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WIMPy Leptogenesis

PATRICK STENGEL

*Department of Physics and Astronomy,
University of Hawaii, Honolulu, HI 96822 USA*

We consider a class of leptogenesis models in which the lepton asymmetry arises from dark matter annihilation processes which violate CP and lepton number. Importantly, a necessary one-loop contribution to the annihilation matrix element arises from absorptive final state interactions. We elucidate the relationship between this one-loop contribution and the CP -violating phase. As we show, the branching fraction for dark matter annihilation to leptons may be small in these models, while still generating the necessary asymmetry.

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1 Introduction

Both the baryon asymmetry of the universe (BAU) and the existence of non-baryonic dark matter (DM) are well motivated by observation at a variety of scales and epochs. Recent observations of the cosmic microwave background (CMB) by the Wilkinson Microwave Anisotropy Probe (WMAP) [1] establish the baryonic and cold dark matter densities in our universe as

$$\begin{aligned}\Omega_b h^2 &\sim 0.022, \\ \Omega_{DM} h^2 &\sim 0.12.\end{aligned}\tag{1}$$

The most common theoretical paradigm used to explain the factor of ~ 5 between the two observed densities has been asymmetric dark matter (ADM) models [2]. While ADM models transfer an asymmetry of the dark sector into the baryonic sector, dynamical generation of the baryon asymmetry from dark matter decays have also been investigated [3]. More recently, weakly interacting massive particle (WIMP) annihilation has been suggested as a mechanism for baryogenesis [4]. In such WIMP models, a departure from thermal equilibrium is guaranteed by the WIMP framework. As such, these models of DM annihilation need only violations of baryon number, C and CP in order to satisfy the Sakharov conditions for baryogenesis [5].

1.1 Matter Asymmetry

In order to demonstrate the mechanism of CP violation when dynamically generating an asymmetry, consider a generic DM annihilation to a multi-particle final state, $XX \rightarrow Y$, and the CP conjugate process $XX \rightarrow \bar{Y}$. We can write the matrix element for such processes as the sum of CP -invariant and CP -violating terms,

$$\begin{aligned}\mathcal{M}_{XX \rightarrow Y} &= \mathcal{M}_{XX \rightarrow Y}^{CP} + \mathcal{M}_{XX \rightarrow Y}^{CPV}, \\ \mathcal{M}_{XX \rightarrow \bar{Y}} &= \pm (\mathcal{M}_{XX \rightarrow Y}^{CP} - \mathcal{M}_{XX \rightarrow Y}^{CPV}).\end{aligned}\tag{2}$$

As a result of the optical theorem, the relative phase between the CP -invariant and CP -violating terms will be $\pm\pi/2$ at tree level, eliminating any asymmetry. Thus, any nonvanishing asymmetry in the production cross section must be an interference of CP -invariant and CP -violating matrix elements with a relative phase difference from an additional imaginary matrix element contribution at loop level,

$$\sigma_{XX \rightarrow Y} - \sigma_{XX \rightarrow \bar{Y}} \propto \text{Re} [\mathcal{M}_{XX \rightarrow Y}^{CP} (\mathcal{M}_{XX \rightarrow \bar{Y}}^{CPV})^*].\tag{3}$$

In standard models of baryogenesis the relative phase is introduced by the interference of a CP -invariant tree-level diagram with a CP -violating loop-level diagram. Alternatively, we separate the CP -violating phase of two tree-level diagrams from an absorptive phase introduced by a final state decay.

1.2 WIMPy Baryogenesis Models

Before discussing a particular WIMPy model, some general features and issues are worth noting. A pair of WIMPs, stabilized by a discrete symmetry, will annihilate to a quark (or lepton) and a heavy exotic field through a set of operators with the C and CP properties necessary to satisfy the Sakharov conditions. This new state needs to also be protected from interactions with standard model by a discrete charge and should transfer its own asymmetry into a hidden sector. In models with a relative phase generated by the interference of tree-level and loop-level diagrams, dangerous washout diagrams arise. Washout processes yield terms in the Boltzmann equations for the rate of the asymmetry production that are proportional to the asymmetry itself. As a consequence, these CP -invariant terms will dampen the generation of an asymmetry.

