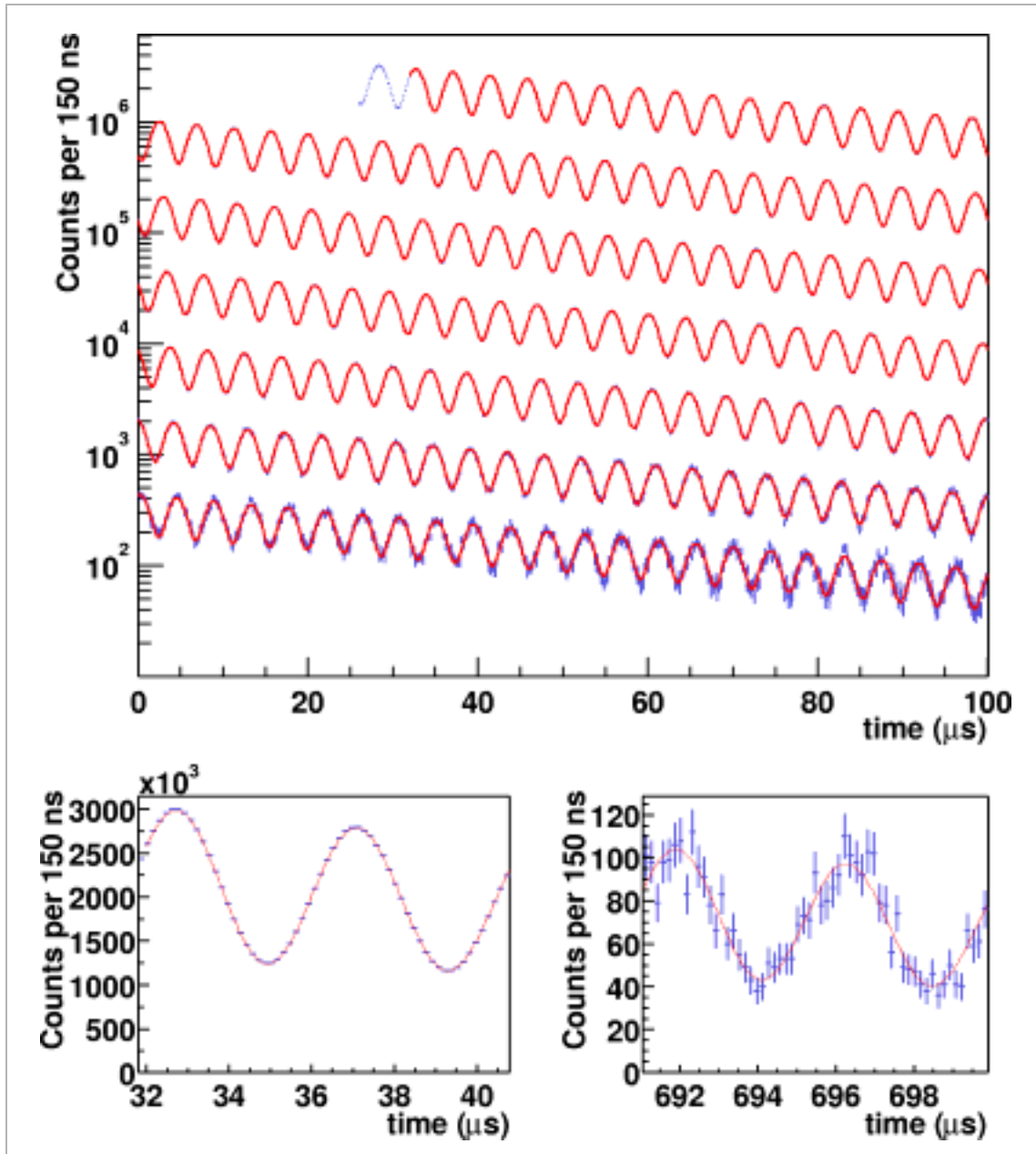
### Aiming at:

- $\sigma_{\text{exp}}(a_\mu) \square 4 \times 10^{-10}$
- $\square \sigma(a_\mu^{\text{had}})$
- $\square a_\mu^{\text{weak [SM]}}$



The fine structure constant at  $M_Z$  is an important ingredient to EW precision fits

# For the Beauty of it: BNL ( $g-2$ ) Measurement



Observed positron rate in successive 100 μs periods

Spin precession frequency:

$$\vec{\omega}_a = \frac{e}{mc} a_\mu \vec{B}$$

obtained from 5-parameter fit to the function:

$$N(t) = N_0 e^{-\lambda t} [1 + A \cos(\omega_a t + \phi)]$$

taken from:

**E821 ( $g-2$ ) [D.W. Hertzog]**

hep-ex/0202024

# 2001: The First Round of BNL Results on $a_\mu$

The E821 ( $g-2$ ) experiment at BNL published early 2001 a value 3× more precise than the previous CERN and BNL expts. combined:

E821 ( $g-2$ )  
hep-ex/0102017

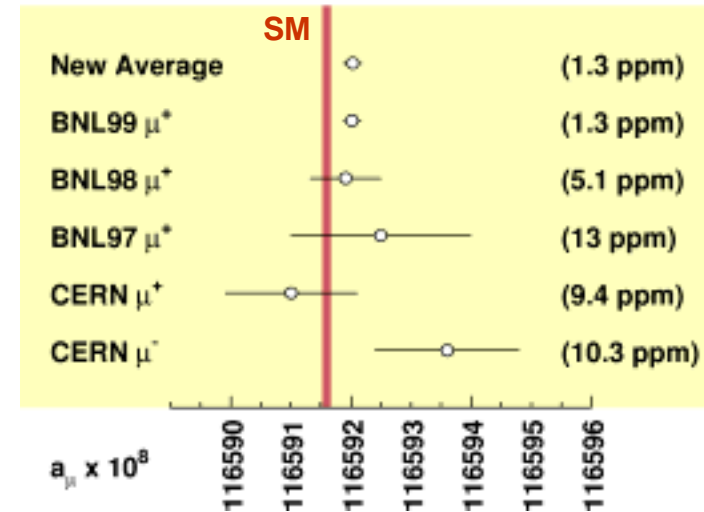
$$a_\mu(\text{exp}) = 11\,659\,202(16) \times 10^{-10}$$

BNL compares with Standard Model prediction:

$$a_\mu(\text{SM}) = 11\,659\,159.6(6.7) \times 10^{-10}$$

Averaging E821 with previous experiments gives:

$$a_\mu(\text{exp}) - a_\mu(\text{SM}) = 43(16) \times 10^{-10} \quad [\rightarrow 2.7 \sigma]$$

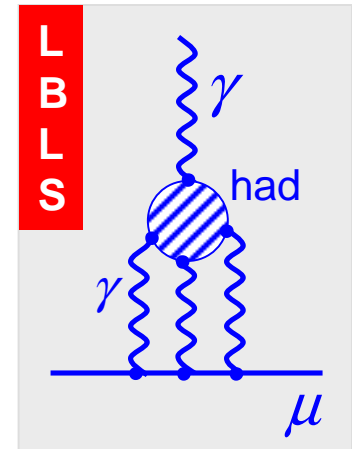


**BUT:** In November 2001, Knecht & Nyffeler have corrected a sign error in the dominant ( $\pi$ -pole) contribution from hadronic light-by-light (LBL) scattering, reducing the above discrepancy to

$$25(16) \times 10^{-10} \quad [\rightarrow 1.6 \sigma]$$

Knecht-Nyffeler, hep-ph/0111058; result approved by:

Hayakawa-Kinoshita, hep-ph/0112102; Bijnens-Pallante-Prades, hep-ph/0112255





# 2002: The Second Round of BNL Results on $a_\mu$

The new analysis, first presented at ICHEP'02, achieves 2× better precision (using 4× more statistics) than the 2001 result:

E821 ( $g-2$ )  
hep-ex/0208001

$$a_\mu(\text{exp}) = 11\,659\,203(7)(5) \times 10^{-10}$$

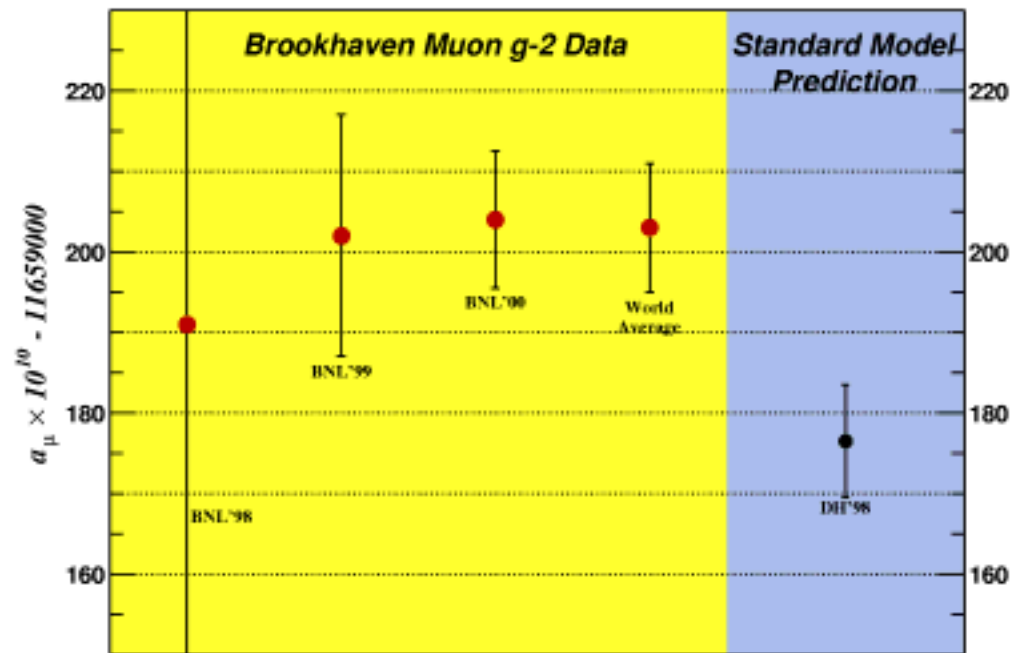
Error dominated by statistics. Systematics:

- $3.6 \times 10^{-10}$  from precession frequency
- $2.8 \times 10^{-10}$  from magnetic field

BNL compares WA with SM prediction (using DH'98 for hadronic vac. pol.):

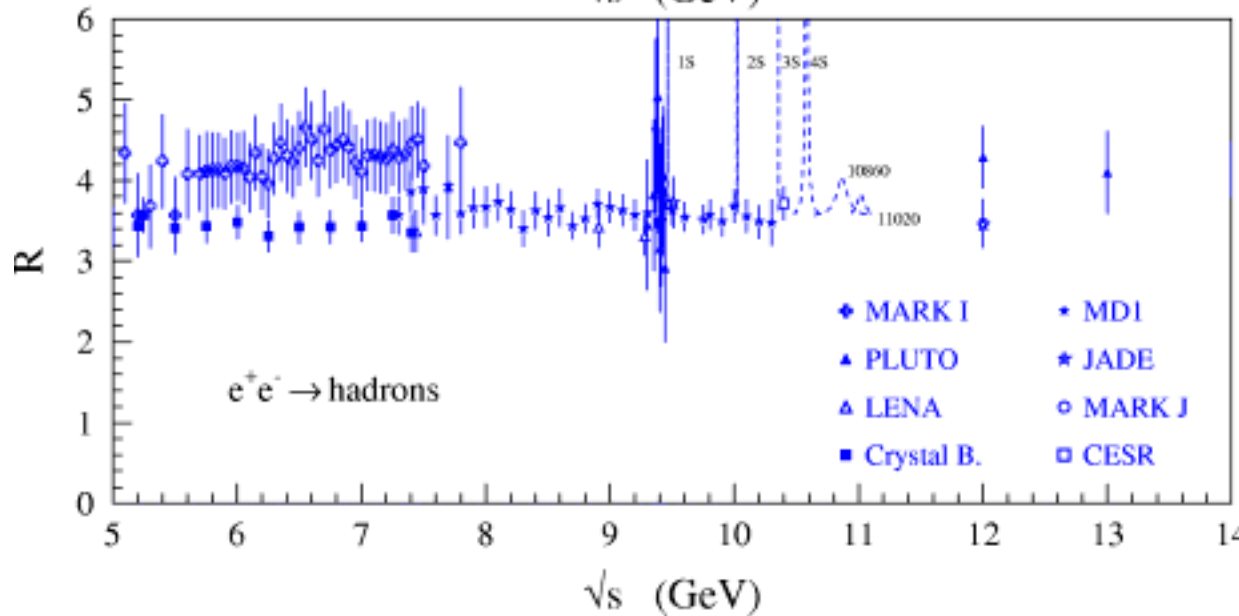
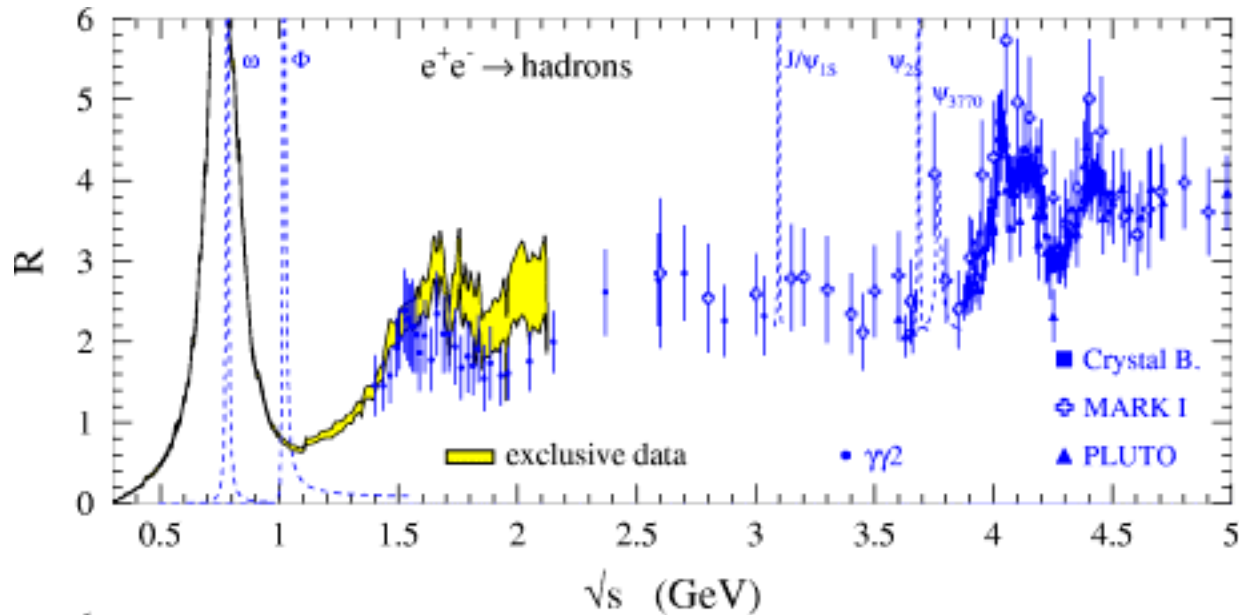
$$a_\mu(\text{exp}) - a_\mu(\text{SM}) = 25(10) \times 10^{-10}$$

Experimental and theoretical uncertainties now of similar order !



