Top quark production

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ABSTRACT

We review the current status of the cross sections measurement of the top-quark at the LHC and at the Tevatron. Total production cross sections, studies using single top quark events and differential $t\bar{t}$ cross sections are discussed. The associated production of top quark pairs with photons, Z and W bosons, including $t\bar{t}Z$ and $t\bar{t}W$ measurements shown for the first time at LHCP2014, are presented.

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1 Introduction

The physics of the top quark is closely related to some of the most interesting issues that still remain open in the Standard Model (SM). Due to its large mass, among all elementary particles the top quark has the largest couplings to other SM particles, e.g. the Higgs boson, and, potentially, to particles in new physics models. Also, the strength of the top quark coupling to the Higgs boson ($\sim 1$) suggests that the top quark may play a special role in the mechanism of electroweak. Associated production of top pairs with the Higgs boson gives direct access to this coupling - a very important SM ingredient still lacking a direct measurement. These are some of the reasons why top quark production is a key topic on the physics programs of the Tevatron and LHC experiments.

Top quarks are produced in pairs, via strong interaction, or via electroweak interactions leading to “single top quark” final states, either in association with jets (in the $t$-channel or $s$-channel) or with a $W$ boson (in the $tW$ channel). In the SM they decay almost exclusively into a $b$ quark and a $W$ boson. Other allowed decay modes are strongly suppressed. Decay channels of the top quarks are named according to the decay of the $W$ boson: either leptonic (into a lepton and a neutrino) or hadronic (into a quark-antiquark pair).

Differences on the incident hadrons and collision energy at the Tevatron and LHC colliders allow different aspects of top quark production to be tested. At the Tevatron, top quark processes from $p\bar{p}$ collisions at 1.96 TeV are mostly initiated by quark-quark interactions. At the LHC, in $pp$ collision at 7 and 8 TeV, gluon-initiated processes are dominant. Therefore, quark-initiated processes, such as $s$-channel single top quark production are relatively more abundant at the Tevatron, while other mostly gluon-initiated processes such as the $tW$ production, or those requiring higher energies, such as $ttZ$ and $ttW$ associated productions, are not accessible at the Tevatron, being their cross sections small even at the LHC s.c.m. energies.

In this contribution, the most recent measurements on top quark production from the ATLAS [1], CDF, CMS [2] and D0 Collaborations are discussed. Measurements of single top quark production, inclusive $tt$ production in nearly all decay modes, differential $tt$ cross sections and associated production are presented. Some topics, such as jet multiplicity in $tt$ and searches for $ttH$ production, are covered in other contributions to this conference, and are not discussed here.

Along with the experimental efforts, precise NNLO calculations have been developed in the last years. In a fruitful synergy between theory and experiment, data observation leads to a finer tune of the theoretical inputs, such as the proton PDF and QCD renormalization and factorization scales, while experiments profit from instance from the development of more accurate Monte Carlo (MC) programs which are used for the estimation of efficiencies and corrections of the data. High precision both at the experimental and the theoretical sides enhances the possibility of signs of new physics to be observed in data, as deviations from the SM.

2 Single top quark production

Single top quark production is characterized by very small signal yields over an overwhelming background. These competing final states come mainly from $tt$, $W$+jets, $Z$+jets, QCD multijets and diboson events. For this reason, analysis techniques for single top quark cross sections measurements involve the use of event topology for signal to background separation, in addition to optimized event selection. Boosted decision trees (BDT), neural networks or multivariate likelihoods are examples such techniques exploiting the full event kinematics to separate signal from background.

Associated production of a single top quark and a $W$ boson is not visible at the Tevatron, due to very small yields. At the LHC, measurements are performed with both $W$ bosons (from the top decay and the prompt $W$) decaying leptonically, leading to a final state with two leptons and a $b$ jet. With only one additional $b$ jet, the final state for $tt$ events decaying dileptonically is found, with a cross section more than 10 times larger. In a very large amount of $tt$ events, one of the $b$ jets misses reconstruction, and the event is selected as $tW$ candidate. CMS measurement [3], based on a sample of integrated luminosity 12.2 fb$^{-1}$ at 8 TeV, uses events containing exactly two leptons ($ee, e\mu, \mu\mu$) and exactly 1 $b$-tagged jet optimized for the signal. The dominant $tt$ background is estimated using a control region containing one additional $b$-tagged jet. The separation of signal and background is achieved on a BDT analysis using 13 event variables, yielding
to collisions is correlated

\[ \sigma_{\text{tW}} = 23.4 \pm 5.4 \text{ pb} \]