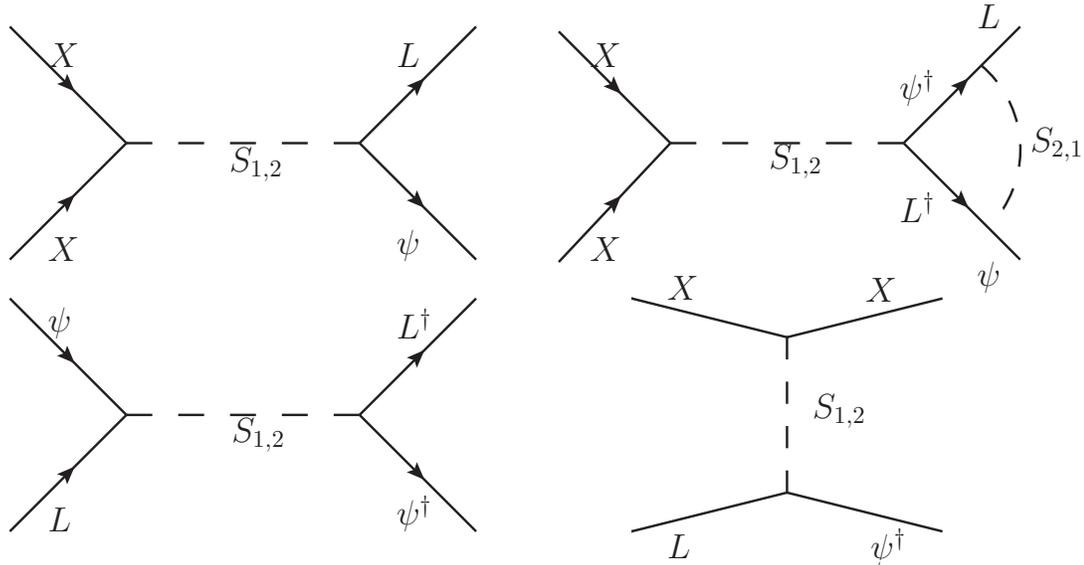


Figure 1: Asymmetry generating DM, X , annihilation diagrams and corresponding washout diagrams in WIMPy models [4]. Note that in order to generate a non-vanishing asymmetry, DM must annihilate to a SM lepton (or quark), L , and an additional weak scale field, ψ , through the interference of tree-level (top left) and loop-level (top right) diagrams. The loop diagram implies the existence of the leading order, “pure”, washout process (bottom left) and crossing the tree-level DM annihilation implies the “mixed” washout diagram (bottom right).

The presence of these dangerous washout processes are clearly implied by the loop-level DM annihilation diagram in Figure 1. Because there is no dark matter in the initial state for these “pure” washout processes, they lack the Boltzmann

suppression present in “mixed” washout processes. Thus, the CP -violating processes must account for a large portion of the total DM annihilation in order to generate a large enough asymmetry. As the mediator, $S_{1,2}$, runs in the loop, the cross section of the asymmetry generating process will also be inversely proportional to the square of the mediator mass. The result is a necessarily small separation of scales between the new external states, X and ψ , and the mediator $S_{1,2}$. This is generally problematic for models constructed within the framework of an effective field theory (EFT) and for WIMPy models which attempt to independently resolve the DM relic density, BAU and ultraviolet (UV) physics.

2 Field Content and Leptogenesis Model

Fields	$SU(2)_L$	$Q_{U(1)_Y}$	$Q_{U(1)_L}$	\mathbb{Z}_2
X	1	0	0	-
$P_L L = l_L$	\square	-1/2	+1	+
H	1	0	0	+
$P_L L = \nu_L$	\square	-1/2	+1	+
H	1	0	0	+

Table 1: Particle Content

We propose a generalized EFT containing generic WIMPy field content [4, 6]. The DM, X , and new heavy exotic, H , are protected by charges shown in Table 1. Unlike previous work, we don’t assume H is prevented from possibly large washout terms (through gauge interactions) by the same discrete symmetry as X . For simplicity, we assume the DM is protected by a Z_2 and annihilates to H and a SM lepton, L , through operators which violate $SU(2)_L$ and, thus, the EFT couplings are proportional to a Higgs vev.

$$\begin{aligned}
\mathcal{O}_1 &= \frac{\lambda_1}{2M_*^2} (i\bar{X}\gamma^5 X)(\bar{H}P_L L) + \frac{\lambda_1^*}{2M_*^2} (i\bar{X}\gamma^5 X)(\bar{L}P_R H) \\
\mathcal{O}_2 &= \frac{\lambda_2}{2M_*^2} (\bar{X}\gamma_\mu\gamma^5 X)(\bar{H}\gamma^\mu P_L L) + \frac{\lambda_2^*}{2M_*^2} (\bar{X}\gamma_\mu\gamma^5 X)(\bar{L}\gamma^\mu P_L H)
\end{aligned} \tag{4}$$