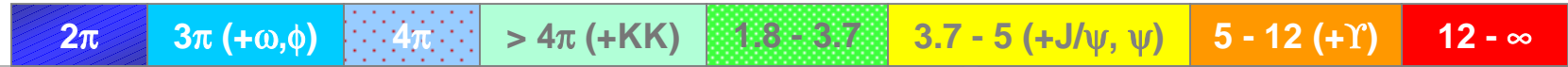
More work and, in particular, better data needed to achieve a more precise prediction of the hadronic contribution

# The Data Situation (around 1995)



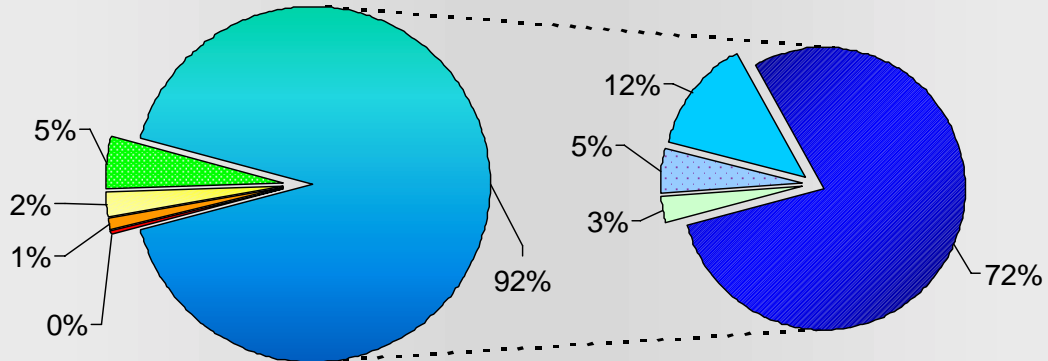
Data density and quality unsatisfactory in some crucial energy regions

# (I) Energy-Dependent Contributions and Errors

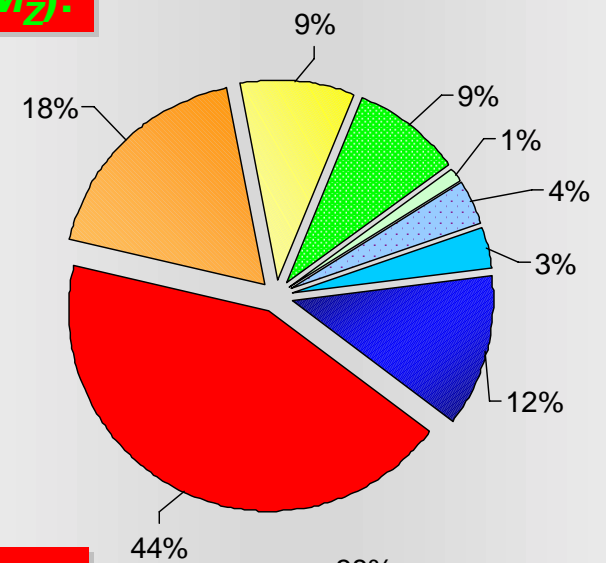


$< 1.8 \text{ GeV}$

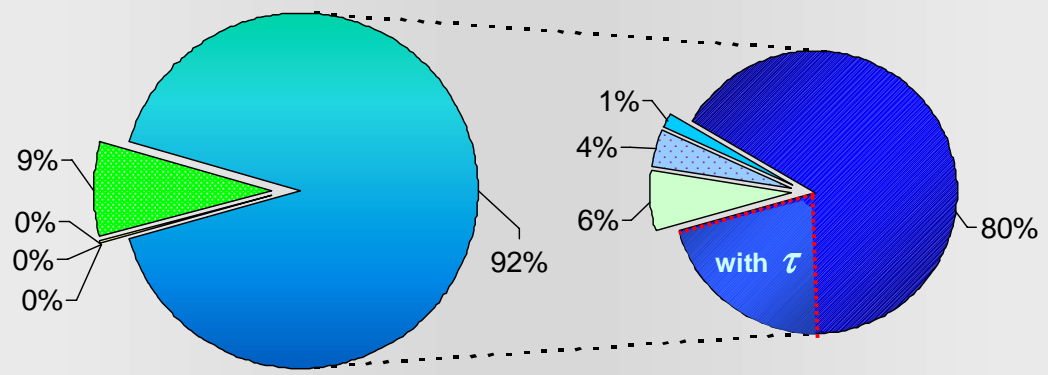
$a_{\mu}^{\text{had}}$



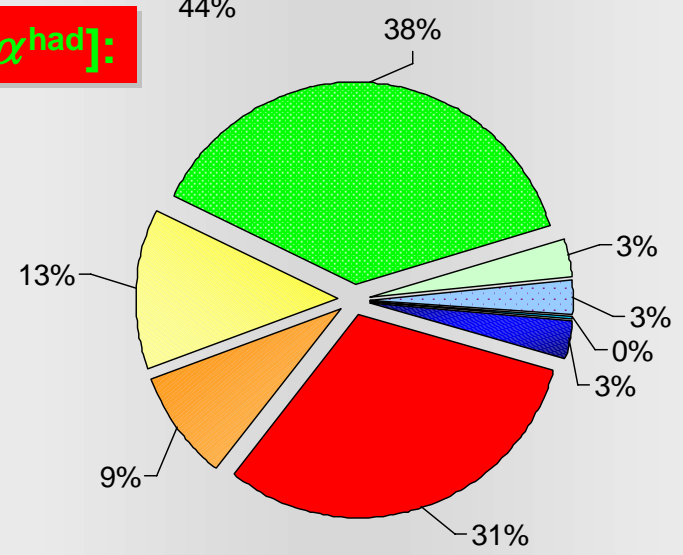
$\Delta\alpha^{\text{had}}(M_Z)$



$\sigma^2[a_{\mu}^{\text{had}}]$



$\sigma^2[\Delta\alpha^{\text{had}}]$



# Improved Determination of the Hadronic Contribution to $(g-2)_\mu$ and $\alpha(M_Z^2)$

Situation 1995 [Eidelman-Jegerlehner] and since ...

**Eidelman-Jegerlehner**

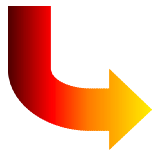
hep-ph/9502298

**Davier-AH**

hep-ph/9805470

Energy range [GeV]	Input 1995	Input after 1998
$2m_\pi - 1.8$	Data	Data ( $e^+e^-$ & $\tau$ ) + QCD
$1.8 - J/\psi$	Data	QCD
$J/\psi - \Upsilon$	Data	Data + QCD
$\Upsilon - 40$	Data	QCD
$40 - \infty$	QCD	QCD

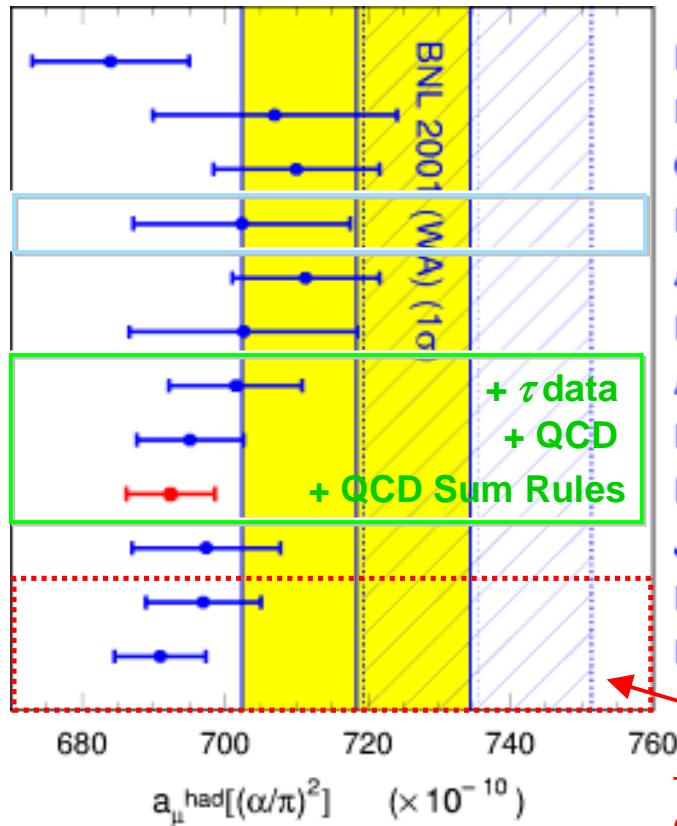
Improvement in 3 Steps:



- Inclusion of precise  $\tau$  data using SU(2) (CVC)
  - Alemany-Davier-AH'97, Narison'01, Trocóniz-Ynduráin'01
- Extended use of (dominantly) perturbative QCD
  - Martin-Zeppenfeld'95, Davier-AH'97, Kühn-Steinhauser'98, Erler'98, + others
- More theoretical constraints from QCD sum rules
  - Groote, Körner, Schilcher, Nasrallah'98, Davier-AH'98, Martin-Outhwaite-Ryskin'00, Cvetič-Lee-Schmidt'01, + others

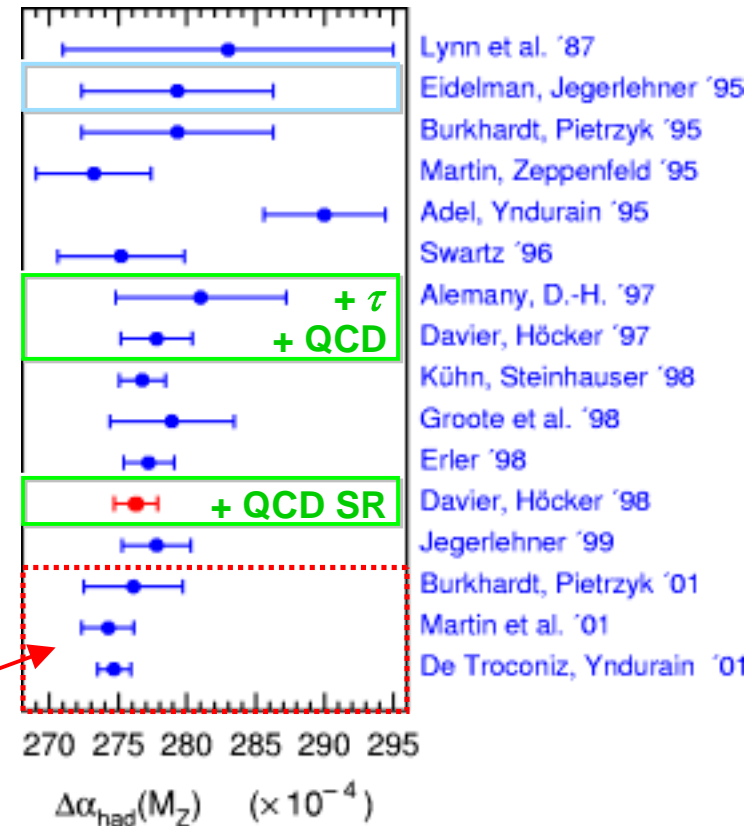
# Situation before Summer 2002

$a_\mu^{\text{had}}$



Barkov et al. '85  
Kinoshita et al. '85  
Casas et al. '85  
Eidelman, Jegerlehner '95  
Adel, Yndurain '95  
Brown, Worstell '96  
Alemany, D.-H. '97  
Davier, Höcker '97  
Davier, Höcker '98  
Jegerlehner '00  
Narison '01  
De Troconiz, Yndurain '01

$\Delta\alpha^{\text{had}}(M_Z)$



Lynn et al. '87  
Eidelman, Jegerlehner '95  
Burkhardt, Pietrzyk '95  
Martin, Zeppenfeld '95  
Adel, Yndurain '95  
Swartz '96  
Alemany, D.-H. '97  
Davier, Höcker '97  
Kühn, Steinhauser '98  
Groote et al. '98  
Eler '98  
Davier, Höcker '98  
Jegerlehner '99  
Burkhardt, Pietrzyk '01  
Martin et al. '01  
De Troconiz, Yndurain '01

These analyses use preliminary  
CMD-2 data (ICHEP'00) which  
suffer from missing rad. corrections

- Reasonable agreement among recent evaluations
- All recent analyses use enhanced theoretical input (compared to EJ'95)
- Most recent analyses employ  $e^+e^-$  and  $\tau$  data

# A New Analysis of $a_\mu^{\text{had}}$

## Motivation for new work:

- **New high precision  $e^+e^-$  results (0.6% sys. error) around  $\rho$  from CMD-2 (Novosibirsk)**
- **New  $\tau$  results from ALEPH using full LEP1 statistics, also: use CLEO data**
- **New  $R$  results from BES between 2 and 5 GeV**
- **New theoretical analysis on SU(2) breaking**

**CMD-2** PL B527, 161 (2002)

**ALEPH** CONF 2002-19

**CLEO** PR D61, 112002; PR D61, 072003 (2000)

**BES** PRL 84 594 (2000); PRL 88, 101802 (2002)

**Cirigliano-Ecker-Neufeld**  
hep-ph/0207310

## Outline of the new analysis:

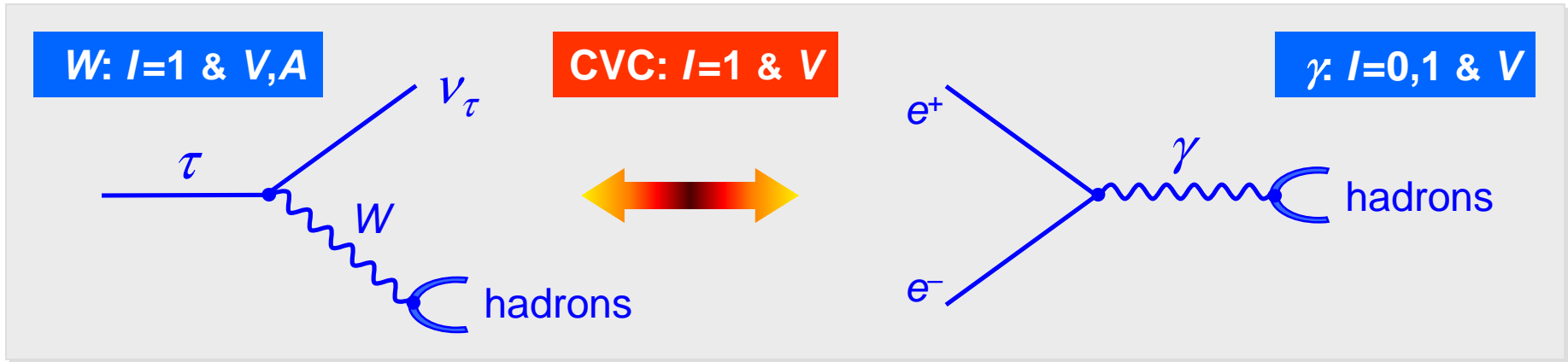
- Include all new Novosibirsk (CMD-2, SND) and ALEPH data
- Apply (revisited) SU(2)-breaking corrections to  $\tau$  data
- Identify application/non-application of radiative corrections
- Recompute all exclusive, inclusive and QCD contributions to dispersion integral; revisit threshold contribution and resonances
- Results, comparisons, discussions...