The measurement corresponds to the first observation of single top quark production in the tW channel, with an observed significance of 6.1\( \sigma \) (with 5.4\( \pm 1.4 \sigma \) expected). ATLAS measurement [4], based on 20.3 fb\(^{-1}\) of 8 TeV data, uses two samples, both containing an \( e\mu \) pair, and either 1 or 2 jets. The sample with one jet is used in a BDT with 19 variables for the signal estimation, while the 2 jets sample is used to determine the \( t\bar{t} \) normalization in a BDT with 20 variables. The measurement, \( \sigma_{\text{tW}} = 27.2 \pm 2.8 \text{ (stat) } \pm 5.4 \text{ (syst.) } \text{ pb} \), is an evidence for tW channel production at 4.2\( \sigma \) (with 4\( \sigma \) expected). Both measurements are in agreement with the SM predictions at approximated NNLO, of 22.2 \( \pm 0.6 \) (scale) \( \pm 1.4 \) (PDF) pb.

Single top quark production in the t-channel has been observed long ago both at the Tevatron and at the LHC. Currently, uncertainties on the most precise cross section measurements are around 20\% at 1.96 TeV and 10\% at 7 [5] and 8 TeV. As more data becomes available at the LHC, t-channel events can be used in more detailed studies. For instance, the measurement of the ratio between tops and anti-tops produced in the t-channel was performed by ATLAS [6] and CMS [7] Collaborations. Since the ratio is mainly driven by the relative proportion of quarks up and charm (leading to t quark production) and quarks down and strange (\( \bar{t} \) quark production) in the proton, the measurement is sensitive to the parametrizations of the proton PDF. Fig. 1(left) shows the results for the CMS Collaboration: the measured ratio is compared to SM predictions using several PDF parametrization. ATLAS measurement yields similar results. ATLAS Collaboration studied the dependence of the model used for acceptance correction in the measurement of t-channel cross section [8]. The t-channel generator choice is, together with jet energy corrections uncertainty, the dominant source of uncertainty on the measurement of the t-channel fiducial cross section. Several NLO MC generators were tested, as well as the LO AcerMC generator with different factorization and renormalization scales. Once the fiducial cross section is extrapolated to the full phase space, NLO generators models are in much better agreement to each other, and other uncertainties on the extrapolation become dominant. Results are presented in Fig. 1(right). Although with current precision no particular PDF parametrization or MC model can be excluded, these studies demonstrate the potential of t-channel events in future analyses as stringent tests of the SM, which may reduce systematic uncertainties on top production measurements.

The single top-quark production mode with smallest cross section in the SM, the s-channel, was first measured at the Tevatron in combination with the t-channel. The typical event selection criteria was as generic as possible, requiring 1 lepton, at least 2 jets (from which at least one b tagged) and some \( E_{\text{T}}^{\text{miss}} \). D0 [9] combines three multivariate techniques, using up to 30 variables to separate signal from background. CDF [10] bases its measurement on NN analyses using up to 14 variables. A third measurement from CDF [11] uses an even looser selection criteria, requiring events with \( E_{\text{T}}^{\text{miss}} \), 2 or three jets, but vetoing isolated leptons. This selection recovers events where the lepton misses reconstruction, keeping the sample orthogonal to that in the previous measurement. The precision on these measurements range from 14 to 48\%.

In addition to the difficult task of separating single top quark s + t-channel events from the background, Tevatron experiments used properties of event topology to separate s- and t-channel events from each other. Events in the t-channel tend to have a distinctive forward jet, whose direction in p\( \bar{p} \) collisions is correlated with the lepton charge, and one b jet. On the other hand, s-channel events are more likely to have central jets, two of them from b quarks. Using these properties, CDF and D0 [12] have formed multivariate discriminants optimized for s- and t-channel separately. These discriminants were used on a Bayesian statistical analysis to combine the results from both collaborations. The three individual measurements above mentioned were combined, yielding the only existing observation of single top quark production in hadron colliders to this date: \( \sigma_s = 1.29^{+0.26}_{-0.24} \) at 6.3 standard deviations (with 5.1 expected).