A model in which both X and H are protected by a Z_4 is phenomenologically very similar, only the DM must be Dirac due to the imaginary charges then assigned to X . Note the DM bilinears in these dimension six operators are the only combination which can yield CP -violating interference terms for Majorana fermion DM [7]. Since \mathcal{O}_1 is CP odd if λ_1 is real and \mathcal{O}_2 is CP odd if λ_2 is imaginary, a non-vanishing CP -violating phase will arise from a term proportional to $Re(\lambda_1\lambda_2^*)$.

2.1 Absorptive Final State Decays

In order to generate an asymmetry given our CP -violating phase, we also need an absorptive phase from a final state H decaying into light hidden sector fields, scalar ϕ and fermion H' . The loop, typically manifested by final state vertex corrections in standard leptogenesis scenarios, is alternatively provided by the external leg correction of the unstable H . The fully corrected propagator for a fermionic intermediate state is given by [8]:

$$S(p) = \frac{\not{p}_H + (m_H - i\Gamma_H/2)}{p_H^2 - m_H^2 - im_H\Gamma_H}. \quad (5)$$

If the width of the H resonance is small there will be regions of phase space where the intermediate H goes on-shell, generating an asymmetry between the cross sections for $XX \rightarrow \phi^*\bar{H}'L$ and $XX \rightarrow \phi\bar{L}H'$. As a consequence of the optical theorem, we can see the imaginary matrix element contribution we need to preserve our CP asymmetry can only come from a term proportional to $\Gamma_H/2$ in the propagator. If we assume both λ_1 and λ_2 are real, \mathcal{O}_1 will maximally violate CP and only couple to the right-handed Weyl spinor H_R , while \mathcal{O}_2 will completely preserve CP and only couple to the left-handed Weyl spinor H_L . We assume H can only decay from the left-handed helicity (as in the standard model) into our hidden sector through the CP -invariant operator

$$\mathcal{O}_H = |g|(\phi^*\bar{H}'P_LH + \phi\bar{H}P_RH'). \quad (6)$$

Thus the relative phase between the CP -violating $XX \rightarrow \bar{H}_R L_L \rightarrow \phi\bar{H}'_R L_L$ and CP -conserving $XX \rightarrow \bar{H}_L L_L \rightarrow \phi\bar{H}'_R L_L$ must arise from the $-i\Gamma_H/2$ contribution to the helicity-flip term of the propagator. Assuming the narrow width approximation, we can then calculate the asymmetry in the full cross sections of the $2 \rightarrow 3$ DM annihilation processes,

$$(\sigma^{XX \rightarrow \phi^*\bar{H}'L} - \sigma^{XX \rightarrow \phi\bar{L}H'})_v = \Gamma_H \frac{Re(\lambda_1\lambda_2^*)m_X}{4\pi M_*^4} \left[1 - \frac{m_H^2}{s}\right]^2. \quad (7)$$

Note \sqrt{s} is the energy in center-of-mass frame and for simplicity we assume $m_{H'}, m_\phi, m_L \ll m_X, m_H \sim \text{TeV}$. In order to keep our UV mediator effectively decoupled from low-energy physics, we set $M_* = 10 \text{ TeV}$. If we define a measure of the asymmetry,

$$\epsilon \equiv \frac{\sigma^{XX \rightarrow \phi^*\bar{H}'L} - \sigma^{XX \rightarrow \phi\bar{L}H'}}{\sigma^{XX \rightarrow \phi^*\bar{H}'L} + \sigma^{XX \rightarrow \phi\bar{L}H'}}, \quad (8)$$

then we find $\epsilon \sim \Gamma_H/m_{H,X}$. This is in contrast with other WIMPy baryogenesis models, where one typically finds $\epsilon \sim m_X^2/M_*^2$ [6] due to the CP -violating loop diagrams. The consequences of this distinction in our model are explained in the next section.