**Davier-Eidelman-AH-Zhang**  
hep-ph/0208177

# The Conserved Vector Current

The CVC property of weak decays follows from the factorization of strong physics produced through the  $\gamma$  and  $W$  propagators out of the QCD vacuum

# The Conserved Vector Current – SU(2)



Hadronic physics factorizes in **Spectral Functions**:

fundamental ingredient relating long distance (resonances) to short distance description (QCD)

Isospin symmetry (CVC) connects  $I=1$   $e^+e^-$  cross section to vector  $\tau$  spectral functions:

$$\sigma^{(I=1)} [e^+e^- \rightarrow \pi^+\pi^-] = \frac{4\pi\alpha^2}{s} v [\tau^- \rightarrow \pi^-\pi^0\nu_\tau]$$

$$v [\tau^- \rightarrow \pi^-\pi^0\nu_\tau] \propto \frac{\text{BR} [\tau^- \rightarrow \pi^-\pi^0\nu_\tau]}{\text{BR} [\tau^- \rightarrow e^-\bar{\nu}_e\nu_\tau]} \cdot \frac{1}{N_{\pi\pi^0}} \frac{dN_{\pi\pi^0}}{ds} \cdot \frac{m_\tau^2}{(1-s/m_\tau^2)^2 (1+s/m_\tau^2)}$$

branching Fractions    mass spectrum    kinematic factor (PS)



# $\tau$ Spectral Functions

Hadronic  $\tau$  decays are a clean probe of hadron dynamics – experimentally in many ways complementary to  $e^+e^- \rightarrow$  hadrons:

$\tau$

- Excellent normalization (branching fractions) due to high statistics, large acceptance, small non- $\tau$  background

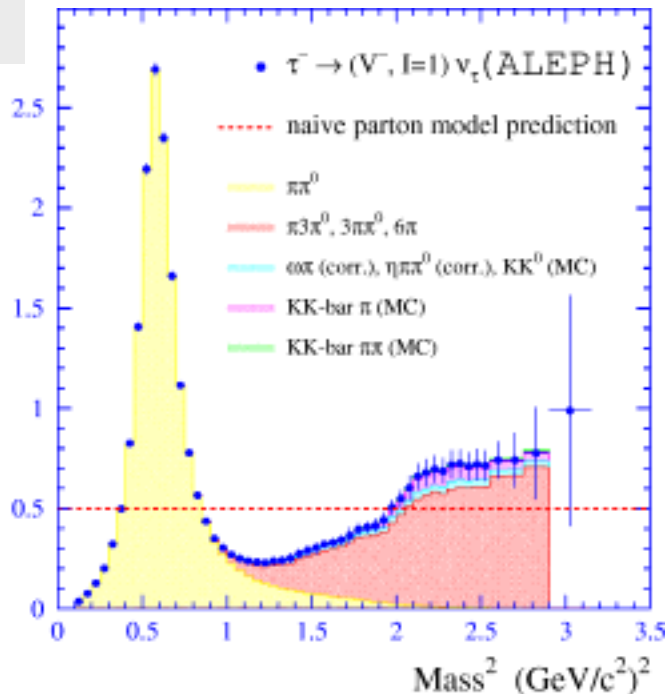
- Shape subject to bin-to-bin resolution corrections (*unfolding*)

- Excellent relative cross sections (correlated systematics)

$e^+e^-$

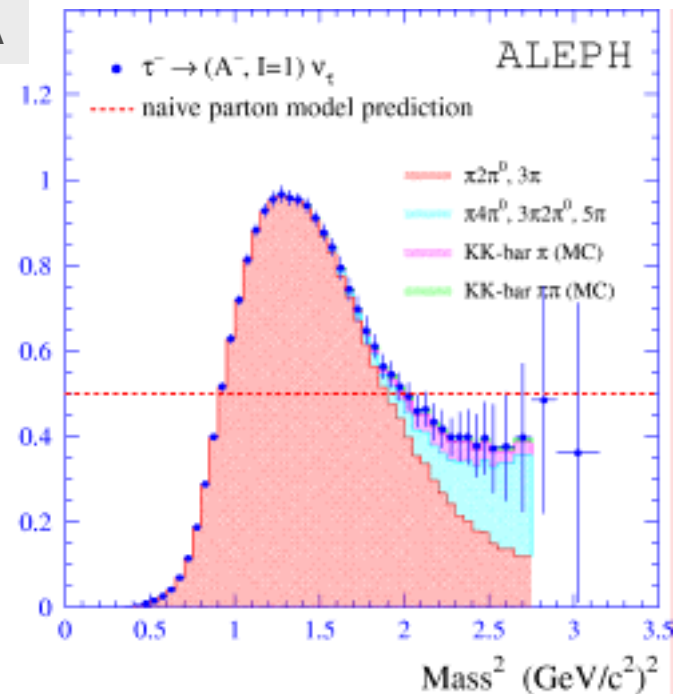
- Overall normalization subject to radiative corrections, systematic uncertainties from acceptance and luminosity

V



ALEPH  
ZP C76, 15  
(1997)

A



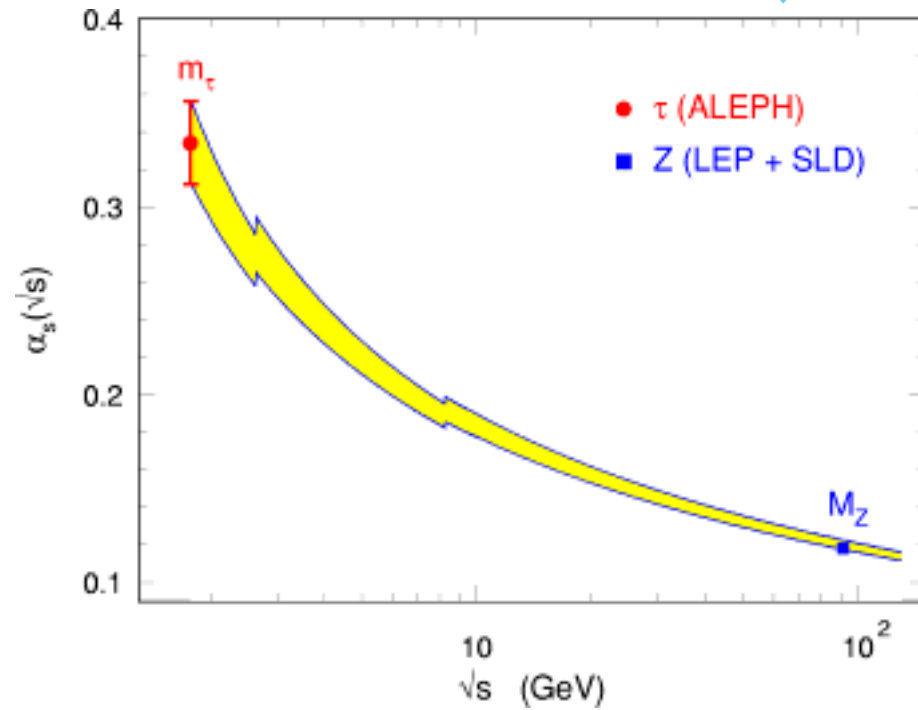
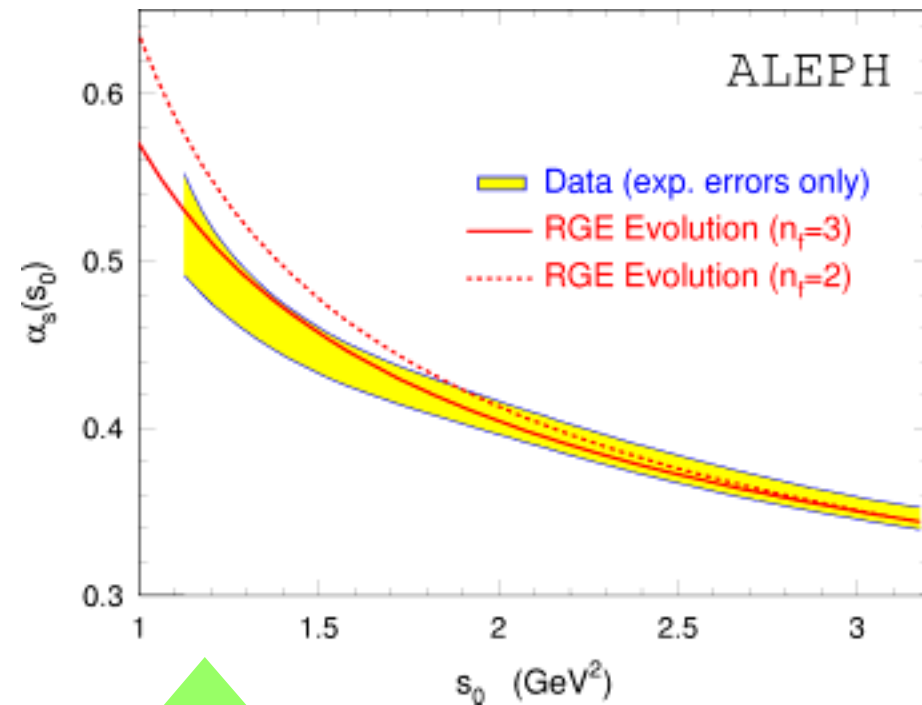
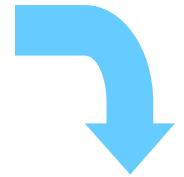
ALEPH  
EPJ C4, 409  
(1998)

# QCD Results from $\tau$ Decays

Evolution of  $\alpha_s(m_\tau)$ , measured using  $\tau$  decays, to  $M_Z$  using RGE (4-loop QCD  $\beta$ -function & 3-loop quark flavor matching) shows the excellent compatibility of  $\tau$  result with EW fit:

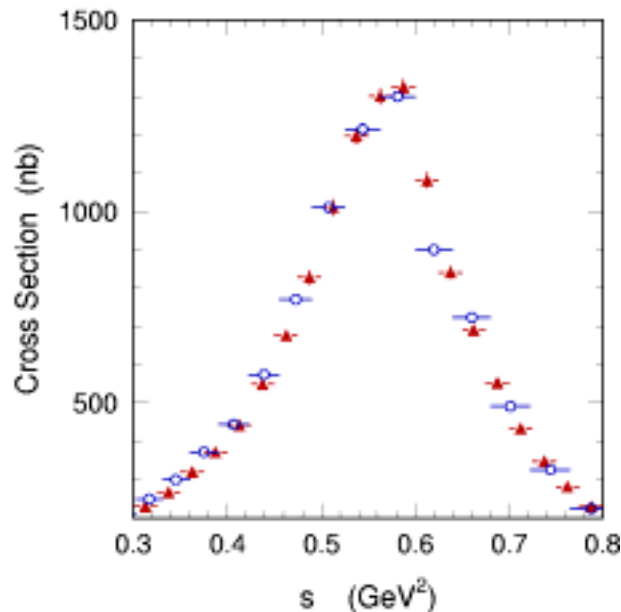
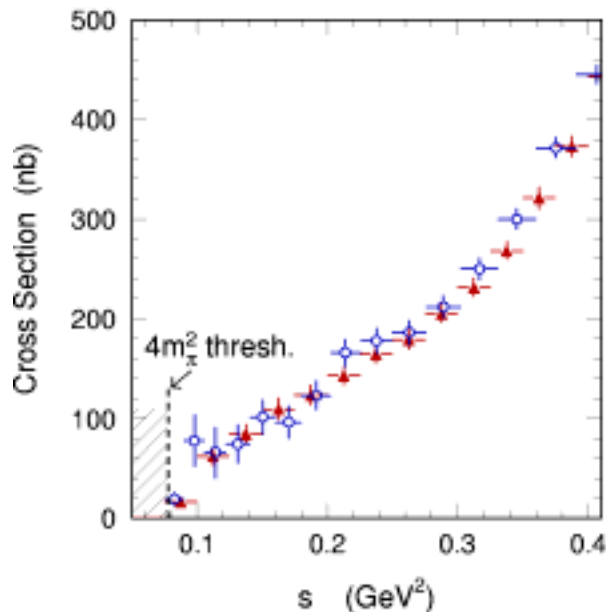
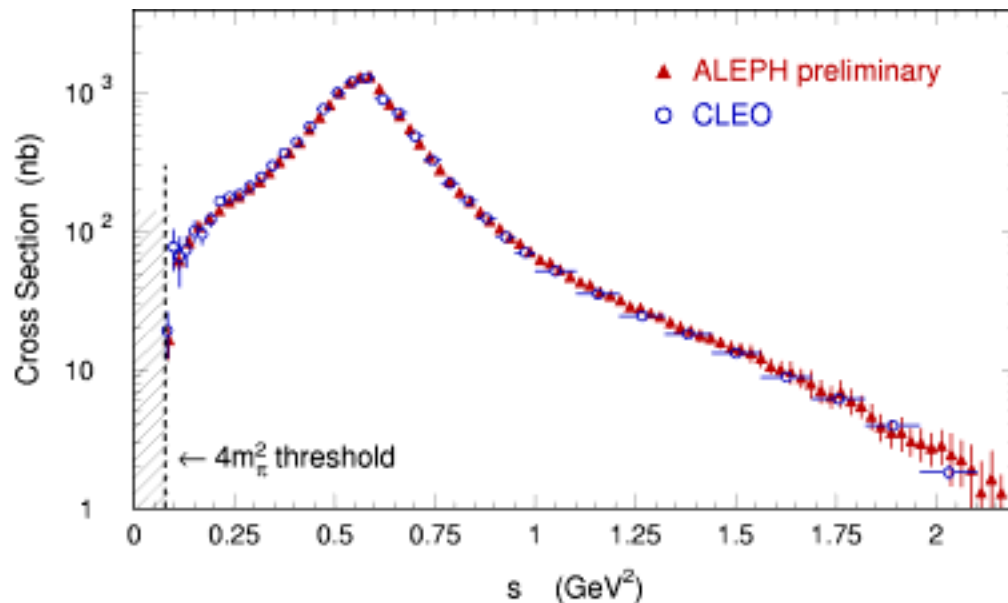
$$\alpha_s(M_Z) = 0.1202 \pm 0.0027 \quad (\text{ALEPH'98, theory dominated})$$

$$\alpha_s(M_Z) = 0.1183 \pm 0.0027 \quad (\text{LEP'00, statistics dominated})$$



The  $\tau$  spectral function allows to directly measure the running of  $\alpha_s(s_0)$  within  $\sqrt{s_0} \in [\sim 1 \dots 1.8 \text{ GeV}]$

# $\tau \rightarrow \pi^- \pi^0 \nu_\tau$ : Comparing ALEPH and CLEO



Spectral functions expressed as cross sections.