With relatively less quark-initiated processes as compared to the Tevatron, LHC has not yet observed the production of s-channel single top quark. With a cut-and-count analysis based on 0.7 pb\(^{-1}\) of 7 TeV data, ATLAS [13] has set an upper limit of 26.5 fb at 95\% CL to the production. CMS Collaboration [14] has performed a more sophisticated BDT analysis based on 19.3 fb\(^{-1}\) of data at 8 TeV, measuring \( \sigma_s = 6.2 \pm 5.4 \pm 5.9 \) pb with 0.7 \( \sigma \) (with 0.9 expected).

### 3 t\( \bar{t} \) production cross section

The production of top quark pairs is an excellent ground to explore the frontiers of the SM. It provides important inputs to improve current limitations of the SM predictions and also plays an important role on
direct searches for physics beyond the SM. Signs of new physics could be evidenced as an enhancement of \( t\bar{t} \) production in a particular decay mode, or have SM \( t\bar{t} \) as background.

The total \( t\bar{t} \) inclusive cross section has been measured both at the Tevatron and at the LHC, at the centre-of-mass energies of 1.96, 7 and 8 TeV. At the Tevatron, CDF and D0 Collaborations [15] reported on a combination of 4 measurements from CDF Collaboration and 2 from D0 Collaboration, substantially improving the precision on the individual measurements. Systematic uncertainties related to the modeling of signal are the largest, and theoretical uncertainties on the background are also important. The most precise measurements at 8 TeV are those using fullleptonic \( t\bar{t} \) decays. In a cut-and-count analysis using ee,\( e\mu \) and \( \mu\mu \) and 2 b-tagged jet final states, and based on integrated luminosity of 5.3 fb\(^{-1}\), CMS [19] measures the cross section with a precision of about 4.5\%. ATLAS [18] uses a different cut-and-count analysis, where by separating events with 1 and 2 b-tagged jets, the efficiency of reconstructing and tagging a b jet can be extracted from the data simultaneously with the cross section, further reducing systematic uncertainties. ATLAS measurement is based on \( e\mu \) final states and 20 fb\(^{-1}\). A summary of the most recent measurements, compared to the latest theoretical predictions, is shown in Figure 2 and Table 1. All possible decay channels, even with the excellent agreement observed between data and theory, these high-precision measurements are able to reveal finer details of the models that still have some ground for improvement. Some examples are given in Fig. 3. MC@NLO predictions are in better agreement with D0 data than Alpgen; ATLAS has shown the dependency of the several PDF parametrizations as a function of the mass of the \( t\bar{t} \) system; the description of the top quark \( p_T \) spectra by the current MC models used by the LHC Collaborations was found to have small deviations.
Table 1: Summary of the most precise measurements of $t\bar{t}(+\gamma,V)$ production. $\sqrt{s}$ are displayed in TeV and cross sections in pb or fb. Experimental uncertainties are displayed as $\pm\text{stat.} \pm\text{syst.} \pm\text{lumi.}$ ATLAS $t\bar{t}$ cross section at 8 TeV quotes an additional uncertainty due to the uncertainty on the LHC beam energy.

from the data, as observed by both ATLAS and CMS independently of the decay mode and collision energy. Approximate NNLO and NLO+NNLL predictions, when available, are able to describe well the data.

Figure 3: Differential $t\bar{t}$ cross sections distributions from D0 [20] (left) ATLAS [21] (center) and CMS [22] (right).

4 Associated production: $t\bar{t}Z$, $t\bar{t}W$, $t\bar{t}\gamma$

Associated production of top pairs with electroweak gauge bosons is one of the most interesting subjects in top quark production, since it can immediately be used to test SM predictions and in the near future provide important information to improve electroweak measurements at the LHC (and even to the next generation of particle colliders). In particular, $t\bar{t}\gamma$ cross section can be reinterpreted in terms of the top quark electrical charge - therefore its measurement can be considered an alternative probe of that physical quantity. The $t\bar{t}Z$ production cross section is proportional to the $t\bar{t}Z$ electroweak gauge coupling, and can be used to impose stringent limits on anomalous top quark couplings, with severe implications on many new physics models. And all of them, together with $t\bar{t}W$, are part of the irreducible background on the measurement of $t\bar{t}H$ production, and share the same final state with many processes beyond the SM.

