

# Nature of the $X(5568)$ : a critical Laplace sum rule analysis at N<sup>2</sup>LO

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In collaboration with R.M. Albuquerque - S. Narison - D. Rabetiarivony

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- We will also extend our results to the charm channel.

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- ⑤ Results an Summary

# Exotic mesons

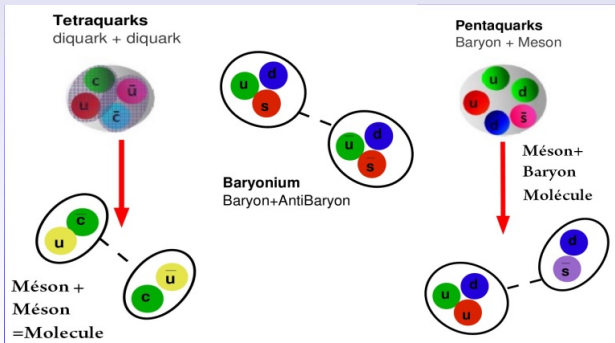
- ① Hadrons = Mesons (1 quark + 1 anti-quark  $q_1\bar{q}_2$ ) or Baryons (3 quarks  $q_1q_2q_3$ )
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- 3 Observables particles must be **color singlet**
- 4 QCD allow other configurations



- 1 High Energy  $\rightarrow$  Perturbative QCD



# Spectral Functions

- ① High Energy  $\rightarrow$  Perturbative QCD
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- 1 High Energy  $\rightarrow$  Perturbative QCD
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QCD spectral Sum Rules
- 3 Two point function :

$$\begin{aligned}\Pi_{mol}^{\mu\nu}(q) &= i \int d^4x e^{iq \cdot x} \langle 0 | T J^\mu(x) J^{\nu\dagger}(0) | 0 \rangle \\ &= -(q^2 g^{\mu\nu} - q^\mu q^\nu) \Pi_{mol}^{(1)}(q^2) + q^\mu q^\nu \Pi_{mol}^{(0)}(q^2)\end{aligned}$$

The 2-point function can be evaluate in two ways :

- QCD side : OPE (Operator Product Expansion)
- Phenomenological side  $\rightarrow f_H$  : the decay constant and it parametrize the coupling of the hadron to the current.

Wilson expansion :

$$\Pi^{\mu\nu}(q) = i \int d^4x e^{iq \cdot x} \langle 0 | J(x) J^\dagger(0) | 0 \rangle = \sum_n C_n(q^2) \langle 0 | : \mathcal{O}_n(0) : | 0 \rangle$$

Wick theorem : Normal ordered product + All contractions

$$: \mathcal{O}_3 : = \bar{q}(0)q(0) \Rightarrow \langle \bar{q}q \rangle$$

$$: \mathcal{O}_4 : = g_s^2 G_{\alpha\beta}^N(0) G_{\alpha\beta}^N(0) \Rightarrow \langle g_s^2 G^2 \rangle$$

$$: \mathcal{O}_5 : = \bar{q}(0) : g_s^2 \sigma^{\alpha\beta} G_{\alpha\beta}^N(0) q(0) \Rightarrow \langle \bar{q}Gq \rangle$$

$$: \mathcal{O}_6^q : = \bar{q}(0)q(0)\bar{q}(0)q(0) \Rightarrow \langle \bar{q}q \rangle^2$$

$$: \mathcal{O}_6^G : = f_{NMK} g_s^3 G_{\alpha\beta}^N(0) G_{\beta\gamma}^N(0) G_{\gamma\alpha}^N(0) \Rightarrow \langle g_s^3 G^3 \rangle$$

$$\rho^{OPE}(t) = \rho_0(t) + \rho_3(t) \langle \bar{q}q \rangle + \rho_4(t) \langle g^2 G^2 \rangle + \rho_5(t) \langle \bar{q}Gq \rangle + \dots$$

Dispersion relation :

$$\Pi^{OPE}(q) = \int_{t_0}^{+\infty} dt \frac{\rho^{OPE}(t)}{t - q^2}$$

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- $\sum H_q \equiv H_0 + \sum H'$  where  $H_0$  represent the ground state and  $H'$  the excited ones.

