The Flat Decay Phase in the Early X-ray Afterglows of Swift GRBs Jonathan Granot KIPAC @ Stanford Collaborators: A. Königl, T. Piran, P. Kumar,

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Includes work from the papers:

- Granot, Königl & Piran 2006 (astro-ph/0601056)
- Granot & Kumar 2006, MNRAS, 366, L13-L16
- Eichler & Granot 2006, ApJ, 641, L5-L8
- Nousek et al. 2006, ApJ, 642, 389-400

GLAST Lunch Seminar, May 25, 2006, SLAC

Outline of the Talk:

Short description of Swift & its main findings Swift finds a flat decay phase in the early X-ray afterglow of many GRBs: F_x < was expected Possible Explanations: Energy injection into the afterglow shock viewing angle slightly outside emitting region Afterglow efficiency increases with time ♦ A two component jet model Possible implications: The efficiency of the gamma-ray emission The kinetic energy in the afterglow shock Conclusions

Swift Observatory launched: November 20, 2004



Orbit 600 km x 28° inclination

Data Downlinks TDRSS rapid (2 kbps) ASI Malindi gnd station

Operations

Ops Center @ Penn State Science Center @ GSFC

BAT New CdZnTe detector ~100 GRBs/yr

XRT

Arcsec GRB positions CCD spectroscopy

UVOT

Sub-arcsec positions Grism spectroscopy

Spacecraft Autonomous slews 20-75s

Swift's Main Discoveries so far: Short-hard GRBs First detection of their afterglows • redshifts ($z \sim 0.2-0.7$) \Rightarrow energy & event rate Host galaxies (different from long-soft GRBs) Constraints on contemporaneous supernova Early afterglow emission ◆ Flat decay phase in many X-ray light curves ♦ X-ray flares in many cases Dim early optical emission (reverse shock) **GRBs** at high redshifts $(z = 6.29) \Rightarrow$ cosmology Resolved rise time of SGR 1806-20 giant flare

Early X-ray Afterglows from Swift:





DIVIC L'APICHAUTURS IVI L'AITY I ICC Decay **1. Energy Injection into the Afterglow Shock Type I**: continuous relativistic wind: • The source must remain active for up to a few hours with $L \propto -t^{-0.5}$ ◆ might be possible for a magnetar (Usov 1992)



The magnetar model is able to reproduce the right energy scale and time scale

$$E \approx 5 \times 10^{52} \left(\frac{\Omega}{10^4 \,\mathrm{s}^{-1}}\right)^2 \,\mathrm{erg}$$

$$\tau \approx 20 \left(\frac{\Omega}{10^4 \,\mathrm{s}^{-1}}\right)^{-2} \left(\frac{B}{3 \times 10^5 \,\mathrm{G}}\right)^{-2} \,\mathrm{s}$$

It has many potential problems:

How to launch a collimated relativistic jet

♦ What sets the GRB duration & energy

Hard to explain the observed diversity

♦ How to avoid collapse into a black hole

Type II Energy Injection:

The outflow is ejected within the duration of the GRB (typically tens of seconds)

It naturally has some spread in Lorentz factors

- The outflow naturally sorts itself out in order of increasing velocity faster ahead of slower
- Slower outflow takes longer to catch up with the afterglow shock & deposit its energy
- gradual energy injection (Sari & Mészáros '00) until the slower outflow that carries most of the energy catches up with the afterglow shock

has predictions for the evolution of the broad band spectrum & the light curves



- 1. Unperturbed ext. medium
- 2. Shocked external medium
- 3. Shocked ejecta
- 4. Freely expanding ejecta





2. Viewing angle outside the emitting region Smaller F, f, E_{peak}? t (days) 10^{2} 10 0.50, 0=0 ,0.50 Hydrodynamic simulation of an -10 $\gamma \sim 3$ initially uniform sharp edged jet F_v (mJy) 1.50₀ $\gamma \sim 50$ 10⁻⁵ 200 θ_0 10-6 300 θ_{obs} 50 10 uniform jet with sharp edges 10⁰ (ylm) 10 10 10 LL. 10 10⁻⁵ 10⁵¹ "thick ring" jet: $\theta_c = \Delta \theta = 0.05$, $\Gamma_0 = 100$ 10° $\theta_{obs} = 0, 0.025, 0.0325, 0.04, 0.05, 0.075, 0.$ 10 0.1. 0.11. 0.1175. 0.125. 0.15. 0.2 Gaussian jet: ε , $\Gamma_0 - 1 \propto e^{-\theta^2/2\theta_0^2}$ 10⁵⁰[10⁰ 10 $\beta_{\rm X} = (p-1)/2$ (vlm) 10 10⁴⁹ 10 10 tail of prompt 10⁻³1 erg/s emission ∝ ť 10^{-6} └-X,tail 10⁴⁸ 10 10⁻⁵ 10⁻² 10^{2} 10^{-1} 10^{0} 10¹ 10 os 10^{47'} X (days) 10 10⁰ Gaussian ϵ + constant Γ_0 10^{46|} GRB 050315: fits to a Gaussian jet (MJV) 10 $(\theta_{obs} = \theta_c = 3^\circ, E = 3 \times 10^{51} \text{ erg}, \Gamma_0 = 300.$ 10 10 10^{45} n = 15 cm⁻³, p = 2, $\varepsilon_{\rm p}$ = 0.3, $\varepsilon_{\rm B}$ = 0.05) and to a 10 thick ring shaped jet ($\theta_c = \Delta \theta = 0.025$, $\theta_{obs} = 0.44\theta_c$, 10⁻⁵ 10⁴⁴ $\rho = Ar^{-2}$, $A_* = 1$, $E = 5 \times 10^{51}$ erg, p = 2, $\varepsilon_{p} = 0.17$, $\varepsilon_{B} = 0.01$ 10^{-2} 10⁻¹ 10⁰ 10¹ 10² 10^{2} 10^{3} 10^{4} 10^{6} 10⁵ 10¹ Granot (Eichler & t/(1+z) in seconds t (days)