Shape comparison only. Both normalized to WA branching fraction.

- Agreement observed
- ALEPH more precise at low  $s$
- CLEO better at high  $s$

# SU(2) Breaking

Electromagnetism does not respect isospin and hence we have to consider isospin breaking when dealing with an experimental precision of 0.5%

Corrections for SU(2) breaking applied to  $\tau$  data for dominant  $\pi^-\pi^+$  contrib.:

## ■ Electroweak radiative corrections:

- ▶ dominant contribution from short distance correction  $S_{EW}$  to effective 4-fermion coupling  $\propto (1 + 3\alpha(m_\tau)/4\pi)(1+2\langle Q \rangle)\log(M_Z/m_\tau)$
- ▶ subleading corrections calculated and small
- ▶ long distance radiative correction  $G_{EM}(s)$  calculated [ add FSR to the bare cross section in order to obtain  $\pi^-\pi^+(\gamma)$  ]

Cirigliano-Ecker-Neufeld, hep-ph/0207310

New  
Development

## ■ Charged/neutral mass splitting:

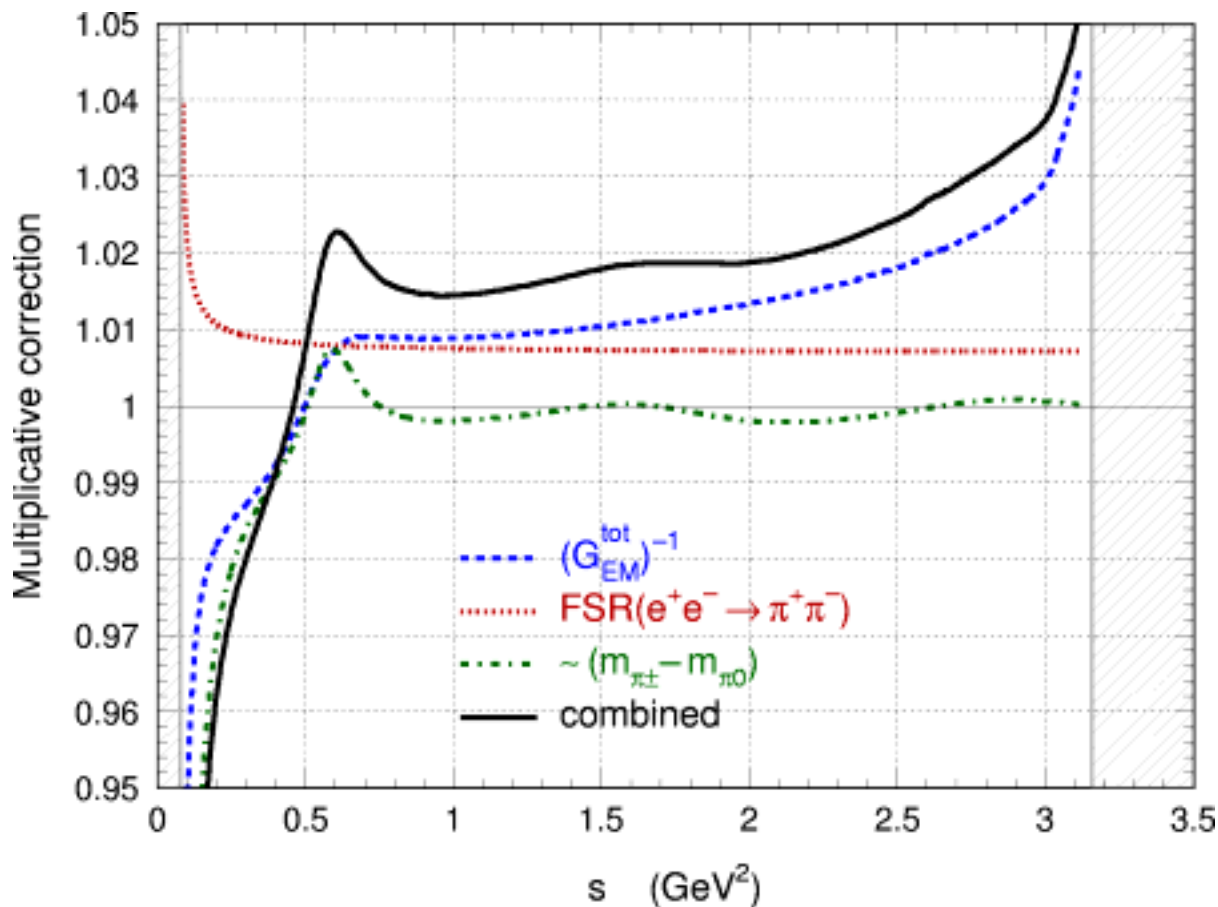
- ▶  $m_{\pi^-} \neq m_{\pi^0}$  leads to phase space (cross sec.) and width (FF) corrections
- ▶  $m_{\rho^-} \neq m_{\rho^0}$  and  $\rho$ - $\omega$  mixing (EM  $\omega \rightarrow \pi^-\pi^+$  decay) corrected using FF model

## ■ Electromagnetic decays, like: $\rho \rightarrow \pi\pi\gamma$ , $\rho \rightarrow \pi\gamma$ , $\rho \rightarrow \eta\gamma$ , $\rho \rightarrow l^+l^-$

## ■ Quark mass difference $m_u \neq m_d$ generating “second class currents” (negligible)

# SU(2) Breaking

Multiplicative SU(2) corrections applied to  $\tau^- \rightarrow \pi^- \pi^0 \nu_\tau$  spectral function:



Only  $\beta^3$  and EW short distance corrections applied to  $4\pi$  spectral functions

# SU(2) Breaking

## Corrections for isospin violation applied to $\tau$ data

Source	$\Delta a_\mu^{\text{had}} (10^{-10})$			
	$\pi^-\pi^+$ (simple)	$\pi^-\pi^+$ (improved)	$\pi^-\pi^+ 2\pi^0$	$2\pi^-2\pi^+$
<b>Short distance radiative corrections to <math>\tau</math> decays</b> ( $S_{\text{EW}} = 1.0267 \pm 0.0027$ ) [Marciano-Sirlin'88, Braaten-Li'90, new evaluation DEHZ'02]	$-13.8 \pm 2.5$	$-13.8 \pm 2.5$	$-0.49 \pm 0.09$	$-0.25 \pm 0.05$
<b>Long distance corrections</b>	-	$-1.0$	-	-
$m_{\pi^-} \neq m_{\pi^0}$ ( $\beta$ in cross section)	$-7.0$	$-7.0$	$+0.6 \pm 0.6$	$-0.4 \pm 0.4$
$m_{\pi^-} \neq m_{\pi^0}$ ( $\beta$ in width)	$+4.0$	$+4.2$	-	-
$m_{\rho^-} \neq m_{\rho^0}$ ( $\pm \approx 1 \text{ MeV}/c^2$ )	$+0 \pm 0.2$	$+0 \pm 2.0$	-	-
$\rho$ - $\omega$ mixing (exp. uncertainty)	$+3.5 \pm 0.6$	$+3.5 \pm 0.6$	-	-
<b>EM decay modes</b>	$-1.4 \pm 1.2$	$-1.4 \pm 1.2$	-	-
<b>Total correction</b>	$-14.7 \pm 2.9$	$-15.7 \pm 2.8$	$+0.1 \pm 0.6$	$-0.7 \pm 0.4$

# $e^+e^-$ Radiative Corrections

Multiple radiative corrections are applied on measured  $e^+e^-$  cross sections

Situation often unclear: whether or not - and if - which corrections were applied

## ■ Vacuum polarization (VP) in the photon propagator:

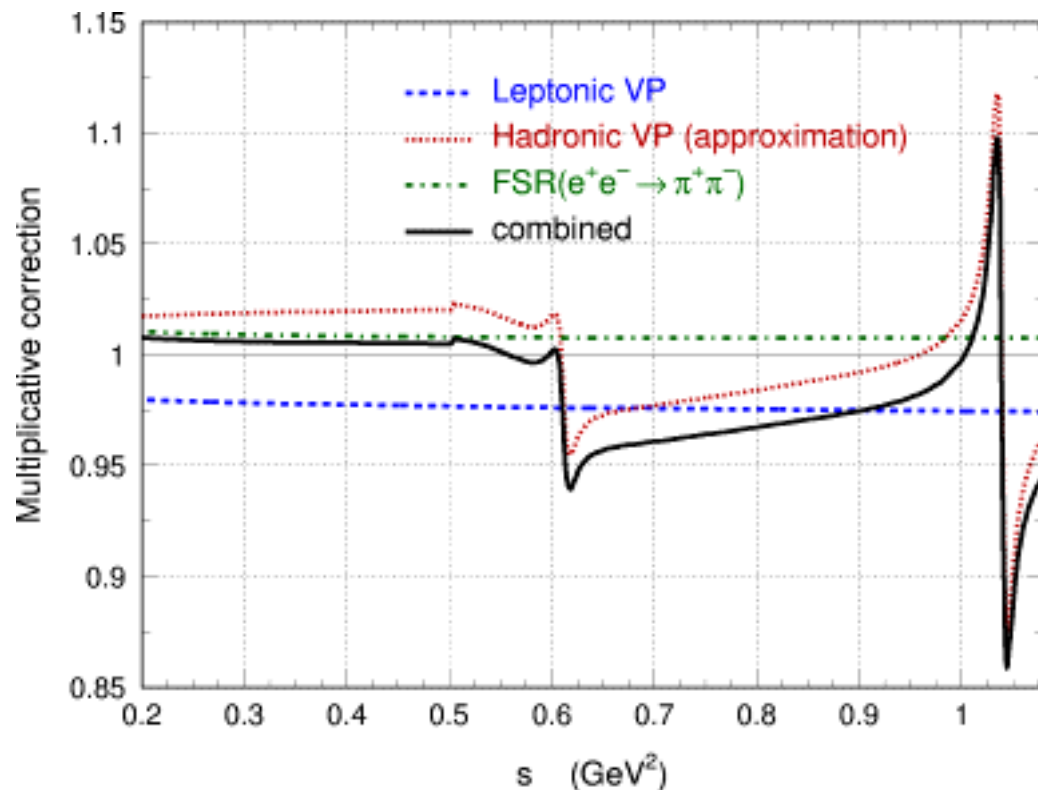
- ▶ leptonic VP mostly corrected
- ▶ hadronic VP not corrected but for CMD-2 (in principle: iterative proc.)

## ■ Initial state radiation (ISR)

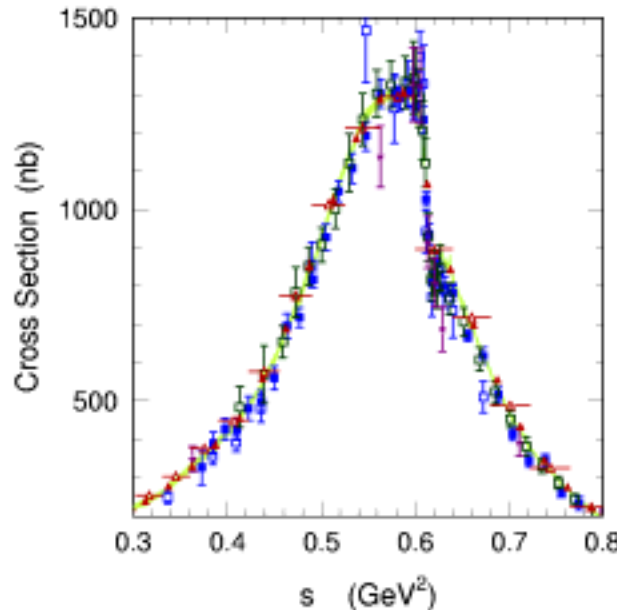
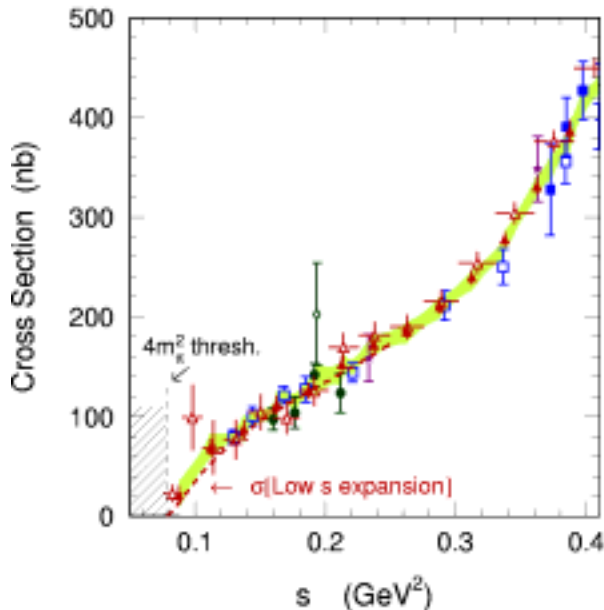
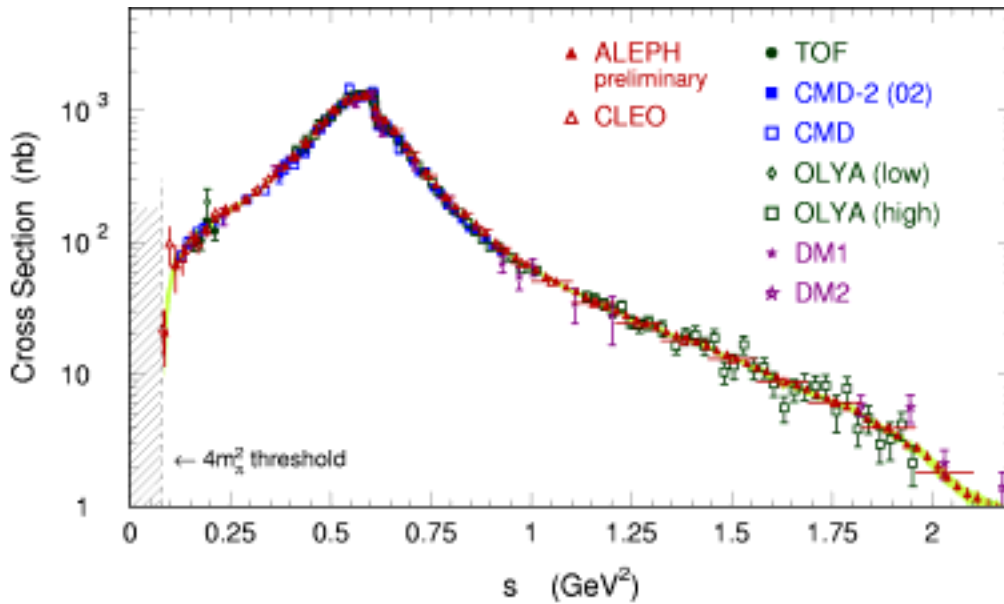
- ▶ corrected by experiments