The dispersion relation become :

$$\Pi^{PHEN}(q) = \frac{\lambda^2}{M_H^2 - q^2} + \int_{t_c}^{+\infty} dt \frac{\rho^{OPE}(t)}{t - q^2}$$

Quark-hadron duality principle :

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Excited states  $\implies$  Inverse Laplace Sum Rules  $\implies$  masses,  
couplings

Molecule :

- $BK$   $0^+$   $(\bar{b}i\gamma^5u)(\bar{d}i\gamma^5s)$
- $B_s\pi$   $0^+$   $(\bar{b}i\gamma^5s)(\bar{d}i\gamma^5u)$
- $B^*K$   $1^+$   $(\bar{b}\gamma^\mu u)(\bar{d}i\gamma^5s)$
- $B_s^*\pi$   $1^+$   $(\bar{b}\gamma^\mu s)(\bar{d}i\gamma^5u)$

Four-quark :

- $1^-$   $(s^T C \gamma^5 u)(\bar{b}\gamma^\mu \gamma^5 C \bar{d}^T) + k(s^T C u)(\bar{b}\gamma^\mu C \bar{d}^T)$
- $1^+$   $(s^T C \gamma^5 u)(\bar{b}\gamma^\mu C \bar{d}^T) + k(s^T C u)(\bar{b}\gamma^\mu \gamma^5 C \bar{d}^T)$

# Mass and Coupling

$$\text{Mass : } M_H^2 = \frac{\int_{4m_Q^2}^{t_c} dt t e^{-t\tau} \frac{1}{\pi} \text{Im}\Pi^{OPE}(t)}{\int_{4m_Q^2}^{t_c} dt e^{-t\tau} \frac{1}{\pi} \text{Im}\Pi^{OPE}(t)}$$

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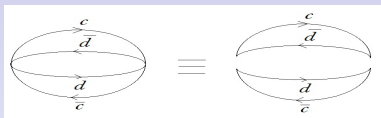
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- Noting that the bilinear (pseudo)scalar current acquires an anomalous dimension due to its normalization, thus the decay constants run to order  $\alpha^2$



The molecular spectral function can be considered as a convolution of two 2-point functions<sup>1</sup>

$$\frac{1}{\pi} \text{Im}\Pi_{mol}^{(0,1)}(t) = \theta(t - 4M_Q^2) \left(\frac{1}{4\pi}\right)^2 t^2 \int_{M_Q^2}^{(\sqrt{t}-M_Q)^2} dt_1 \times \int_{M_Q^2}^{(\sqrt{t}-\sqrt{t_1})^2} dt_2 \times \dots$$

$$\text{For spin 0 : } \dots = \lambda^{1/2} \left[ \left( \frac{t_1}{t} + \frac{t_2}{t} - 1 \right)^2 \right] \frac{1}{\pi} \text{Im}\Pi^{(0)}(t_1) \frac{1}{\pi} \text{Im}\Pi^{(0)}(t_2)$$

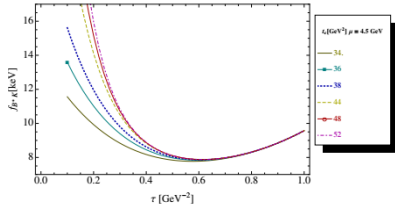
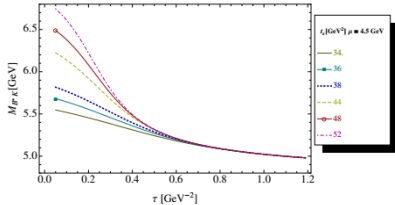
$$\dots = \lambda^{3/2} \frac{1}{\pi} \text{Im}\Pi^{(1)}(t_1) \frac{1}{\pi} \text{Im}\Pi^{(1)}(t_2)$$

$$\text{For spin 1 : } \dots = \lambda^{1/2} \left[ \left( \frac{t_1}{t} + \frac{t_2}{t} - 1 \right)^2 + \frac{8t_1 t_2}{t^2} \right] \frac{1}{\pi} \text{Im}\Pi^{(0)}(t_1) \frac{1}{\pi} \text{Im}\Pi^{(1)}(t_2)$$

$$\rho_{OPE} = \rho_{pert} \left( 1 + \frac{NLO+N2LO}{LO} \right) + \rho_{\langle q\bar{q} \rangle} + \rho_{\langle g^2 G^2 \rangle} + \rho_{\langle qGq \rangle} + \rho_{\langle qq \rangle^2} + \rho_{\langle g^3 G^3 \rangle}$$

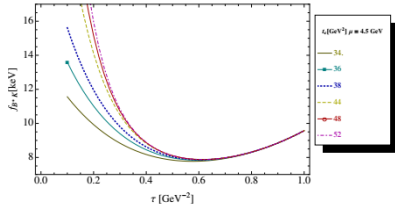
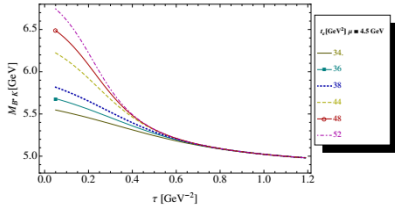
1. A. Pich and E. de Rafael/S. Narison and A. Pivovarov.