Implications of viewing angle interpretation:

 If the regions of prominent gamma-ray and afterglow emission coincide, then a smaller typical photon energy, fluence and peak flux are expected: X-ray flashes (Yamazaki et al. '02,03,04)

This does **not show up** in the **data**: it might suggest that **these two regions do not coincide**

 Along some lines of sight there is a bright gamma-ray emission but dim afterglow
 emission ⇒ high gamma-ray efficiency and/or low afterglow efficiency 3. Increase in afterglow Efficiency: (Granot, Königl & Piran 2006)

The energy in the afterglow shock is constant The emission is from along our line of sight The afterglow efficiency initially increases with time, due to a change in one or more of: $\bullet \epsilon_{e} =$ fraction of energy in relativistic electrons $\bullet \epsilon_{\rm B}$ = fraction of energy in the magnetic field $\mathbf{\xi}_{e} =$ fraction of electrons that are relativistic The shock microphysics parameters eventually saturate at some asymptotic values This may have interesting implications for relativistic collisionless shocks

Inferring the Kinetic Energy: $F_{\nu>\max(\nu_m,\nu_c)} \propto \xi_e^{2-p} (1+Y)^{-1} \varepsilon_e^{p-1} \varepsilon_B^{(p-2)/4} E_{k,iso}^{(p+2)/4} t^{(2-3p)/4} v^{-p/2}$ $E_{k,iso} \propto \xi_e^{4(p-2)/(p+2)} (1+Y)^{4/(p+2)} \varepsilon_e^{-4(p-1)/(p+2)} \varepsilon_B^{-(p-2)/(p+2)}$ Flux independent of external density ($v > v_m, v_c$) **Weak dependence on \varepsilon_{\rm B} which is very uncertain E**_{k.iso} inferred from $F_X(10hr)$ pre-*Swift* (for $\xi_e=1$) **Problems with this method: Typically** $\varepsilon_{\rm R} \ll \varepsilon_{\rm e} \Longrightarrow (1+Y) \sim (\varepsilon_{\rm e}/\varepsilon_{\rm R})^{1/2}$, $F_v \propto (\varepsilon_e)^{-2(2p-3)/(p+2)} (\varepsilon_B)^{-p/(p+2)} (\text{larger } \varepsilon \text{-dependence})$ **There is a degeneracy: the observed** F_v is not effected by $(E,n) \rightarrow (E,n)/\xi_e$, $(\varepsilon_e,\varepsilon_B) \rightarrow \xi_e(\varepsilon_e,\varepsilon_B)$ for $m_e/m_p < \xi_e \le 1$ (Eichler & Waxman 2005)

Implications for γ -ray Efficiency $= \varepsilon_{\gamma} = E_{\gamma}/E_0, \varepsilon_{\gamma}/(1 - \varepsilon_{\gamma}) = \kappa f; \kappa = E_{\gamma}/E_k(t), f = E_k(t)/E_{k,0}$ $\mathbf{k} \sim 1$ from the X-ray afterglow flux at t = 10 hr **f** \geq 10 if the early flat decay phase is interpreted as energy injection into afterglow shock: $\varepsilon_{\gamma} \ge 0.9$ Such high efficiencies are very hard to produce If the flat decay phase is due to an increase in the afterglow efficiency then $f \sim 1 \& \epsilon_{y} \sim 0.5$ If also $E_k(t = 10 \text{ hr})$ is underestimated (e.g., $\xi_e \sim$ 0.1 instead of 1) then possibly $\kappa \sim 0.1 \& \varepsilon_{\gamma} \sim 0.1$ $\blacksquare \Rightarrow$ a typical afterglow kinetic energy $\ge 10^{52}$ erg $(\geq 10^{53} \text{ erg})$ for a uniform (structured) jet

4. A Two Component Jet: (Granot, Königl & Piran 2006) wide jet: Γ₀ ~ 20-50 → narrow jet: Γ₀ > 100 > observer

- In the collapsar model the wide jet is produced by the cocoon (Ramirez-Ruiz et al. 2002)
- Neutron decoupling in a hydromagnetically driven neutron rich jet (Vlahakis et al. 2003)

t_{dec} ∝ Γ₀-^{2(4-k)/(3-4)} for ρ_{ext} ∝ r^{-k} ⇒ t_{dec,n} ≪ t_{dec,w}
 E_{iso,w} < E_{iso,n} & E_w > E_n ⇒ lowers the required γ-ray efficiency ε_v (Peng, Königl & Granot 2005)

Afterglow light curves for a two component jet:

 $E_w > E_n$ is required for the wide jet to dominate the flux at late times

 E_{iso,w} < E_{iso,n} is in order to lower the required ε_γ



(Peng, Königl & Granot 2005)



(Granot, Königl & Piran 2006)

Simultaneous Optical light curves
 The optical light curves do not show a similar flat decay phase or a similar steepening at its end
 This is problematic for most models



Conclusions:

There are different possible explanations for the early flat decay phase in the X-ray afterglows the different explanations may potentially be distinguished via multi-wavelength observations There are potentially interesting implications for The gamma-ray efficiency ◆ The afterglow kinetic energy Collisionless relativistic shocks A lot of work still remain to be done