## ■ Final state radiation (FSR) [we need $e^+e^- \rightarrow \text{hadrons} (\gamma)$ in dispersion integral]

- ▶ mostly, experiments obtain bare cross section so that FSR has to be added “by hand”; done for CMD-2, (supposedly) not done for others



# Comparing $e^+e^- \rightarrow \pi^+\pi^-$ and $\tau \rightarrow \pi^-\pi^0\nu_\tau$



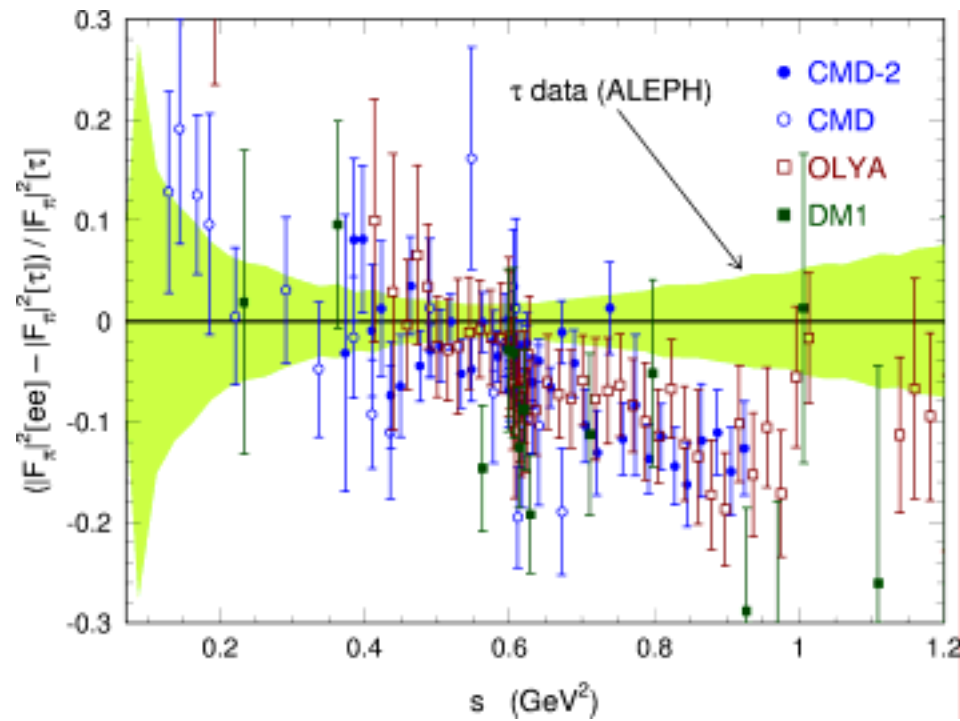
Correct  $\tau$  data for missing  $\rho$ - $\omega$  mixing (taken from BW fit) and all other SU(2)-breaking sources

- Remarkable agreement
- But: is it good enough ?

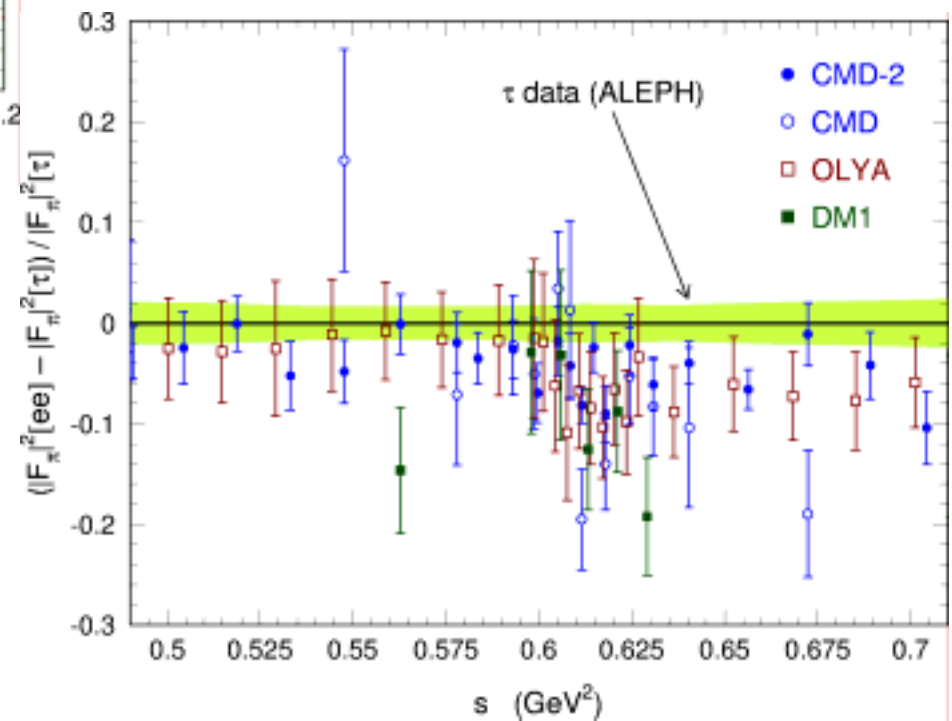
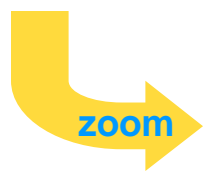




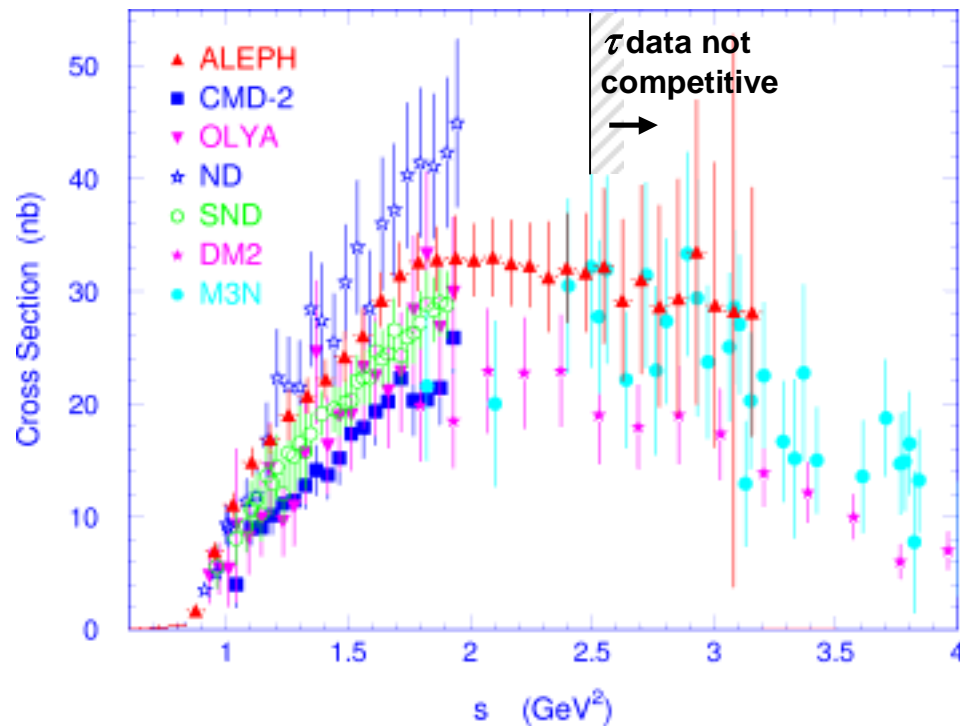
# Comparing $e^+e^- \rightarrow \pi^+\pi^-$ and $\tau^- \rightarrow \pi^-\pi^0\nu_\tau$



Relative difference between  $\tau$  and  $e^+e^-$  data



# Comparing the $4\pi$ Spectral Functions



$e^+e^- \rightarrow \pi^+\pi^-2\pi^0$

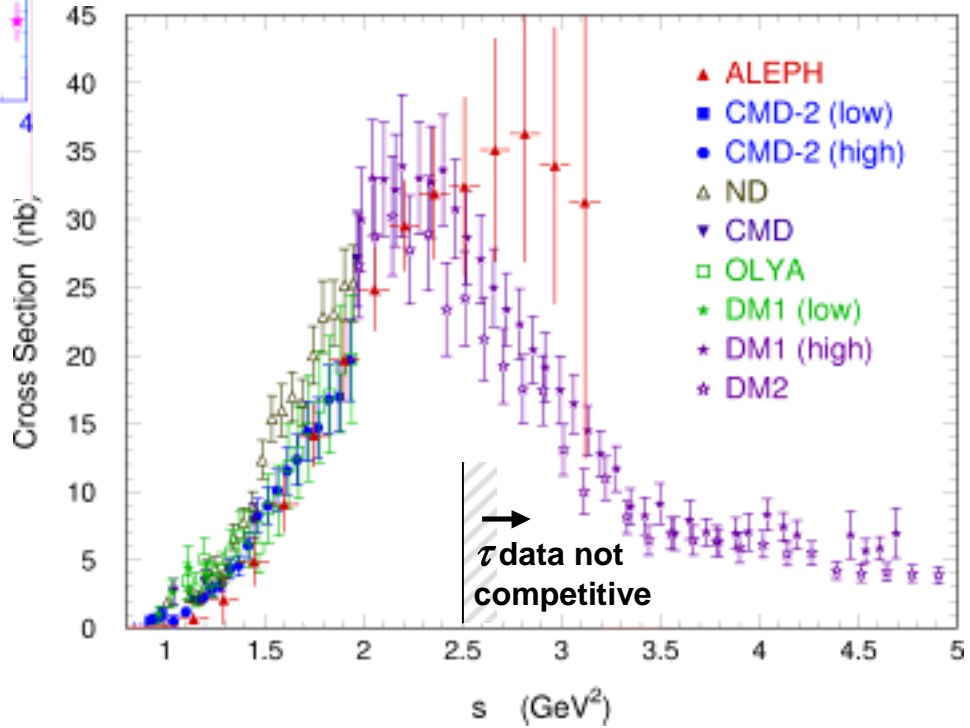
Large discrepancies between experiments

$e^+e^- \rightarrow 2\pi^+2\pi^-$

CVC relations (Isospin rotation):

$$\sigma_{2\pi^+2\pi^-}^{(I=1)} = 2 \cdot \frac{4\pi\alpha^2}{s} v_{\pi^-3\pi^0\nu_\tau}$$

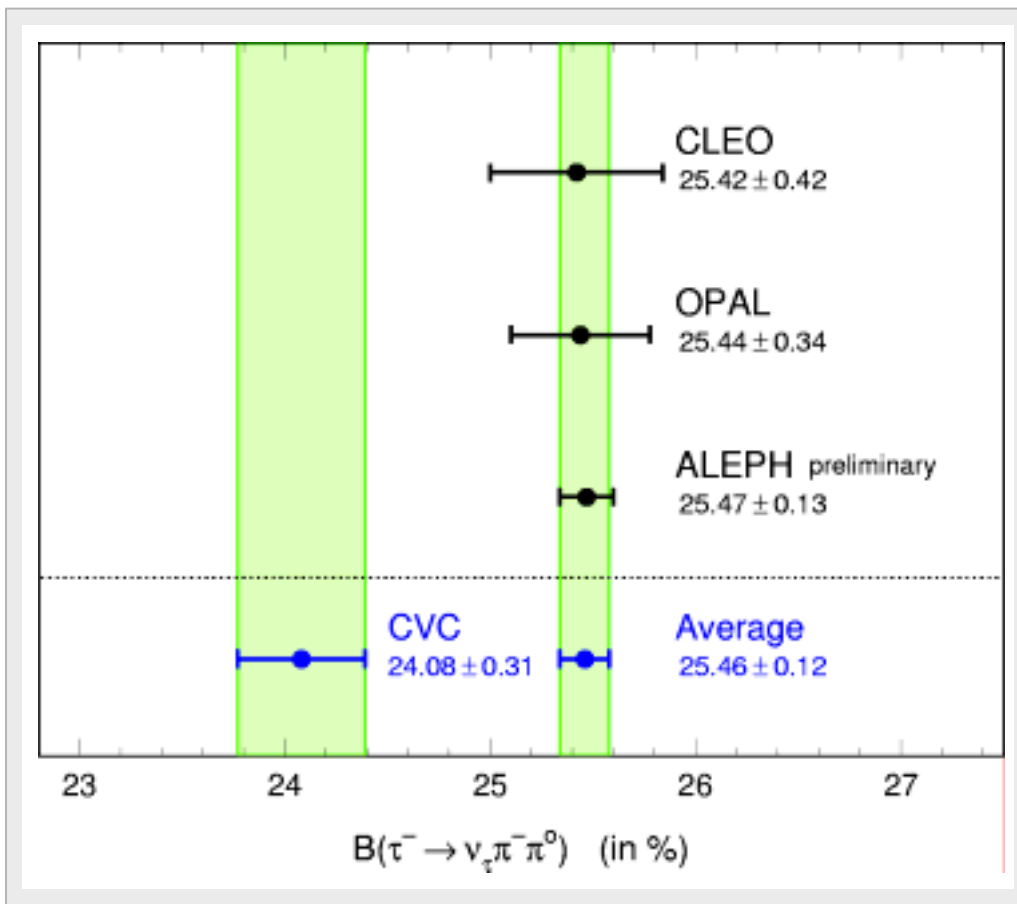
$$\sigma_{\pi^+\pi^-2\pi^0}^{(I=1)} = \frac{4\pi\alpha^2}{s} \left( v_{2\pi^-\pi^+\pi^0\nu_\tau} - v_{\pi^-3\pi^0\nu_\tau} \right)$$