# Result for $B^*K$ molecule at LO

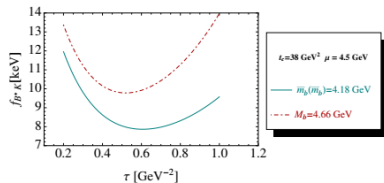
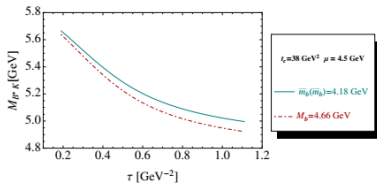


Stability	$t_c$ (GeV $^2$ )	$\tau$ (GeV $^{-2}$ )	Mass (GeV)	Coupling (keV)
	$\geq 34$	0.45 - 0.6	5.19 - 5.20	7.77 - 7.88
	44-48	0.6		

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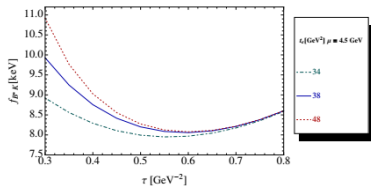
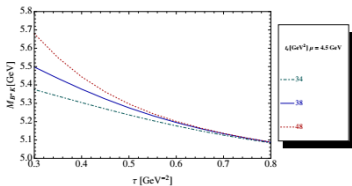
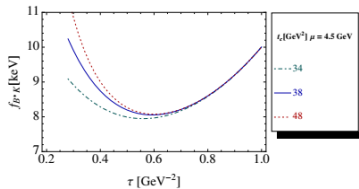
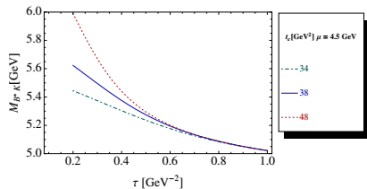


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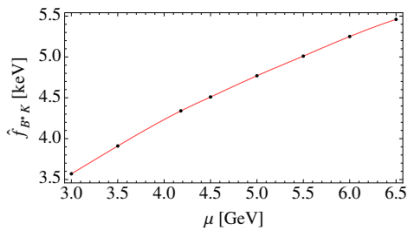
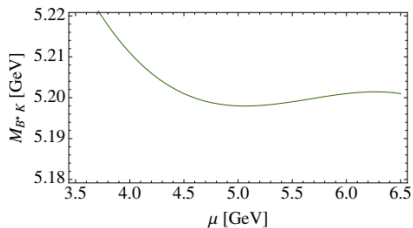
The effect of the definitions (running and pole) of the heavy quark mass used should be added as errors in the LO analysis

# $B^*K$ results at NLO and N2LO

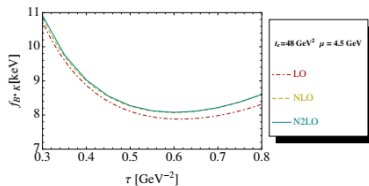
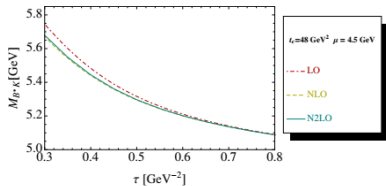


	$t_c$ ( $\text{GeV}^2$ )	$\tau$ ( $\text{GeV}^{-2}$ )	Mass (GeV)	Coupling (keV)
NLO	$\geq 34 - 48$	0.56 - 0.60	5.200 - 5.201	7.95 - 8.07
N2LO	34 - 48	0.58 - 0.62	5.185 - 5.187	7.95 - 8.09

