Radiative correction from the Wigner's little group perspective

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Radiative calculation

Outline

▷ Introduction

- Heavy ion collision
- Aim of the work
- \triangleright Wigner's Little group W
 - Action of a W on a state
 - Helicity representation
- Radiative calculation
 - Kinematic setup
 - Soft factor radiation

⊳ Conclusion



Today

Life on earth



4 hillion years



Introduction





▷ Heavy Ion Collision:

- Study collectivity and thermodynamics of QCD.
- Collide heavy nuclei at relativistic energies (RHIC, LHC).
- Creating a small volume of matter with extreme density and temperature.

> Experimental Observation:

- Thermalized system: collective response to pressure gradients.
- High transverse momentum processes: predicted from pQCD.





▷ Hard Probes:

- The outgoing partons propagates through the medium.
- Understand the properties of hot and dense QCD matter.
- A dramatic decrease of the jet energy, jet quenching.
- \triangleright Nuclear modification R_{AA}
 - Ratio: σ_{AuAu} and σ_{pp} .
 - O Deviation from 1?
 - R_{AA}^{γ} vs R_{AA}^{jet}





▷ Hard Probes:

- The outgoing partons propagates through the medium.
- Understand the properties of hot and dense QCD matter.
- A dramatic decrease of the jet energy, jet quenching.
- ▷ Energy loss mechanisms:
 - Collisionnal energy loss: multiple scattering center, drag coefficient.
 - Radiative energy loss: soft/collinear photon and gluon radiation.





Radiative calculation:

- Effects of radiation (γ , g) on a jet through the medium.
- Radiation without explicit details on the hard collisions!
- Soft contributions can be factorized from hard collisions.



 $\circ\,$ Hard collision $\sim 2 \rightarrow 2$ collisions between the probe and the constituents of the bulk.

Wigner's Little Group



- ▷ Little group W: is a subgroup of the Lorentz group that leave a momentum p invariant (Wp = p).
- \triangleright Single particle state (m=0): $\hat{U}(W) \ket{p,h} = e^{i\theta h} \ket{p,h}$.
- ▷ Massless: $W \cong U(1)$ generated by $\hat{L} | p, h \rangle = h | p, h \rangle$.
- Scattering amplitudes





Massless spinor representation

$$p_{\mu} \xrightarrow{p^2 = 0} \lambda_a \bar{\lambda}_{\dot{a}}$$

 \triangleright The little group W becomes a simple scaling of λ and $\bar{\lambda}$

$$W: \begin{cases} \lambda \longrightarrow t \, \lambda \\ \bar{\lambda} \longrightarrow t^{-1} \bar{\lambda} \end{cases}$$

 \triangleright Spinor representation of the generator \hat{L}

$$\hat{L} = -\frac{1}{2} \left(\lambda_a \frac{\partial}{\partial \lambda_a} - \bar{\lambda}_{\dot{a}} \frac{\partial}{\partial \bar{\lambda}_{\dot{a}}} \right)$$

 $Dash \hat{L}$ act on any function $f(\lambda, \bar{\lambda})$ as a simple derivative.



Little group action on the factorized amplitude

$$\hat{L}^{(j)} \xrightarrow{p_3 \ p_1}_{p_1 \ p_2 \ p_1 \ p_2 \ p_2 \ p_2 \ p_3 \ p_4 \ p_4$$

Helicity constraints on the soft factor

$$\begin{cases} \hat{L}^{(s)}H(s;\{i\}) = h_s H(s;\{i\})\\ \\ \hat{L}^{(i)}H(s;\{i\}) = 0 \end{cases}$$

 \triangleright Solving system of (n+1) linear equations.



> Little group action on the factorized amplitude

$$\hat{L}^{(j)}_{p_2} \xrightarrow{p_1}{p_1 \atop k_s} \simeq \hat{L}^{(j)} H(s;\{i\}) \times \stackrel{p_3}{\underset{p_2}{\longrightarrow}} \stackrel{p_4}{\underset{p_1}{\longrightarrow}} + H(s;\{i\}) \times \hat{L}^{(j)}_{p_2} \xrightarrow{p_3} \stackrel{p_4}{\underset{p_1}{\longrightarrow}} + H(s;\{i\}) \times \hat{L}^{(j)}_{p_2}$$

Helicity constraints on the soft factor

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 \triangleright Solving system of (n+1) linear equations.

Radiative Calculation

Radiative calculation

▷ Kinematic setup

• Parton medium interactions are dominated by *t*-channel exchange.

Radiative calculation

- The interactions are characterized by
- $H(s; \{i\})$ of an helicity -1 radiation depends only on λ , i.e.

 $\lambda_a^{(3)} \sim \lambda_a^{(4)} \sim \lambda_a^{(q)}$.

- $\langle ij\rangle = \lambda^{(i)}\cdot\lambda^{(j)}$
- Solving the helicity constraint for a single and a double radiations.







gluon exchange





▷ Single soft radiation

$$H(s) = A \frac{\langle 14 \rangle}{\langle 1s \rangle \langle s4 \rangle} + B \frac{\langle 42 \rangle}{\langle 4s \rangle \langle s2 \rangle}$$

• A and B contain information on the non-kinematic dof.

- A and B can be fixed from the symmetry of the theory.
- Abelian solution: $A = B \sim e$

$$H^{\gamma}(s) \sim e \frac{\langle 12 \rangle}{\langle 1s \rangle \langle s2 \rangle} \Longleftrightarrow e \left(\frac{p_1^{\mu}}{p_1 \cdot k_s} - \frac{p_2^{\mu}}{p_2 \cdot k_s} \right) \mathcal{E}_{\mu}(k_s)$$

Non-Abelian solution: color degrees of freedom

$$A(a_s) \sim gT_{a_s}\hat{C}_R$$
 and $B(a_s) \sim g\hat{C}_R T_{a_s}$

Wigner's Little Group

Radiative calculation

Conclusion

Radiative calculation

▷ Two soft radiation

$$H(s_1, s_2) = H(s_1)H(s_2) + \frac{\varepsilon_{\mathrm{NA}}}{\langle s_1 s_2 \rangle} R(s_1, s_2)$$

- $\circ \varepsilon_{\scriptscriptstyle\rm NA}$ is a non-Abelian effect between the radiations.
- $\circ R(s_1, s_2)$ leads to the correlation between the two radiations.
- Suggestive to the correlation from central p + Pb collisions measured by the ALICE collaboration.





Conclusion

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Summary

- Jet energy loss are characterized by collisional and radiative mechanisms.
- ▷ Study the induced radiative effect on parton through a medium.
- \triangleright From the scale factorization, we introduce the helicity constraints from the *W* action on *H*(*s*).
- Presented some early result for one and two spin-1 radiations two compare to the ALICE correlation measurement.

Outlook

- \triangleright Consider massive parton with $W \cong SU(2)$.
- ▷ Deep investigation on the kinematics for multiple exchange.
- ▷ Combine with the existing model (GLV, ASW, ...)