# Testing CVC

Infer  $\tau$  branching fractions from  $e^+e^-$  data:

$$\text{BR}_{\text{CVC}}(\tau^- \rightarrow \pi^- \pi^0 \nu_\tau) = \frac{6\pi |V_{ud}|^2 S_{EW}}{m_\tau^2} \int_0^{m_\tau^2} ds \text{kin}(s) \cdot v^{\text{SU}(2)\text{-corrected}}(s)$$



←  $\tau^- \rightarrow \pi^- \pi^0 \nu_\tau$

Difference:  $\text{BR}[\tau] - \text{BR}[e^+e^- \text{ (cvc)}]$ :

Mode	$\Delta(\tau - e^+e^-)$	„Sigma“
$\tau^- \rightarrow \pi^- \pi^0 \nu_\tau$	$+1.38 \pm 0.33$	4.0
$\tau^- \rightarrow \pi^- 3\pi^0 \nu_\tau$	$-0.09 \pm 0.11$	0.8
$\tau^- \rightarrow 2\pi^- \pi^+ \pi^0 \nu_\tau$	$+0.88 \pm 0.25$	3.5

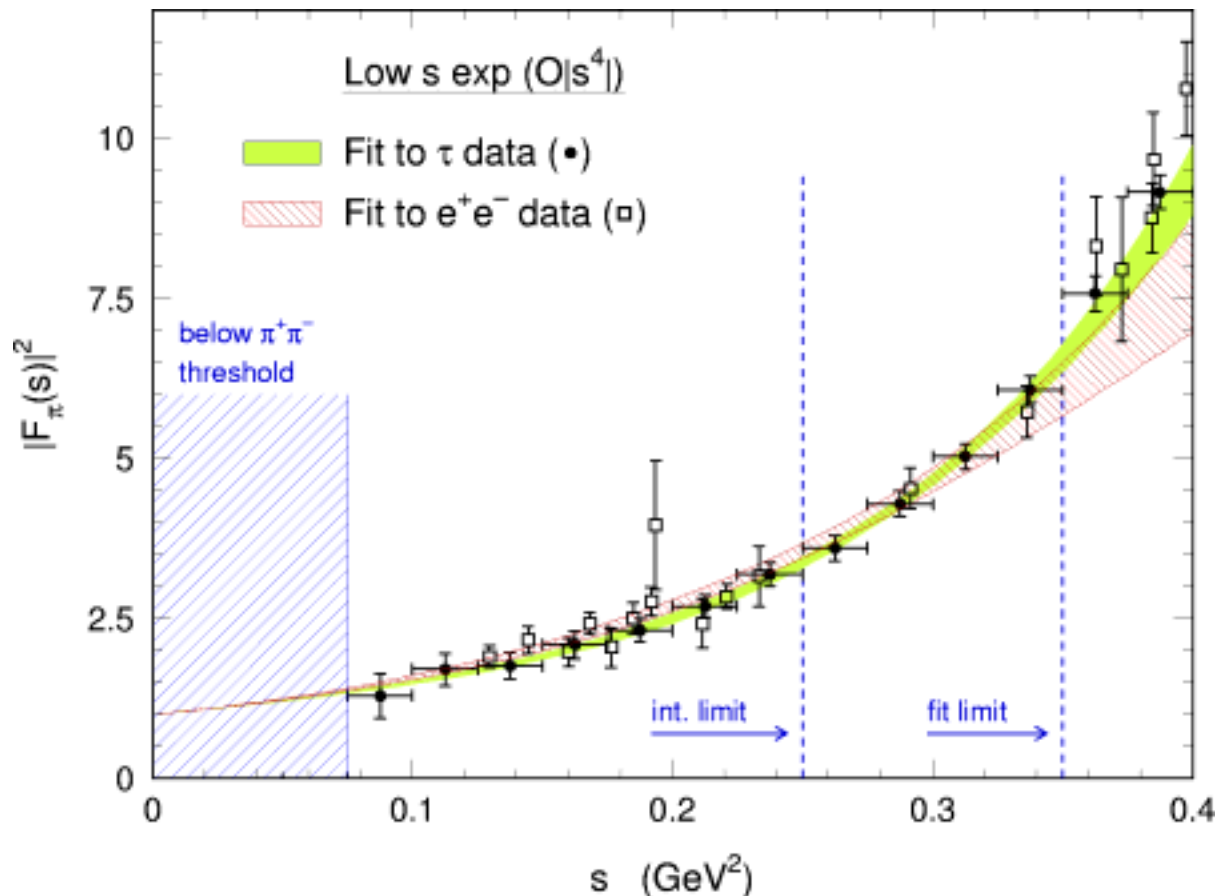
# **Specific Contributions**

# Specific Contributions: Low s Threshold

Use Taylor expansion for  $\pi^+\pi^-$  threshold:

$$\sigma_{\pi\pi} = \frac{\pi\alpha^2\beta^3}{3s} |F_\pi|^2 \quad \text{and:} \quad F_\pi = 1 + \frac{1}{6} \langle r^2 \rangle_\pi s + c_1 s^2 + c_2 s^3 + O(s^4)$$

- exploiting precise **space-like data**,  $\langle r^2 \rangle_\pi = (0.439 \pm 0.008) \text{ fm}^2$ , and **fitting**  $c_1$  and  $c_2$



Fit range:

- up to  $0.35 \text{ GeV}^2$

Integration range:

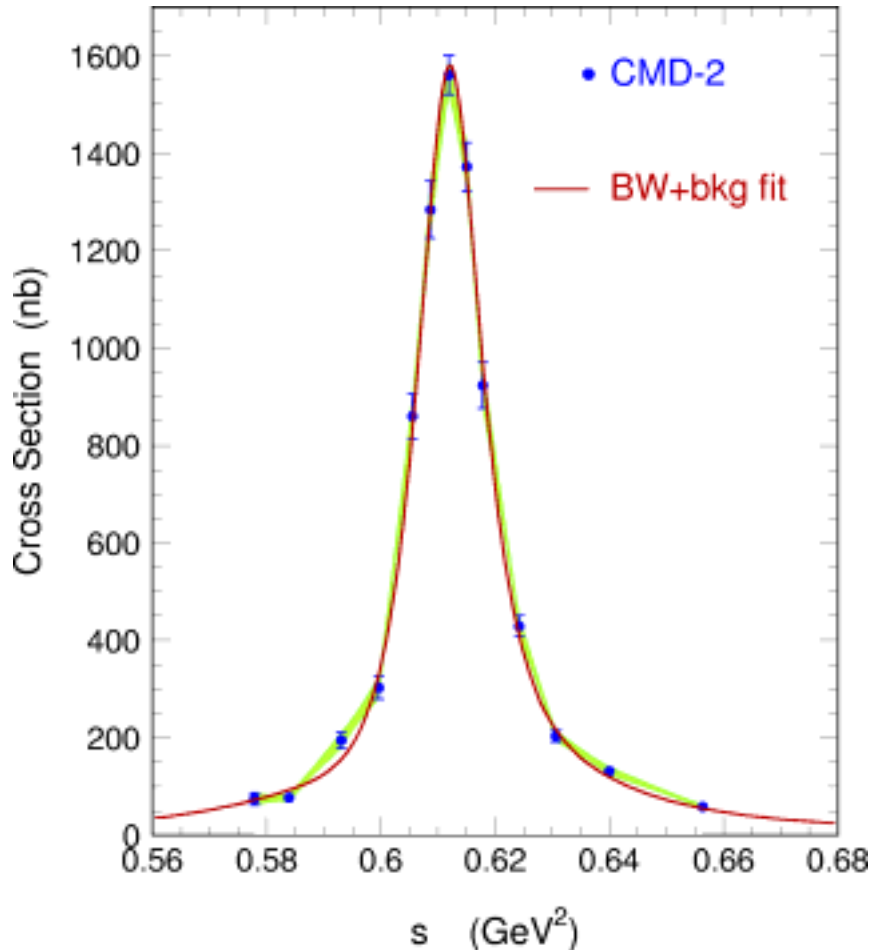
- up to  $0.25 \text{ GeV}^2$

- Excellent  $\chi^2$  for both  $\tau$  and  $e^+e^-$  data
- strong anti-correlation between  $c_1$  and  $c_2$

# Specific Contributions: Narrow Resonances

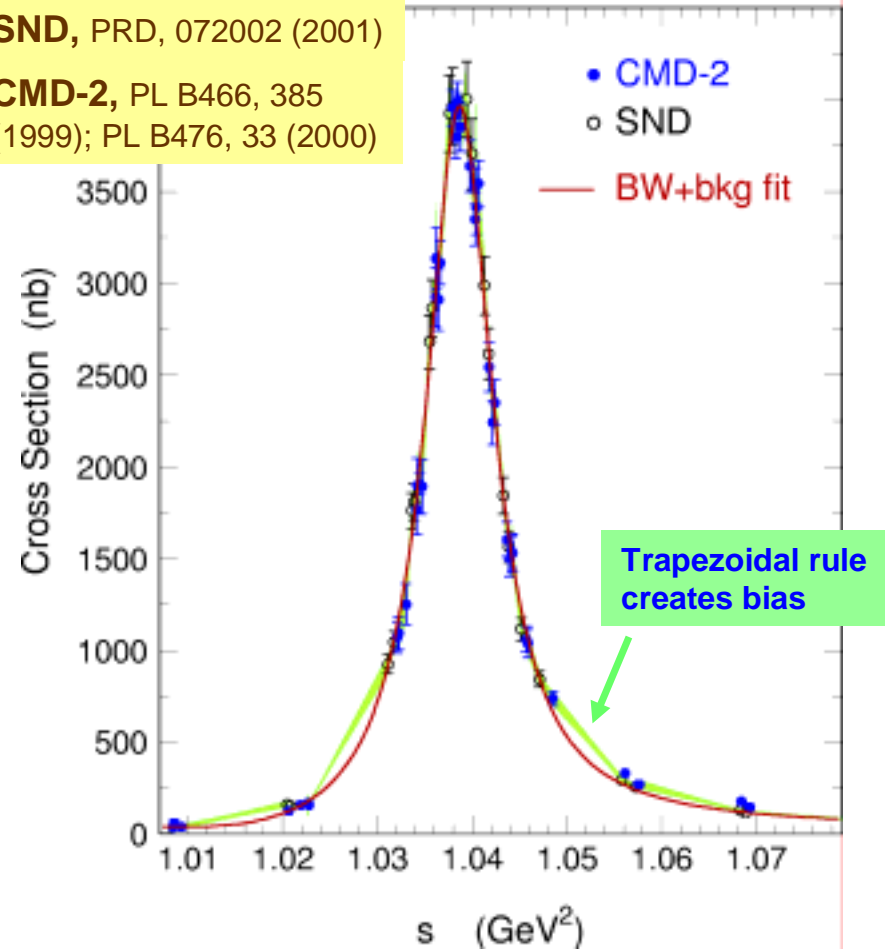
Use direct data integration for  $\omega(782)$  and  $\phi(1020)$  to account for non-resonant contributions. However, careful integration necessary:

- trapezoidal rule creates systematics for functions with strong curvature
- use phenomenological fit



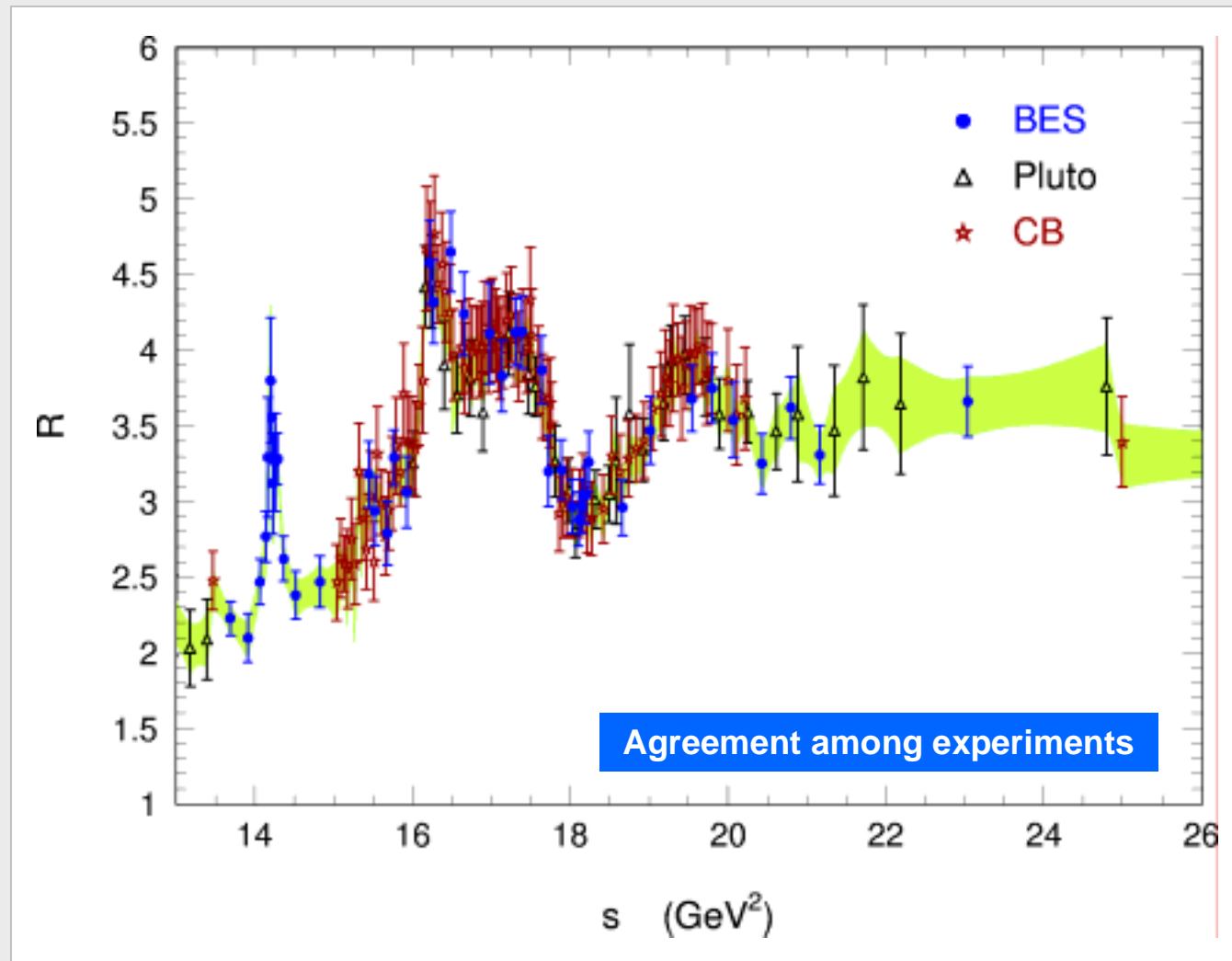
SND, PRD, 072002 (2001)