Among these three associated production processes, only $t\bar{t}\gamma$ is accessible at the Tevatron. The dominant background to this process is instrumental, when a QCD jet satisfies all selection criteria normally passed by a photon. It has been measured by the CDF Collaboration [23] using semileptonic $t\bar{t}$ decays on a selected sample of 30 events. The main background is estimated using jets passing most of photon selection criteria, but failing at least one of them. At the LHC, ATLAS [25] and CMS measurements [27] at 7 and 8 TeV, respectively, use template fits on variables with large discriminating power to separate real from fake photons. Results are presented in Table 1 in comparison to the predictions from the SM.
Measurements of $t\bar{t}V$ typically use leptonic decays of top quarks and $V$ bosons, leading to final states containing 2 to 4 leptons, with characteristic lepton charge combination for each process. Given that SM processes with multilepton final states are rare, backgrounds to these processes are also mostly instrumental, from QCD jets misidentified as leptons or misidentified lepton charge.

CMS measurement [28], presented at this Conference for the first time, is based on 19.5 fb$^{-1}$ of data at 8 TeV, and uses events with three final states containing electrons and/or muons: 2 same-sign leptons, aiming to tag semileptonic $t\bar{t}$+ W boson events, and 3 or 4 leptons to identify $t\bar{t}Z$ events. The instrumental background is estimated via data-driven techniques: jet-enriched samples used to control the rate of fakes and charge misassignment is studied on opposite sign ee or $e\mu$ events, predominantly from Drell-Yan or $t\bar{t}$ processes. Other backgrounds are estimated on MC. All final states are combined on a profile likelihood analysis. A simultaneous fit to all channels yields a combined cross section $\sigma_{t\bar{t}V} = 380^{+100}_{-90}(\text{stat.})^{+100}_{-90}(\text{syst})$ fb with 3.7 $\sigma$ significance. Then, fits from using same-sign (the remaining) channels $t\bar{t}W$ ($t\bar{t}Z$) cross section is extracted. Two-dimensional fits where $t\bar{t}W$ and $t\bar{t}Z$ are extracted simultaneously yield the same results. Results for the individual channels are given in Table 1.

5 Summary and conclusions

The era of observations and evidences of the top quark production was initiated at the Tevatron collider in 1995, when a few dozen of $t\bar{t}$ events were observed for the first time. Tevatron also holds the first evidence for single top quark production, and its $s$-channel observation. At the LHC, single top quark production was observed at the $t$-channel and $tW$ channels. Most recently, $t\bar{t}$ processes in association to bosons, at much lower cross sections, have also been observed.

By the time Tevatron ceased operations, CDF and D0 collaborations had each collected samples of almost 70,000 top pairs. Meanwhile, LHC ramped up activities as a “top factory”, collecting more than 5.5 million top pairs and 2.7 million single top events per experiment available for detailed studies. These millions of top quark events have brought us to the precision era today, where theory and experiment challenge each other to reach better and better accuracy.

Measurements of total inclusive cross sections of top quark pairs have reached an experimental precision of about 5% at the three probed s.c.m. energies (1.96, 7 and 8 TeV), while single top quark processes have been measured with precisions of 9-30%, depending on the production channel and collision energy. This accuracy is mainly limited by the uncertainties on the modelling of the top signal MC. For single top quarks, the statistical precision is also important, specially in the measurements using Tevatron data.

Differential cross sections have the potential to constrain some theoretical uncertainties in the signal modelling. For instance, with current precision, differential cross sections data, such as jet multiplicity, can rule out extreme variations of the QCD factorization and renormalization scales; ratios of single top quark-antiquark production in the $t$-channel is sensitive to the choice of the proton PDF, although with the current precision no set of PDFs is favoured. Reducing theoretical uncertainties, the dominant systematic uncertainties in most of the top quark measurements, is a prospect for improving future measurements of top quark production.

The associated production of top pairs with photons, Z’s and W bosons have also been measured. These production processes at extremely low cross sections have finally been observed, yielding significant tests of the SM: firstly, the cross sections proportional the magnitude of the electroweak couplings to the top quark are very sensitive to new physics; secondly they are an important milestone on the measurement of the SM $t\bar{t}H$ production. Since these processes are statistically limited, their precision at the planned LHC high luminosity runs at 13 TeV is expected to improve. The precision of the direct measurements of the top-Higgs Yukawa coupling achieved at the LHC in the next years may have a deep impact on the future generations of colliders, since they are, along with the discovery or further investigation of new physics, at the core of physics goals for planned experiments such as the International Linear Collider.

So far, no sign of new physics have been seen in top quark production measurements at 1.96, 7 and 8 TeV. These processes will continue to be tested on a new energy regime when the LHC resumes operation in 2015. If new physics exists at an scale reachable at the LHC, and couples to the top quark, it may become evident on deviations between the very high precision data and theory. In this case, the era of new physics...
discovery may be started soon.

References

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