2.2 Boltzmann Equations

We can write the Boltzmann equations in terms of dimensionless variables $x = m_X/T$ and $Y = n/s$, where n is the number density and s is the entropy density. Assuming an adiabatic process, the entropy S should be constant, and Y is essentially a comoving number density. As the hidden sector fields, ϕ and H' , are effectively massless, they should remain approximately in equilibrium throughout the relevant cosmological epoch. L is also light but we want to track even a small departure from equilibrium in order to generate the small observed baryon (or, through sphalerons, lepton) asymmetry $Y_{\Delta L} \equiv Y_L - Y_{\bar{L}} \sim 10^{-10}$. Thus, we can assume $Y_L + Y_{\bar{L}} \simeq 2Y_{L_{eq}}$. The coupled Boltzmann equations for the primary DM annihilation to any fermion/antifermion pair, as well as the subdominant DM annihilation that gives us the injection of our lepton asymmetry are [9]:

$$\begin{aligned}
\frac{x^2 H(m_X)}{s(m_X)} \frac{dY_X}{dx} &= -\langle \sigma_{Av} \rangle (Y_X^2 - Y_{X_{eq}}^2), & (9) \\
\frac{x^2 H(m_X)}{s(m_X)} \frac{dY_{\Delta L}^{inj}}{dx} &= \frac{1}{2} [\langle \sigma_{XX \rightarrow \phi^* \bar{H}' L} v \rangle] (Y_X^2 - Y_{X_{eq}}^2 Y_L / Y_{L_{eq}}) \\
&\quad - \frac{1}{2} [\langle \sigma_{XX \rightarrow \phi \bar{L} H'} v \rangle] (Y_X^2 - Y_{X_{eq}}^2 Y_{\bar{L}} / Y_{\bar{L}_{eq}}) \\
&\quad - \langle \sigma_{XL \rightarrow \phi X H'} v \rangle Y_X (Y_L - Y_{L_{eq}}) + \langle \sigma_{X\bar{L} \rightarrow \phi^* X \bar{H}'} v \rangle Y_X (Y_{\bar{L}} - Y_{\bar{L}_{eq}}) \\
&\quad + \dots & (10)
\end{aligned}$$

The “+...” terms involved suppressed processes in which an on-shell resonance is kinematically forbidden. In order to achieve the correct relic DM density, given a weak scale m_X , we always set $\langle \sigma_{Av} \rangle = 1$ pb. Note that $H(T)$ is the Hubble parameter at temperature T given a flat, radiation-dominated early universe. While the equilibrium rates for the $3 \rightarrow 2$ processes are given by detailed balance with the rates for the $2 \rightarrow 3$ processes, the actual $3 \rightarrow 2$ rates are calculated by rescaling the equilibrium rates by the ratio of the actual mass densities to the equilibrium mass densities. We can rewrite the the equation for $dY_{\Delta L}^{inj}/dx$ as

$$\frac{x^2 H(m_X)}{s(m_X)} \frac{dY_{\Delta L}^{inj}}{dx} \sim \langle \sigma_{XX}^{CPV} v \rangle (Y_X^2 - Y_{X_{eq}}^2) - \langle \sigma_{XX}^{CP} v \rangle Y_{X_{eq}}^2 Y_{\Delta L} / Y_{L_{eq}} - \langle \sigma_{XL}^{CP} v \rangle Y_X Y_{\Delta L},$$

where

$$\begin{aligned}
\langle \sigma_{XX}^{CP} v \rangle &\equiv \frac{1}{2} \left[\langle \sigma_{XX \rightarrow \phi^* \bar{H}' L} v \rangle + \langle \sigma_{XX \rightarrow \phi \bar{L} H'} v \rangle \right], \\
\langle \sigma_{XX}^{CPV} v \rangle &\equiv \frac{1}{2} \left[\langle \sigma_{XX \rightarrow \phi^* \bar{H}' L} v \rangle - \langle \sigma_{XX \rightarrow \phi \bar{L} H'} v \rangle \right], \\
\langle \sigma_{XL}^{CP} v \rangle &\equiv \frac{1}{2} \left[\langle \sigma_{XL \rightarrow \phi X H'} v \rangle + \langle \sigma_{X\bar{L} \rightarrow \phi^* X \bar{H}'} v \rangle \right].
\end{aligned} \tag{11}$$

Before the electroweak phase transition (EWPT) any asymmetry generated within leptons will be shared with baryons through electroweak sphalerons. Recent discovery of the Higgs mass allows for lattice calculation of the sphaleron rate through the EWPT until sphalerons are effectively decoupled [10]. The expression for the lepton injection rate is the source term in the coupled differential equations for the evolution of the baryon and lepton number densities.

$$\begin{aligned}
xH(T) \frac{dY_{\Delta B}}{dx} &= -\gamma(T) [Y_{\Delta B} + 3\eta(T)Y_{\Delta L}] \\
xH(T) \frac{dY_{\Delta L}}{dx} &= -\frac{1}{3}\gamma(T) [Y_{\Delta B} + \eta(T)Y_{\Delta L}] + xH(T) \frac{dY_{\Delta L}^{inj}}{dx}
\end{aligned} \tag{12}$$

The functions $\eta(T)$ and $\gamma(T)$ are defined in terms of the temperature T , the temperature-dependent Higgs field expectation value v_{min} , and the Chern-Simons diffusion rate $\Gamma_{diff}(T)$, plotted in Figure 2.