CMD-2, PL B466, 385 (1999); PL B476, 33 (2000)



# Specific Contributions: the Charm Region

New precise BES data improve  $c\bar{c}$  resonance region:

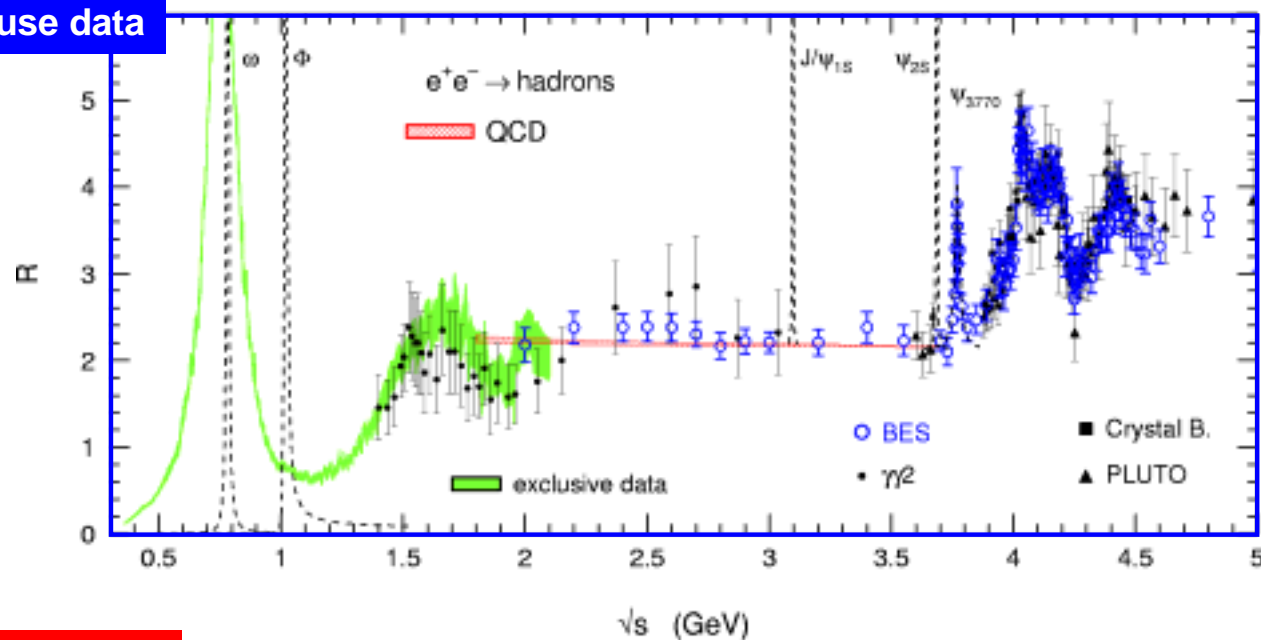


# Results

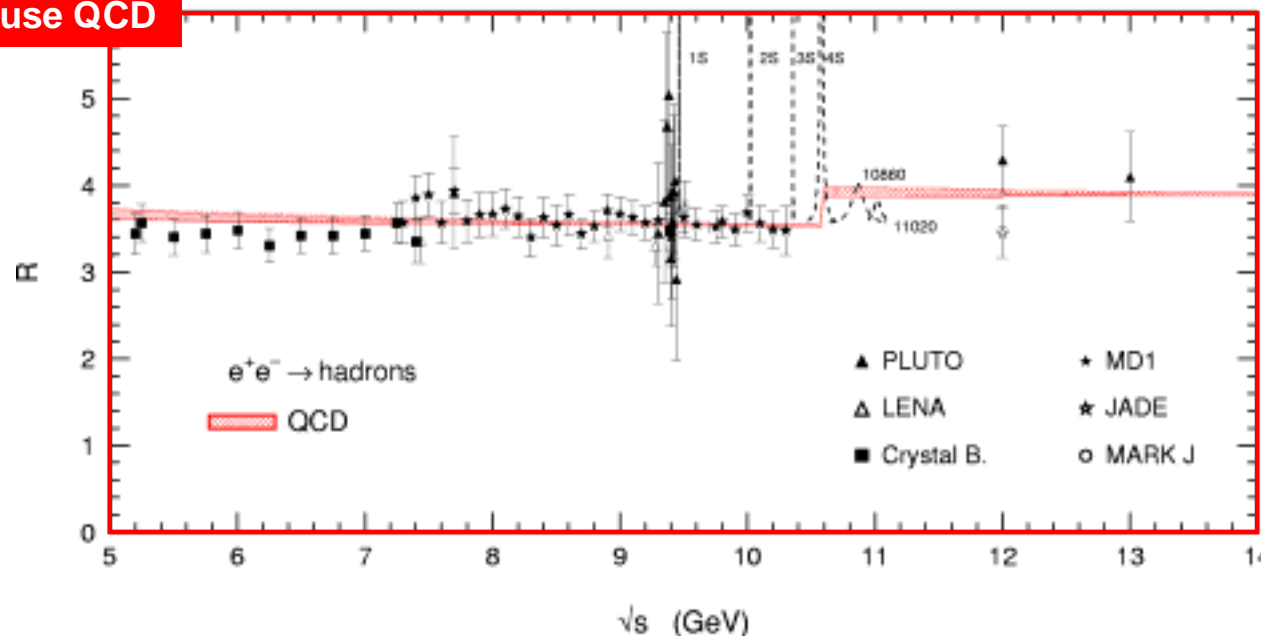


# Results: the Data & the Theory

use data



use QCD



■ Agreement between Data (BES) and pQCD

■ Better agreement between exclusive and inclusive ( $\gamma\gamma^2$ ) data than in previous analysis

# Results: the Compilation

Contributions to  $a_\mu^{\text{had}}$  from the different energy domains:

Modes	Energy range [GeV]	$a_\mu^{\text{had}}$ ( $10^{-10}$ )	
		$e^+e^-$	$\tau$
Low $s$ expansion	$2m_\pi - 0.5$	$58.0 \pm 1.7 \pm 1.1_{\text{rad}}$	$54.0 \pm 1.7 \pm 0.3_{\text{SU}(2)}$
$\pi^+\pi^-$	$2m_\pi - 1.8$	$440.8 \pm 4.7 \pm 1.5_{\text{rad}}$	$459.0 \pm 2.9 \pm 2.5_{\text{SU}(2)}$
$\pi^+\pi^- 2\pi^0$	$2m_\pi - 1.8$	$16.7 \pm 1.3 \pm 0.2_{\text{rad}}$	$21.4 \pm 1.1 \pm 0.6_{\text{SU}(2)}$
$2\pi^+ 2\pi^-$	$2m_\pi - 1.8$	$14.0 \pm 0.9 \pm 0.2_{\text{rad}}$	$12.3 \pm 1.0 \pm 0.4_{\text{SU}(2)}$
$\omega$ (782)	0.3 – 0.81	$36.9 \pm 0.8 \pm 0.8_{\text{rad}}$	-
$\phi$ (1020)	1.0 – 1.055	$34.8 \pm 0.9 \pm 0.6_{\text{rad}}$	-
Other exclusive	$2m_\pi - 2.0$	$32.2 \pm 1.6 \pm 0.3_{\text{rad}}$	-
$J/\psi, \psi(2S)$	3.08 – 3.11	$7.4 \pm 0.4 \pm 0_{\text{rad}}$	-
$R$ [data]	2.0 – 5.0	$33.9 \pm 1.7_{\text{exp}} \pm 0_{\text{rad}}$	-
$R$ [QCD]	5.0 – $\infty$	$9.9 \pm 0.2_{\text{theo}}$	-
Sum	$2m_\pi - \infty$	$684.7 \pm 6.0 \pm 3.6_{\text{rad}}$	$701.9 \pm 4.7 \pm 1.2_{\text{rad}} \pm 3.8_{\text{SU}(2)}$

# Discussion

The problem of the  $\pi^+\pi^-$  contribution [shifts given in units of  $10^{-10}$ ]:

- **Experimental conspiracy:**

- ▶ new CMD-2 data produce downward shift  $[-1.9]$ , with much better precision
- ▶ new ALEPH BRs produce upward shift  $[+3.5]$
- ▶ CLEO spectral functions produce upward shift  $[+2.2]$
- ▶ previous difference was:  $\Delta[\tau - e^+e^-] = (11 \pm 15) 10^{-4} \rightarrow$  we could use average

- **Who is wrong ?**

- ▶  $e^+e^-$  is consistent with among experiments, but error dominated by CMD-2; large radiative corrections applied
- ▶  $\tau$  is consistent with among experiments, but error dominated by ALEPH
- ▶ **SU(2) corrections:** basic contributions identified and stable since long; overall correction applied to  $\tau$  is  $(-2.2 \pm 0.5)\%$ , dominated by uncontroversial short distance piece; additional long-distance corrections found to be small

**At present, we believe that it is inappropriate to combine  $\tau$  and  $e^+e^-$**

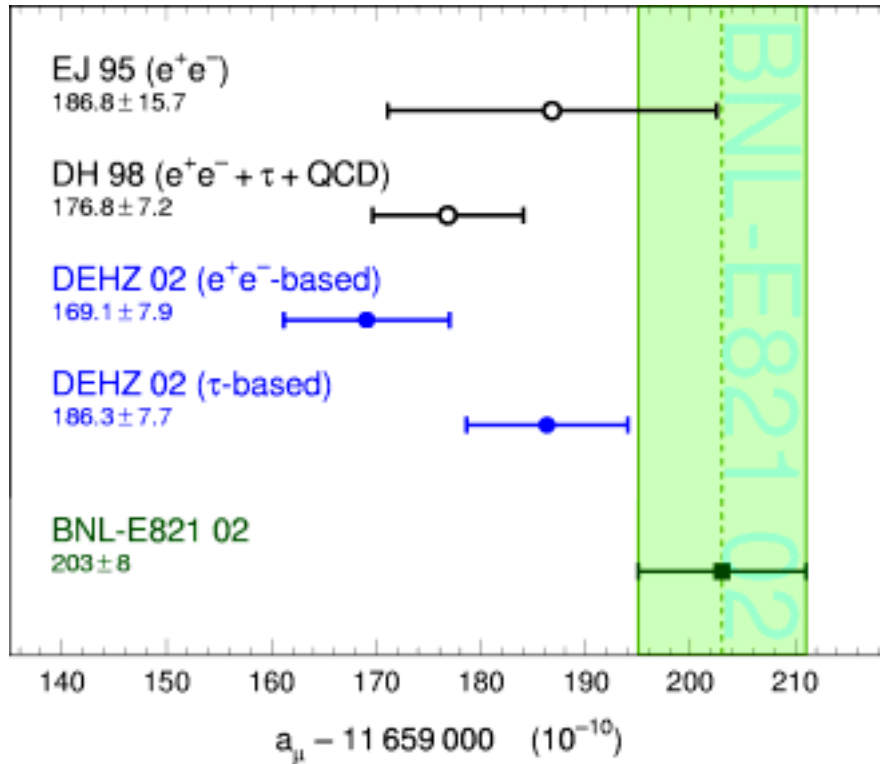
Other changes with respect to DH'98 analysis:

- **Experimental Contribution from  $\phi$  resonance much smaller**  $[-4.3]$
- **Constraints from isospin lead to reduced contributions from  $6\pi$  modes**

# Final Results

[DEHZ'02]	$a_\mu^{\text{had}} [ee] = (684.7 \pm 7.0) 10^{-10}$	(exp and theo errors added in quadrature)	$692.4 \pm 6.2$ [DH'98]
	$a_\mu^{\text{had}} [\tau] = (701.9 \pm 6.2) 10^{-10}$	$\Delta[\tau - e^+e^-] = (0.62 \pm 0.29)\%$	
	$a_\mu [ee] = (11\,659\,169.1 \pm 7.0_{\text{had}} \pm 3.5_{\text{LBL}} \pm 0.3_{\text{QED+EW}}) 10^{-10}$		
	$a_\mu [\tau] = (11\,659\,186.3 \pm 6.2_{\text{had}} \pm 3.5_{\text{LBL}} \pm 0.3_{\text{QED+EW}}) 10^{-10}$		

including:	<b>Hadronic contribution from higher order</b> : $a_\mu^{\text{had}} [(\alpha_s/\pi)^3] = -(10.0 \pm 0.6) 10^{-10}$
	<b>Hadronic contribution from LBL scattering</b> : $a_\mu^{\text{had}} [\text{LBL}] = + (8.6 \pm 3.5) 10^{-10}$



## Observed Discrepancy:

$$a_\mu [\text{exp}] - a_\mu [\text{SM}] = \begin{cases} 34 \pm 11 & [e^+e^-] \\ 17 \pm 11 & [\tau] \end{cases} (10^{-10})$$

Effect on  $\Delta\alpha_{\text{had}}(M_Z^2)$ :