# $B^*K$ results ( $\mu$ ), PT-series

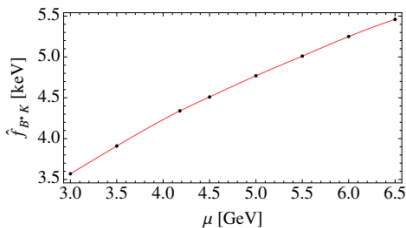
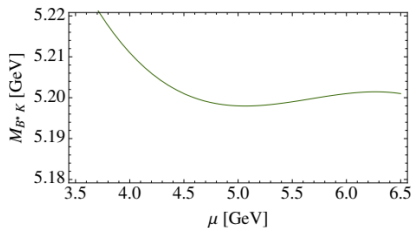


$\mu$ (GeV)	Mass (GeV)	Coupling (keV)
4.5 - 5.0	5.201 - 5.198	4.51 - 4.77

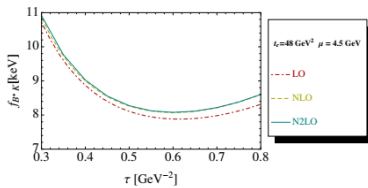
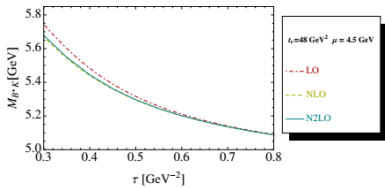




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Good convergence of PT-series

# Final results

Nature	$J^P$	Mass [MeV]	$\hat{f}_X$ [keV]	$f_X(4.5)$ [keV]
<b><i>b</i>-quark channel</b>				
<i>Molecule</i>				
$B^*K$	$1^+$	$5186 \pm 13$	$4.48 \pm 1.45$	$8.02 \pm 2.60$
$BK$	$0^+$	$5195 \pm 15$	$2.57 \pm 0.75$	$8.26 \pm 2.40$
$B_s^*\pi$	$1^+$	$5200 \pm 18$	$5.61 \pm 0.87$	$10.23 \pm 1.59$
$B_s\pi$	$0^+$	$5199 \pm 24$	$3.15 \pm 0.70$	$10.5 \pm 2.30$
<i>Four-quark (su)(<math>\bar{b}\bar{d}</math>)</i>				
$A_b$	$1^+$	$5186 \pm 16$	$5.05 \pm 1.32$	$9.04 \pm 2.37$
$S_b$	$0^+$	$5196 \pm 17$	$2.98 \pm 0.70$	$9.99 \pm 2.36$
<b><i>c</i>-quark channel</b>				
<i>Molecule</i>				
$D^*K$	$1^+$	$2395 \pm 48$	$155 \pm 36$	$226 \pm 52$
$DK$	$0^+$	$2402 \pm 42$	$139 \pm 26$	$254 \pm 48$
$D_s^*\pi$	$1^+$	$2395 \pm 48$	$215 \pm 35$	$308 \pm 49$
$D_s\pi$	$0^+$	$2404 \pm 37$	$160 \pm 22$	$331 \pm 46$
<i>Four-quark (su)(<math>\bar{c}\bar{d}</math>)</i>				
$A_c$	$1^+$	$2400 \pm 47$	$192 \pm 41$	$260 \pm 55$
$S_c$	$0^+$	$2395 \pm 68$	$122 \pm 26$	$221 \pm 47$

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- We do not include higher contributions ( $\geq 7$ ) in our estimate but only consider them as a source of the errors.
- Our previous analysis within stability criteria with respect to  $\tau$ ,  $t_c$  and  $\mu$  have done some successful predictions in different hadronic channels. However, we did not predict a mass of a pure exotic  $BK, B^*K, B_s\pi$  molecule or four-quark state around the D0's  $X(5568)$  wich is not confirmed by LHCb.

- We can see a good convergence after including higher correction, the existence of this convergence confirm the veracity of our results.
- Our analysis including high order PT corrections has given a more meaning on the input value and the definition of heavy quark mass. The ill-defined heavy quark mass definition used at LO is not enough to have better results.
- We do not include higher contributions ( $\geq 7$ ) in our estimate but only consider them as a source of the errors.
- Our previous analysis within stability criteria with respect to  $\tau$ ,  $t_c$  and  $\mu$  have done some successful predictions in different hadronic channels. However, we did not predict a mass of a pure exotic  $BK, B^*K, B_s\pi$  molecule or four-quark state around the D0's  $X(5568)$  which is not confirmed by LHCb.
- From our analysis, one may suggest to scan the regions (2327 - 2444) MeV and (5173 - 5226) MeV for detecting these unmixed exotic hadrons.

# Thank you

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**Nature of the X(5568) — A critical Laplace sum rule analysis at N<sup>2</sup>LO**

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