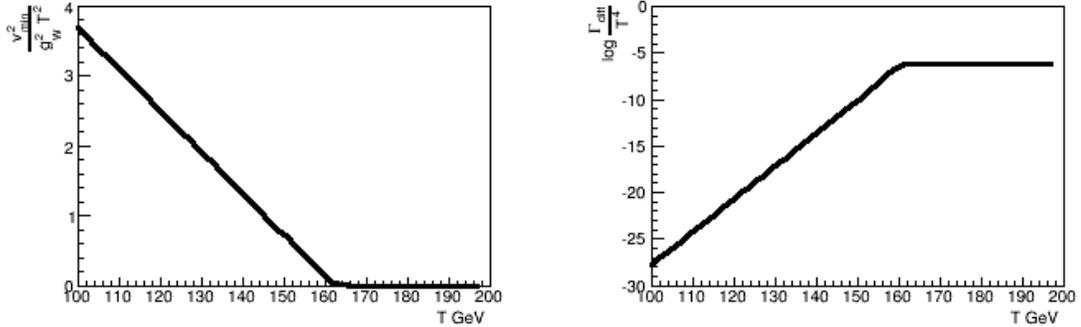


Figure 2: Lattice calculation of $v_{min}^2 / (T^2 g_{weak}^2)$ (left) and $\log(\Gamma_{diff} / T^4)$ (right) through the transition region [10]. For $T < 140$ GeV, $\log(\Gamma_{diff} / T^4)$ is extrapolated from analytical calculations deep in the broken phase [11], with a constant rescaling to provide consistency with the lattice calculation for $T = 140 - 155$ GeV.

3 Results

We assume $m_H \leq 2m_X$, so that H can go on-shell. We also assume the DM is in equilibrium up until $x = 1$, after which we numerically solve the coupled Boltzmann/sphaleron equations. As we only consider cases with $\langle\sigma_{Av}\rangle \gg \langle\sigma_{\bar{X}X \rightarrow \bar{H}L}v\rangle$, the annihilation process $\bar{X}X \rightarrow \bar{H}L, \bar{L}H$ will not significantly affect the WIMP miracle. Thus, the relevant parameters of our model are m_X , m_H/m_X , Γ_H/m_H and $Re(\lambda_1\lambda_2^*)$.

As scans over our parameter space are phenomenologically similar to [6], we can distinguish our results with several benchmark points. For all benchmark models, the parameters are chosen to both generate the correct baryon asymmetry and the observed DM density. In Figure 3, we plot the thermally-averaged cross sections for the processes $XX \rightarrow \phi^*\bar{H}'L$ (both CP -invariant and CP -violating terms) and $XL \rightarrow \phi H'X$ (the CP -invariant term) as a function of $x = m_X/T$. We also plot the contribution of these terms to the lepton source injection rate, as well as Y_B , Y_X and $Y_{X_{eq}}$. We have chosen the ‘‘high-mass’’ benchmark parameters $m_X = 5$ TeV, $m_H = 7$ TeV, $\lambda_1 = \lambda_2 = 0.5$, $\langle\sigma_{Av}\rangle = 1$ pb, $\Gamma_H/m_H = 0.1$.