$\Delta[\tau - e^+e^-] = (2.37 \pm 0.62) 10^{-4}$

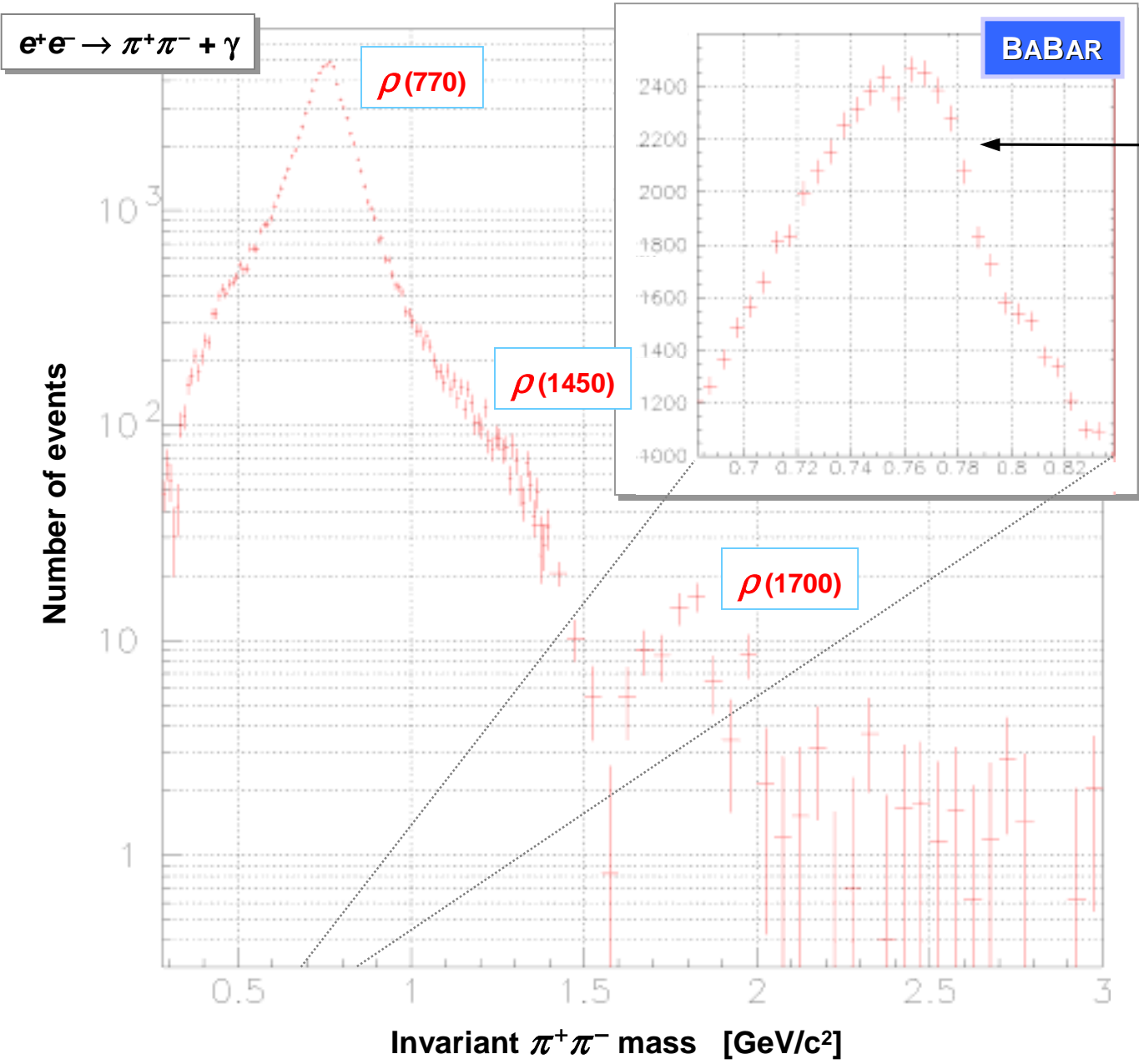
$\rightarrow \Delta M_{\text{Higgs}} \approx -15 \text{ GeV}/c^2$  for  $\tau$

# Conclusions/Perspectives

- Hadronic vacuum polarization creates dominant systematics for SM predictions of many precision measurements
- New analysis of leading hadronic contribution motivated by new, precise  $e^+e^-$  (0.6% systematic error for  $e^+e^-$ ) and  $\tau$  (0.5% error on normalization) data
- New theoretical analysis confirmed the rules to correct for SU(2) breaking
- Radiative (VP and FSR) corrections in  $e^+e^-$  are major source of systematics
- We have re-evaluated all exclusive and inclusive as well as resonance contributions
- We conclude with two incompatible numbers from  $e^+e^-$  and (mainly)  $\tau$ , leading to SM predictions that differ by **3.0  $\sigma$  [ $e^+e^-$ ]** and **1.6  $\sigma$  [ $\tau$ ]** from the experiment

- The key problem is the quality of the experimental data...
- Future experimental projects are:
  - PEP-N (SLAC):  $e^+e^- \rightarrow$  hadrons between 1.4-3.1 GeV ★
  - ISR production  $e^+e^- \rightarrow$  hadrons +  $\gamma$  @ KLOE, BABAR, CLEO & BES as  $\tau$  /charm factories ➔ (systematics?)

# Proof of Principle of ISR Method



Finite mass resolution visible in sharp  $\rho$ - $\omega$  interference. Requires unfolding.

- The data correspond to an integrated luminosity of only 22 fb<sup>-1</sup>.
- Background from  $e^+e^- \rightarrow \mu^+\mu^- + \gamma$  events is at the level of 1%.
- The present statistics (>80 fb<sup>-1</sup>) is more than competitive with the latest results from CMD-2.
- The main concern is the control of the systematics, in particular for particle identification.

# Backup Slides

how could we further improve  $a_{\mu}^{\text{had}}$ ,  
if there were no incompatibilities ?

# Extend the Use of QCD

- The inclusion of  $\tau$  data has **reduces** the error on  $a_\mu^{\text{had}}$  by 40%
- However, **no** significant improvement for  $\Delta\alpha^{\text{had}}(M_Z)$

...how can we further improve – in particular:  $\Delta\alpha^{\text{had}}(M_Z)$  ?

**HINT:**

Inclusive hadronic  $\tau$  decays have shown that QCD is safely applicable at  $m_\tau \sim 1.8 \text{ GeV}/c^2$

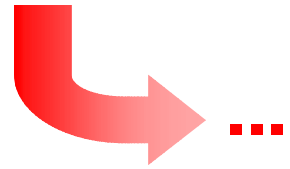
ALEPH (1993, 1998)  
CLEO (1994)  
OPAL (1998)

⇒ **Why not also for  $e^+e^-$  ?**

**TEST:**

Do equivalent QCD analysis as for  $\tau$  decays using spectral moments of  $\sigma[e^+e^- \rightarrow \text{hadrons}](s)$

$$R_{e^+e^-}(s_0) = \frac{1}{2\pi i} \oint_{|s|=s_0} \frac{ds}{s} D(s)$$



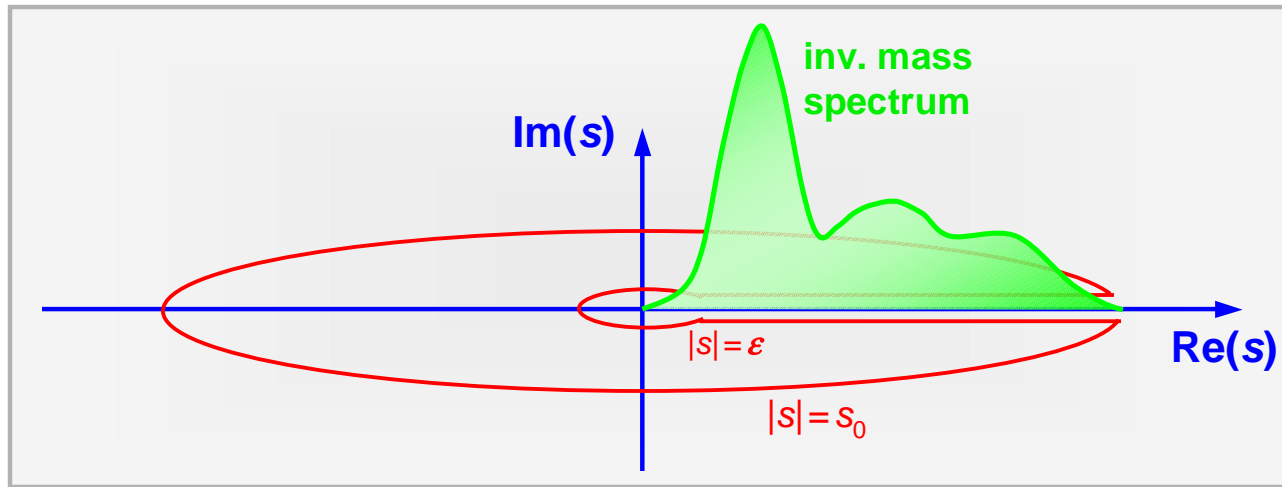
The combined fit of spectral moment to  $e^+e^-$  data showed consistent OPE prediction and revealed small non-perturbative contributions at  $s_0=1.8 \text{ GeV}$



# Spectral Functions and QCD

- (1) Optical theorem  $v(s) \propto \text{Im} \Pi(s)$
- (2) Apply Cauchy's theorem for "save" (i.e., sufficiently large)  $s_0$ :

$$R(s_0) \propto \int_0^{s_0} ds \underbrace{f(s)}_{\text{kinematic factor}} \text{Im} \Pi(s) \quad \Leftrightarrow \quad -\frac{1}{2i} \oint_{|s|=s_0} ds f(s) \Pi(s)$$



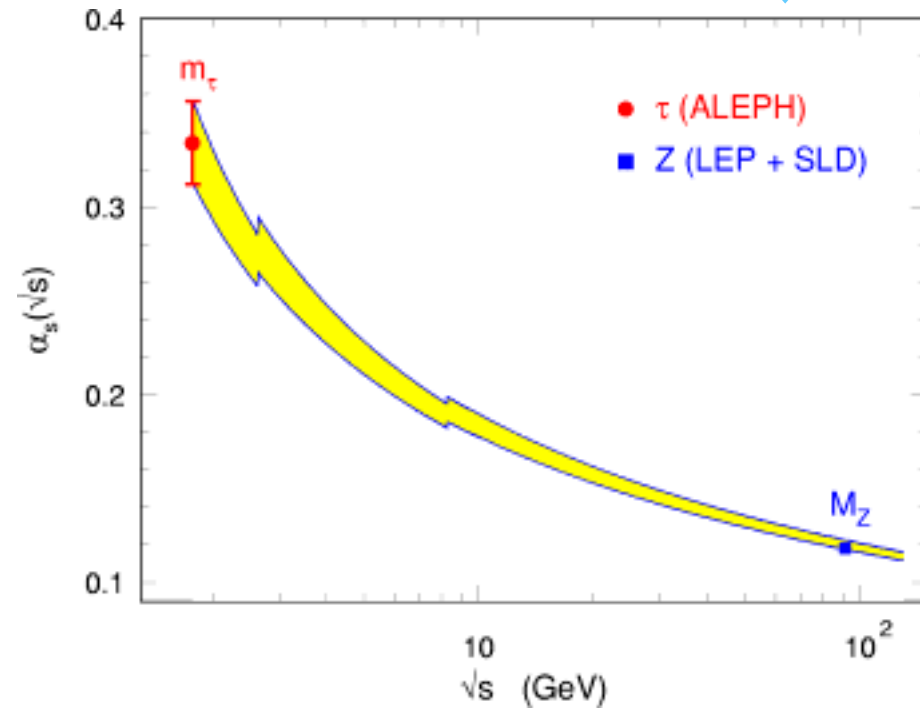
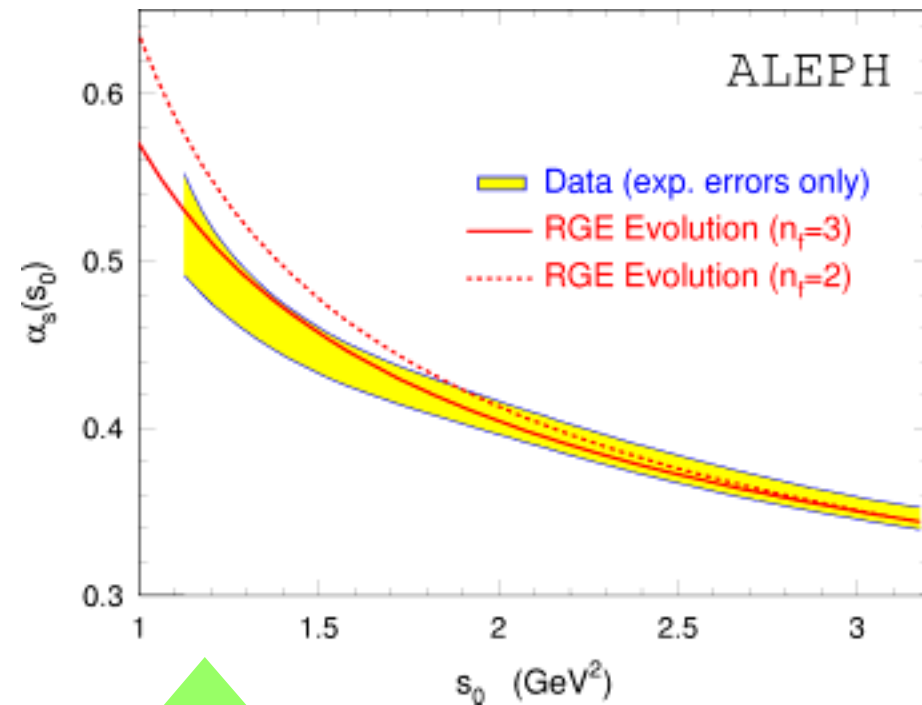
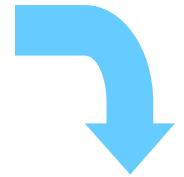
- (3) Use the Adler function to remove unphysical subtractions:  $D(s) = s \frac{d\Pi(s)}{ds}$
- (4) Use global quark-hadron duality in the framework of the Operator Product Expansion (OPE) to predict:  $D(s) \sim D_{\text{pert}}(s) + D_{\text{q-mass}}(s) + D_{\text{non-pert}}(s)$
- (5) Use analytical moments  $f_n(s) = f(s) \cdot \text{poly}_n(s)$  to fix non-perturbative parameters of the OPE ... and then fit  $\alpha_s(m_\tau)$

# QCD Results from $\tau$ Decays

Evolution of  $\alpha_s(m_\tau)$  to  $M_Z$  using RGE (4-loop QCD  $\beta$ -function & 3-loop quark flavor matching) shows the excellent compatibility of  $\tau$  result with EW fit:

$$\alpha_s(M_Z) = 0.1202 \pm 0.0027 \quad (\text{ALEPH'98, theory dominated})$$

$$\alpha_s(M_Z) = 0.1183 \pm 0.0027 \quad (\text{LEP'00, statistics dominated})$$



The  $\tau$  spectral function allows to directly measure the running of  $\alpha_s(s_0)$  within  $\sqrt{s_0} \in [\sim 1 \dots 1.8 \text{ GeV}]$

# Data Driven QCD Sum Rules

- The use of  $\tau$  data and extended QCD **reduces** the error on  $a_\mu^{\text{had}}$  by **51%**
- and provides a **60% (!)** reduction of the error on  $\Delta\alpha^{\text{had}}(M_Z)$

...however, there exist domains for which no theoretical constraints are used so far... we can thus still increase our information budget !

**USE:** Modified Dispersion Integral:

$$\int_0^{s_0} ds f \cdot \text{Im} \Pi = \int_0^{s_0} ds [f - P_n] \text{Im} \Pi + \frac{i}{2} \oint_{|s|=s_0} ds P_n \cdot \Pi$$

The diagram illustrates the Modified Dispersion Integral equation. The left side is the full integral. The right side is split into two parts: a blue-shaded box containing the integral of  $[f - P_n] \text{Im} \Pi$  from 0 to  $s_0$ , with a blue arrow pointing down to the word "Data"; and a red-shaded box containing the contour integral  $\frac{i}{2} \oint_{|s|=s_0} ds P_n \cdot \Pi$ , with a red arrow pointing down to the word "Theory".

where the  $P_n$  are (analytic) polynomials that approximate the kernel  $f(s)$  and thus reduce the data piece of the integral replaced by known theory. No new assumptions !

**An optimization procedure minimizes the total experimental and (conservative) theoretical error**