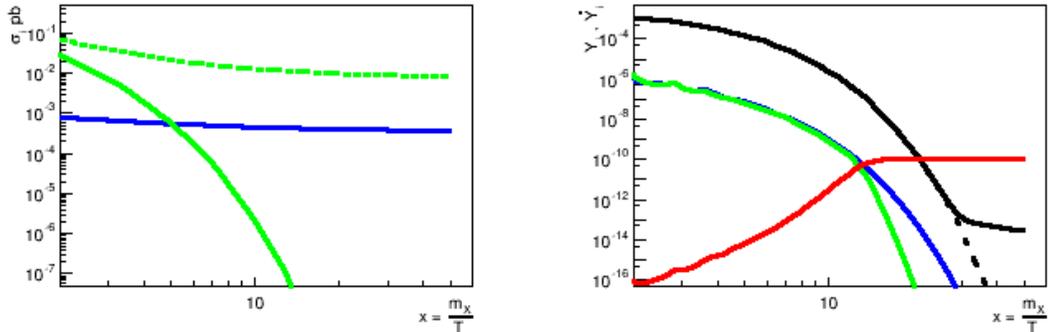


Figure 3: The left panel shows the thermally-averaged cross sections $\langle\sigma_{XX}^{CP}v\rangle$ (green dashed), $\langle\sigma_{XX}^{CPV}v\rangle$ (blue) and $\langle\sigma_{XL}v\rangle$ (green) (as defined in eq. 11). The right panel shows the corresponding contributions to the lepton source rate (sum of washout terms in green and source term in blue) as well as Y_B (red), Y_X (black) and $Y_{X_{eq}}$ (black dashed). We have chosen parameters $m_X = 5$ TeV, $m_H = 7$ TeV, $\lambda_1 = \lambda_2 = 0.5$, $\langle\sigma_{Av}\rangle = 1$ pb, $\Gamma_H/m_H = 0.1$.

Note that the CP -violating part of the $XX \rightarrow \phi^*\bar{H}'L$ process starts to drive the asymmetry as soon as the DM becomes non-relativistic, when $x > 1$. Although the departure from equilibrium is small at $x \sim 1$ relative to when DM freezes out (typically $x \sim 20$), the DM density is much higher and even a small deviation from equilibrium will generate an asymmetry. The asymmetry won’t become large until washout processes freeze out, usually around $x \sim 10$ in WIMP models [12].

Other benchmarks include a narrower-width model with $m_X = 5$ TeV, $m_H = 7$ TeV, $\lambda_1 = \lambda_2 = 0.5$, $\langle\sigma_{Av}\rangle = 1$ pb, $\Gamma_H/m_H = 0.05$ and a low-mass model with $m_X = 1.5$ TeV, $m_H = 2.2$ TeV, $\lambda_1 = \lambda_2 = 1$, $\langle\sigma_{Av}\rangle = 1$ pb, $\Gamma_H/m_H = 0.1$. These cases require stronger couplings in order to compensate for weaker asymmetry production. The narrower-width benchmark is nearly identical to the high mass benchmark, only washout processes freeze out slightly later. For the low-mass benchmark, sphalerons begin to decouple around when washout processes freeze out, thus forcing a sharper freeze out of baryon number. The parameters of these benchmark models are summarized in Table 2.

benchmark	m_X	m_H	Γ_H/m_H	$\lambda_1 = \lambda_2$	ϵ	$\langle\sigma_{XX\rightarrow\phi^*\bar{H}'L^v}\rangle/\langle\sigma_{Av}\rangle$
low-mass	1.5 TeV	2.2 TeV	0.10	1.0	0.045	0.002
high-mass	5.0 TeV	7.0 TeV	0.10	0.5	0.045	0.008
narrower-width	5.0 TeV	7.0 TeV	0.05	1.0	0.022	0.033

Table 2: Benchmarks

The lack of “pure” washout terms in our model allows us to generate the observed BAU with smaller ϵ independently of the mediator scale, M_* . We can then have $m_X = 1.5$ TeV while still separating the UV physics from our new weak scale interactions. We also can have our CP -violating DM annihilation channel very small relative to any other channel one might consider in a more complete WIMP framework.

4 Conclusion

We have shown that the CP -violating process required by the Sakharov conditions for the dynamical generation of the BAU can arise from tree-level diagrams and an absorptive final state interaction. This mechanism sequesters the one-loop suppression of the asymmetry and eliminates dangerous tree-level washout processes which are not Boltzmann suppressed. As a result, we can effectively decouple our dark sector from the UV and address the DM relic density and BAU more independently. We also have carefully treated nonperturbative sphaleron effects through the electroweak crossover.

For simplicity, we have focused only on the approximation where the intermediate H is a narrow resonance and can be produced on-shell. Although one would expect an even more efficient asymmetry generation if we disregarded this assumption, additional washout processes incapable of producing a narrow resonance would have to be considered. Also, baryogenesis through DM annihilation directly into quarks would avoid the sphaleron shutoff, but the mass of H would be highly constrained by collider searches for DM production by quarks. In principal, the application of

final state absorptive interactions could also be used to eliminate dangerous washout processes from more traditional baryogenesis or leptogenesis scenarios.

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